Matrix product states for the absolute beginner

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Brief overview: Why tensor networks?

Graphical notation

Matrix Product States and Matrix Product Operators

Compressing Matrix Product States

Energy optimization

Time evolution

Periodic and infinite MPS

Focus on basic computations and algorithms with MPS

not covered: entanglement area laws, RG, topological aspects, symmetries etc.

Quantum mechanics is complex



Dirac

The fundamental laws necessary for the mathematical treatment of a large part of physics and the whole of chemistry are thus completely known ...

the difficulty lies only in the fact that application of these laws leads to equations that are too complex to be solved.

$$\Psi(r_1,r_2,...r_n)$$
 $\Psi(s_1,s_2,...s_L)$ $\Psi(n_1,n_2,...n_L)$ n electron positions L spins L particle occupancies

Exponential complexity to represent wavefunction

This view of QM is depressing

[The Schrodinger equation] cannot be solved accurately when the number of particles exceeds about 10. No computer existing, or that will ever exist, can break this barrier because it is a catastrophe of dimension ...

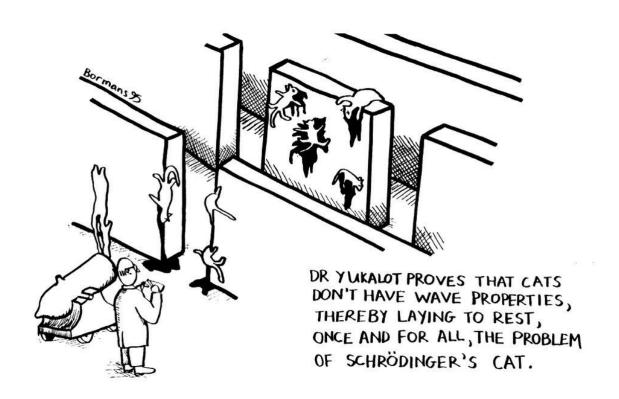
Pines and Laughlin (2000)

in general the many-electron wave function Ψ ... for a system of N electrons is not a legitimate scientific concept [for large N]

Kohn (Nobel lecture, 1998)

illusion of complexity

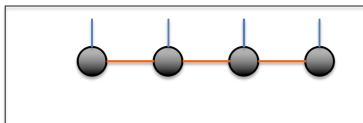
nature does not explore all possibilities



Nature is local: ground-states have low entanglement

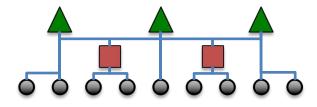
Language for low entanglement states is tensor networks

different tensor networks reflect geometry of entanglement



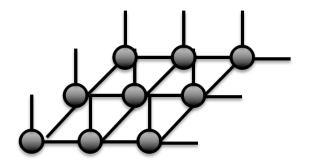
Matrix Product State

1D entanglement for **gapped** systems (basis of DMRG - often used in quasi-2D/3D)



MERA

1D/nD entanglement for gapless systems



Tensor Product State (PEPS)

nD entanglement for gapped systems

Graphical language

$$|\Psi\rangle = \sum_{n_1 n_2 n_3} \Psi^{n_1 n_2 n_3} |n_1 n_2 n_3\rangle$$

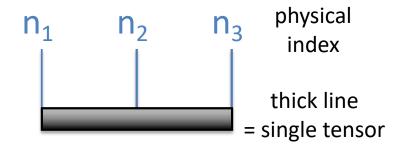
e.
$$|n\rangle=\{|\uparrow\rangle,|\downarrow\rangle\}$$
 spin 1/2

g.
$$|n\rangle=\{|0\rangle,|1\rangle\}$$
 particles

Algebraic form

 $\Psi^{n_1 n_2 n_3}$

Graphical form



General state

Graphical language, cont'd

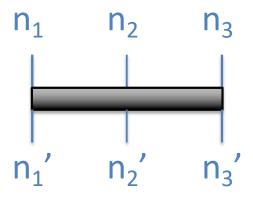
$$\hat{O} = \sum_{nn'} O_{n'_1 n'_2 n'_3}^{n_1 n_2 n_3} |n_1 n_2 n_3\rangle \langle n'_1 n'_2 n'_3|$$

