

Caleb Logemann,
James
Rossmanith

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Discontinuous Galerkin Method for Solving Thin Film Equations

Caleb Logemann James Rossmanith

Mathematics Department,
Iowa State University

logemann@iastate.edu

May 9, 2019

Overview

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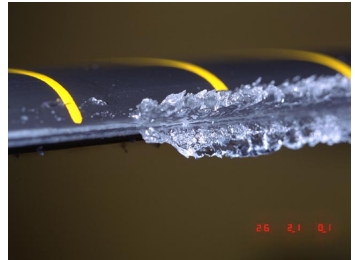
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- Aircraft Icing
- Runback



- Industrial Coating

Model Equations

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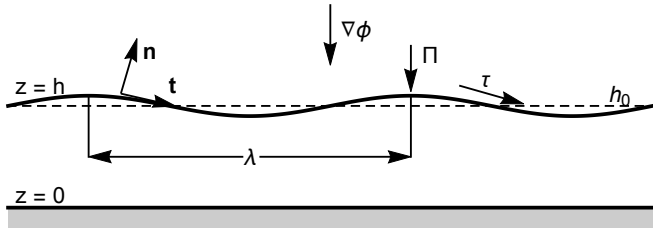
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■ Incompressible Navier-Stokes Equation

$$u_x + w_z = 0$$

$$\rho(u_t + uu_x + wu_z) = -p_x + \mu\Delta u - \phi_x$$

$$\rho(w_t + uw_x + ww_z) = -p_z + \mu\Delta w - \phi_z$$

$$w = 0, u = 0 \quad \text{at } z = 0$$

$$w = h_t + uh_x \quad \text{at } z = h$$

$$\mathbf{T} \cdot \mathbf{n} = (-\kappa\sigma + \Pi)\mathbf{n} + \left(\frac{\partial\sigma}{\partial s} + \tau\right)\mathbf{t} \quad \text{at } z = h$$

Nondimensionalization

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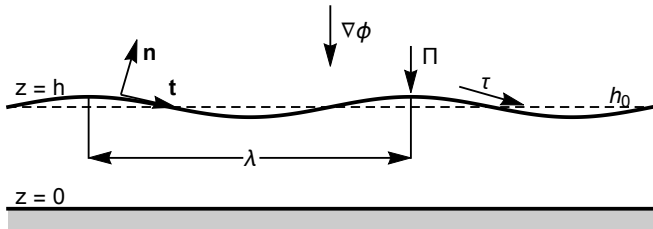
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$$\varepsilon = \frac{h_0}{\lambda} \ll 1$$

$$Z = \frac{z}{h_0}$$

$$X = \frac{\varepsilon x}{h_0}$$

$$U = \frac{u}{U_0}$$

$$W = \frac{w}{\varepsilon U_0}$$

$$T = \frac{\varepsilon U_0 t}{h_0}$$

Nondimensionalization

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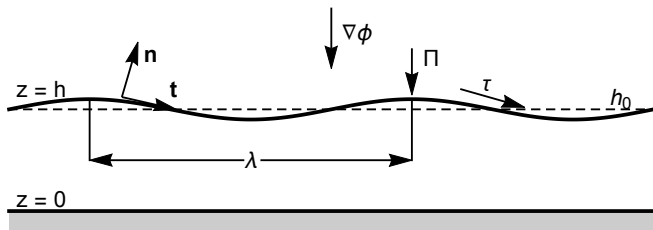
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$$U_X + W_Z = 0$$

$$\varepsilon \text{Re}(U_T + UU_X + WW_Z) = -P_X + U_{ZZ} + \varepsilon^2 U_{XX} - \Phi_X$$

$$\varepsilon^3 \text{Re}(W_T + WW_X + WW_Z) = -P_Z + \varepsilon^2 (W_{ZZ} + \varepsilon^2 W_{XX}) - \Phi_Z$$

$$W = 0, U = 0 \quad \text{at } Z = 0$$

$$W = H_T + UH_X \quad \text{at } Z = H$$

$$U_Z + \varepsilon^2 W_X - 4\varepsilon^2 H_X U_X = \tau + \Sigma_X \quad \text{at } Z = H$$

