

Caleb Logemann,
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Nonconservative Discontinuous Galerkin Method for Generalized Shallow Water Equations

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Overview

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Generalized Shallow Water

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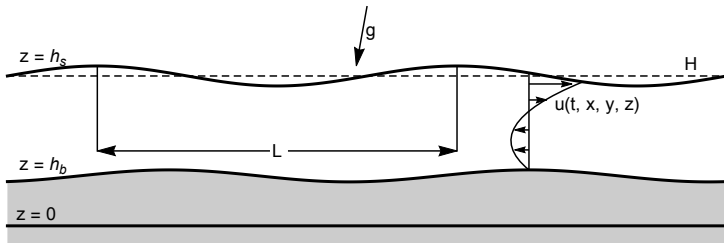
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Navier Stokes Equations with a free surface

$$\nabla \cdot \mathbf{u} = 0$$

$$\mathbf{u}_t + \nabla \cdot (\mathbf{u}\mathbf{u}) = -\frac{1}{\rho}\nabla p + \frac{1}{\rho}\nabla \cdot \boldsymbol{\sigma} + \mathbf{g}$$

$$(h_s)_t + [u(t, x, y, h_s), v(t, x, y, h_s)]^T \cdot \nabla h_s = w(t, x, y, h_s)$$

$$(h_b)_t + [u(t, x, y, h_b), v(t, x, y, h_b)]^T \cdot \nabla h_b = w(t, x, y, h_b)$$

Polynomial Ansatz

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$$\begin{aligned}\tilde{u}(t, x, y, \zeta) &= u_m(t, x, y) + u_d(t, x, y, \zeta) \\ &= u_m(t, x, y) + \sum_{j=1}^N (\alpha_j(t, x, y) \phi_j(\zeta))\end{aligned}$$

$$\begin{aligned}\tilde{v}(t, x, y, \zeta) &= v_m(t, x, y) + v_d(t, x, y, \zeta) \\ &= v_m(t, x, y) + \sum_{j=1}^N (\beta_j(t, x, y) \phi_j(\zeta))\end{aligned}$$

Orthogonality Condition

$$\int_0^1 \phi_j(\zeta) \phi_i(\zeta) d\zeta = 0 \quad \text{for } j \neq i$$

$$\phi_0(\zeta) = 1, \quad \phi_1(\zeta) = 1 - 2\zeta, \quad \phi_2(\zeta) = 1 - 6\zeta + 6\zeta^2$$

Constant Moments

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$$\begin{aligned}
 & h_t + (hu_m)_x + (hv_m)_y = 0 \\
 & (hu_m)_t + \left(h \left(u_m^2 + \sum_{j=1}^N \frac{\alpha_j^2}{2j+1} \right) + \frac{1}{2} g e_z h^2 \right)_x \\
 & + \left(h \left(u_m v_m + \sum_{j=1}^N \frac{\alpha_j \beta_j}{2j+1} \right) \right)_y = -\frac{\nu}{\lambda} \left(u_m + \sum_{j=1}^N \alpha_j \right) + hg(e_x - e_z(h_b)_x) \\
 & (hv_m)_t + \left(h \left(v_m^2 + \sum_{j=1}^N \frac{\alpha_j \beta_j}{2j+1} \right) + \frac{1}{2} g e_z h^2 \right)_y \\
 & + \left(h \left(u_m v_m + \sum_{j=1}^N \frac{\alpha_j \beta_j}{2j+1} \right) \right)_x = -\frac{\nu}{\lambda} \left(v_m + \sum_{j=1}^N \beta_j \right) + hg(e_y - e_z(h_b)_y)
 \end{aligned}$$

Higher Order Moments

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$$\begin{aligned}
 (h\alpha_i)_t + \left(2hu_m\alpha_i + h \sum_{j,k=1}^N A_{ijk}\alpha_j\alpha_k \right)_x + \left(hu_m\beta_i + hv_m\alpha_i + h \sum_{j,k=1}^N A_{ijk}\alpha_j\beta_k \right)_y \\
 = u_mD_i - \sum_{j,k=1}^N B_{ijk}D_j\alpha_k - (2i+1)\frac{\nu}{\lambda} \left(u_m + \sum_{j=1}^N \left(1 + \frac{\lambda}{h}C_{ij} \right) \alpha_j \right) \\
 (h\beta_i)_t + \left(hu_m\beta_i + hv_m\alpha_i + h \sum_{j,k=1}^N A_{ijk}\alpha_j\beta_k \right)_x + \left(2hv_m\beta_i + h \sum_{j,k=1}^N A_{ijk}\beta_j\beta_k \right)_y \\
 = v_mD_i - \sum_{j,k=1}^N B_{ijk}D_j\beta_k - (2i+1)\frac{\nu}{\lambda} \left(v_m + \sum_{j=1}^N \left(1 + \frac{\lambda}{h}C_{ij} \right) \beta_j \right)
 \end{aligned}$$

