

# Chp 25: Current, Resistance, and Electromotive Force

## Goals for Chapter 25

- To understand current and how charges move in a conductor
- To understand resistivity and conductivity
- To calculate the resistance of a conductor
- To learn how an emf causes current in a circuit
- To calculate energy and power in circuits

# Current

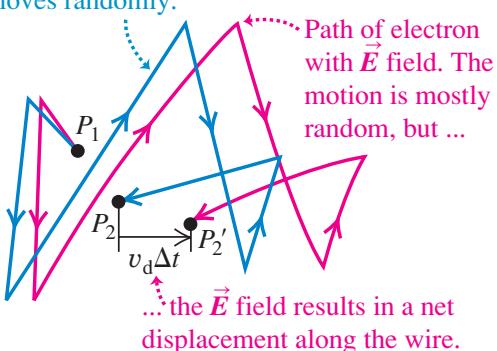
Electric circuits contain *charges in motion*. Circuits are at the heart of modern devices such as computers, televisions, and industrial power systems.

Current is any motion of charge from one region to another and is defined as  $I = dQ/dt$ .

Conductor without internal  $\vec{E}$  field



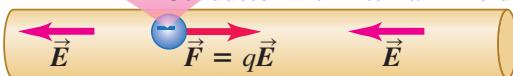
Path of electron without  $\vec{E}$  field. Electron moves randomly.



...the  $\vec{E}$  field results in a net displacement along the wire.

Path of electron with  $\vec{E}$  field. The motion is mostly random, but ...

Conductor with internal  $\vec{E}$  field



An electron has a negative charge  $q$ , so the force on it due to the  $\vec{E}$  field is in the direction opposite to  $\vec{E}$ .

# Current

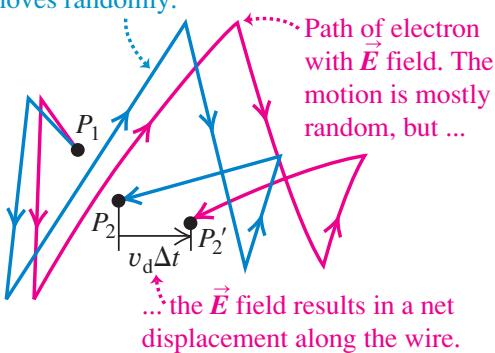
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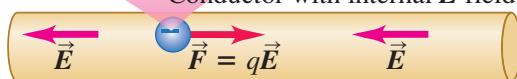


Path of electron without  $\vec{E}$  field. Electron moves randomly.



In electrostatic situations the electric field is zero everywhere within the conductor, and there is *no current*. However, this does not mean that all charges within the conductor are at rest. In an ordinary metal some of the electrons are free to move within the conducting material in all directions, somewhat like the molecules of a gas but with much greater speeds, of the order of  $10^6$  m/s. The electrons nonetheless do not escape from the conducting material, because they are attracted to the positive ions of the material. The motion of the electrons is random, so there is no *net flow of charge* in any direction and hence no current.

Conductor with internal  $\vec{E}$  field



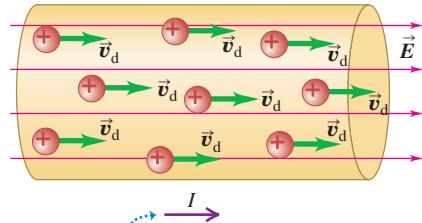
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If a constant, steady electric field  $\vec{E}$  is established inside a conductor, free electrons inside the conducting material are then subjected to a steady force  $\vec{F} = q\vec{E}$ . A charged particle moving in a conductor undergoes frequent collisions with the massive, nearly stationary ions of the material.

This motion described in terms of **drift velocity**  $v_d$

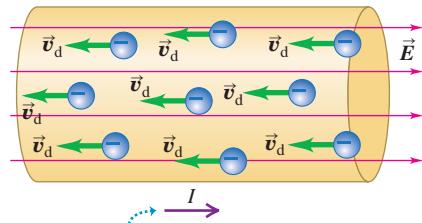
# The Direction of Current Flow

(a)



A conventional current is treated as a flow of positive charges, regardless of whether the free charges in the conductor are positive, negative, or both.

(b)



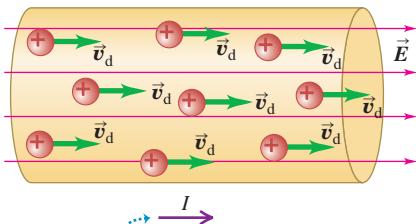
In a metallic conductor, the moving charges are electrons — but the *current* still points in the direction positive charges would flow.

$$I = \frac{dQ}{dt} \quad (\text{definition of current})$$

The SI unit of current is the ampere; one ampere is defined to be one coulomb per second  
1 A = 1 C/s

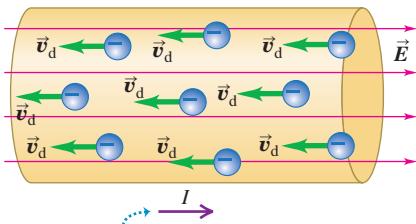
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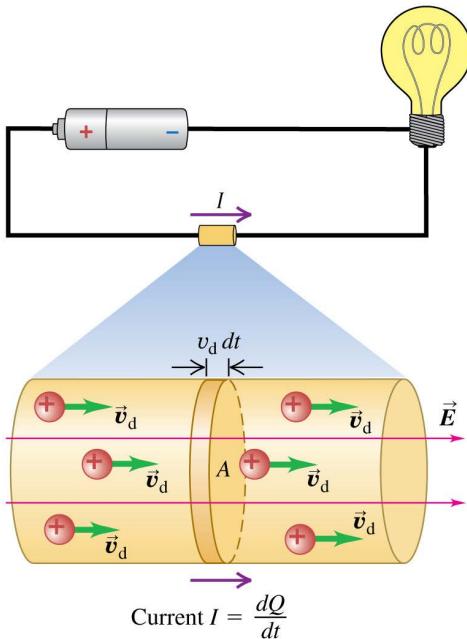
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The drift of moving charges through a conductor can be interpreted in terms of work and energy. The electric field  $E$  does work on the moving charges. The resulting kinetic energy is transferred to the material of the conductor by means of collisions with the ions, which vibrate about their equilibrium positions in the crystalline structure of the conductor. This energy transfer increases the average vibrational energy of the ions and therefore the temperature of the material. Thus much of the work done by the electric field goes into heating the conductor, *not* into making the moving charges move ever faster and faster. This heating is sometimes useful, as in an electric toaster, but in many situations is simply an unavoidable by-product of current flow.

We describe currents as though they consisted entirely of positive charge flow, even in cases in which we know that the actual current is due to electrons. Hence the current is to the right in both Figures. This choice or convention for the direction of current flow is called **conventional current**. While the direction of the conventional current is *not* necessarily the same as the direction in which charged particles are actually moving, we'll find that the sign of the moving charges is of little importance in analyzing electric circuits.

# Current, drift velocity, and current density



Suppose there are  $n$  moving charged particles per unit volume. We call  $n$  the concentration of particles; its SI unit is  $\text{m}^{-3}$ .

In a time interval  $dt$ , each particle moves a distance  $v_d dt$ .

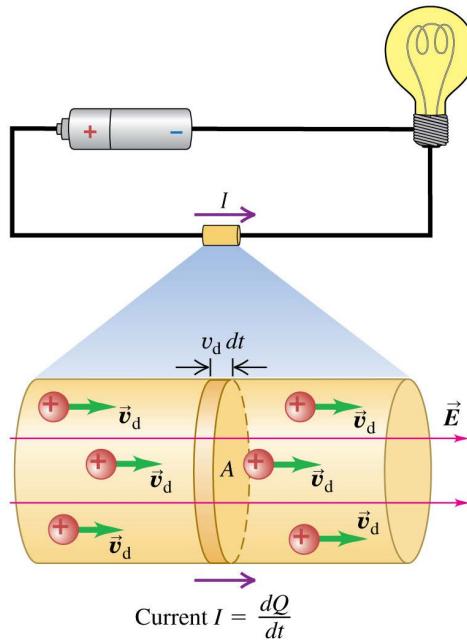
The volume of the shaded cylinder is  $A v_d dt$ , and the number of particles within it is  $n A v_d dt$ .

the charge  $dQ$  that flows out of the end of the cylinder during time  $dt$  is

$$dQ = q(n A v_d dt) = n q v_d A dt$$

$$\rightarrow I = \frac{dQ}{dt} = n |q| v_d A \quad (\text{general expression for current})$$

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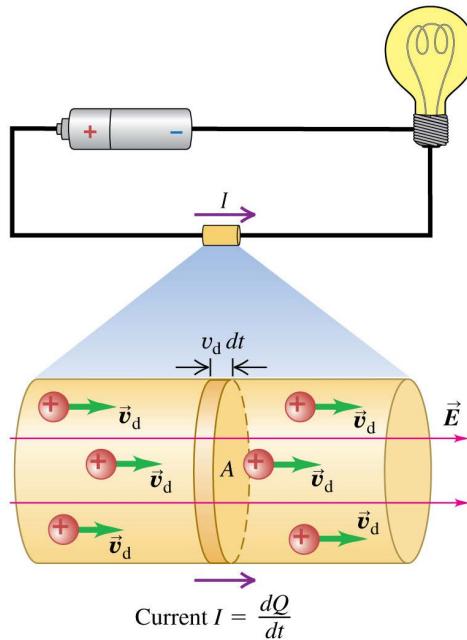
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The current in the solution is carried by  $\text{Na}^+$  and  $\text{Cl}^-$  ions

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The current *per unit cross-sectional area* is called the current density  $J$ :

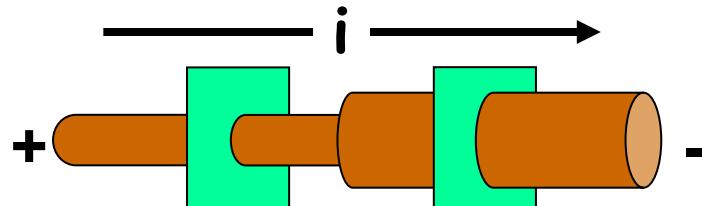
$$J = \frac{I}{A} = n|q|v_d \quad (\text{A/m}^2)$$

$$\vec{J} = nq\vec{v}_d \quad (\text{vector current density})$$

There are *no* absolute value signs here. If  $q$  is positive,  $\vec{v}_d$  is in the same direction as  $\vec{E}$ ; if  $q$  is negative,  $\vec{v}_d$  is opposite to  $\vec{E}$ . In either case,  $\vec{J}$  is in the same direction as  $\vec{E}$ .

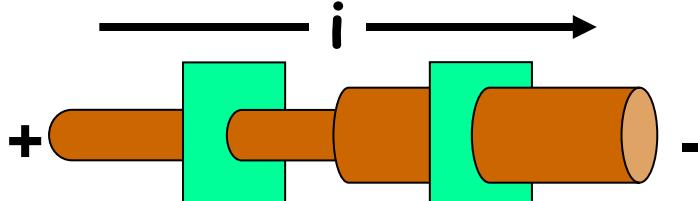
Note that current density  $\vec{J}$  is a vector but the current  $I$  is not!