$$-P - \Pi + \varepsilon^2 U_X (\varepsilon^2 H_X^2 - 1) = \varepsilon^2 H_X (U_Z + \varepsilon^2 W_X) + C^{-1} \varepsilon^3 H_{XX} \quad \text{at } Z = H$$

$$U_X + W_Z = 0$$

$$U_{ZZ} = P_X + \Phi_X$$

$$0 = -P_Z - \Phi_Z$$

$$W = 0 \quad \text{at } Z = 0$$

$$U = 0$$

$$W = H_T + UH_X \quad \text{at } Z = H$$

$$U_Z = \tau_0 + \Sigma_X$$

$$-\Pi_0 - P = \overline{C}^{-1} H_{XX}$$

Operator Splitting

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■ Simplified Model

$$q_t + (q^2 - q^3)_x = -(q^3 q_{xxx})_x \quad (0, T) \times \Omega$$

■ Operator Splitting

$$q_t + (q^2 - q^3)_x = 0$$

$$q_t + (q^3 u_{xxx})_x = 0$$

■ Strang Splitting

$\frac{1}{2}\Delta t$ step of Convection

$$q_t + (q^2 - q^3)_x = 0$$

Δt step of Diffusion

$$q_t + (q^3 u_{xxx})_x = 0$$

$\frac{1}{2}\Delta t$ step of Convection

$$q_t + (q^2 - q^3)_x = 0$$

Convection

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■ Convection Equation

$$\begin{aligned}q_t + f(q)_x &= 0 & (0, T) \times \Omega \\f(q) &= q^2 - q^3\end{aligned}$$

■ Weak Form

Find q such that

$$\int_{\Omega} (q_t v - f(q) v_x) dx + \hat{f} v \Big|_{\partial\Omega} = 0$$

for all test functions v

Notation

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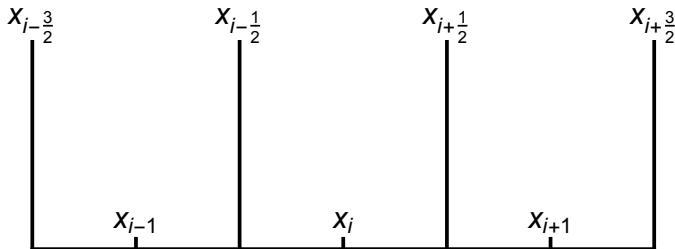
References

- Partition the domain, $[a, b]$ as

$$a = x_{1/2} < \cdots < x_{j-1/2} < x_{j+1/2} < \cdots < x_{N+1/2} = b$$

- $I_j = [x_{j-1/2}, x_{j+1/2}]$

- $x_j = \frac{x_{j+1/2} + x_{j-1/2}}{2}$.



Runge Kutta Discontinuous Galerkin

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- Find $Q(t, x)$ such that for each time $t \in (0, T)$,
 $Q(t, \cdot) \in V_h = \left\{ v \in L^1(\Omega) : v|_{I_j} \in P^k(I_j) \right\}$

$$\int_{I_j} Q_t v \, dx = \int_{I_j} f(Q) v_x \, dx \\ - \left(\mathcal{F}_{j+1/2} v^-(x_{j+1/2}) - \mathcal{F}_{j-1/2} v^+(x_{j-1/2}) \right)$$

for all $v \in V_h$

- Rusanov/Local Lax-Friedrichs Numerical Flux

$$\mathcal{F}_{j+1/2} = \frac{1}{2} \left(f(Q_{j+1/2}^-) + f(Q_{j+1/2}^+) \right) + \frac{1}{2} \max_q \{ |f'(q)| \} (Q_{j+1/2}^- - Q_{j+1/2}^+)$$

- Solve this system of ODEs with any Explicit Strong Stability Preserving (SSP) Runge-Kutta Method.

