

# Lecture 22: Balance equations – Momentum and energy balances

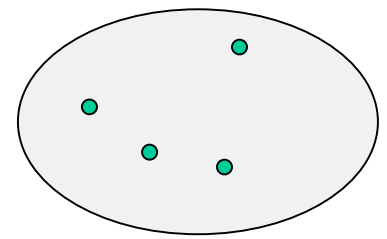
- Recap balance laws
- The momentum balance
- The energy balance
- Differential balance laws

Book: Ch. 11.2, 11.4

# Process modeling and balance laws

- The balance laws are formulated for «conserved quantities»:
  - Mass (or other quantities that are «equivalent» to mass, such as moles, particles, etc.)
  - Momentum
  - Energy
- Process modeling is done by
  1. formulating the relevant balance laws, and
  2. finding the «closure relations» that is used to determine the flows in a balance law, as function of the state («inventory») of the balance law
- The state («inventory») of a balance law is what is used as a measure for the conserved quantity
  - Such as mass, moles, concentration, level, pressure, ... for mass balance,
  - velocity or flows for momentum balance, and
  - temperature for energy balance

# The basic physical principles



Consider a volume consisting of a **fixed** number of fluid particles, with total mass  $m$ , total momentum  $\vec{p}$  and total energy  $E$ . From basic physics (conservation laws), we know the following principles hold:

- Conservation of mass (mass balance):

$$\frac{dm}{dt} = 0$$

- Newton's second law (momentum balance)

$$\frac{d\vec{p}}{dt} = \vec{F}$$

Also holds for angular momentum,  $\vec{h} = \vec{r} \times \vec{p}$ :

$$\frac{d\vec{h}}{dt} = \vec{r} \times \vec{F} = \vec{T}$$

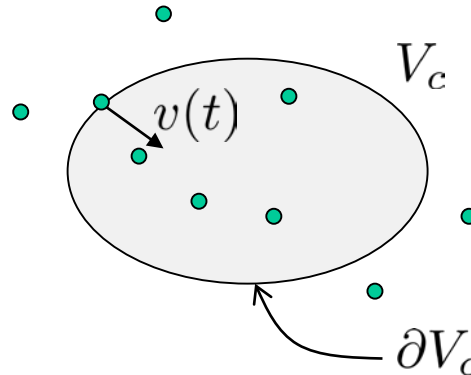
- First law of thermodynamics (conservation of energy, energy balance):

$$\frac{dE}{dt} = \dot{Q} - \dot{W}$$

Rate of heat flowing into volume from surroundings  $\rightarrow \dot{Q}$   
 $\dot{W}$   $\rightarrow$  Rate at which work is done by the body at surroundings

# The balance laws

- Assume a **fixed** control volume (of arbitrary size and shape), where fluid flows across the control volume

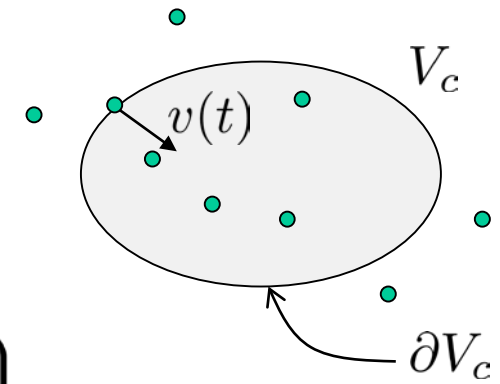


$$B = \iiint_{V_c} \rho \beta(\mathbf{x}, t) dV$$

- The general integral (macroscopic) balance law for  $B$  is

$$\frac{d}{dt} B = \left\{ \begin{array}{l} \text{transfer of } B \text{ through} \\ \text{surface } \partial V_c \text{ by} \\ \text{fluid flow (convection)} \end{array} \right\} + \left\{ \begin{array}{l} \text{other effects that} \\ \text{transfer } B \text{ into } V_c \\ \text{(indep. of fluid flow)} \end{array} \right\}$$

# The integral balance laws



- **Mass balance** (without reactions/phase transfer)

$$\frac{d}{dt}m = \left\{ \begin{array}{l} \text{transfer of mass into} \\ V_c \text{ by fluid flow} \\ \text{across surface } \partial V_c \end{array} \right\}$$