**Current is the same across each cross-section of a wire**



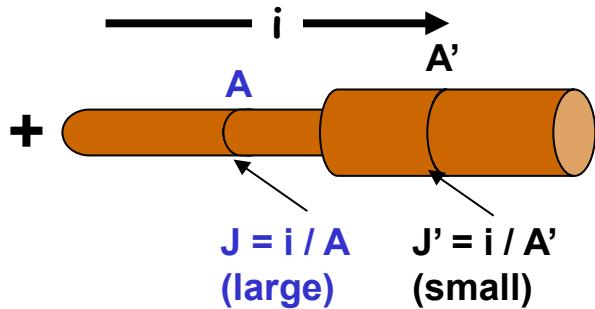
**Current density  $J$  may vary**  
 $[J] = \text{current}/\text{area}$

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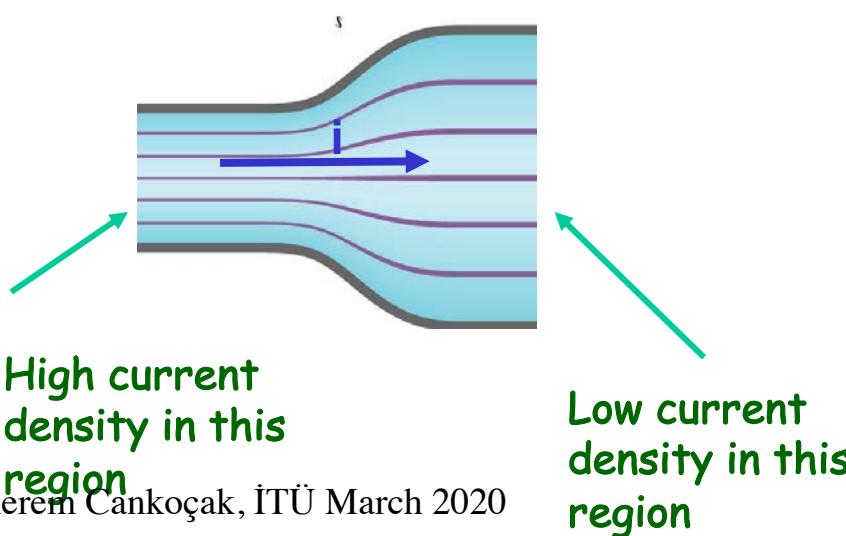


Current density  $\underline{J}$  may vary  
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## Current density $J$ : Current / Unit Area (Vector)



Same current crosses larger or smaller Surfaces,  
current density  $J$  varies



For uniform density  $i = J/A$  or  $J = i/A$

If not uniform:  $di \equiv \vec{J} \circ \hat{n} dA$

$$i = \int_{\text{area}} \vec{J} \circ d\vec{A}$$

## Resistivity ( $\rho$ )

The current density  $J$  in a conductor depends on the electric field  $E$  and on the properties of the material. In general, this dependence can be quite complex. But for some materials, especially metals, at a given temperature,  $J$  is nearly *directly proportional* to  $E$ , and the ratio of the magnitudes of  $E$  and  $J$  is constant. This relationship, called Ohm's law, was discovered in 1826 by the German physicist Georg Simon Ohm (1787-1854).

The **resistivity** of a material is the ratio of the electric field in the material to the current density it causes:  $\rho = E/J$ .

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## Resistivities at Room Temperature (20°C)

Substance	$\rho$ ( $\Omega \cdot m$ )	Substance	$\rho$ ( $\Omega \cdot m$ )
Silver	$1.47 \times 10^{-8}$	Semiconductors	
Copper	$1.72 \times 10^{-8}$	Pure carbon (graphite)	$3.5 \times 10^{-5}$
Gold	$2.44 \times 10^{-8}$	Pure germanium	0.60
Aluminum	$2.75 \times 10^{-8}$	Pure silicon	2300
Tungsten	$5.25 \times 10^{-8}$	Insulators	
Steel	$20 \times 10^{-8}$	Amber	$5 \times 10^{14}$
Lead	$22 \times 10^{-8}$	Glass	$10^{10}-10^{14}$
Mercury	$95 \times 10^{-8}$	Lucite	$>10^{13}$
Manganin (Cu 84%, Mn 12%, Ni 4%)	$44 \times 10^{-8}$	Mica	$10^{11}-10^{15}$
Constantan (Cu 60%, Ni 40%)	$49 \times 10^{-8}$	Quartz (fused)	$75 \times 10^{16}$
Nichrome	$100 \times 10^{-8}$	Sulfur	$10^{15}$
		Teflon	$>10^{13}$
		Wood	$10^8-10^{11}$

$$\rho = \frac{E}{J}$$

$$(\text{V/m})/(\text{A/m}^2) = \text{V} \cdot \text{m/A}$$

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$$\rho = \frac{E}{J}$$

$$(V/m)/(A/m^2) = V \cdot m/A$$

1 V/A is called one ohm ( $\Omega$ )  $\rightarrow$  SI units for  $\rho$  is  $\Omega \cdot m$ .

The reciprocal of resistivity is **conductivity ( $\sigma$ )**. Its units are  $(\Omega \cdot m)^{-1}$

# Resistivity and temperature

The resistivity of a *metallic* conductor nearly always increases with increasing temperature. As temperature increases, the ions of the conductor vibrate with greater amplitude, making it more likely that a moving electron will collide with an ion; this impedes the drift of electrons through the conductor and hence reduces the current. Over a small temperature range (up to 100 C° or so), the resistivity of a metal can be represented approximately by the equation

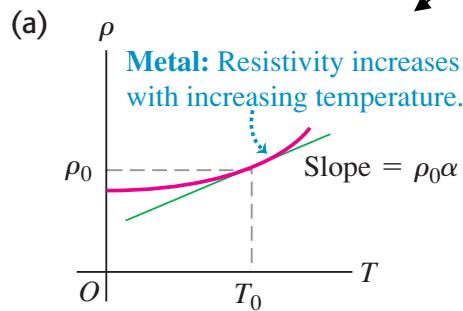
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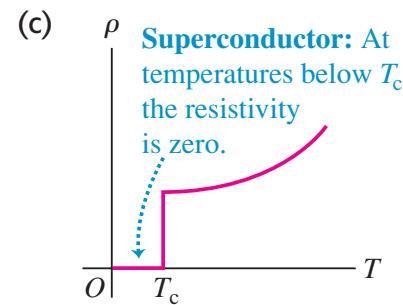
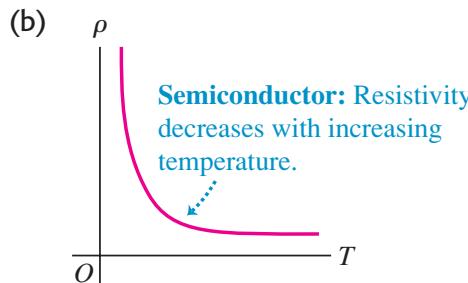
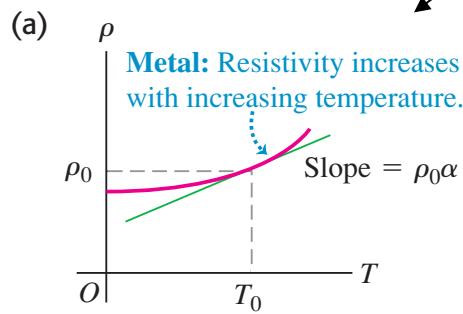


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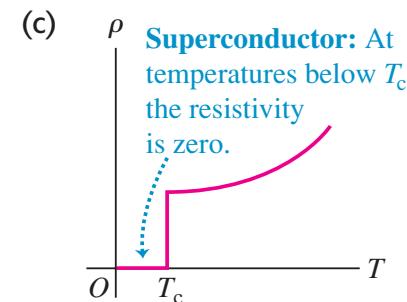
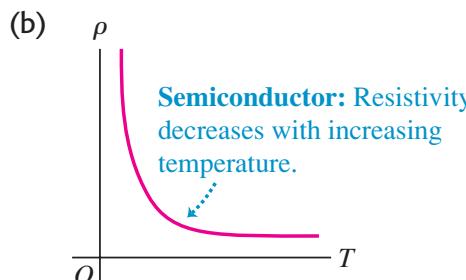
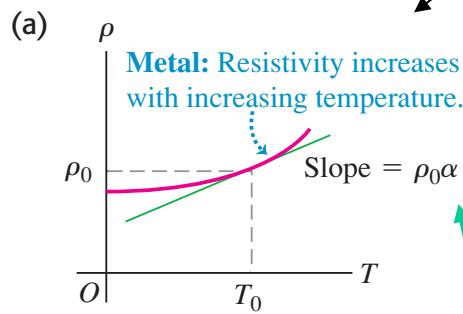


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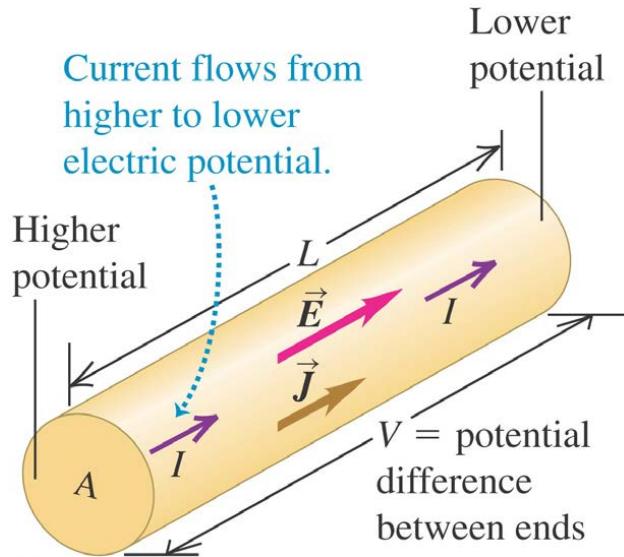
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Aluminum	0.0039	Lead	0.0043
Brass	0.0020	Manganin	0.00000
Carbon (graphite)	-0.0005	Mercury	0.00088
Constantan	0.00001	Nichrome	0.0004
Copper	0.00393	Silver	0.0038
Iron	0.0050	Tungsten	0.0045

# Resistance ( $R$ )

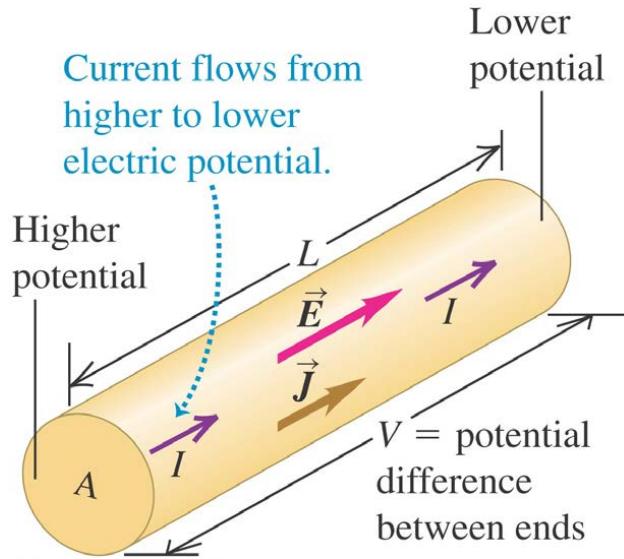


The resistance of a conductor is  $R = \rho L/A$

The potential across a conductor is  $V = IR$ .

If  $V$  is directly proportional to  $I$  (that is, if  $R$  is constant), the equation  $V = IR$  is called *Ohm's law*.

# Resistance ( $R$ )



Current flows from higher to lower electric potential.