Algebraic form

General operator

 $O_{n'_1 n'_2 n'_3}^{n_1 n_2 n_3}$

Graphical form



Ex: overlap, expectation

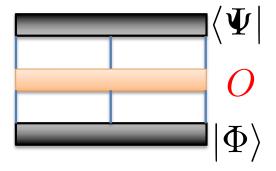
Overlap

$$\langle \Psi | \Phi \rangle = \sum_{n}^{n} \Psi_{n_1 n_2 n_3} \Phi^{n_1 n_2 n_3}$$



Expectation value

$$\langle \Psi | O | \Phi \rangle = \sum_{n,n'} \Psi_{n_1 n_2 n_3} O_{n'_1 n'_2 n'_3}^{n_1 n_2 n_3} \Phi^{n_1 n_2 n_3}$$



Low entanglement states

What does it mean for a state to have low entanglement?

Consider system with two parts, 1 and 2

No entanglement

$$\Psi^{n_1 n_2} = A^{n_1} A^{n_2}$$

local measurements on separated system 1, system 2 can be done independently. **Local realism** (classical)

Entangled state

$$\Psi^{n_1 n_2} = \sum_{i} A_{i}^{n_1} A_{i}^{n_2}$$

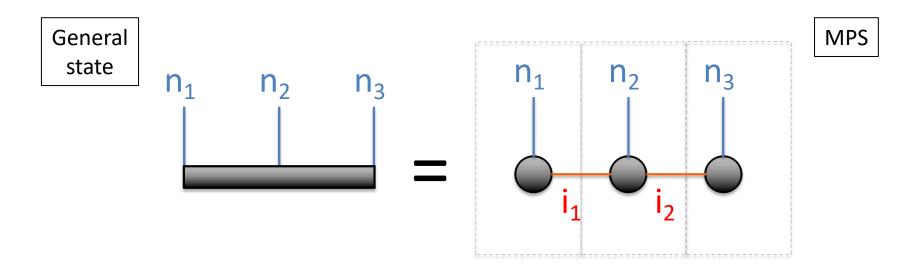
Low entanglement : small number of terms in the sum

Matrix product states

first and last tensors have one fewer auxiliary index

$$m{\Psi}^{n_1n_2n_3...n_l} = \sum_i \stackrel{\downarrow}{A_{i_1}^{n_1}A_{i_2i_2}^{n_2}A_{i_2i_3}^{n_3}...A_{i_l}^{n_l}}^{1D ext{ structure}}$$
 of entanglement $^{ ext{"M" or "D" or "}\chi"} = \mathbf{A}^{n_1}\mathbf{A}^{n_2}\ldots\mathbf{A}^{n_l}$

amplitude is obtained as a product of matrices

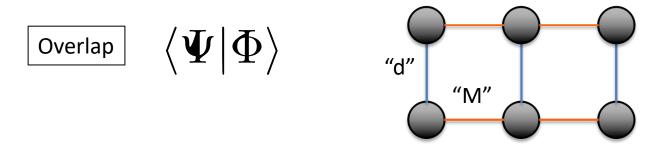


MPS gauge

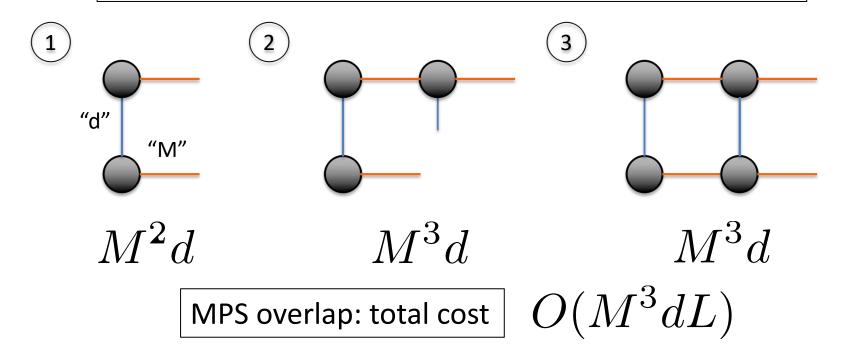
MPS are not unique: defined up to gauge on the auxiliary indices

