Basics / Notation

Signature (1, -1, -1, -1)

$$\gamma = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}}$$

$$d' = d/\gamma$$

$$E^2 = \vec{p}^2 c^2 + m^2 c^4$$

$$x^{\mu} = (t, x, y, z)$$

$$x_{\mu} = (t, -x, -y, -z)$$

$$p = (E, -p_1, -p_2, -p_3)$$

$$p^2 = p^{\mu} p_{\mu} = E^2 - \vec{p}^2$$

Don't fuck up boosts:

$$x = \gamma(x^* - vt)$$

$$t = \gamma(t^* - vx^*)$$

$$x^* = \gamma(x + vt)$$

$$t^* = \gamma(t + vx)$$

$$p_x = \gamma(p_x^* - vE^*)$$

$$E = \gamma(E^* - vp_x^*)$$

$$p_x^* = \gamma(p_x + vE)$$

$$E^* = \gamma(E + vp_x)$$

Rel.
$$K: K = E - mc^2$$

Mandelstam: $1+2 \rightarrow 3+4$

$$s = (p_1 + p_2)^2$$
$$t = (p_1 - p_3)^2$$
$$u = (p_1 - p_4)^2$$

Conservation Laws

Conservation of energy, momentum, charge, spin, baryon number, lepton number, colour charge, charge conjugation $C\psi = i\gamma^2\psi^{\dagger}$

Lepton flavour conserved at γ , Z but not W vertices.

Symm. \rightarrow conserved:

- transform under symm, Taylor expand, write as $\psi' = (1 + i\epsilon \hat{A})\psi$, $\Longrightarrow \hat{A}$ conserved
- Translation \leftrightarrow momentum
- Flavour symm. in strong interaction

Spin Algebra Stuff

$$\sigma_x = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}, \sigma_y = \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix}, \sigma_z = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \quad \begin{array}{l} \text{Dirac: } (i\gamma^\mu\partial_\mu - m)\psi = 0 \\ \gamma^\mu\gamma^\nu + \gamma^\nu\gamma^\mu = 2g^{\mu\nu} \end{array}$$

$$S_{-} = S_x - iS_y, S_{+} = S_x + iS_y$$

Isospin, $\hat{T} = \frac{1}{2}\sigma$, same algebra, $T_+d = u$ $T_3u = \frac{1}{2}u, T_3d = -\frac{1}{2}d$

$$T_3\bar{u} = -\frac{1}{2}\bar{u}, T\bar{d} = \frac{1}{2}\bar{d}$$

$$T_3s = 0$$

$$T^2 = \frac{4}{3}I$$

$$\lambda_{1} = \begin{bmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}, \lambda_{2} = \begin{bmatrix} 0 & -i & 0 \\ i & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

$$\lambda_{3} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 0 \end{bmatrix}, \lambda_{4} = \begin{bmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ 1 & 0 & 0 \end{bmatrix}$$

$$\lambda_{5} = \begin{bmatrix} 0 & 0 & -i \\ 0 & 0 & 0 \\ i & 0 & 0 \end{bmatrix}, \lambda_{6} = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{bmatrix}$$

$$\lambda_{7} = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & -i \\ 0 & i & 0 \end{bmatrix}, \lambda_{8} = \frac{1}{\sqrt{3}} \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -2 \end{bmatrix}$$

Fermi's Golden Rule

$$\Gamma_{fi} = 2\pi |T_{fi}|^2 \rho(E_i)$$

Lifetime: $\tau = \frac{\hbar}{\Gamma}$

Lor. Inv. T: M

$$M_{fi}^{ab} = 4(E_a E_b E_c E_d)^{\frac{1}{2}} T_{fi}^{ab}$$

If we have an $a \to 1 + 2$ decay,

$$\Gamma_{fi} = \frac{p^*}{32\pi^2 m_a^2} \int |M_{fi}|^2 d\Omega$$

Cross Sections

For $1 + 2 \to 3 + 4$:

$$\sigma = \frac{1}{64\pi^2 s} \frac{p_f^*}{p_i^*} \int |M_{fi}|^2 d\Omega^*$$

$$\frac{d\sigma}{d\Omega^*} = \frac{1}{64\pi^2 s} \frac{p_f^*}{p_i^*} |M_{fi}|^2$$