Example Systems

1D model with h_b constant, $e_x = e_y = 0$, and $e_z = 1$
Constant System

$$\begin{bmatrix} h \\ hu_m \end{bmatrix}_t + \begin{bmatrix} hu_m \\ hu_m^2 + \frac{1}{2}gh^2 \end{bmatrix}_x = -\frac{\nu}{\lambda} \begin{bmatrix} 0 \\ u_m \end{bmatrix}$$

Flux Jacobian Eigenvalues, $u_m \pm \sqrt{gh}$
Linear System, $\tilde{u} = u_m + \alpha_1 \phi_1$

$$\begin{bmatrix} h \\ hu_m \\ h\alpha_1 \end{bmatrix}_t + \begin{bmatrix} hu_m \\ hu_m^2 + \frac{1}{2}gh^2 + \frac{1}{3}h\alpha_1^2 \\ 2hu_m\alpha_1 \end{bmatrix}_x = Q \begin{bmatrix} h \\ hu_m \\ h\alpha_1 \end{bmatrix}_x - \mathbf{s}$$

$$Q = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & u_m \end{bmatrix} \quad \mathbf{s} = \frac{\nu}{\lambda} \begin{bmatrix} 0 \\ u_m + \alpha_1 \\ 3(u_m + \alpha_1 + 4\frac{\lambda}{h}\alpha_1) \end{bmatrix}$$

Flux Jacobian Eigenvalues, $u_m \pm \sqrt{gh + \alpha_1^2}$, u_m

Example Systems

1 dimensional with h_b constant, $e_x = e_y = 0$, and $e_z = 1$ Quadratic Vertical Profile, $\tilde{u} = u + \alpha_1 \phi_1 + \alpha_2 \phi_2$

$$\begin{bmatrix} h \\ hu \\ h\alpha_1 \\ h\alpha_2 \end{bmatrix}_t + \begin{bmatrix} hu \\ hu^2 + \frac{1}{2}gh^2 + \frac{1}{3}h\alpha_1^2 + \frac{1}{5}h\alpha_2^2 \\ 2hu\alpha_1 + \frac{4}{5}h\alpha_1\alpha_2 \\ 2hu\alpha_2 + \frac{2}{3}h\alpha_1^2 + \frac{2}{7}h\alpha_2^2 \end{bmatrix}_x = Q \begin{bmatrix} h \\ hu \\ h\alpha_1 \\ h\alpha_2 \end{bmatrix}_x - \mathbf{s}$$

$$Q = \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & u - \frac{\alpha_2}{5} & \frac{\alpha_1}{5} \\ 0 & 0 & \alpha_1 & u + \frac{\alpha_2}{7} \end{bmatrix}, P = \frac{\nu}{\lambda} \begin{bmatrix} 0 \\ u + \alpha_1 + \alpha_2 \\ 3(u + \alpha_1 + \alpha_2 + 4\frac{\lambda}{h}\alpha_1) \\ 5(u + \alpha_1 + \alpha_2 + 12\frac{\lambda}{h}\alpha_2) \end{bmatrix}$$

Flux Jacobian Eigenvalues, $u \pm c\sqrt{gh}$

$$c^4 - \frac{10\alpha_2}{7}c^3 - \left(1 + \frac{6\alpha_2^2}{35} + \frac{6\alpha_1^2}{5}\right)c^2 + \left(\frac{22\alpha_2^3}{35} - \frac{6\alpha_2\alpha_1^2}{35} + \frac{10\alpha_2}{7}\right)c - \frac{\alpha_2^4}{35} - \frac{6\alpha_2^2\alpha_1^2}{35} - \frac{3\alpha_2^2}{7} + \frac{\alpha_1^4}{5} + \frac{\alpha_1^2}{5} = 0$$

Nonconservative Products

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Model Equation

$$\mathbf{q}_t + \nabla \cdot \mathbf{f}(\mathbf{q}) + Q_i(\mathbf{q})\mathbf{q}_{x_i} = \mathbf{s}(\mathbf{q}) \quad \text{for } (\mathbf{x}, t) \in \Omega \times [0, T]$$