$$\mathbf{G}\mathbf{G}^{-1} = \mathbf{1}$$
insert gauge matrices
$$=$$

Ex: MPS contraction



Efficient computation: contract in the correct order!



MPS from general state

recall singular value decomposition (SVD) of matrix

$$\Psi^{nm} = \sum_{i} L_{i}^{n} \sigma_{i} R_{i}^{m}$$
 singular values

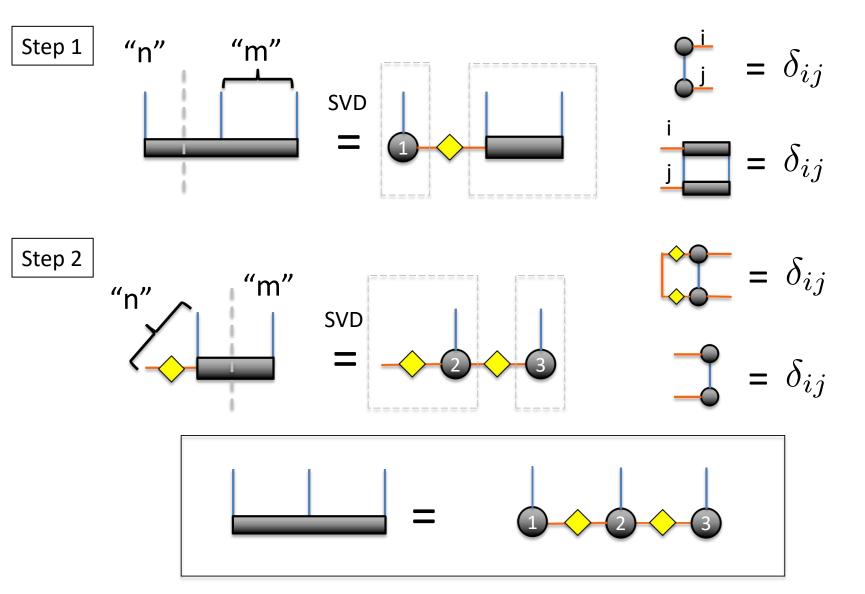
$$\sum_{i=1}^{n}L_{i}^{n}L_{j}^{n}=\delta_{ij}$$

$$\sum_{i=1}^{n}R_{i}^{n}R_{j}^{n}=\delta_{ij}$$
 orthogonality conditions

$$\int_{j}^{i} = \delta_{ij}$$

$$= \delta_{ij}$$
orthogonality conditions

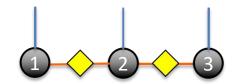
MPS from general state, cont' d



Common canonical forms

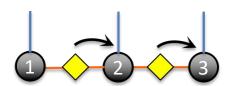
"Vidal" form





different canonical form: absorb singular values into the tensors

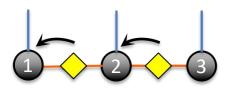
left canonical



all tensors contract to unit matrix from left

$$\int_{\underline{j}}^{\underline{i}} = \delta_{ij} \quad \boxed{ \int_{\underline{j}}^{\underline{i}} = \delta_{ij} } = \text{etc}$$

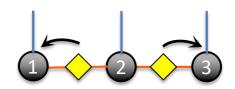
right canonical



all tensors contract to unit matrix from right

$$\int_{j}^{i} = \delta_{ij} \quad \int_{j}^{i} = \delta_{ij} \text{ etc}$$

mixed canonical
around site 2
(DMRG form)