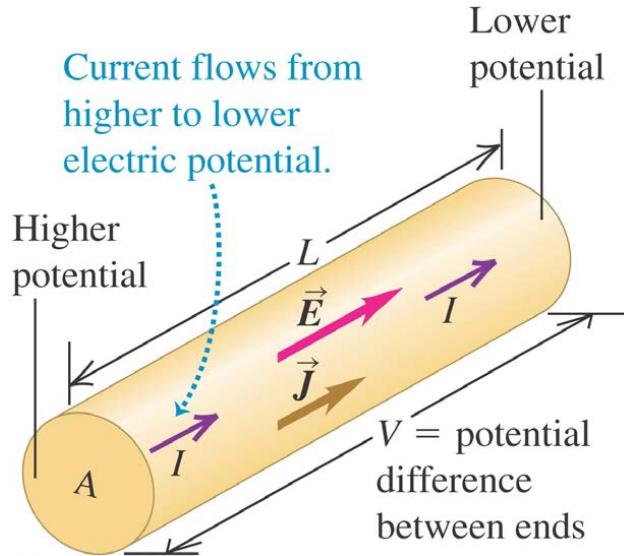
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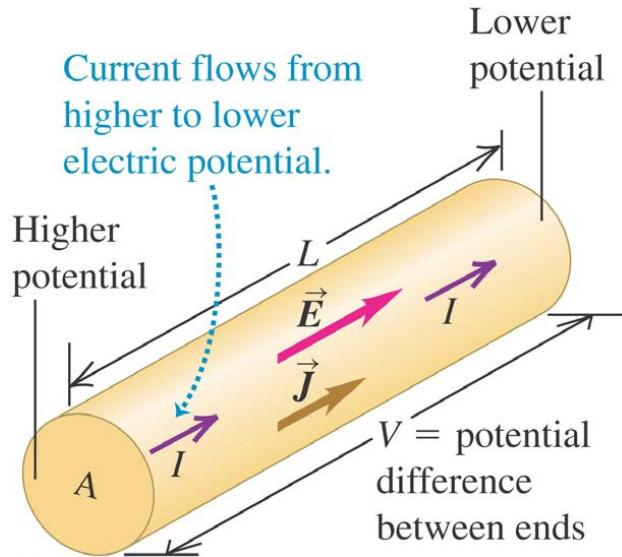
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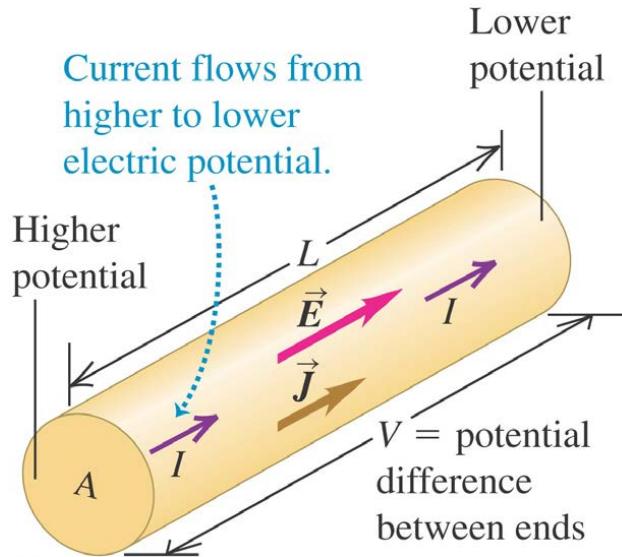


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$$R = \rho L / A$$

The SI unit of resistance is the ohm, equal to one volt per ampere ( $1 \Omega = 1 V/A$ ).  
The kilohm ( $1 k\Omega = 10^3 \Omega$ ) and the megohm ( $1 M\Omega = 10^6 \Omega$ ) are also in common use.

# Resistance ( $R$ )



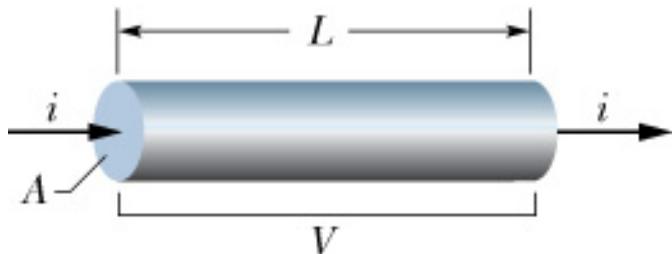
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A 100-m length of 12-gauge copper wire, the size usually used in household wiring, has a resistance at room temperature of about  $0.5 \Omega$ . A 100-W, 120-V light bulb has a resistance (at operating temperature) of  $140 \Omega$ . If the same current  $I$  flows in both the copper wire and the light bulb, the potential difference  $V = IR$  is much greater across the light bulb, and much more potential energy is lost per charge in the light bulb. This lost energy is converted by the light bulb filament into light and heat. You don't want your household wiring to glow white-hot, so its resistance is kept low by using wire of low resistivity and large cross-sectional area.

# Calculating resistance, given the resistivity

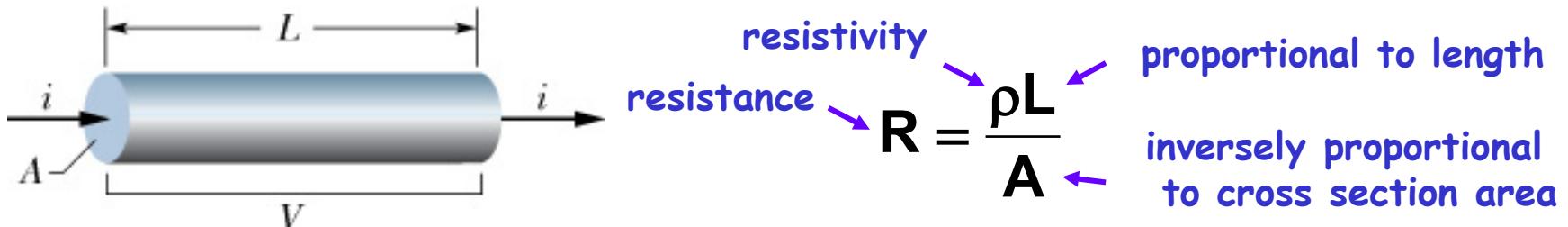

$$R = \frac{\rho L}{A}$$

The diagram shows a cylindrical conductor with length  $L$  and cross-sectional area  $A$ . Current  $i$  flows through the conductor from left to right. The resistivity  $\rho$  is proportional to the length  $L$ , and the resistance  $R$  is inversely proportional to the cross-sectional area  $A$ .

resistivity  $\rho$  is proportional to length  $L$

resistance  $R$  is inversely proportional to cross section area  $A$

# Calculating resistance, given the resistivity



EXAMPLE:

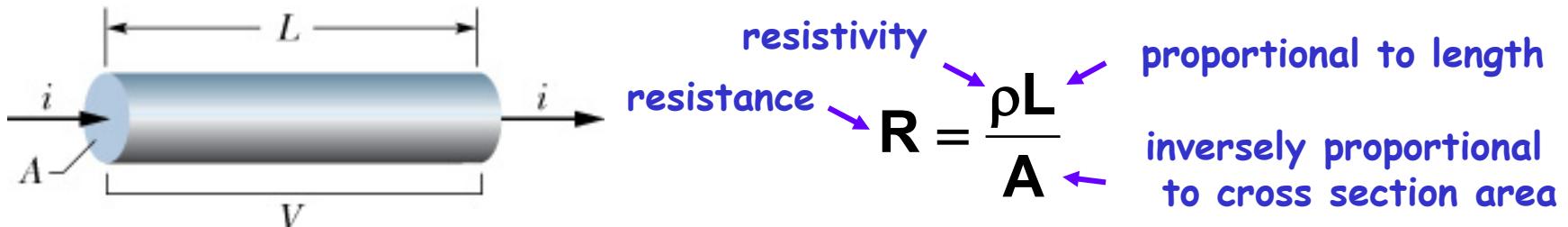
Find  $R$  for a 10 m long iron wire, 1 mm in diameter

$$R = \frac{\rho L}{A} = \frac{9.7 \times 10^{-8} \Omega \cdot \text{m} \times 10 \text{ m}}{\pi \times (10^{-3}/2)^2 \text{ m}^2} = 1.2 \Omega$$

Find the potential difference across  $R$  if  $i = 10 \text{ A}$ . (Amperes)

$$\Delta V = iR = 12 \text{ V}$$

# Calculating resistance, given the resistivity



EXAMPLE:

Find  $R$  for a 10 m long iron wire, 1 mm in diameter

$$R = \frac{\rho L}{A} = \frac{9.7 \times 10^{-8} \Omega \cdot \text{m} \times 10 \text{ m}}{\pi \times (10^{-3}/2)^2 \text{ m}^2} = 1.2 \Omega$$

Find the potential difference across  $R$  if  $i = 10 \text{ A}$ . (Amperes)

$$\Delta V = iR = 12 \text{ V}$$

EXAMPLE:

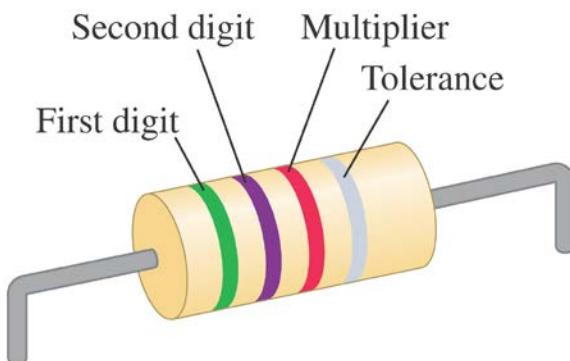
Find resistivity of a wire with  $R = 50 \text{ m}\Omega$ ,

diameter  $d = 1 \text{ mm}$ , length  $L = 2 \text{ m}$

$$\rho = \frac{RA}{L} = \frac{50 \times 10^{-3} \Omega \times 10 \text{ m}}{2 \text{ m}} \times \pi \left(10^{-3}/2\right)^2 = 1.96 \times 10^{-8} \Omega \cdot \text{m}$$

Use a table to identify material. Not Cu or Al, possibly an alloy

# Resistors are color-coded for easy identification



This resistor has a resistance of  $5.7 \text{ k}\Omega$  with a tolerance of  $\pm 10\%$ .

## Color Codes for Resistors

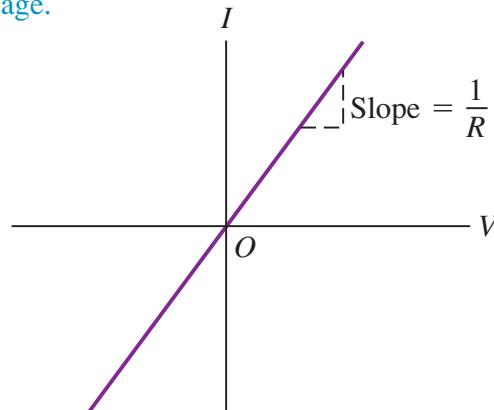
Color	Value as Digit	Value as Multiplier
Black	0	1
Brown	1	$10$
Red	2	$10^2$
Orange	3	$10^3$
Yellow	4	$10^4$
Green	5	$10^5$
Blue	6	$10^6$
Violet	7	$10^7$
Gray	8	$10^8$
White	9	$10^9$

Fourth band: no band means  $\pm 20\%$ , a silver band  $\pm 10\%$ , and a gold band  $\pm 5\%$ .

# Ohmic and nonohmic resistors

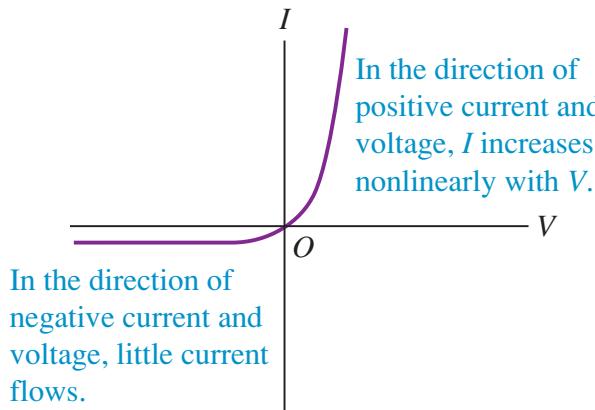
(a)

**Ohmic resistor** (e.g., typical metal wire): At a given temperature, current is proportional to voltage.



(b)

**Semiconductor diode: a nonohmic resistor**



Because the resistivity of a material varies with temperature, the resistance of a specific conductor also varies with temperature. For temperature ranges that are not too great, this variation is approximately a linear relationship,

$$R(T) = R_0[1 + \alpha(T - T_0)]$$

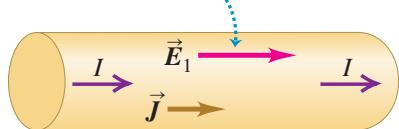
$R_0$  is the resistance at temperature  $T_0$

temperature coefficient of resistance

# Electromotive force ( $\varepsilon$ ) and circuits

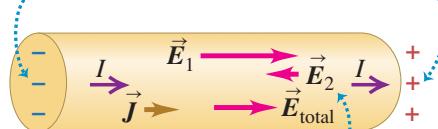
## Open circuit:

- (a) An electric field  $\vec{E}_1$  produced inside an isolated conductor causes a current.



