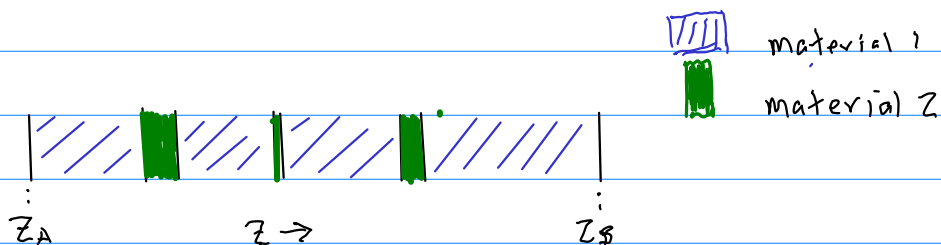


Fields institute mini-course : Lecture #3

Numerical methods for the 1D Schrodinger - Poisson equation

Equations



$\chi(z)$ = dielectric factor - piecewise constant, constant over each layer

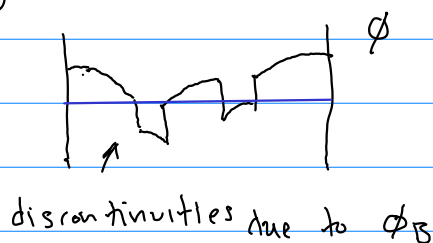
The potential $\phi(z) = \tilde{\phi}(z) + \phi_D(z) + \phi_B(z)$ where

$\phi_B(z)$ = "bandshift" - piecewise constant, constant over each layer

$$\frac{d}{dz} \chi(z) \frac{d}{dz} \phi_D = -\rho_D(z) \quad \phi_D(z_A) = q_A \quad \phi_D(z_B) = q_B$$

$$\frac{d}{dz} \chi(z) \frac{d}{dz} \tilde{\phi} = -\rho_e(z) \quad \tilde{\phi}(z_A) = 0 \quad \tilde{\phi}(z_B) = 0$$

Sample ϕ plot:



$$\rho_e(z) = 2 \sum_{\lambda_j < E_f} \psi_j^*(z) \psi_j(z) W(\lambda_j, E_f)$$

$$\left\{ (E_f - \lambda_j) \frac{m_r(z)}{2\pi \hbar^2} \right\}$$

where

$$\frac{\hbar^2}{2} \left[\frac{d}{dz} \left(\frac{1}{m_e(z)} \frac{d}{dz} \right) + [\tilde{\phi}(z) + \phi_D(z) + \phi_B(z)] \right] \psi_j = \lambda_j \psi_j$$

$m_e(z)$ = effective mass - piecewise constant, constant over each layer

First computational task.

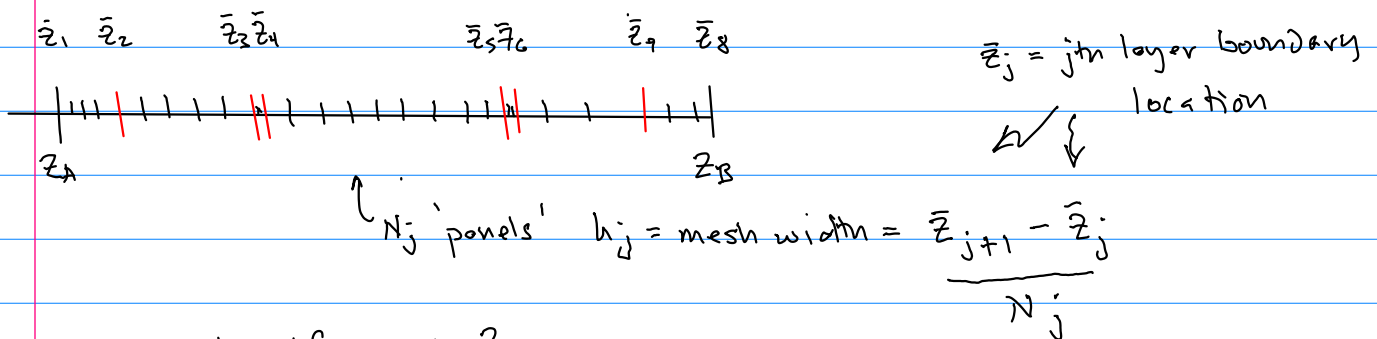
Solve problems of the form

$$\frac{d}{dz} \left(a(z) \frac{du}{dz} \right) = f(z) \quad u(z_A) = g_A \quad u(z_B) = g_B \quad a(z) \text{ piecewise constant.}$$