For elastic $e^-p \to e^-p$ scattering:

$$\frac{d\sigma}{d\Omega} = \frac{1}{64\pi^2} \left(\frac{E_3}{m_p E_1}\right)^2 |M_{fi}|^2$$

$$(E_3 = E_3(\theta))$$

Relativistic QM

Klein-Gordon: $(\partial^{\mu}\partial_{\mu} + m^2)\psi = 0$

Dirac:
$$(i\gamma^{\mu}\partial_{\mu} - m)\psi = 0$$

$$\gamma^{\mu}\gamma^{\nu} + \gamma^{\nu}\gamma^{\mu} = 2g^{\mu\nu}$$

$$\gamma^{0} = \beta = \begin{pmatrix} I & 0 \\ 0 & -I \end{pmatrix}$$
$$\gamma^{i} = \beta \alpha_{i} = \begin{pmatrix} 0 & \sigma_{k} \\ -\sigma_{k} & 0 \end{pmatrix}$$

Adjoints:

$$\begin{split} \bar{\psi} &= \psi^{\dagger} \gamma^0 \\ &= (\psi_0^*, \psi_1^*, -\psi_2^*, -\psi_3^*) \end{split}$$

Positive E solutions:

$$u_1(p) = \sqrt{E+m} \begin{pmatrix} 1\\0\\\frac{p_z}{E+m}\\\frac{p_x+ip_y}{E+m} \end{pmatrix}$$

$$u_2(p) = \sqrt{E+m} \begin{pmatrix} 0 \\ 1 \\ \frac{p_x - ip_y}{E+m} \\ \frac{-p_z}{E+m} \end{pmatrix}$$

Negative E solutions: antiparticles

$$v_1(p) = \sqrt{E+m} \begin{pmatrix} \frac{p_x - ip_y}{E+m} \\ -\frac{p_z}{E+m} \\ 0 \\ 1 \end{pmatrix}$$

$$v_2(p) = \sqrt{E+m} \begin{pmatrix} \frac{p_z}{E+m} \\ \frac{p_x+ip_y}{E+m} \\ 1 \\ 0 \end{pmatrix}$$

Current:

$$j^{\mu} = \psi^{\dagger} \gamma^0 \gamma^{\mu} \psi = \bar{\psi} \gamma^{\mu} \psi$$

Helicity

$$h = \frac{\mathbf{S} \cdot \mathbf{p}}{p} = \begin{pmatrix} \sigma \cdot \hat{p} & 0\\ 0 & \sigma \cdot \hat{p} \end{pmatrix}$$
$$s = \sin(\theta/2), c = \cos(\theta/2)$$

$$u^{\uparrow} = \sqrt{E + m} \begin{pmatrix} c \\ se^{i\phi} \\ \frac{p}{E + m} c \\ \frac{p}{E + m} se^{i\phi} \end{pmatrix}$$

$$u^{\downarrow} = \sqrt{E+m} \begin{pmatrix} -s \\ ce^{i\phi} \\ \frac{p}{E+m} s \\ -\frac{p}{E+m} ce^{i\phi} \end{pmatrix}$$

$$v^{\uparrow} = \sqrt{E+m} \begin{pmatrix} -\frac{p}{E+m} s \\ -\frac{p}{E+m} c e^{i\phi} \\ -s \\ c e^{i\phi} \end{pmatrix}$$

$$u^{\downarrow} = \sqrt{E+m} \begin{pmatrix} \frac{p}{E+m}c\\ \frac{p}{E+m}se^{i\phi}\\ c\\ se^{i\phi} \end{pmatrix}$$