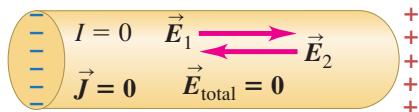
An *electromotive force (emf)* makes current flow. In spite of the name, an emf is *not a force*.

- (b) The current causes charge to build up at the ends.



The charge buildup produces an opposing field  $\vec{E}_2$ , thus reducing the current.

- (c) After a very short time  $\vec{E}_2$  has the same magnitude as  $\vec{E}_1$ ; then the total field is  $\vec{E}_{\text{total}} = \mathbf{0}$  and the current stops completely.

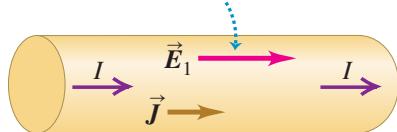


$$\vec{E} = \vec{E}_1 + \vec{E}_2 = \mathbf{0}$$

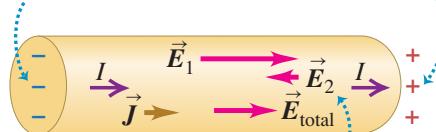
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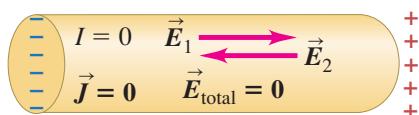


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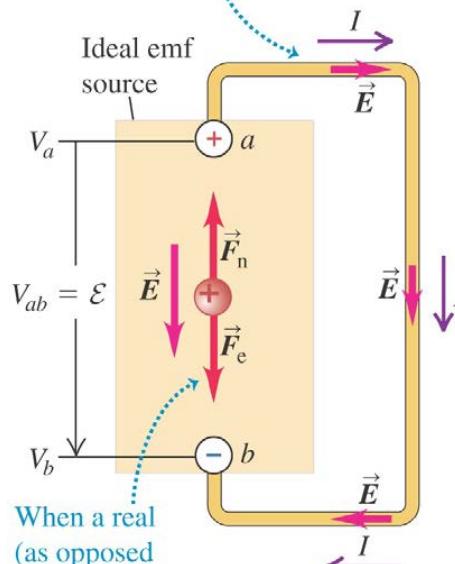


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Potential across terminals creates electric field in circuit, causing charges to move.



When a real (as opposed to ideal) emf source is connected to a circuit,  $V_{ab}$  and thus  $F_e$  fall, so that  $F_n > F_e$  and  $\vec{F}_n$  does work on the charges.

$$\mathcal{E} = V_{ab} = IR$$

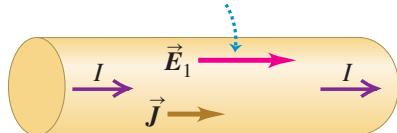


Emf ( $\epsilon$ ) acts like a water pump in a water fountain! 29

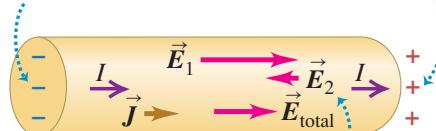
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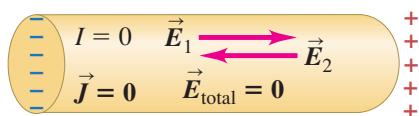


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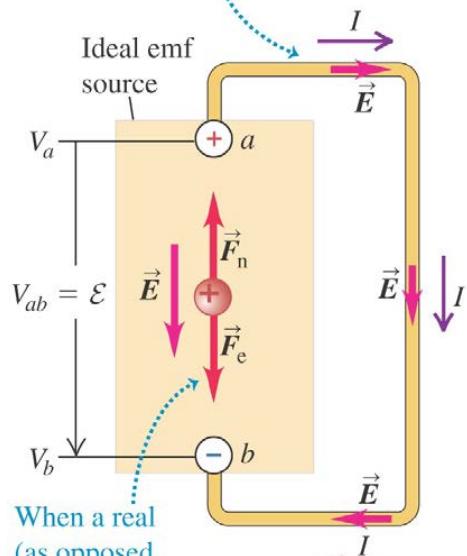


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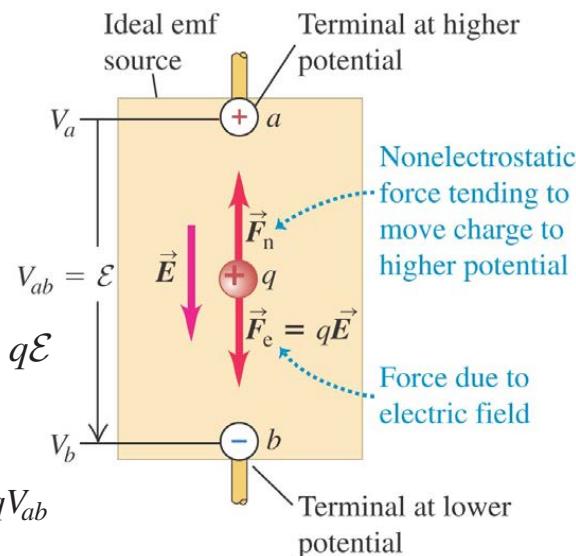
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When the emf source is not part of a closed circuit,  $F_n = F_e$  and there is no net motion of charge between the terminals.

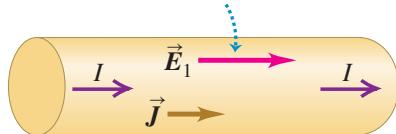


*Emf ( $\varepsilon$ ) acts like a water pump in a water fountain!*

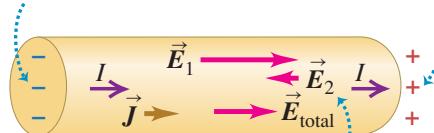
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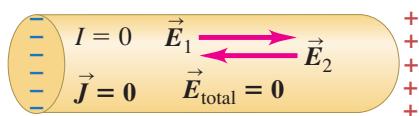


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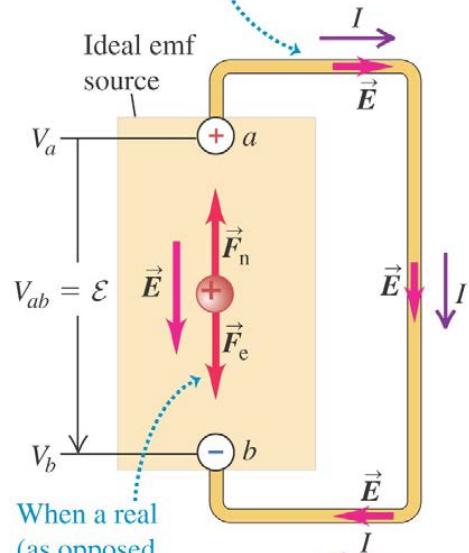
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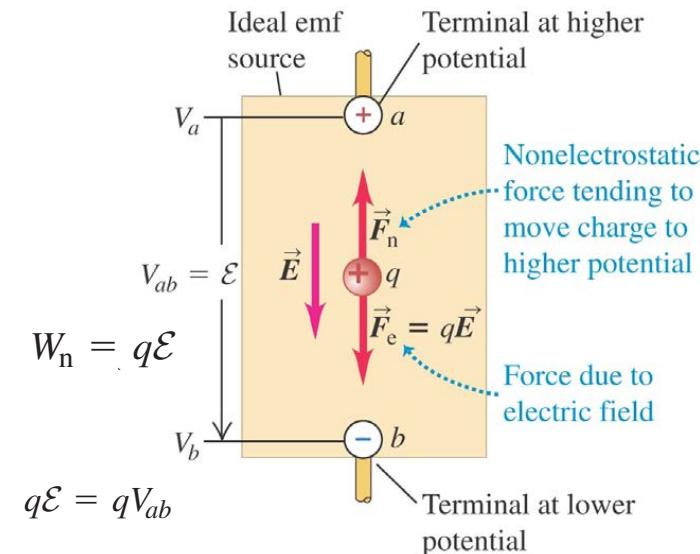
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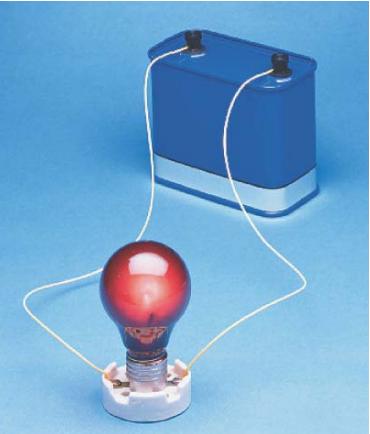
Emf ( $\epsilon$ ) acts like a water pump in a water fountain!

Unit of emf ( $\epsilon$ ) is the same as potential (1 V = 1 J/C)

A battery with an emf of 1.5 V does 1.5 J of work on every coulomb of charge that pass through it

# Internal resistance

Real sources of emf actually contain some *internal resistance*  $r$ .  
The *terminal voltage* of an emf source is  $V_{ab} = \mathcal{E} - Ir$ .



# Internal resistance

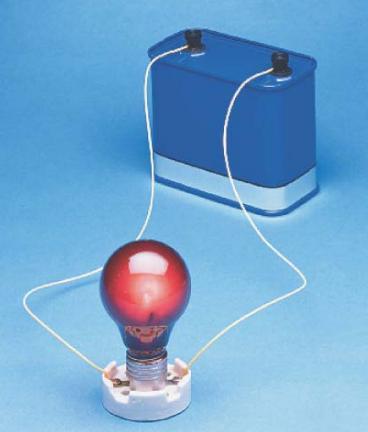
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The terminal voltage of an emf source is  $V_{ab} = \mathcal{E} - Ir$ .

The terminal voltage of the 12-V battery shown at the right is less than 12 V when it is connected to the light bulb.

The potential  $V_{ab}$ , called the **terminal voltage**, is less than the emf  $\mathcal{E}$  because of the term  $Ir$  representing the potential drop across the internal resistance  $r$ .

$$\rightarrow \quad \mathcal{E} - Ir = IR \qquad I = \frac{\mathcal{E}}{R + r}$$



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## Symbols for circuit diagrams



Conductor with negligible resistance



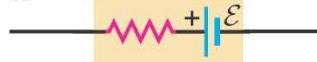
Resistor



Source of emf (longer vertical line always represents the positive terminal, usually the terminal with higher potential)



or



Source of emf with internal resistance  $r$  ( $r$  can be placed on either side)

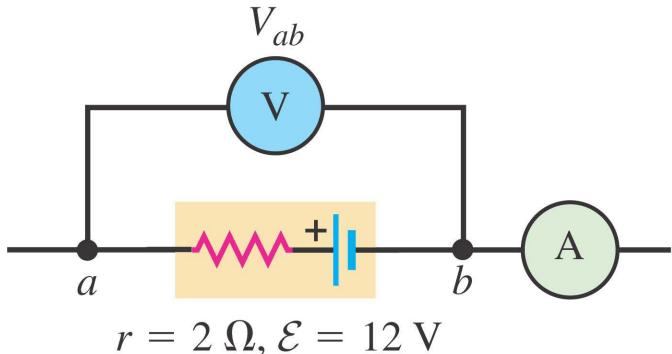


Voltmeter (measures potential difference between its terminals)



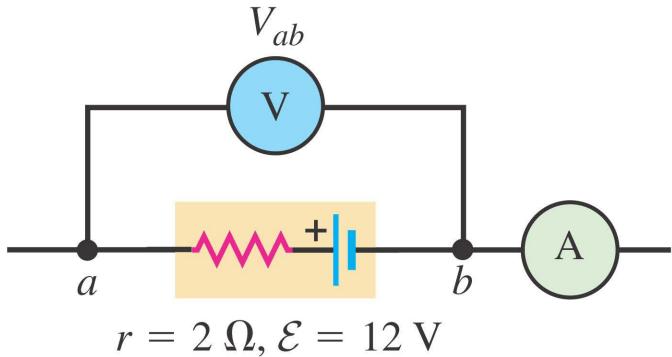
Ammeter (measures current through it)

## A source in an open circuit



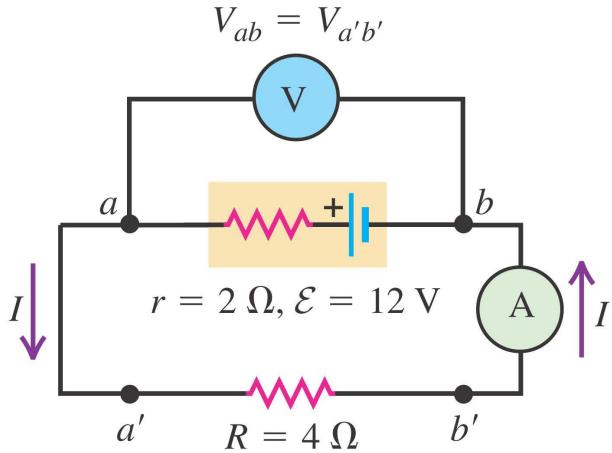
There is **zero current** because there is no complete circuit. Hence the **ammeter reads  $I = 0$** . Because there is no current through the battery, there is no potential difference across its internal resistance. The potential difference  $V_{ab}$  across the battery terminals is equal to the emf. So the voltmeter reads  $V_{ab} = \mathcal{E} = 12 \text{ V}$ .