Explicit SSP Runge Kutta Methods

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■ Forward Euler

$$q^{n+1} = q^n + \Delta t L(q^n)$$

■ Second Order

$$\begin{aligned} q^* &= q^n + \Delta t L(q^n) \\ q^{n+1} &= \frac{1}{2}(q^n + q^*) + \frac{1}{2}\Delta t L(q^*) \end{aligned}$$

Diffusion

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■ Diffusion Equation

$$q_t = -(q^3 q_{xxx})_x \quad (0, T) \times \Omega$$

■ Linearize operator at $t = t^n$, let $f(x) = q^3(t = t^n, x)$

$$q_t = -(f(x) q_{xxx})_x \quad (0, T) \times \Omega$$

Local Discontinuous Galerkin

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References

Find $Q(t, x), R(x), S(x), U(x)$ such that for all $t \in (0, T)$

$$Q(t, \cdot), R, S, U \in V_h = \left\{ v \in L^1(\Omega) : v|_{I_j} \in P^k(I_j) \right\}$$

$$\int_{I_j} Rv \, dx = - \int_{I_j} Qv_x \, dx + \left(\hat{Q}_{j+1/2} v_{j+1/2}^- - \hat{Q}_{j-1/2} v_{j-1/2}^+ \right)$$

$$\int_{I_j} Sw \, dx = - \int_{I_j} R w_x \, dx + \left(\hat{R}_{j+1/2} w_{j+1/2}^- - \hat{R}_{j-1/2} w_{j-1/2}^+ \right)$$

$$\begin{aligned} \int_{I_j} Uy \, dx = & \int_{I_j} S_x f y \, dx - \left(S_{j+1/2}^- f_{j+1/2}^- y_{j+1/2}^- - S_{j-1/2}^+ f_{j-1/2}^+ y_{j-1/2}^+ \right) \\ & + \left(\hat{S}_{j+1/2} \hat{f}_{j+1/2} y_{j+1/2}^- - \hat{S}_{j-1/2} \hat{f}_{j-1/2} y_{j-1/2}^+ \right) \end{aligned}$$

$$\int_{I_j} Q_t z \, dx = - \int_{I_j} U z_x \, dx + \left(\hat{U}_{j+1/2} z_{j+1/2}^- - \hat{U}_{j-1/2} z_{j-1/2}^+ \right)$$

for all $I_j \in \Omega$ and all $v, w, y, z \in V_h$.

Numerical Fluxes

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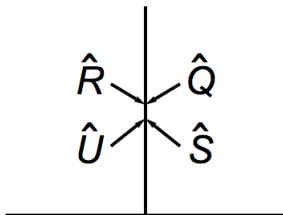
$$\hat{f}_{j+1/2} = \frac{1}{2} \left(f_{j+1/2}^+ + f_{j+1/2}^- \right)$$

$$\hat{Q}_{j+1/2} = Q_{j+1/2}^+$$

$$\hat{R}_{j+1/2} = R_{j+1/2}^-$$

$$\hat{S}_{j+1/2} = S_{j+1/2}^+$$

$$\hat{U}_{j+1/2} = U_{j+1/2}^-$$



LDG Complications

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- Explicit time step scales with h^4
- Implicit System is difficult to solve efficiently
 - GMRES iterations scale with size of system
 - Preconditioned GMRES

$$P = A_0^{-1}$$

$$PAx = Pb$$

- Geometric Multigrid fails to converge

Finite Difference Approach

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- Let cell centers, x_i , form finite difference grid.
- Finite difference space, \mathbb{R}^N .
- $Q_{DG} \in V_h \rightarrow Q_{FD} \in \mathbb{R}^N$

$$(Q_{FD})_i = \frac{1}{h} \int_{K_i} Q_{DG} \, dx$$

- $Q_{FD} \in \mathbb{R}^N \rightarrow Q_{DG} \in V_h$

$$Q_{DG}|_K \in P^1(K)$$

$$\frac{1}{h} \int_{K_i} Q_{DG} \, dx = (Q_{FD})_i$$

$$\partial_x Q_{DG}|_{K_i} = \frac{(Q_{FD})_{i+1} - (Q_{FD})_{i-1}}{2h}$$

Finite Difference Approximation

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■ First derivative approximation

$$(-(f(x)q_{xxx})_x)_i \approx -\frac{q_{i+1/2}^3(q_{xxx})_{i+1/2} - q_{i-1/2}^3(q_{xxx})_{i-1/2}}{h}$$

