

# Computational examples in transport phenomena with Python

## 1 Steady-state conduction in a lengthwise-insulated rod with internal heat production

This is a derivation of the temperature distribution function  $T(x)$  for a steady-state heat conduction in a straight rod of length  $L$ . We assume that the internal heat production  $Q_p$  is present in every point inside the rod volume (which may simulate heating up of an electrical wire). The rod is perfectly insulated along its length and it loses heat only through its endpoints which in a steady-state case are kept at a fixed temperature  $T_0$ .

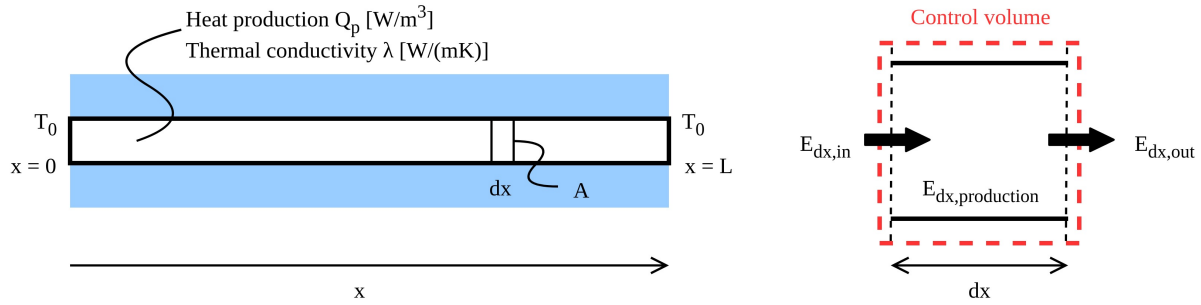


Figure 1: Conduction in a rod with internal heat production

We will take for the control volume a slice  $dx$  from the rod.

The energy balance for the rod element  $dx$  reads

$$\frac{dE_{dx}}{dt} = E_{dx,in} - E_{dx,out} + E_{dx,production} \quad (1)$$

(note here that  $E_{dx,in}$ ,  $E_{dx,out}$  and  $E_{dx,production}$  are energies per unit time and so have the units of W)

The heat flow can be described by the Fourier's law for one-dimensional heat conduction

$$\phi = \lambda A \left( -\frac{dT}{dx} \right) \quad (2)$$

Hence

$$E_{dx,in} = \lambda A \left( -\frac{dT}{dx} \right)_x \quad (3)$$

$$E_{dx,out} = \lambda A \left( -\frac{dT}{dx} \right)_{x+dx} \quad (4)$$

The energy per unit time coming from the production can be written as  $Q_p$  multiplied by the volume of the slice  $dx$

$$E_{dx,production} = Q_p A dx \quad (5)$$

In the steady-state  $\frac{dE}{dt} = 0$  and the energy balance becomes

$$\lambda A \left( -\frac{dT}{dx} \right)_x - \lambda A \left( -\frac{dT}{dx} \right)_{x+dx} + Q_p A dx = 0 \quad (6)$$

Simplifying the above energy balance we get

$$\frac{\left( \frac{dT}{dx} \right)_{x+dx} - \left( \frac{dT}{dx} \right)_x}{dx} = -\frac{Q_p}{\lambda}$$

It is interesting to note here that we have lost the dependence on the cross-sectional surface area of the rod.

If we now substitute some function  $f(x) = \frac{dT}{dx}$  we notice that we have

$$\frac{f(x+dx) - f(x)}{dx} = -\frac{Q_p}{\lambda}$$

in other words

$$\frac{df(x)}{dx} = -\frac{Q_p}{\lambda} \quad (7)$$

With the above substitution, the differential equation that we are about to solve becomes

$$\frac{d^2T}{dx^2} = -\frac{Q_p}{\lambda} \quad (8)$$

Applying the boundary conditions from both ends of the rod, the solution to the above differential equation is

$$T(x) = -\frac{Q_p}{2\lambda}(x^2 - Lx) + T_0 \quad (9)$$

## Computational example

As a computational example we will draw the graph of the temperature distribution in a copper rod  $200m$  long. We take that the thermal conductivity for this rod is  $400 \frac{W}{m \cdot K}$ . The internal heat production in the entire rod is  $20W$ .

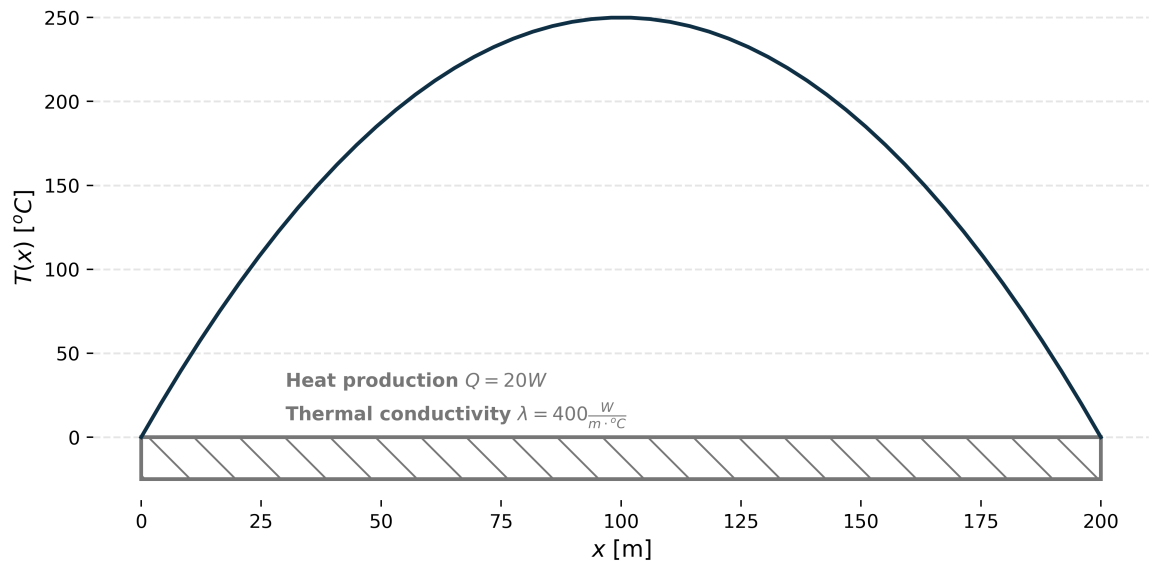


Figure 2: Temperature distribution in a rod with internal heat production of  $20W$

## 2 Evaporating sphere

## 3 The Stefan problem

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