Rel.Lim.: fraction=1, E + m = E

Bethe-Bloch

$$\frac{dE}{dx} \approx -4\pi\hbar^2 c^2 \alpha^2 \frac{nZ}{m_2 v^2} \left(\ln \left(\frac{2\beta^2 \gamma^2 c^2 m_e}{I_e} \right) - \beta^2 \right) \left\langle |M_{fi}|^2 \right\rangle = \frac{m_p^2 e^4}{E^2 \sin^4(\theta/2)} \cos^2(\theta/2)$$

$$\frac{d\sigma}{d\Omega} = \frac{\alpha^2}{4E^2 \sin^4(\theta/2)} \cos^2(\theta/2)$$

Parity

$$\hat{P} = \gamma^{0}$$

$$\hat{P}u(m,0) = +u(m,0)$$

$$\hat{P}v(m,0) = -v(m,0)$$

$$\hat{P}u_{\uparrow}(p,\theta,\phi) = u_{\downarrow}(p,\pi-\theta,\pi+\phi)$$

(not conserved in weak interactions)

Chirality

$$\gamma^5 = i\gamma^0 \gamma^1 \gamma^2 \gamma^3 = \begin{pmatrix} 0 & I \\ I & 0 \end{pmatrix}$$

$$P_R = \frac{1}{2} (1 + \gamma^5), P_R u_R = u_R$$

$$P_L = \frac{1}{2} (1 - \gamma^5), P_L u_L = u_L$$

$$(\gamma^5)^2 = 1, \gamma^5 \gamma^\mu = \gamma^\mu \gamma^5$$

In rel. limit (but not generally!),

$$u_R = u_{\uparrow}, u_L = u_{\downarrow}, v_R = v_{\uparrow}, v_L = v_{\downarrow}$$

I.e. $u_R = \lim_{rel} u_{\uparrow}$

For massless particles, helicity = chirality

Matrix Elements

$$\langle |M_{fi}|^2 \rangle = \frac{1}{4} \sum_{\text{spins}} |M|^2$$

$$e^+ e^- \to \mu^+ \mu^-$$

$$\langle |M_{fi}|^2 \rangle = 2e^4 \left(\frac{t^2 + u^2}{s^2}\right)$$

$$e^- q \to \mu^- q$$

$$\langle |M_{fi}|^2 \rangle = 2Q_q^2 e^4 \left(\frac{s^2 + u^2}{t^2}\right)$$

Deep Inelastic Scattering Rutherford: non-rel. e^- , no proton recoil

$$\left\langle |M_{fi}|^2 \right\rangle = \frac{m_p^2 m_e^2 e^4}{p^4 \sin^4(\theta/2)}$$
$$\frac{d\sigma}{d\Omega} = \frac{\alpha^2}{16E_K^2 \sin^4(\theta/2)}$$

Mott: rel. e^- , no proton recoil

$$\frac{\langle |M_{fi}|^2 \rangle = \frac{m_p \sigma}{E^2 \sin^4(\theta/2)} \cos^2(\theta/2)}{\frac{d\sigma}{d\Omega} = \frac{\alpha^2}{4E^2 \sin^4(\theta/2)} \cos^2(\theta/2)}$$

Form factors: p finite size, some recoil

$$\begin{split} M_{fi} &= M_{fi}^{pt} F(q^2) \\ F(q^2) &= \int \rho(r) e^{iq \cdot r} d^3 r \\ \frac{d\sigma}{d\Omega} &\to \left(\frac{d\sigma}{d\Omega}\right)_{\text{Mott}} |F(q^2)|^2 \\ \frac{d\sigma}{d\Omega} &= \frac{\alpha^2}{4E_1^2 \sin^4(\theta/2)} \frac{E_3}{E_1} \\ \left(\frac{G_e^2 + \tau G_M^2}{1 + \tau} + 2\tau G_M^2 \tan^2 \frac{\theta}{2}\right) \end{split}$$