Traditionally searching for weak solutions, find \mathbf{q} such that

$$\int_0^T \int_{\Omega} (\mathbf{q}_t + \nabla \cdot \mathbf{f}(\mathbf{q}) + Q_i(\mathbf{q})\mathbf{q}_{x_i}) v \, d\mathbf{x} \, dt = \int_0^T \int_{\Omega} \mathbf{s}(\mathbf{q}) v \, d\mathbf{x} \, dt$$

for all $v \in C_0^1(\Omega \times [0, T])$

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Manufactured Solution

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Effect of Higher Moments

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Model Equation

$$\mathbf{q}_t + \mathbf{f}(\mathbf{q})_x = Q(\mathbf{q})\mathbf{q}_x - \mathbf{P}(\mathbf{q}) \quad \text{for } (x, t) \in [a, b] \times [0, T]$$

Weak Form, find \mathbf{q} such that

$$\int_a^b \mathbf{q}_t v \, dx + \int_a^b \mathbf{f}(\mathbf{q})_x v \, dx = \int_a^b Q(\mathbf{q})\mathbf{q}_x v \, dx - \int_a^b \mathbf{P}(\mathbf{q}) v \, dx$$

for all $v \in L^2([a, b] \times [0, T])$

Notation

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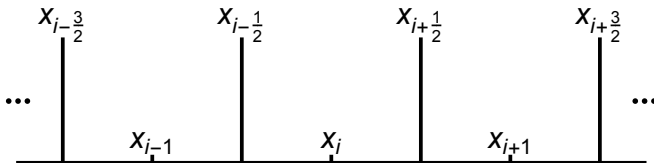
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- Partition the domain, $[a, b]$ as

$$a = x_{1/2} < \cdots < x_{j-1/2} < x_{j+1/2} < \cdots < x_{N+1/2} = b$$

- $I_j = [x_{j-1/2}, x_{j+1/2}]$
- $x_j = \frac{x_{j+1/2} + x_{j-1/2}}{2}$
- $\Delta x_j = x_{j+1/2} - x_{j-1/2}$
- $\Delta x_j = \Delta x$ for all j .



Discontinuous Galerkin Space

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Finite Dimensional DG Space

$$V^k = \left\{ v \in L^2([a, b]) \mid v|_{I_j} \in P^k(I_j) \right\}$$

Basis for V^k

$$\{\phi_j^\ell\} \text{ where } \phi_j^\ell(x)|_{I_j} = \phi^\ell(\xi_j(x)) \text{ and } \phi_j^\ell(x)|_{\bar{I}_j} = 0$$

for $j = 1, \dots, N$ and $\ell = 1, \dots, k$.

Legendre Polynomials

$$\phi^k \in P^k([-1, 1]) \text{ with } \frac{1}{2} \int_{-1}^1 \phi^k(\xi) \phi^\ell(\xi) d\xi = \delta_{k\ell}$$

and

$$\xi_j(x) = \frac{2}{\Delta x_j}(x - x_j)$$

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Find $\mathbf{q}_h \in V^k$ such that

$$\begin{aligned} \int_{I_j} (\mathbf{q}_h)_t \phi_j^\ell(x) \, dx &= \int_{I_j} \mathbf{f}(\mathbf{q}_h)_x \phi_j^\ell \, dx \\ &\quad - F_{j+1/2} \phi_j^\ell(x_{j+1/2}) + F_{j-1/2} \phi_j^\ell(x_{j-1/2}) \\ &\quad + \int_{I_j} Q(\mathbf{q}_h)(\mathbf{q}_h)_x \phi_j^\ell \, dx - \int_{I_j} \mathbf{P}(\mathbf{q}_h) \phi_j^\ell \, dx \end{aligned}$$

for all ϕ_j^ℓ .

Local Lax-Friedrichs Flux

$$\mathbf{q}_h^+ = \lim_{x \rightarrow x_{j+1/2}^+} (\mathbf{q}_h(x))$$

$$\mathbf{q}_h^- = \lim_{x \rightarrow x_{j+1/2}^-} (\mathbf{q}_h(x))$$

$$a = \max_{\mathbf{q} \in [\mathbf{q}_h^-, \mathbf{q}_h^+]} \{\rho(\mathbf{f}'(\mathbf{q}) - Q(\mathbf{q}))\}$$

$$F_{j+1/2} = \frac{1}{2}(\mathbf{f}(\mathbf{q}_h^+) + \mathbf{f}(\mathbf{q}_h^-)) - \frac{1}{2}a(\mathbf{q}_h^+ - \mathbf{q}_h^-)$$

