# Hybridized Discontinuous Galerkin Method

These notes are intended to give background on Hyridized Discontinuous Galerkin methods and then explore how to apply this method to Thin Film Equations

# 1 Introduction/Main Idea

To start we will consider Poisson's equation with Dirichlet boundary conditions

$$-\Delta u = f \qquad \text{in } \Omega$$
$$u = g \qquad \text{on } \partial \Omega$$

We will consider this in mixed form by introducing the auxilliary variable  $q = -\nabla u$ , then the equation becomes

$$\mathbf{q} + \nabla u = 0$$
 in  $\Omega$   
 $\nabla \cdot \mathbf{q} = f$  in  $\Omega$   
 $u = g$  on  $\partial \Omega$ 

This problem has an exact solution that can be found analytically assuming some nice properties.

I will introduce a triangulation of  $\Omega$ ,  $\mathcal{T}_h$ , and reformulate the problem on this triangulation that will give the same exact solution.

### 1.1 Notation

First some notation,

$$\partial \mathcal{T}_h = \{ \partial K : K \in \mathcal{T}_h \}$$

$$F = \partial K \cap \partial \Omega \text{ for } K \in \mathcal{T}_h$$

$$F = \partial K^+ \cap \partial K^- \text{ for } K^+, K^- \in \mathcal{T}_h$$

Let  $\varepsilon_h$  be the set of all faces, F, and  $\varepsilon_h^0$  be interior faces, and  $\varepsilon_h^{\partial}$  be boundary faces.

Let  $n^+$  and  $n^-$  be the outward unit normals of  $\partial K^+$  and  $\partial K^-$  respectively, and  $(q^{\pm}, u^{\pm})$  be the interior values of (q, u) on F for  $K^{\pm}$ . Define

$$[\mathbf{q} \cdot n] = \mathbf{q}^+ \cdot \mathbf{n}^+ + \mathbf{q}^- \cdot \mathbf{n}^-$$

$$[\mathbf{u}\mathbf{n}] = u^+ \mathbf{n}^+ + u^- \mathbf{n}^-$$

$$\{\mathbf{q}\} = \frac{\mathbf{q}^+ + \mathbf{q}^-}{2}$$

$$\{u\} = \frac{u^+ + u^-}{2}$$

### 1.2 Reformulation

Now we can reformulate the original Poisson's problem on  $\mathcal{T}_h$  as a local problem for each K

$$\mathbf{q} + \nabla u = 0$$
$$\nabla \cdot \mathbf{q} = f$$

a transmission condition on each interior face,  $F \in \mathcal{E}_h^0$ 

$$[\![u\boldsymbol{n}]\!] = \mathbf{0}$$
$$[\![\boldsymbol{q} \cdot \boldsymbol{n}]\!] = 0$$

and the boundary condition on each boundary face,  $F \in \varepsilon_h^{\partial}$ 

$$u = g$$

This problem is equivalent to the original problem on  $\Omega$ . The (q, u) that satisfies this problem also solve the original problem.

We would like to be able to solve the local problem locally, but this requires boundary conditions on each element K for the local problem to be solved. Therefore consider the local problem

$$\mathbf{q} + \nabla u = 0$$
 in  $K$   
 $\nabla \cdot \mathbf{q} = f$  in  $K$   
 $u = \hat{u}$  on  $\partial K$ 

We have introduced another unknown  $\hat{u}$  on each interior face  $F \in \varepsilon_h^0$ . This unknown automatically makes us satisfy ||un|| = 0, so the transmission condition becomes

$$[\boldsymbol{q} \cdot \boldsymbol{n}] = 0$$
 on  $F \in \varepsilon_h^0$ 

and we still have the boundary condition

$$u = g$$
 on  $F \in \varepsilon_h^{\partial}$ 

Now solving for  $(q, u, \hat{u})$  will give the same solution as the original problem, however q and u can be solved locally and only  $\hat{u}$  needs to be solved globally.