(last one = Rosenbluth)

$$\left\langle R^2 \right\rangle = -6 \left[\frac{dF(q^2)}{dq^2} \right]_{q^2=0}$$

For high Q^2 elastics

$$\frac{d\sigma}{d\Omega}_{\rm elastic} \propto \frac{1}{Q^6} \frac{d\sigma}{d\Omega}_{\rm Mott}$$

Bjorken x: fraction of momentum transferred to struck quark

$$x = \frac{Q^2}{2p_2 \cdot q}$$

Elasticity (fract. energy loss): $y = \frac{p_2 \cdot q}{p_2 \cdot p_1}$

$$v = \frac{p_2 \cdot q}{m_p}$$

Any 2 of 4 variables are indep, x, y, v, Q^2 Deep inelastic scattering $Q^2 >> m_n^2 y^2$:

$$\begin{split} \frac{d^2\sigma}{dxdQ^2} &\approx \frac{4\pi\alpha^2}{Q^4} \left[(1-y) \frac{F_2(x,Q^2)}{x} + y^2 F_1(x,Q^2) \right] \\ Q^2 &= 4E_1 E_3 \sin^2 \frac{\theta}{2} \\ x &= \frac{Q^2}{2m_p(E_1 - E_3)} \end{split} \qquad \begin{aligned} &\text{Weak} \\ & y = 1 - \frac{E_3}{E_1} \end{aligned} \qquad &\text{If } q^2 \text{ small: } \frac{ig_{\mu\nu}}{m_W^2} \end{split}$$

 $F_i = \text{structure functions.}$

Bjorken scaling: $F_i \approx \text{flat over } Q^2 \text{ and }$ $F_2(x) = 2xF_1(x).$

Can rewrite σ in terms of q^2 :

$$\frac{d\sigma}{dq^2} = \frac{2\pi\alpha^2Q_q^2}{q^4}\left[1+\left(1+\frac{q^2}{s}\right)^2\right]$$

Feynman Rules

 $-iM = \prod [\text{terms}]$

- Put adjoints first

- Put polarizations out front

- Same index on all things involved in a ver-

QED

In lepton: u(p)

Out lepton: $\bar{u}(p)$

In anti-lepton: $\bar{v}(p)$

Out anti-lepton: v(p)

In photon: $\epsilon_{\mu}(p)$

Out photon: $\epsilon_{\mu}(p)^*$

Lepton prop: $\frac{i\gamma^{\mu}q_{\mu}+m}{q^2-m^2}$

Photon prop: $\frac{-ig_{\mu\nu}}{a^2}$

QCD

Quarks: same as QED leptons

Gluons: same as QED photons

Gluon prop: $-\frac{ig_{\mu\nu}}{a^2}\delta^{ab}$

Vertices (qqg, ggg, gggg): $-ig_s \frac{1}{2} \lambda_{ii}^a \gamma^{\mu}$

$$\psi = \psi_{\text{flavor}} \chi_{\text{spin}} \xi_{\text{colour}} \eta_{\text{space}}$$

Colour confinement: colour neutral final states

Colour factors $C(ik \to jl) = \frac{1}{4} \sum_{a=1}^{8} \lambda_{ii}^{a} \lambda_{lk}^{a}$ (r=1, g=2, b=3). Remember colour conservation when applicable!

$$\left\langle |C|^2 \right\rangle = \frac{1}{9} \sum_{ijkl=1}^3 |C(ij \to kl)|^2$$

$$M_{QCD} = CM_{QED}$$

Fermions: same as before

W prop: $-\frac{i[g_{\mu\nu}-q_{\mu}q_{\nu}/m_W^2]}{q^2-m_W^2}$

If q^2 small: $\frac{ig_{\mu\nu}}{m^2}$

Vertex: $\frac{-ig_w}{\sqrt{2}}\frac{1}{2}\gamma^{\mu}(1-\gamma^5)$