Nonconservative Flux

Need to evaluate

$$\int^{I_j} Q \mathbf{q}_x \phi_j^\ell dx$$

$$\mathbf{q}|_{I_j} = \sum_{\ell=1}^k (Q_j^\ell \phi_j^\ell(x)), \quad \mathbf{q}_x|_{I_j} = \sum_{\ell=1}^k (Q_x^\ell \phi_j^\ell(x))$$

where

$$\begin{bmatrix} Q_x^1 \\ Q_x^2 \\ Q_x^3 \\ Q_x^4 \\ Q_x^5 \end{bmatrix} = \frac{1}{2\Delta x} \begin{bmatrix} \Delta Q^1 - 2\sqrt{5}\Delta Q^3 + 78\Delta Q^5 \\ \Delta Q^2 - \frac{10}{3}\sqrt{3}\sqrt{7}\Delta Q^4 \\ \Delta Q^3 - 14\sqrt{5}\Delta Q^5 \\ \Delta Q^4 \\ \Delta Q^5 \end{bmatrix}$$

$$\Delta Q^\ell = Q_{i+1}^\ell - Q_{i-1}^\ell$$

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Inviscid Example

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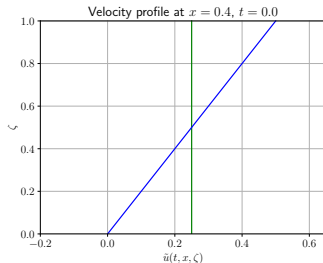
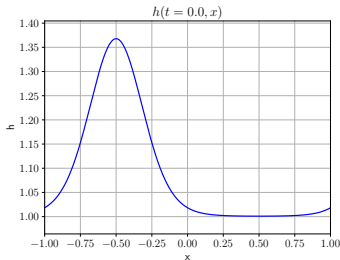
$$x \in [-1, 1] \quad t \in [0, 2.0]$$

$$h(t = 0, x) = 1 + e^{3 \cos(\pi(x+0.5)) - 4}$$

$$\tilde{u}(t = 0, x, \zeta) = \begin{cases} 0.25 & \text{constant} \\ 0.5\zeta & \text{linear} \end{cases}$$

$$u_m = 0.25$$

$$s = -0.25$$



Inviscid Example

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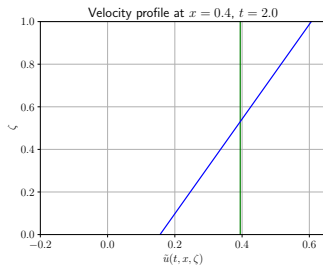
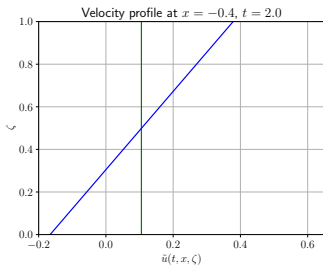
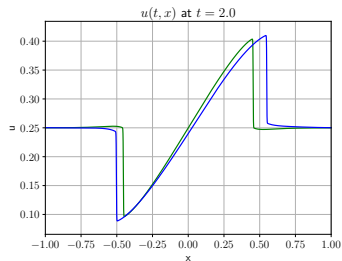
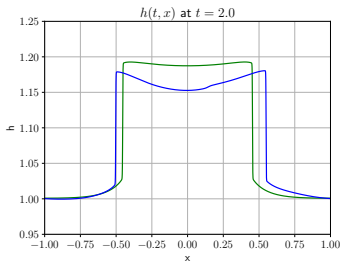
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Higher Moment Equations

1 dimensional with h_b constant, $e_x = e_y = 0$, and $e_z = 1$ Quadratic Vertical Profile, $\tilde{u} = u_m + s\phi_1 + \kappa\phi_2$

$$\begin{bmatrix} h \\ hu \\ hs \\ h\kappa \end{bmatrix}_t + \begin{bmatrix} hu \\ hu^2 + \frac{1}{2}gh^2 + \frac{1}{3}hs^2 + \frac{1}{5}h\kappa^2 \\ 2hus + \frac{4}{5}hs\kappa \\ 2huk + \frac{2}{3}hs^2 + \frac{2}{7}h\kappa^2 \end{bmatrix}_x = Q \begin{bmatrix} h \\ hu \\ hs \\ h\kappa \end{bmatrix}_x - P$$

$$Q = \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & u_m - \frac{\kappa}{5} & \frac{s}{5} \\ 0 & 0 & s & u_m + \frac{\kappa}{7} \end{bmatrix} \quad P = \frac{\nu}{\lambda} \begin{bmatrix} 0 \\ u_m + s + \kappa \\ 3(u_m + s + \kappa + 4\frac{\lambda}{h}s) \\ 5(u_m + s + \kappa + 12\frac{\lambda}{h}\kappa) \end{bmatrix}$$