■ Third derivative approximation

$$(q_{xxx})_{i+1/2} \approx \frac{-Q_{i-1} + 3Q_i - 3Q_{i+1} + Q_{i+2}}{h^3}$$

■ Value of Q^3 at boundary

$$q_{i+1/2}^3 = \left(\frac{Q_i + Q_{i+1}}{2} \right)^3$$

Implicit L-Stable Runge Kutta

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■ Backward Euler

$$q^{n+1} = q^n + \Delta t L(q^{n+1})$$

■ 2nd Order

$$q^* = q^n + \frac{1}{4} \Delta t (L(q^n) + L(q^*))$$
$$3q^{n+1} = 4q^* - q^n + \Delta t L(q^{n+1})$$

Nonlinear Solvers

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■ Picard Iteration

$$L(q) = A(f \approx q^3)q$$

$$q_0^{n+1} = q^n$$

$$q_{m+1}^{n+1} = q^n + \Delta t A(q_m^{n+1}) q_{m+1}^{n+1}$$

$$q_{m+1}^* = q^n + \frac{1}{4} \Delta t (L(q^n) + A(q_m^*) q_{m+1}^*)$$

$$3q_{m+1}^{n+1} = 4q^* - q^n + \Delta t A(q_m^{n+1}) q_{m+1}^{n+1}$$

■ Newton's Method

$$q_{m+1}^{n+1} = q_m^{n+1} - J(q_m^{n+1})^{-1} F(q_m^{n+1})$$

$$F(q) = q - q^n - \Delta t L(q)$$

$$J(q) = I - \Delta t L'(q)$$

Manufactured Solution

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$$q_t = -(q^3 q_{xxx})_x + s(x, t)$$
$$q(x, t) = 0.1 * \sin(2\pi(x - t)) + 0.15$$

Backward Euler				
1 Iteration			2 Iterations	
N	error	order	error	order
100	0.0131	—	0.0053	—
200	0.0064	1.0264	0.0026	1.0466
400	0.0033	0.96	0.0013	0.9704
800	0.0016	1.0069	0.0007	1.0134

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$$q_t = -(q^3 q_{xxx})_x + s(x, t)$$
$$q(x, t) = 0.1 * \sin(2\pi(x - t)) + 0.15$$

2nd Order IRK						
1 Iteration			2 Iterations		3 Iterations	
N	error	order	error	order	error	order
50	0.0075	—	0.00047	—	0.0004901	—
100	0.0041	0.8601	0.00012	1.9844	0.0001209	2.0194
200	0.0020	1.0391	0.0000312	1.9451	0.0000305	1.9887
400	0.0010	0.9652	0.0000082	1.9244	0.0000078	1.9641

Manufactured Solution

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$$q_t = -(q^3 q_{xxx})_x + s(x, t)$$
$$q(x, t) = \frac{2}{10} e^{-10t} e^{-300(x-\frac{1}{2})^2} + \frac{1}{10}$$

Backward Euler				
1 Iteration			2 Iterations	
N	error	order	error	order
100	0.0097	—	0.0933	—
200	0.0050	0.95	0.0421	1.1494
400	0.0027	0.87	3.756	-6.48
800	33.21	-13.5	16.51	-2.14

Manufactured Solution with Newton's Method

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$$q_t = -(q^3 q_{xxx})_x + s(x, t)$$
$$q(x, t) = \frac{2}{10} e^{-10t} e^{-300(x-\frac{1}{2})^2} + \frac{1}{10}$$