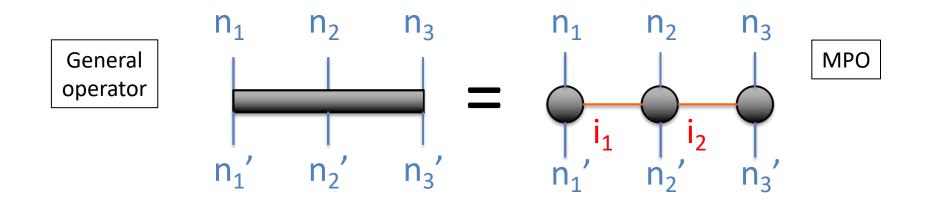


$$\delta_{ij} \quad \delta_{ij} \quad \int_{\mathbf{j}}^{\mathbf{i}} = \delta_{ij}$$

Matrix product operators

$$O_{n'_{1}n'_{2}n'_{3}...n'_{l}}^{n_{1}n_{2}n_{3}...n'_{l}} = \sum_{i} W_{i_{1}}^{n_{1},n'_{1}} W_{i_{1}i_{2}}^{n_{2},n'_{2}} W_{i_{2}i_{3}}^{n_{3}n'_{3}} \dots W_{i_{l}}^{n_{l}n'_{l}}$$

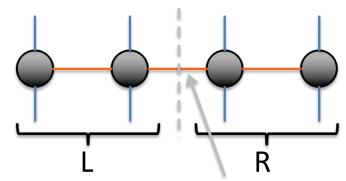
each tensor has a bra and ket physical index



Typical MPO's

$$H = \sum_{\langle ij \rangle} S_i \cdot S_j$$

What is bond dimension as an MPO?



joins pairs of operators on both sides

$$H_L\otimes \mathbf{1}_R+\mathbf{1}_L\otimes H_R+\sum_{lpha=x,y,z}S_2^lpha\cdot S_3^lpha$$

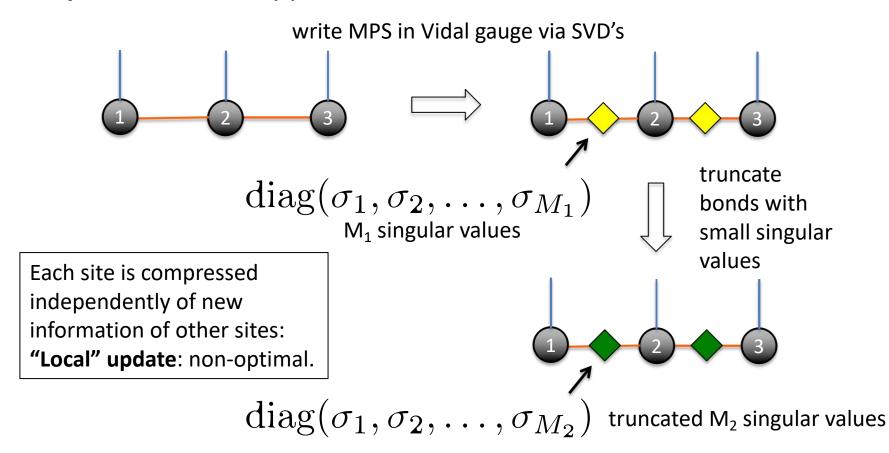
MPO acting on MPS

MPO on MPS leads to new MPS with **product** of bond dimensions

MPS compression: SVD

many operations (e.g. MPOxMPS, MPS+MPS) increase bond dim.

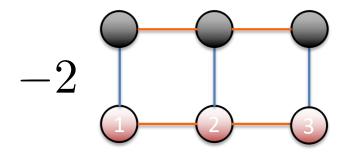
compression: best approximate MPS with smaller bond dimension.