General approach: set up a linear system of equations whose solution gives values that approximate the exact solution values,

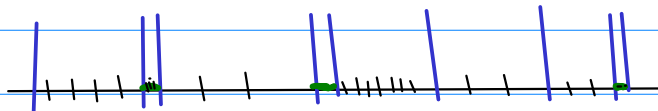
Derive the discrete equations using a simple, but very useful technique "finite volume discretization".

Grid: Semi-uniform, a collection of M uniform grids, one for each material layer with N_j panels the j th layer.



Why a semi-uniform grid?

Some problems consist of a mixture of very thin layers and very thick layers; we don't want a small mesh width in the thin layers to dictate the mesh widths in all the other layers.



Background: Finite Volume Discretization.

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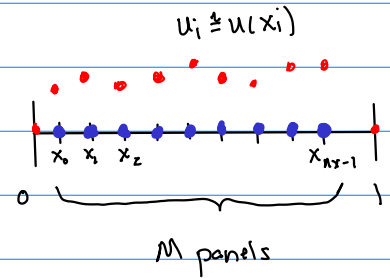
Single uniform grid

$$(1D) \frac{\partial}{\partial x} (a(x) \frac{\partial u}{\partial x}) = f(x) \quad x \in [0,1] \quad u(0) = u(1) = 0 \quad (*)$$

Seek approximate values $u_i \approx u(x_i)$

at interior grid points x_i .

$$x_i = (i+1)h_x \quad i = 0, 1, 2, \dots, n_x-1 \quad \text{and} \quad h_x = \frac{1}{n_x}.$$



$$h_x = \frac{1}{n_x}$$

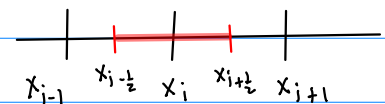
$n_x = M-1$ interior points.

Obtain an equation for each u_i by requiring an integral approximation to $(*)$ hold in a finite volume

$$[x_{i-1/2}, x_{i+1/2}] \quad \left(x_{i+1/2} = x_i + \frac{h_x}{2} \quad x_{i-1/2} = x_i - \frac{h_x}{2} \right)$$

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$$\frac{d}{dx} (a(x) \frac{du}{dx}) = f(x) \Rightarrow \int_{x_{i-1/2}}^{x_{i+1/2}} \frac{d}{dx} (a(x) \frac{du}{dx}) ds = \int_{x_{i-1/2}}^{x_{i+1/2}} f(s) ds.$$



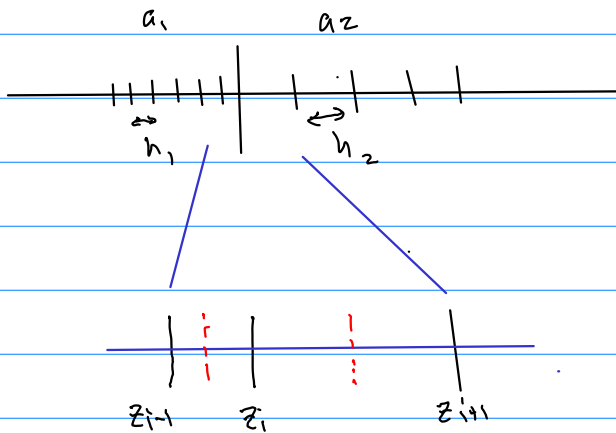
$$\Rightarrow a(x) \frac{du}{dx} \Big|_{x_{i-1/2}}^{x_{i+1/2}} = \int_{x_{i-1/2}}^{x_{i+1/2}} f(s) ds. \quad \text{Obtain equations by approximating this exact relation, with } \frac{du}{dx} \Big|_{x_{i+1/2}} \approx \frac{u_{i+1} - u_i}{h_x}, \quad \frac{du}{dx} \Big|_{x_{i-1/2}} \approx \frac{u_i - u_{i-1}}{h_x} \quad \text{and the midpoint rule for the integral of } f.$$