Backward Euler

N	error	order
50	0.0280	—
100	0.0153	0.8765
200	0.0080	0.9249
400	5.5e75	-258

Hyperbolic Wave Structure

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■ Conservation Law

$$q_t + f(q)_x = 0$$

■ Riemann Problem Initial Data

$$q(x, 0) = \begin{cases} q_l & x < d \\ q_r & x > d \end{cases}$$

■ Rankine-Hugoniot Condition

$$s = \frac{f(q_l) - f(q_r)}{q_l - q_r}$$

Convex Flux Function

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- Shock Wave

$$f'(q_l) > s > f'(q_r)$$

- Rarefaction

$$f'(q_l) < s < f'(q_r)$$

Nonconvex Flux Function

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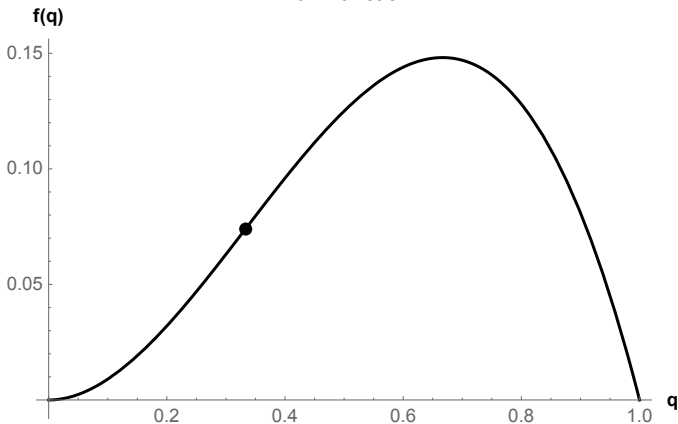
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$$q_t + (q^2 - q^3)_x = 0$$

Flux Function



Nonconvex Flux Function

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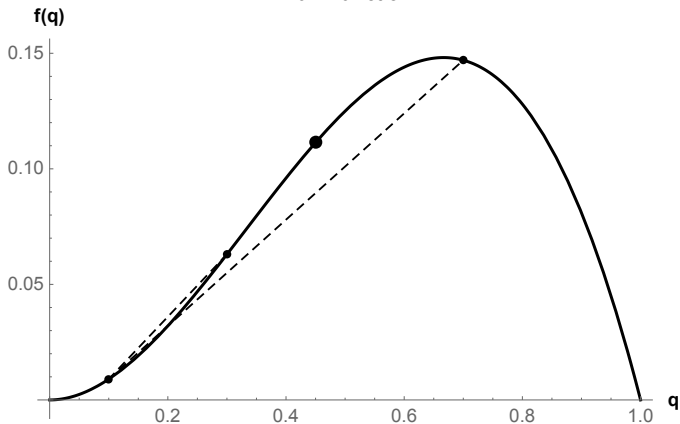
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$$f'(q_b) = s$$

$$q_b = (1 - q_r)/2$$

Flux Function



Compressive Shock

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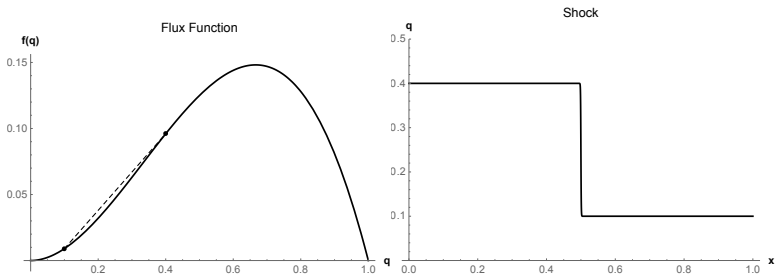
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$$q_l < q_b$$

$$q(x, t) = \begin{cases} q_l & x \leq st \\ q_r & x > st \end{cases}$$



Rarefaction-Compressive Shock

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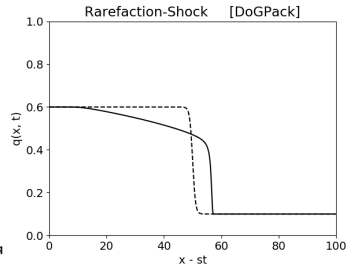
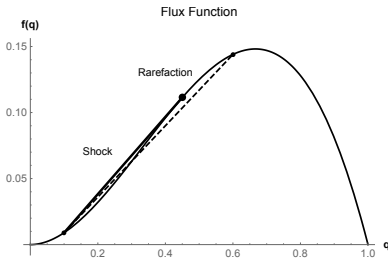
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$$q_l > q_b$$

$$q(x, t) = \begin{cases} q_l & x < f'(q_l)t \\ h_r(x) & f'(q_l)t < x < f'(q_b)t \\ q_r & x > f'(q_b)t \end{cases}$$