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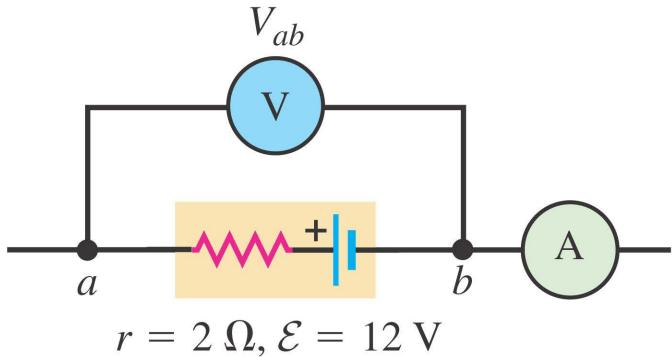
## Source in a complete circuit



The ideal ammeter has zero resistance, so the total resistance external to the source is  $R = 4 \Omega$ . The current through the circuit  $aa'bb'$  is then

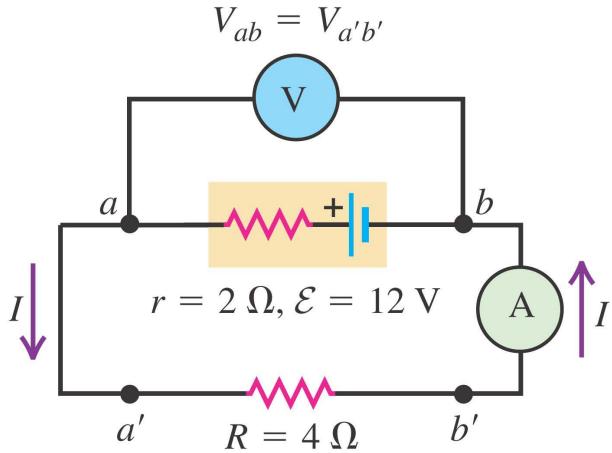
$$I = \frac{\mathcal{E}}{R + r} = \frac{12 \text{ V}}{4 \Omega + 2 \Omega} = 2 \text{ A}$$

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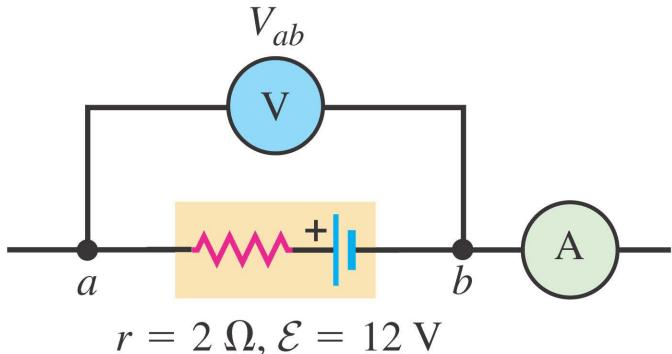
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There is no potential difference between points  $a$  and  $a'$  or between points  $b$  and  $b'$ ; that is,  $V_{ab} = V_{a'b'}$ .

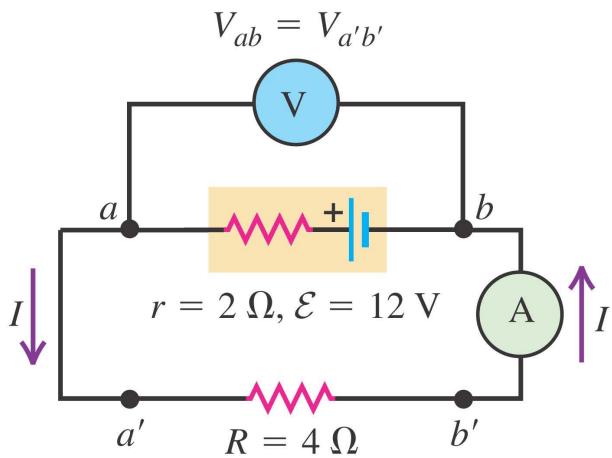
$$V_{a'b'} = IR = (2 \text{ A})(4 \Omega) = 8 \text{ V}$$

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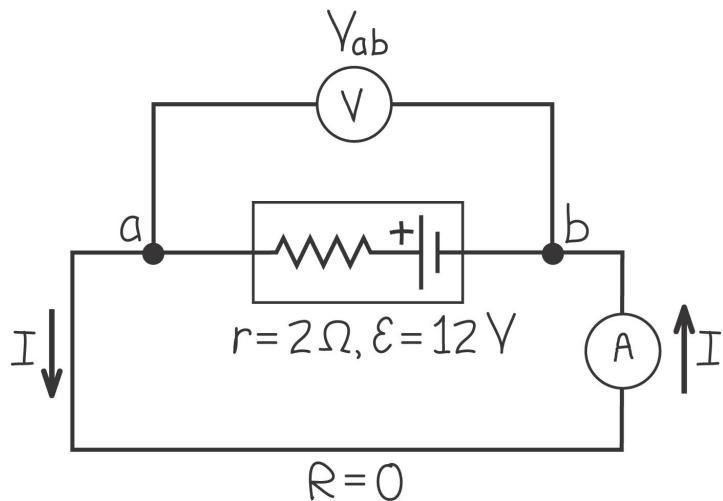
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$$V_{ab} = \mathcal{E} - Ir = 12 \text{ V} - (2 \text{ A})(2 \Omega) = 8 \text{ V}$$

## A source with a short circuit

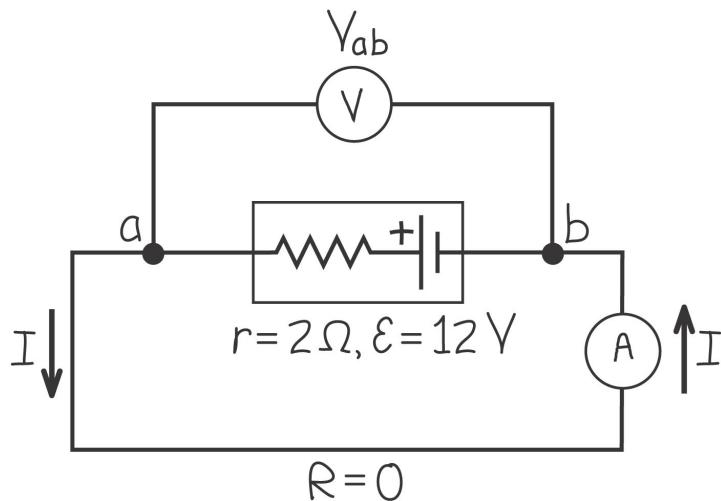


There is now a zero-resistance path between points  $a$  and  $b$ , through the lower loop, so the potential difference between these points must be zero:

$$V_{ab} = IR = I(0) = 0.$$

$$V_{ab} = \mathcal{E} - Ir = 0 \quad I = \frac{\mathcal{E}}{r} = \frac{12 \text{ V}}{2 \Omega} = 6 \text{ A}$$

## A source with a short circuit

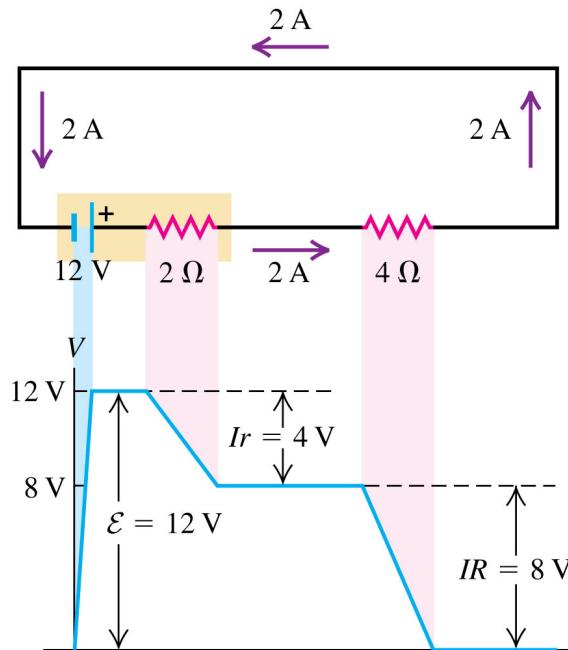


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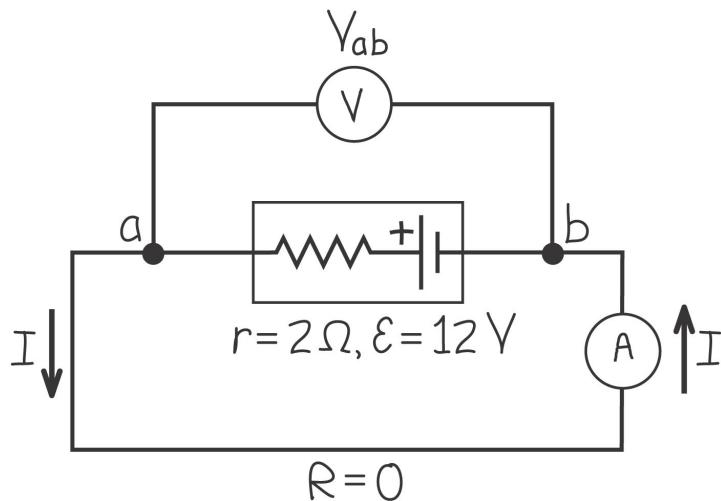
## Potential changes around a circuit



$$\mathcal{E} - Ir - IR = 0$$

A potential gain of  $\mathcal{E}$  is associated with the *emf*, and potential drops of  $Ir$  and  $IR$