Flux Jacobian Eigenvalues, $u_m \pm c\sqrt{gh}$

$$c^4 - \frac{10\kappa}{7}c^3 - \left(1 + \frac{6\kappa^2}{35} + \frac{6s^2}{5}\right)c^2 + \left(\frac{22\kappa^3}{35} - \frac{6\kappa s^2}{35} + \frac{10\kappa}{7}\right)c - \frac{\kappa^4}{35} - \frac{6\kappa^2 s^2}{35} - \frac{3\kappa^2}{7} + \frac{s^4}{5} + \frac{s^2}{5} = 0$$

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- Higher Order Numerical Methods
- Slope Limiters
- Two Dimensional Meshes
- Icosahedral Spherical Mesh
- Positivity Preserving Limiters

Icosahedral Mesh

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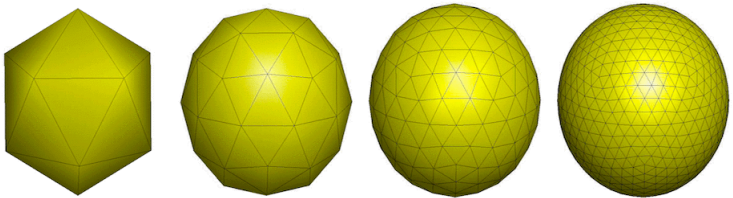
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Subdivide each edge
Project vertices onto sphere

Spherical Test Cases

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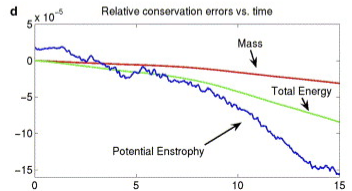
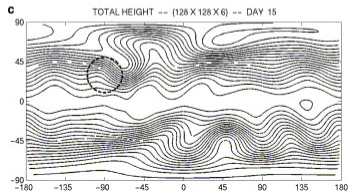
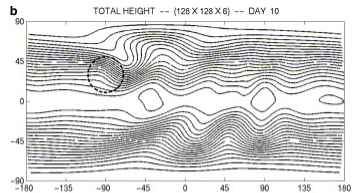
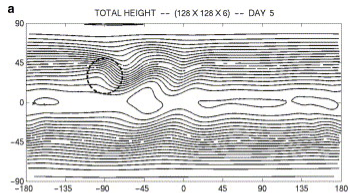
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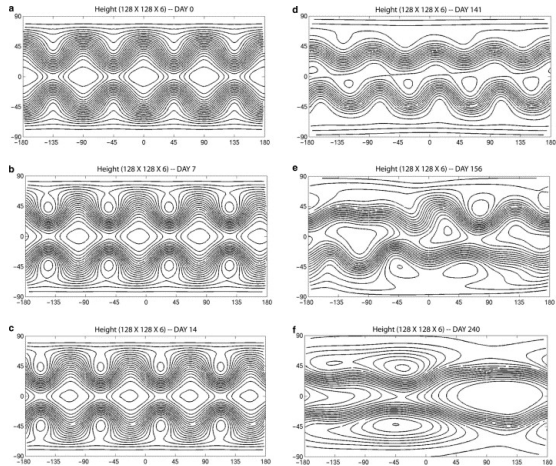
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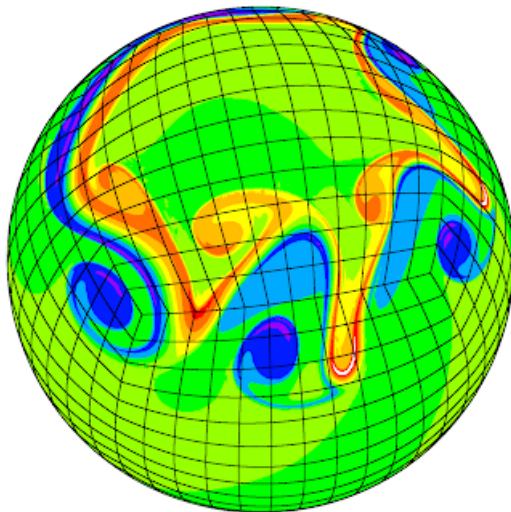
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