- **Momentum** (note: momentum is a vector)

$$\frac{d}{dt}\mathbf{p} = \left\{ \begin{array}{l} \text{transfer of momentum into} \\ V_c \text{ by fluid flow} \\ \text{across surface } \partial V_c \end{array} \right\} + \left\{ \begin{array}{l} \text{generation of momentum} \\ \text{in } V_c \text{ due to forces} \\ \text{acting on } V_c \end{array} \right\}$$

- **Energy**

$$\frac{d}{dt}E = \left\{ \begin{array}{l} \text{transfer of energy into} \\ V_c \text{ by fluid flow} \\ \text{across surface } \partial V_c \end{array} \right\} + \left\{ \begin{array}{l} \text{transfer of energy into} \\ V_c \text{ by heat transfer} \\ \text{and by work} \end{array} \right\}$$

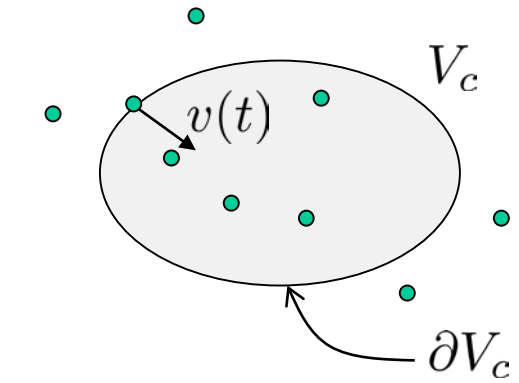
# The mass balance

- In words

$$\frac{d}{dt}m = \left\{ \begin{array}{l} \text{transfer of mass into} \\ V_c \text{ by fluid flow} \\ \text{across surface } \partial V_c \end{array} \right\}$$

- Mathematically

$$\frac{d}{dt}m = \frac{d}{dt} \iiint_{V_c} \rho dV = - \iint_{\partial V_c} \rho \vec{v} \cdot \vec{n} dA$$

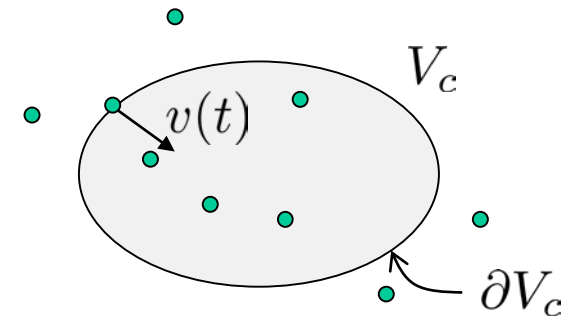


“Convection”

- Often, we have one (or more) «point inflows»  $w_{\text{in},i}$ , and outflows  $w_{\text{out},i}$ . Then mass balance can be formulated as

$$\frac{d}{dt}m = \sum_i w_{\text{in},i} - \sum_i w_{\text{out},i}$$

# The momentum balance



- In words

$$\frac{d}{dt}\mathbf{p} = \left\{ \begin{array}{c} \text{transfer of momentum into} \\ V_c \text{ by fluid flow} \\ \text{across surface } \partial V_c \end{array} \right\} + \left\{ \begin{array}{c} \text{generation of momentum} \\ \text{in } V_c \text{ due to forces} \\ \text{acting on } V_c \end{array} \right\}$$

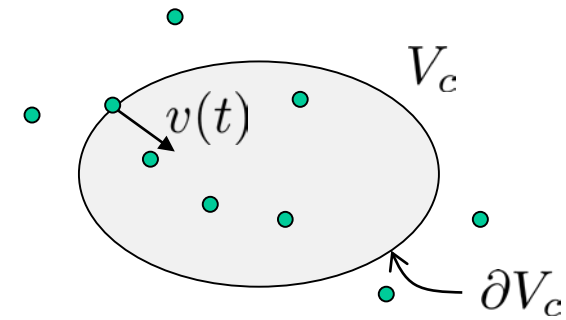
- Mathematically

$$\frac{d}{dt}\vec{p} = \frac{d}{dt} \iiint_{V_c} \rho \vec{v} dV = - \iint_{\partial V_c} \rho \vec{v} \vec{v} \cdot \vec{n} dA + \vec{F}^{(r)}$$

where  $\vec{F}^{(r)}$  is resultant force on fluid in control volume

(often: gravity (hydrostatic) and/or friction (hydrodynamic))