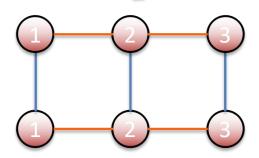


MPS: variational compression

solve minimization problem

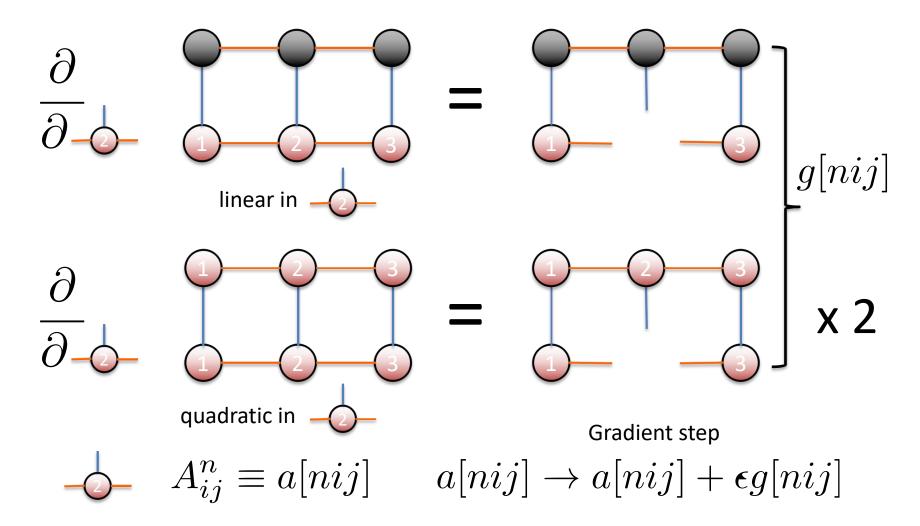
$$\min_{\Phi \text{ original new MPS }} \langle \Psi - \Phi | \Psi - \Phi \rangle = \min_{\Phi} [-2 \langle \Phi | \Psi \rangle + \langle \Phi | \Phi \rangle]$$
 (fixed) MPS





Gradient algorithm

To minimize quantity, follow its gradient until it vanishes



Sweep algorithm (DMRG style)

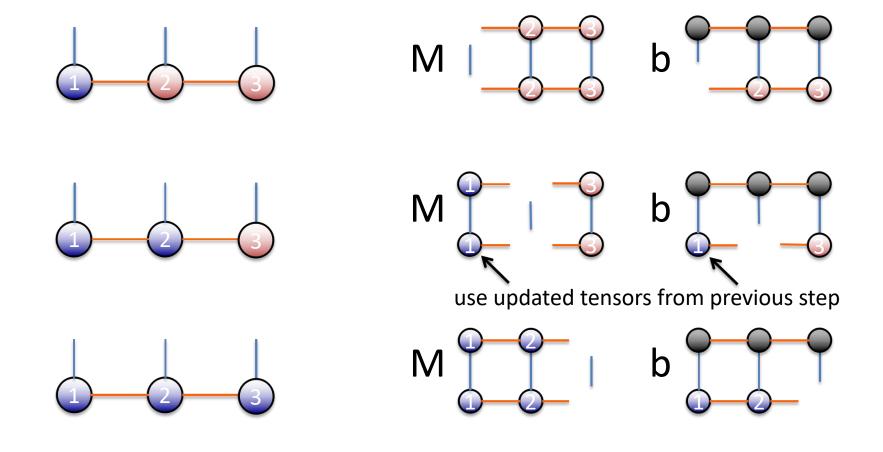
$$\langle \Psi - \Phi | \Psi - \Phi \rangle$$
 bilinear in $\bullet - \bullet - \bullet$ as vector $A^n_{ij} \equiv a[nij]$
$$= b^T a \quad \text{where } b^T = \bullet \bullet \bullet$$

$$= a^T M a \quad \text{where } M = \bullet \bullet \bullet$$

$$\min(a^T M a - b^T a) \Rightarrow M a = b$$

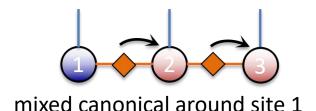
Sweep algorithm cont' d

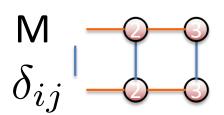
Minimization performed site by site by solving $\,Ma=b\,$