Wave Structure with Nonlinear Hyper Diffusion

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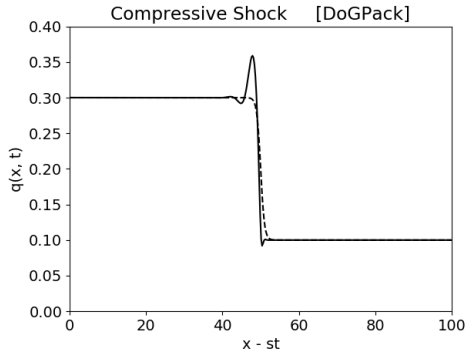
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$$q_t + (q^2 - q^3)_x = -(q^3 q_{xxx})_x$$

$$q_r = 0.1 \quad q_l = 0.3$$



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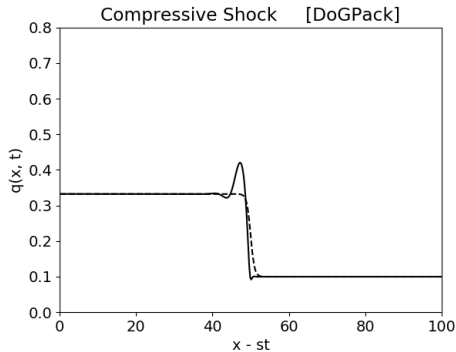
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$$q_r = 0.1 \quad q_l = 0.3323$$

$$q(x, 0) = (-\tanh(x - 50) + 1) \frac{q_l - q_r}{2} + q_r$$



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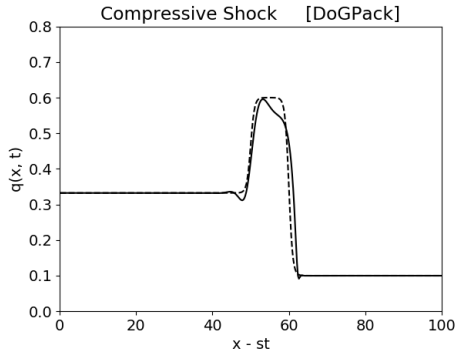
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$$q_r = 0.1 \quad q_l = 0.3323 \quad q_m = 0.6$$

$$q(x, 0) = \begin{cases} \frac{q_m - q_l}{2} \tanh(x - 50) + \frac{q_m + q_l}{2} & x < 55 \\ -\frac{q_m - q_r}{2} \tanh(x - 60) + \frac{q_m + q_r}{2} + q_r & x > 55 \end{cases}$$



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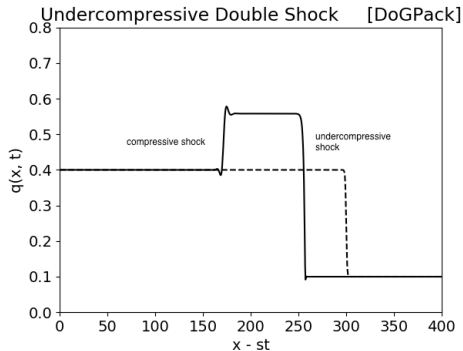
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$$q_r = 0.1 \quad q_l = 0.4$$

$$q(x, 0) = (-\tanh(x - 50) + 1) \frac{q_l - q_r}{2} + q_r$$



Wave Structure with Nonlinear Hyper Diffusion

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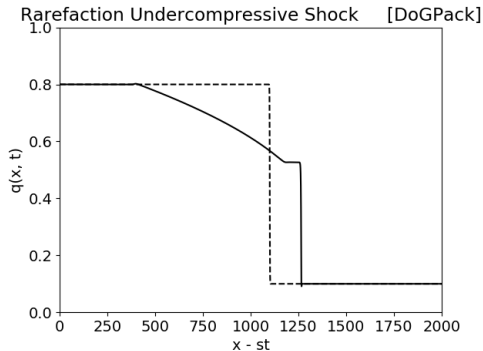
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$$q_r = 0.1 \quad q_l = 0.8$$

$$q(x, 0) = (-\tanh(x - 1100) + 1) \frac{q_l - q_r}{2} + q_r$$



Conclusion

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Observations

- Nonlinear Hyper Diffusion has subtle instabilities

Future Work

- Higher Order Convergence
 - Runge Kutta IMEX
 - Local Discontinuous Galerkin Method
 - Hybridized Discontinuous Galerkin Method

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