# The energy balance



- In words

$$\frac{d}{dt}E = \left\{ \begin{array}{c} \text{transfer of energy into} \\ V_c \text{ by fluid flow} \\ \text{across surface } \partial V_c \end{array} \right\} + \left\{ \begin{array}{c} \text{transfer of energy into} \\ V_c \text{ by heat transfer} \\ \text{and by work} \end{array} \right\}$$

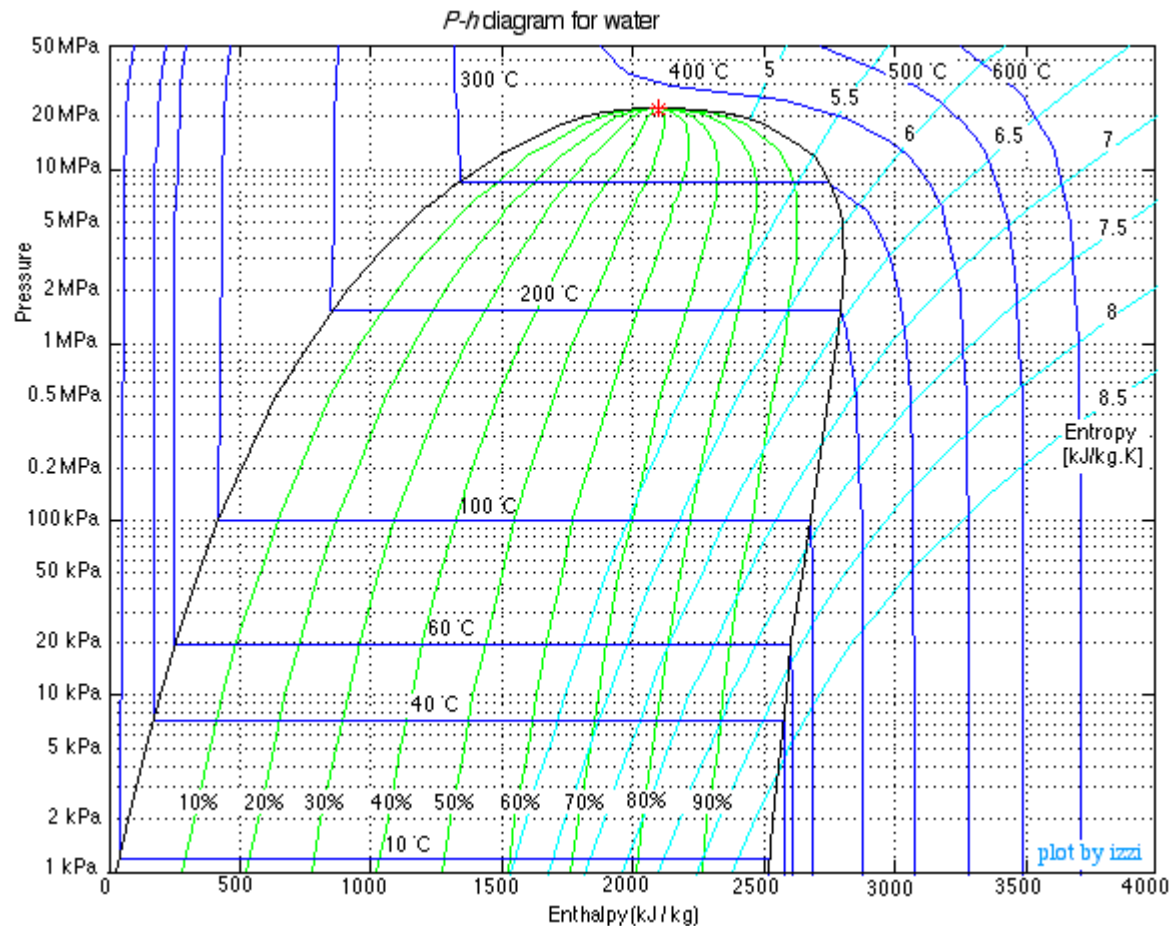
- Mathematically

$$\frac{d}{dt}E = \frac{d}{dt} \iiint_{V_c} \rho e dV = \underbrace{- \iint_{\partial V_c} \rho e \vec{v} \cdot \vec{n} dA}_{\text{Energy flow by convection}} + \dot{Q} - \dot{W}$$

- What is the energy of a fluid?