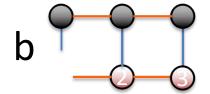


Sweep and mixed canonical form

in mixed canonical form $\,Ma=b\Rightarrow a=b\,$



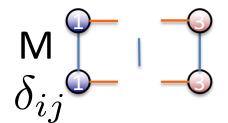


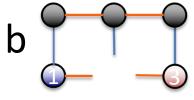


change canonical form



mixed canonical around site 2 note: updated singular values





SVD vs. variational compression

variational algorithms – optimization for each site depends on all other sites. Uses "full environment"

SVD compression: "local update". Not as robust, but cheap!

MPS: full environment / local update same computational scaling, only differ by number of iterations.

General tensor networks (e.g. PEPS): full environment may be expensive to compute or need further approximations.

Energy optimization

$$\min_{\mathbf{\Psi}} \left[\langle \mathbf{\Psi} | H | \mathbf{\Psi} \rangle - E \langle \mathbf{\Psi} | \mathbf{\Psi} \rangle \right]$$

bilinear form: similar to compression problem

commonly used algorithms

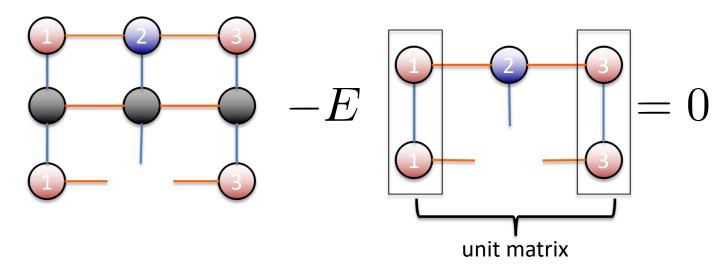
DMRG: variational sweep with full environment

imag. TEBD: local update, imag. time evolution + SVD compression

DMRG energy minimization

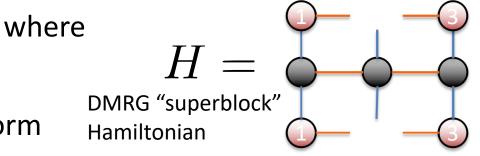
$$\frac{\partial}{\partial \mathbf{P}} \langle \mathbf{\Psi} | H | \mathbf{\Psi} \rangle - E \langle \mathbf{\Psi} | \mathbf{\Psi} \rangle = 0$$

use mixed canonical form around site 2



$$Ha = Ea$$

eigenvalue problem for each site, in mixed canonical form



Time evolution

real time evolution

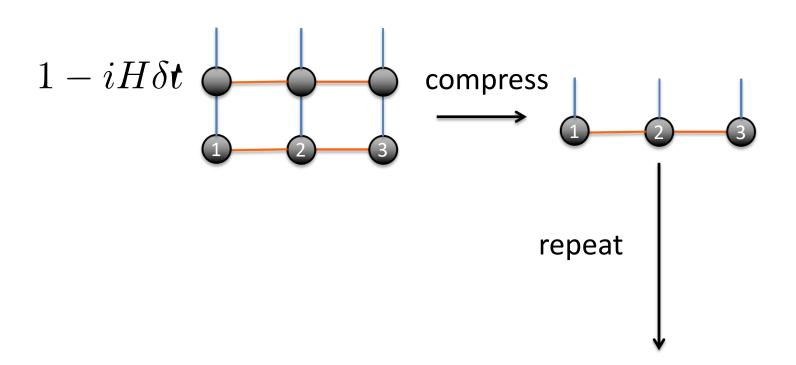
$$|\Psi(\delta t)\rangle = \exp(-iH\delta t)|\Psi(0)\rangle$$

imaginary time evolution: replace i by 1.

projects onto ground-state at long times.