### 1.3 General Algorithm

Here is the outline for solving for q, u, and  $\hat{u}$ . First split the local problem in two, so that one part depends on f and the other part depends on  $\hat{u}$ , that is let  $q = Q_f + Q_{\hat{u}}$  and  $u = U_f + U_{\hat{u}}$ , where

$$Q_f + \nabla U_f = 0$$
 in  $K$   
 $\nabla \cdot Q_f = f$  in  $K$   
 $U_f = 0$  on  $\partial K$ 

and

$$\mathbf{Q}_{\hat{u}} + \nabla U_{\hat{u}} = 0 \quad \text{in } K$$

$$\nabla \cdot \mathbf{Q}_{\hat{u}} = 0 \quad \text{in } K$$

$$U_{\hat{u}} = \hat{u} \quad \text{on } \partial K$$

Now the transmission condition becomes

$$\llbracket oldsymbol{Q}_{\hat{u}} 
rbracket = - \llbracket oldsymbol{Q}_f 
rbracket$$

First solve for  $Q_f$  exactly, then solve for  $Q_{\hat{u}}$  in terms of  $\hat{u}$ . Now the transmission condition gives a global linear algebra problem

$$\llbracket oldsymbol{Q}_{\hat{u}} 
Vert = - \llbracket oldsymbol{Q}_f 
Vert$$

since  $[Q_{\hat{u}}]$  is a linear system in  $\hat{u}$  and  $-[Q_f]$  is known.

After this linear algebra problem is solved, the values of  $U_{\hat{u}}$  and  $Q_{\hat{u}}$  can be found/reconstructed locally. The full solution is then  $q = Q_f + Q_{\hat{u}}$  and  $u = U_f + U_{\hat{u}}$ .

#### 1.4 1D example

In 1D, the problem becomes

$$Q_f + U_f' = 0 \quad \text{in } K$$
 
$$Q_f' = f \quad \text{in } K$$
 
$$U_f = 0 \quad \text{on } \partial K$$

and

$$\begin{aligned} Q_{\hat{u}} + U_{\hat{u}}' &= 0 & \text{in } K \\ Q_{\hat{u}}' &= 0 & \text{in } K \\ U_{\hat{u}} &= \hat{u} & \text{on } \partial K \end{aligned}$$

and

$$Q_{\hat{u}}(x_{j+1/2}^{-}) - Q_{\hat{u}}(x_{j+1/2}^{+}) = -Q_f(x_{j+1/2}^{-}) + Q_f(x_{j+1/2}^{+})$$

where the uniform mesh is given by  $x_i$  at cell centers,  $x_{i+1/2}$  at cell interfaces, and spacing h. Solving the  $\hat{u}$  system we see that

$$Q'_{\hat{u}} = 0$$
$$Q_{\hat{u}} = c$$

and

$$Q_{\hat{u}} + U'_{\hat{u}} = 0$$

$$U'_{\hat{u}} = -c$$

$$U_{\hat{u}} = -cx + b$$

with the boundary conditions, we know  $U_{\hat{u}}$  is a line from  $\hat{u}_{j-1/2}$  to  $\hat{u}_{j+1/2}$ , and  $Q_{\hat{u}}$  is the opposite of the slope of this line.

$$U_{\hat{u}} = \frac{\hat{u}_{j+1/2} - \hat{u}_{j-1/2}}{h} (x - x_{j-1/2}) - \hat{u}_{j-1/2}$$

$$Q_{\hat{u}} = -\frac{\hat{u}_{j+1/2} - \hat{u}_{j-1/2}}{h}$$

Now we can form the linear system given by the transmission condition

$$Q_{\hat{u}}(x_{j+1/2}^{-}) - Q_{\hat{u}}(x_{j+1/2}^{+}) = -Q_f(x_{j+1/2}^{-}) + Q_f(x_{j+1/2}^{+})$$
$$-\frac{\hat{u}_{j+1/2} - \hat{u}_{j-1/2}}{h} + \frac{\hat{u}_{j+3/2} - \hat{u}_{j+1/2}}{h} = -Q_f(x_{j+1/2}^{-}) + Q_f(x_{j+1/2}^{+})$$
$$\frac{\hat{u}_{j-1/2} - 2\hat{u}_{j+1/2} + \hat{u}_{j+3/2}}{h} = -Q_f(x_{j+1/2}^{-}) + Q_f(x_{j+1/2}^{+})$$

After solving this linear system we already have expressions for  $Q_{\hat{u}}$  and  $U_{\hat{u}}$  in terms of  $\hat{u}$ .

# 2 Hybridizable Discontinuous Galerkin Method

### 2.1