# P-h-diagram for water



# Energy

$$\frac{d}{dt}E = \frac{d}{dt} \iiint_{V_c} \rho e dV = - \iint_{\partial V_c} \rho e \vec{v} \cdot \vec{n} dA + \dot{Q} - \dot{W}$$

- The energy of a fluid of mass  $m$ , moving with a velocity  $v$  at a height  $z$  in a gravitational field:

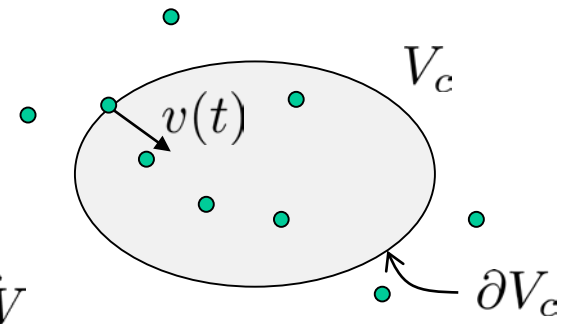
$$E = \underbrace{U}_{\text{internal energy}} + \underbrace{\frac{1}{2}mv^2}_{\text{kinetic energy}} + \underbrace{mgz}_{\text{potential energy}}$$

- Specific energy:

$$e = u + \frac{1}{2}v^2 + gz$$

# Heat and work flow

$$\frac{d}{dt}E = \frac{d}{dt} \iiint_{V_c} \rho e dV = - \iint_{\partial V_c} \rho e \vec{v} \cdot \vec{n} dA + \dot{Q} - \dot{W}$$



- Heat flow

$$\dot{Q} = \iint_{\partial V_c} \vec{j}_Q \cdot \vec{n} dA$$

- Work flow

$$\dot{W} = \underbrace{\iint_{\partial V_c} p \vec{v} \cdot \vec{n} dA}_{\text{flow work}} + \underbrace{\dot{W}_s}_{\text{shaft work}}$$

# Enthalpy

- The energy balance can be written

$$\frac{d}{dt} \iiint_{V_c} \rho e dV = - \iint_{\partial V_c} \rho \left( e + \frac{p}{\rho} \right) \vec{v} \cdot \vec{n} dA - \dot{W}_s + \dot{Q}$$

where the first term on the RHS is convection and flow work

- Define **enthalpy** as

$$h = u + \frac{p}{\rho}$$

- Then

$$\frac{d}{dt} \iiint_{V_c} \rho \left( u + \frac{1}{2} v^2 + gz \right) dV = - \iint_{\partial V_c} \rho \left( h + \frac{1}{2} v^2 + gz \right) \vec{v} \cdot \vec{n} dA - \dot{W}_s + \dot{Q}$$

# Internal energy and enthalpy

- Specific heat capacities:

$$c_v := \left. \frac{\partial u}{\partial T} \right|_{\text{constant volume}} \qquad c_p := \left. \frac{\partial h}{\partial T} \right|_{\text{constant pressure}}$$

(found in tables for different fluids, often assumed constant)