## A source with a short circuit

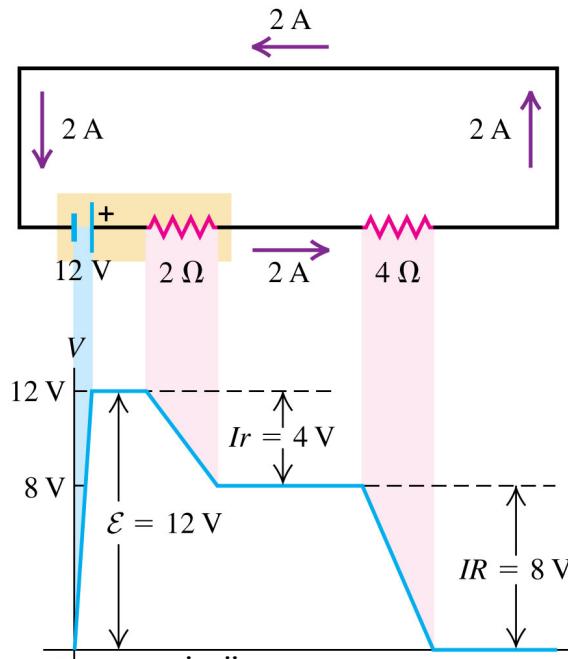


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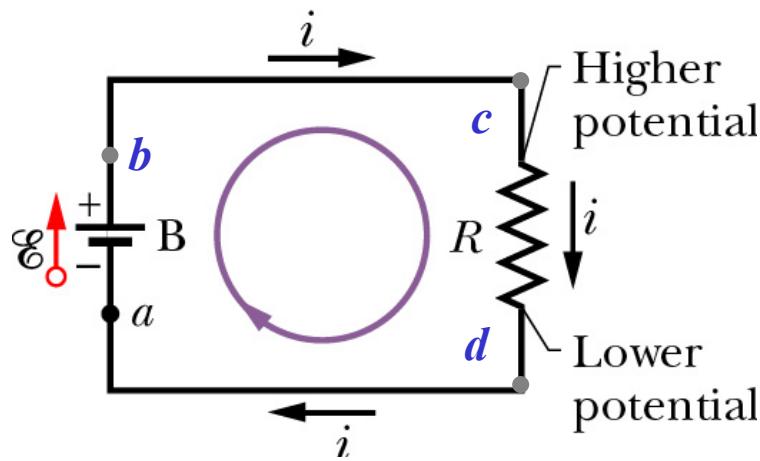
$$\mathcal{E} - Ir - IR = 0$$

A potential gain of  $\mathcal{E}$  is associated with the *emf*, and potential drops of  $Ir$  and  $IR$

The principal difference between a **fresh flashlight** battery and an **old one** is NOT in the *emf*, which decreases only slightly with use, but in the **internal resistance (r)**, which may increase from less than an ohm when the battery is fresh to as much as  $1000 \Omega$  or more after long use.

# Example: CW or CCW around a single-loop circuit

Assume current direction as shown



- **Traverse clockwise from a:**

$$\Delta V_{ba} = V_b - V_a = +\epsilon$$

$$\Delta V_{cb} = 0$$

$$\Delta V_{dc} = V_d - V_c = -iR$$

$$\Delta V_{ad} = 0$$

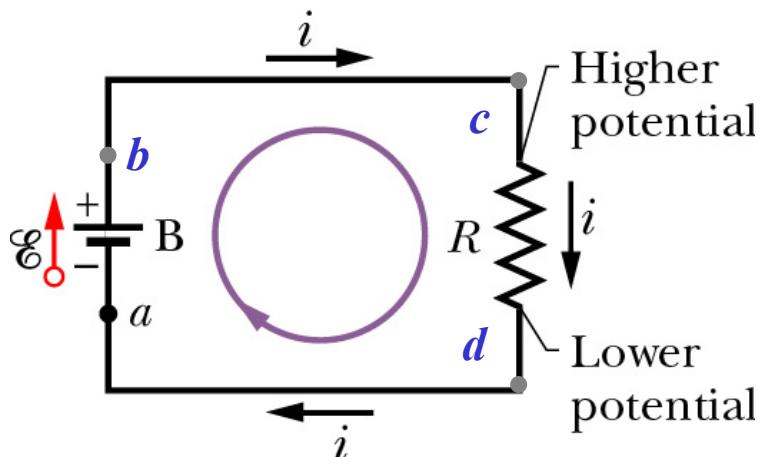
$$\epsilon - iR = 0$$

$$i = \frac{\epsilon}{R}$$

$$\sum_{\text{closed loop}} \Delta V = 0 = \epsilon + 0 - iR + 0$$

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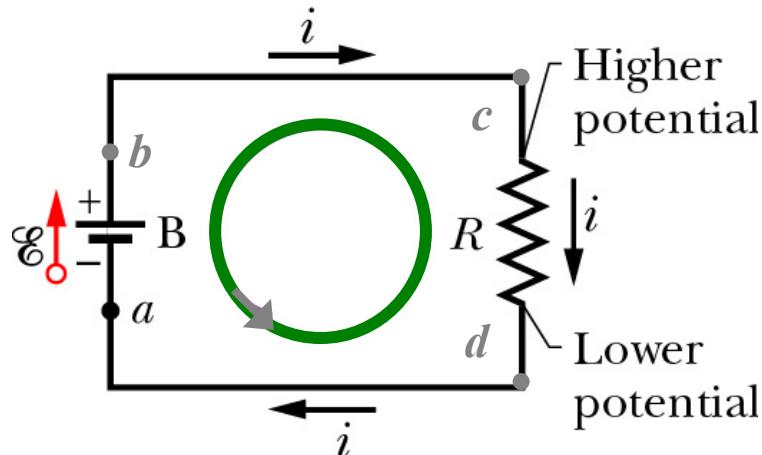
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$$\epsilon - iR = 0$$

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$$\sum_{\text{closed loop}} \Delta V = 0 = \epsilon + 0 - iR + 0$$



- Traverse counterclockwise from a:

$$\Delta V_{da} = 0$$

$$\Delta V_{cd} = V_c - V_d = +iR$$

$$\Delta V_{bc} = 0$$

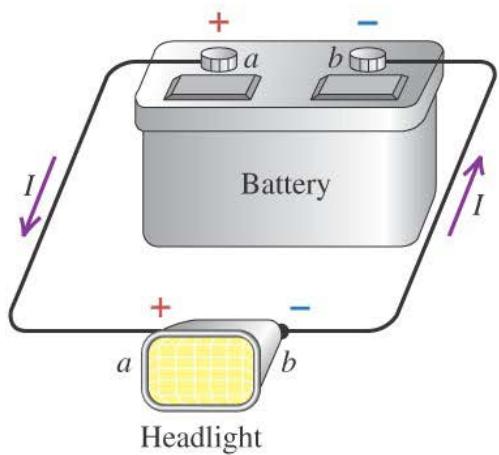
$$\Delta V_{ab} = V_a - V_b = -\epsilon$$

$$\sum_{\text{closed loop}} \Delta V = 0 = 0 + iR + 0 - \epsilon$$

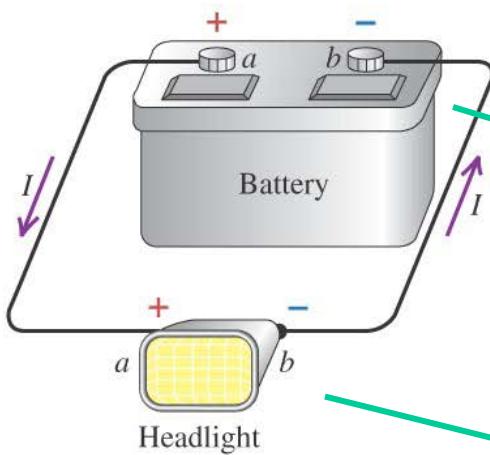
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Same result

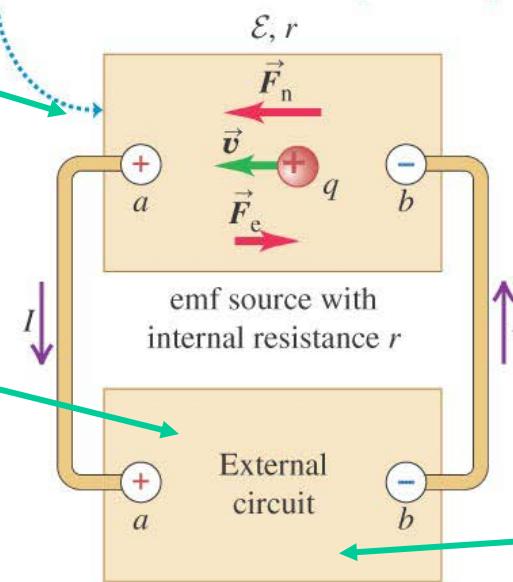
# Energy and power in electric circuits



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- The emf source converts nonelectrical to electrical energy at a rate  $\mathcal{E}I$ .
- Its internal resistance *dissipates* energy at a rate  $I^2r$ .
- The difference  $\mathcal{E}I - I^2r$  is its power output.



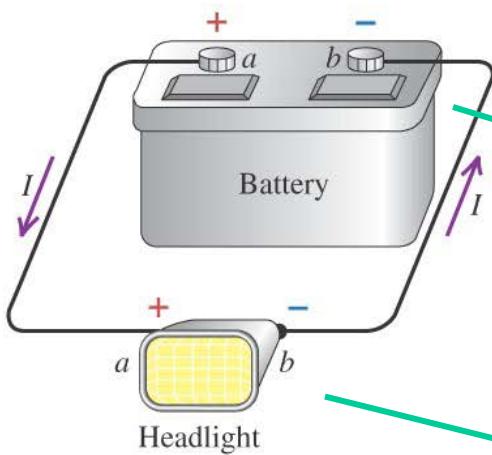
Chemical energy converted to electrical energy

Internal resistance  $R$

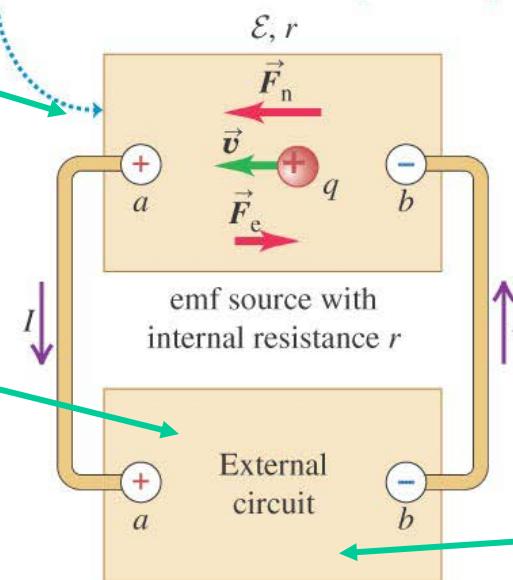
The power input to the circuit element between  $a$  and  $b$  is

$$P = (V_a - V_b)I = V_{ab}I$$

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Chemical energy converted to electrical energy

Internal resistance  $R$

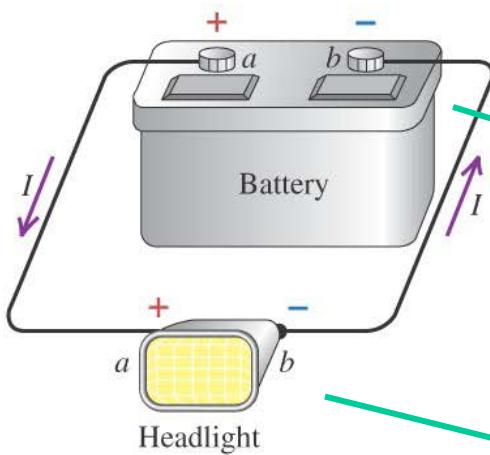
The power input to the circuit element between  $a$  and  $b$  is

$$P = (V_a - V_b)I = V_{ab}I$$

$$P = V_{ab}I \quad (\text{rate at which energy is delivered to or extracted from a circuit element})$$

$$1 \text{ W} = 1 \text{ J/s} = (1 \text{ J/C}) \times (1 \text{ C/s})$$

# Energy and power in electric circuits

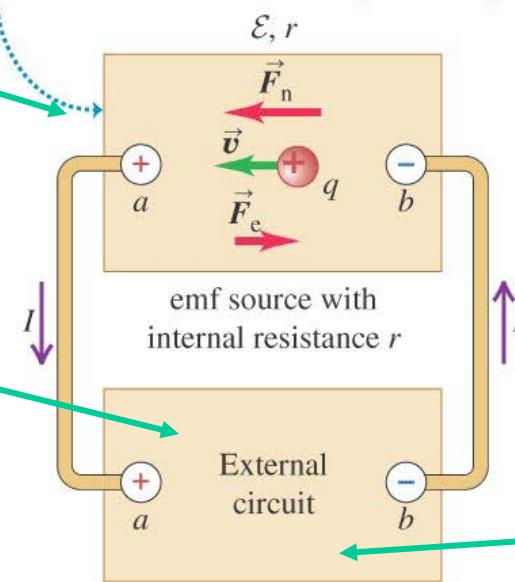


$$V_{ab} = IR$$

$$P = V_{ab}I = I^2R = \frac{V_{ab}^2}{R}$$

(power delivered to a resistor)

- The emf source converts nonelectrical to electrical energy at a rate  $\mathcal{E}I$ .
- Its internal resistance *dissipates* energy at a rate  $I^2r$ .
- The difference  $\mathcal{E}I - I^2r$  is its power output.