$$\frac{a_{i+1/2} [u_{i+1} - u_i]}{h_x} - \frac{a_{i-1/2} [u_i - u_{i-1}]}{h_x} = f_i h_x \quad u_0, u_{n_x} \text{ given}$$

$$\begin{bmatrix} (a_{1/2} + a_{3/2}) & a_{3/2} & & & \\ a_{3/2} & (a_{3/2} + a_{5/2}) & a_{5/2} & & \\ & a_{5/2} & \ddots & \ddots & \\ & & \ddots & \ddots & a_{n_x-1/2} \\ & & & & a_{n_x-1/2} \end{bmatrix} \begin{bmatrix} u_1 \\ u_2 \\ \vdots \\ u_{n_x-1} \end{bmatrix} = (h_x) \begin{bmatrix} f_1 \\ f_2 \\ \vdots \\ f_{n_x-1} \end{bmatrix} + \begin{bmatrix} -(a_{1/2} u_0)/h_x \\ \\ \\ -(a_{n_x-1/2} u_{n_x})/h_x \end{bmatrix}$$

"symmetric tri diagonal"

What changes have to be made when one is using a semi-uniform mesh?
 The only change has to do with the creation of the equations at points that are lower boundaries. :



$$\int_{z_{i-1/2}}^{z_{i+1/2}} \frac{d}{dz} a(z) \frac{d}{dz} dz = \int_{z_{i-1/2}}^{z_{i+1/2}} f(z) dz$$

$$\Rightarrow a_2 \left. \frac{du}{dz} \right|_{z_{i+1/2}} - a_1 \left. \frac{du}{dz} \right|_{z_{i-1/2}} = \int_{z_{i-1/2}}^{z_{i+1/2}} f(z) dz$$

$$a_2 \left(\frac{u_{i+1} - u_i}{h_2} \right) - a_1 \left(\frac{u_i - u_{i-1}}{h_1} \right) \approx f(z_i) \left(\frac{h_1}{2} + \frac{h_2}{2} \right)$$

This leads to a system of equations that is symmetric tri-diagonal, which, when one includes the modification of the right hand side that is due to the boundary conditions, has the form

$$\begin{bmatrix} 0 & & \\ & \ddots & \\ 0 & & 1 \end{bmatrix} \vec{u} = D \vec{f} - \vec{f}_{\text{bary}} \quad D = \begin{bmatrix} n_1 & & \\ & (n_1+h_1)/2 & \\ & h_2 & \\ & & (n_2+h_2)/2 \end{bmatrix}$$

due to the change in mesh sizes

So, solution procedure for

$$\frac{d}{dz} \kappa(z) \frac{d}{dz} \phi = f \quad \phi(z_a) = g_a \quad \phi(z_b) = g_b$$

- (1) Construct matrices associated with the discrete approximation obtained using finite volume based discretizations.

Grid specification $\rightsquigarrow L$ (tri-diagonal matrix)
 $+ \rightsquigarrow D$ (diagonal matrix of mesh weights).
 $\kappa(z)$

(2) Given f , form $f^* = Df - \begin{bmatrix} \kappa_1 \phi'(z_a) \\ h_1 \\ 0 \\ \vdots \\ 0 \\ \kappa_m \phi(z_b) \\ h_m \end{bmatrix}$ \swarrow Boundary value forcing term.

(3) Solve $L\vec{\phi} = f^*$ to obtain $\phi_i \approx \phi(z_i)$.

Note: L is symmetric tri-diagonal so using a standard tri-diagonal solver, $L\vec{\phi} = f^*$ can be obtained in $O(n)$ operations where $n = \text{dimension of } L$.

\swarrow (calls the solver for you.)

In Matlab or Octave, use $\phi = L \setminus f^*$ for
 or, you can write your own tri-diagonal solver (a typical exercise given in numerical analysis classes).

Generally, when creating discrete approximations to self-adjoint operators one strives to obtain symmetric matrices. Why? Symmetric matrices have properties that facilitate the construction of efficient solution procedures.