- If assumed constant, implies that energy and enthalpy is (linear) function of temperature only:

$$\frac{du}{dt} = c_v \frac{dT}{dt} \qquad u(T_2) - u(T_1) = c_v (T_2 - T_1)$$

$$\frac{dh}{dt} = c_p \frac{dT}{dt} \qquad h(T_2) - h(T_1) = c_p (T_2 - T_1)$$

- For ideal gases:

$$c_v = c_p + R$$

- For incompressible fluids (often assumed for liquids):

$$c_v = c_p$$

# Examples energy balances...

# Differential mass balance

- Recall the integral mass balance:

$$\underbrace{\frac{d}{dt} \iiint_{V_c} \rho dV}_{\text{Mathematics (obvious?)}} = - \underbrace{\iint_{\partial V_c} \rho \vec{v} \cdot \vec{n} dA}_{\text{Divergence theorem}}$$

$$\frac{d}{dt} \iiint_{V_c} \rho dV = \iiint_{V_c} \frac{\partial \rho}{\partial t} dV$$

$$\iint_{\partial V_c} \rho \vec{v} \cdot \vec{n} dA = \iiint_{V_c} \vec{\nabla} \cdot (\rho \vec{v}) dV$$

- That is:

$$\iiint_{V_c} \frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot (\rho \vec{v}) dV = 0$$

- This must hold for arbitrary control volumes, which implies

$$\frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot (\rho \vec{v}) = 0$$

Differential mass balance, also called *continuity equation* or *advection equation*

# Alternative formulations

- The differential mass balance

$$\frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot (\rho \vec{v}) = 0$$

- From definition of nabla operator, this is the same as

$$\frac{\partial \rho}{\partial t} + \sum_{i=1}^3 \frac{\partial}{\partial x_i} (\rho v_i) = 0, \quad \mathbf{v} = (v_1, v_2, v_3)^T$$

- If we introduce the *material derivative*,

$$\frac{D\phi}{Dt} := \frac{\partial \phi}{\partial t} + \mathbf{v}^T \nabla \phi = \frac{\partial \phi}{\partial t} + \sum_{i=1}^3 \frac{\partial \phi}{\partial x_i} v_i$$

The material derivative is the derivative following a particle (as opposed to the derivative at a fixed point in space)

and use product rule, we can write

$$\frac{D\rho}{Dt} + \rho \nabla \cdot \vec{v} = 0$$



# Differential momentum and energy balances

- Differential momentum balance for inviscid fluid (*Euler's equation*)

$$\rho \frac{D\vec{v}}{Dt} = -\vec{\nabla} p + \rho \vec{f}, \quad \text{where } \rho \vec{f} \text{ is the mass force (e.g. gravity)}$$

- For viscous (Newtonian) fluids, the differential momentum balance is the famous *Navier-Stokes* equation:

$$\rho \frac{D\vec{v}}{Dt} = -\vec{\nabla} p + \mu \vec{\nabla}^2 \vec{v} + \rho \vec{f}$$

- Differential energy balance (for example)

$$\rho \frac{D}{Dt} \left( \frac{1}{2} \vec{v}^2 + u \right) = -\vec{\nabla} \cdot (p\vec{v}) - \vec{\nabla} \cdot \vec{j}_Q + \rho \vec{v} \cdot \vec{f}$$

# Computational fluid dynamics

- CFD = solving momentum + mass balances (that is, Navier-Stokes + continuity equation) for different setups



<http://physbam.stanford.edu/~fedkiw/>

# Example of differential energy balances: The heat equation of a solid

- The energy balance:

$$\rho \frac{D}{Dt} \left( \frac{1}{2} \vec{v}^2 + u \right) = -\vec{\nabla} \cdot (p\vec{v}) - \vec{\nabla} \cdot \vec{j}_Q + \rho \vec{v} \cdot \vec{f}$$

- Solid: Disregard kinetic and potential energy, no velocity:

$$\rho \frac{\partial u}{\partial t} = -\vec{\nabla} \cdot \vec{j}_Q$$

- We need a «closure relation». Here in the form of *Fourier's law*:

$$\vec{j}_Q = -\alpha \vec{\nabla}(\rho c_p T)$$

- Combined with

$$\frac{\partial u}{\partial t} = c_p \frac{\partial T}{\partial t}$$

we get

$$\frac{\partial T}{\partial t} - \alpha \vec{\nabla} \cdot \vec{\nabla} T = 0$$

In one dimension:

$$\frac{\partial T(x, t)}{\partial t} - \alpha \frac{\partial^2 T(x, t)}{\partial x^2} = 0$$