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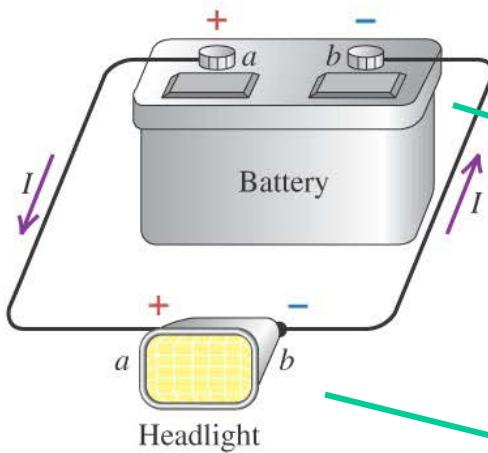
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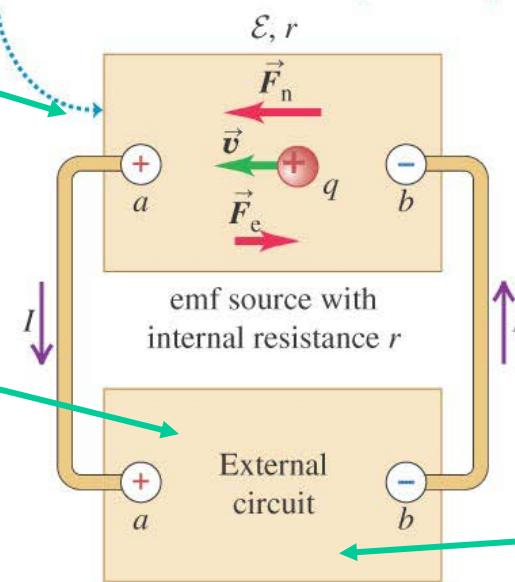
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## Power Output of a Source

$$V_{ab} = \mathcal{E} - Ir$$

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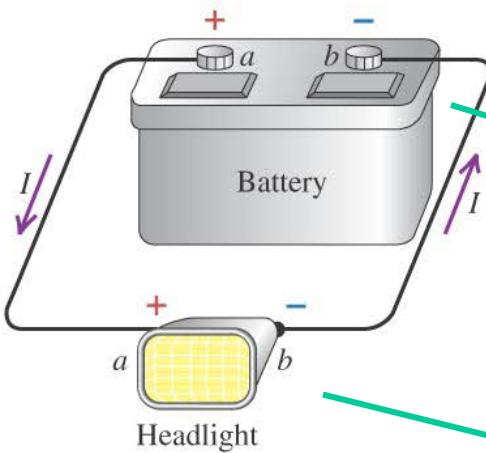
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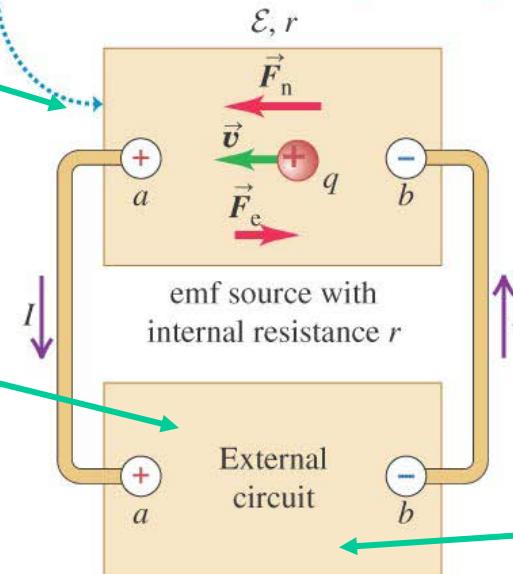
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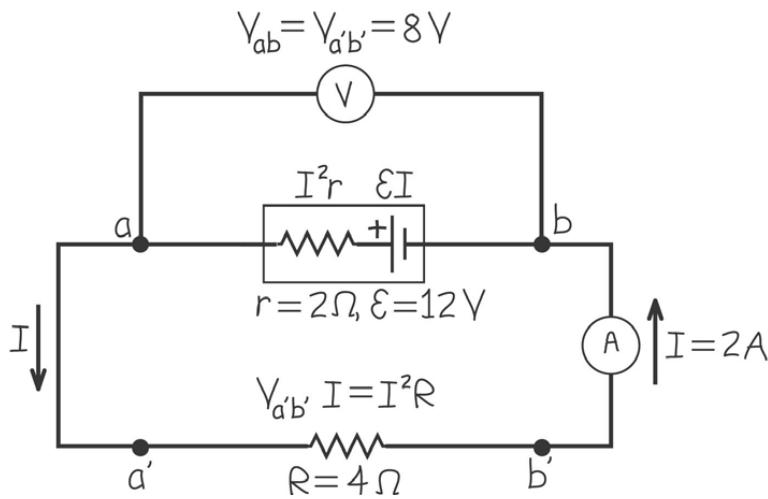
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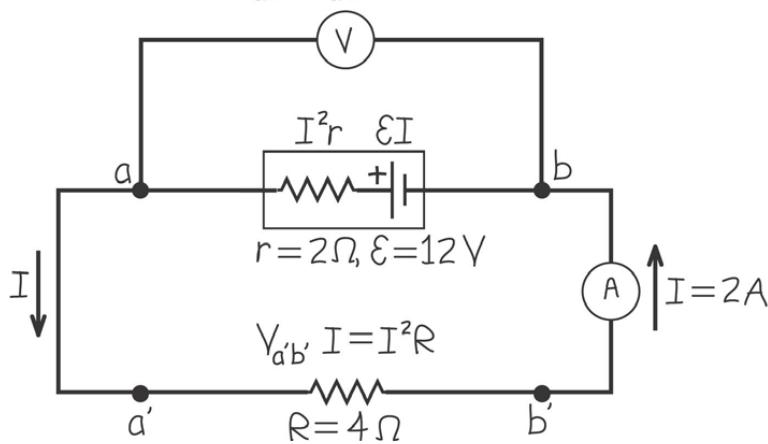
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# Power input and output



# Power input and output

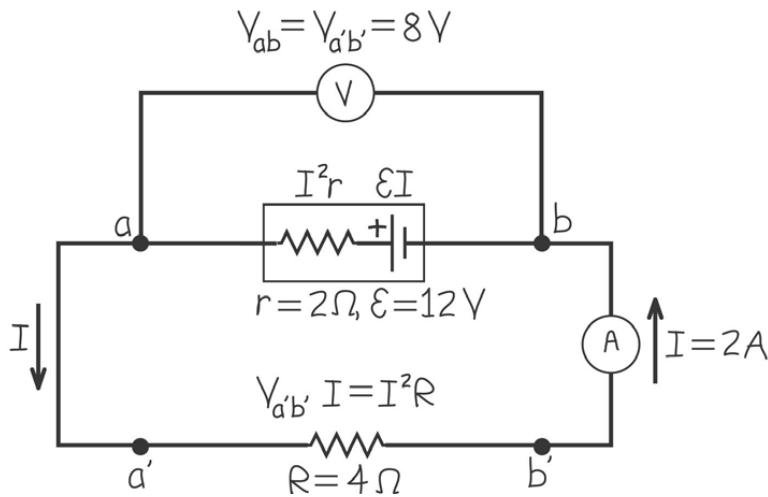
$$V_{ab} = V_{a'b'} = 8 \text{ V}$$



the rate of energy conversion (chemical to electrical) in the battery is

$$\epsilon I = (12 \text{ V})(2 \text{ A}) = 24 \text{ W}$$

# Power input and output



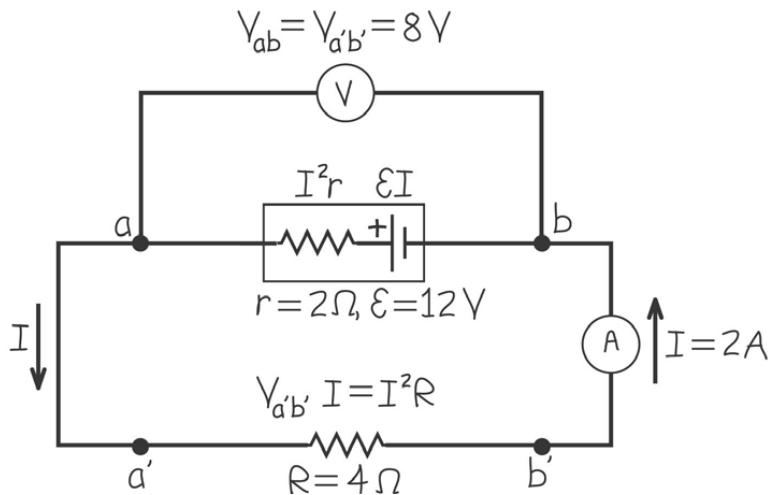
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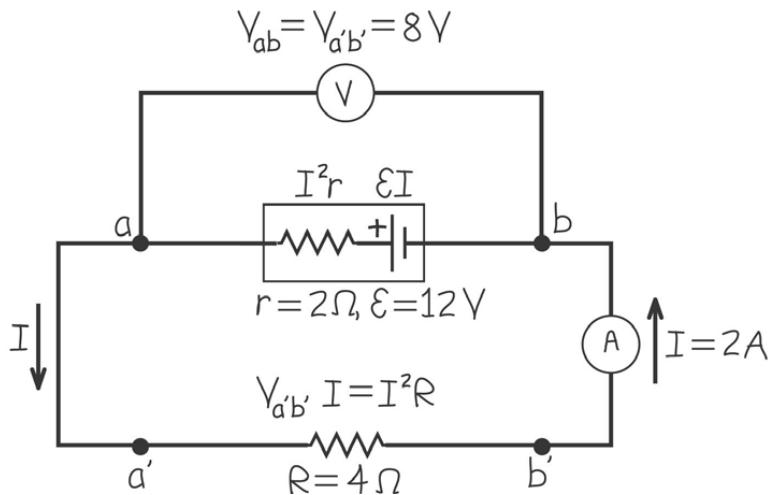
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$$\epsilon I - I^2 r = 16 \text{ W}$$