General time-evolution

$$1 - iH\delta t = -iH\delta t$$

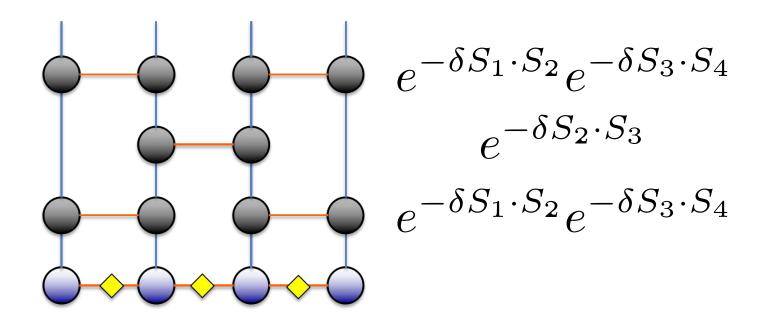


Short range H: Trotter form

$$H = \sum_{ij} S_i \cdot S_j$$

$$e^{-\delta H} = e^{-\delta S_1 \cdot S_2} e^{-\delta S_2 \cdot S_3} \dots + O(\delta^2)$$
 evolution on pairs of bonds

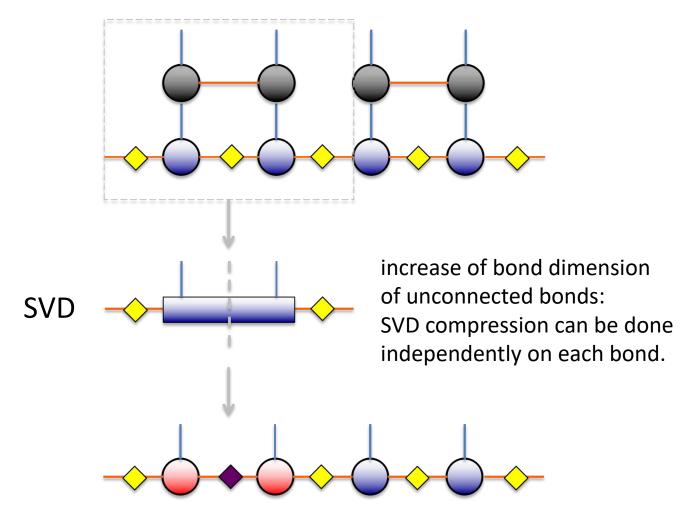
Even-odd evolution



time evolution can be broken up into even and odd bonds

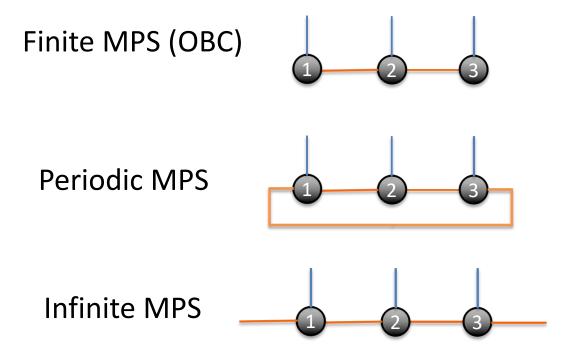
Time-evolving block decimation

even/odd evolution easy to combine with SVD compression: TEBD



Periodic and infinite MPS

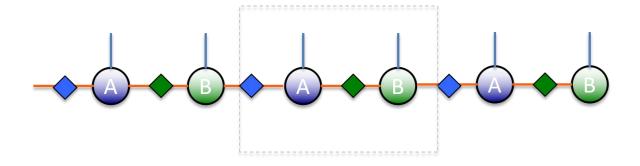
MPS easily extended to PBC and thermodynamic limit



Infinite TEBD

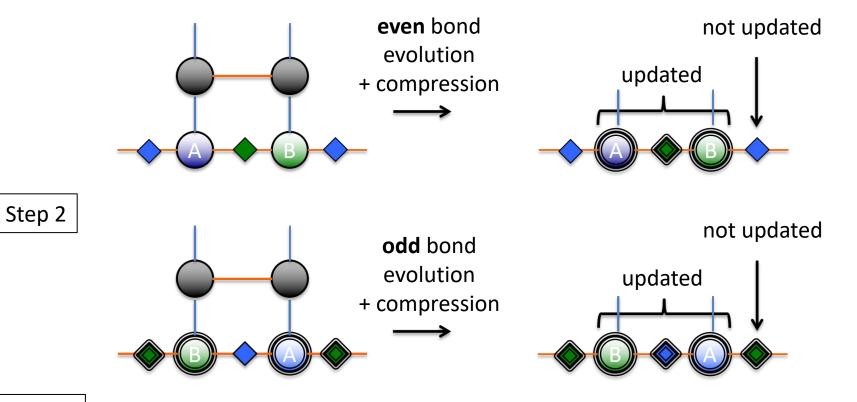
local algorithms such as TEBD easy to extend to infinite MPS

Unit cell = 2 site infinite MPS



i-TEBD cont' d

Step 1



Repeat

Symmetries

Given global symmetry group, local site basis can be labelled by irreps of group – quantum numbers

U(1) – site basis labelled by integer n (particle number)
$$\left| n \right\rangle$$

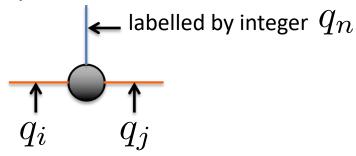
SU(2) symmetry – site basis labelled by j, m (spin quanta) |jm
angle

Total state associated with good quantum numbers

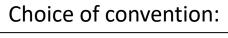
$$|\Psi\rangle = |\Psi(n, j, m...)\rangle$$

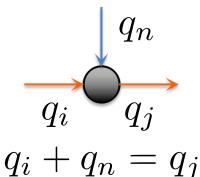
MPS and symmetry

bond indices can be labelled by same symmetry labels as physical sites e.g. particle number symmetry

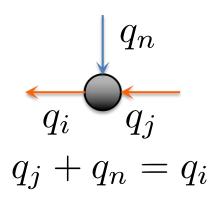


MPS: well defined Abelian symmetry, each tensor fulfils rule

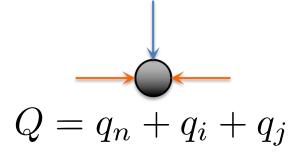




$$\sum q_{in} = \sum q_{out}$$



tensor with no arrows leaving gives total state quantum number



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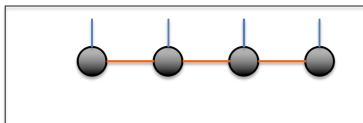
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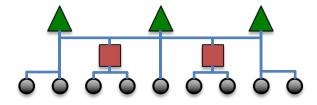
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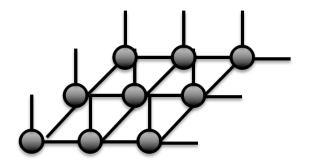
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MERA

1D/nD entanglement for gapless systems



Tensor Product State (PEPS)

nD entanglement for gapped systems

Questions

- 1. What is the dimension of MPS (M_1) + MPS (M_2) ?
- 2. How would we graphically represent the DM of an MPS, (tracing out sites n_3 to n_L ?)
- 3. What is the dimension of the MPO of an electronic Hamiltonian with general quartic interactions?
- 4. What happens when we use an MPS to represent a 2D system?
- 5. What happens to the bond-dimension of an MPS as we evolve it in time? Do we expect the MPS to be compressible? How about for imaginary time evolution?
- 6. How would we alter the discussion of symmetry for non-Abelian symmetry e.g. SU(2)?