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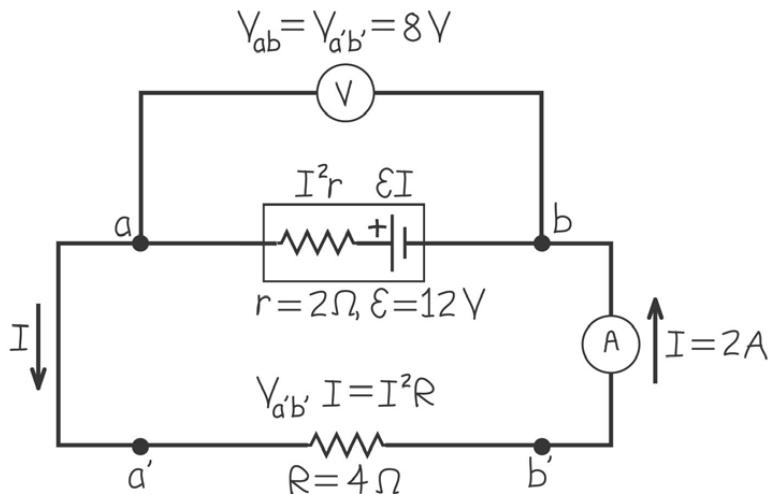
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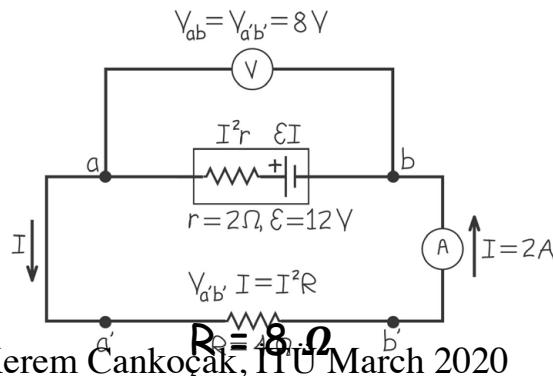
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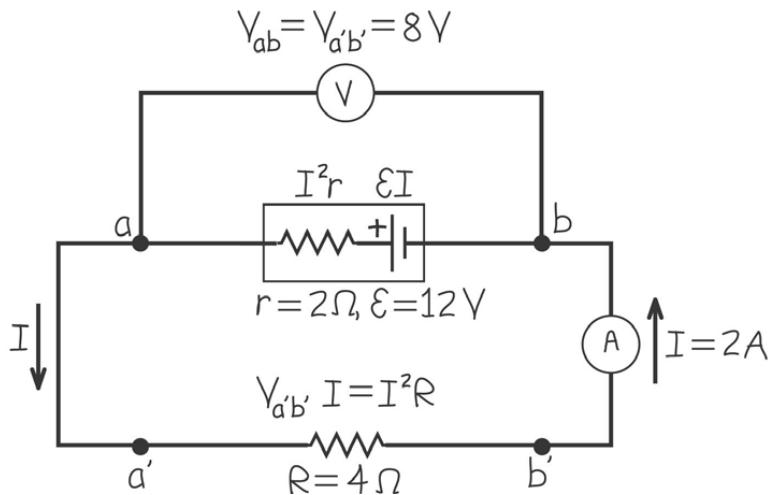
Increasing the resistance ( $4 \Omega \rightarrow 8 \Omega$ )



$$I = \frac{\epsilon}{R + r} = \frac{12 \text{ V}}{8 \Omega + 2 \Omega} = 1.2 \text{ A} \quad \text{Current decreases}$$

$$V_{ab} = IR = (1.2 \text{ A})(8 \Omega) = 9.6 \text{ V} \quad \text{Potential diff. increases}$$

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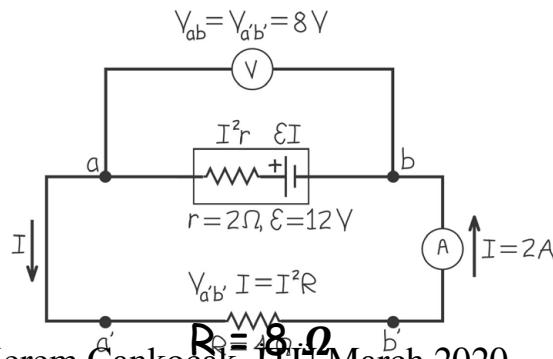
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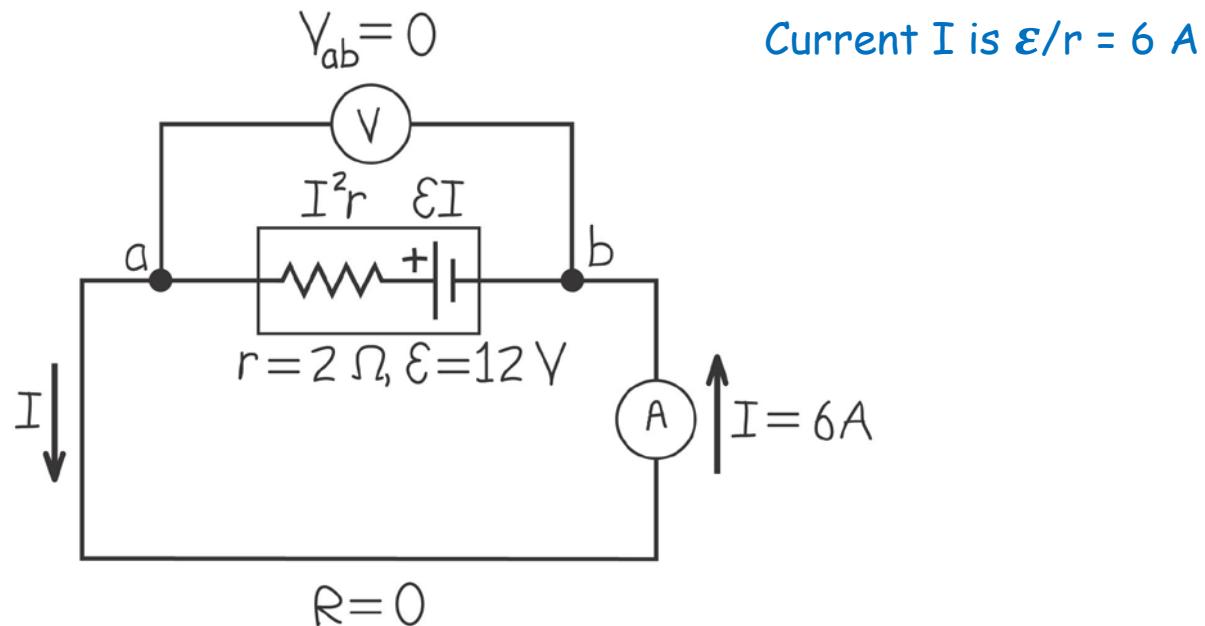
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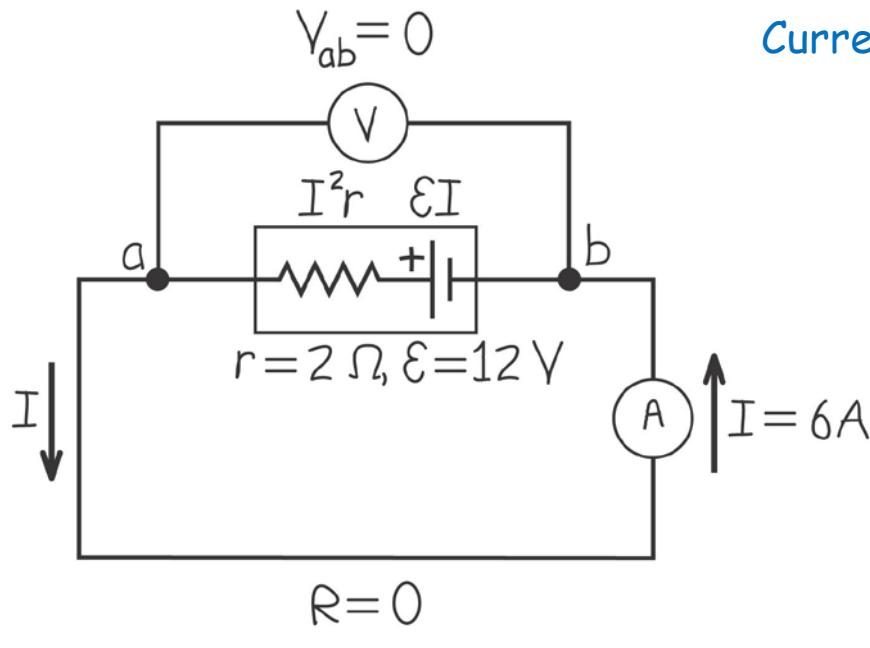
$$P = I^2 R = (1.2 \text{ A})^2 (8 \Omega) = 12 \text{ W}$$

Power dissipated in  $R$  decreases (50 W light bulb has a greater resistance than a 100 W light bulb)

# Power in a short circuit



# Power in a short circuit



Current  $I$  is  $\mathcal{E}/r = 6 \text{ A}$

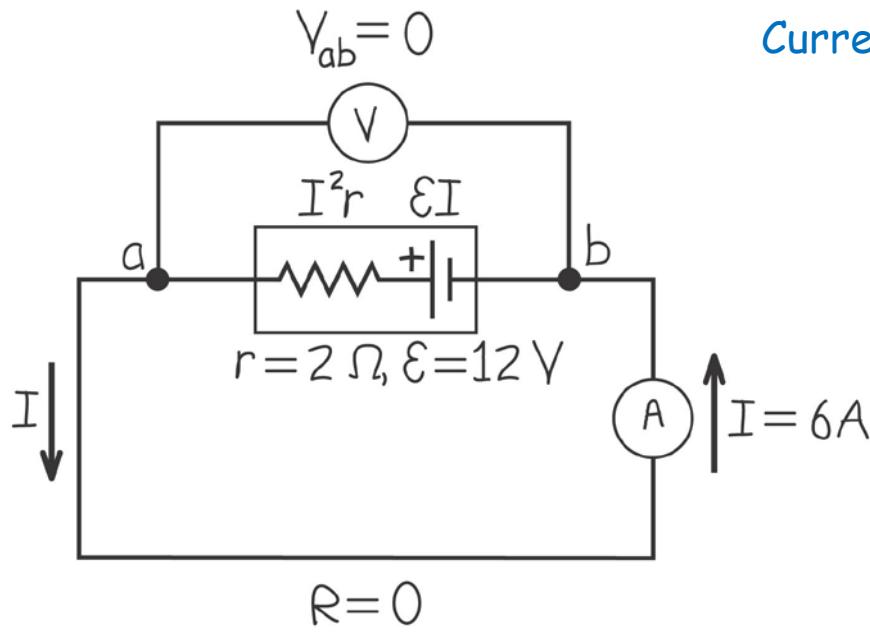
the rate of energy conversion (chemical to electrical) in the battery is then

$$\mathcal{E}I = (12 \text{ V})(6 \text{ A}) = 72 \text{ W}$$

the rate of dissipation of energy in the battery is

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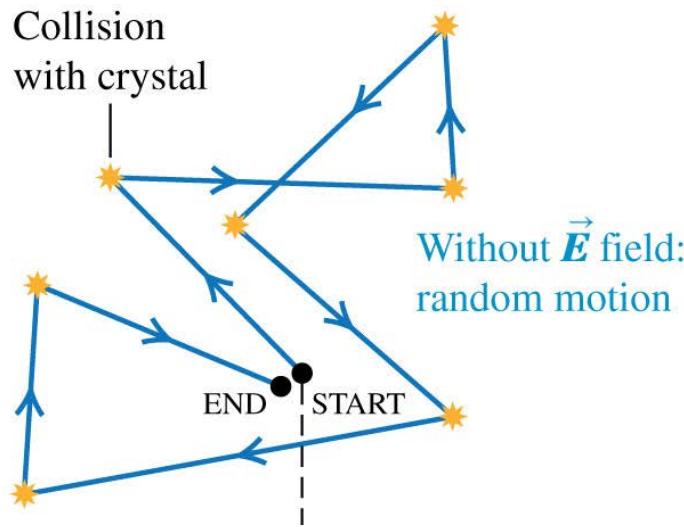
$$I^2 r = (6 \text{ A})^2 (2 \Omega) = 72 \text{ W}$$

The net power is  $\epsilon I - I^2 r = 0$

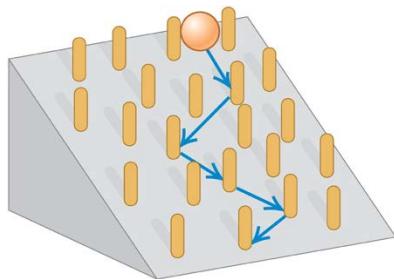
Same result from  $P = V_{ab}I = 0$  (since  $V_{ab} = 0$ )

# Theory of metallic conduction

(a)



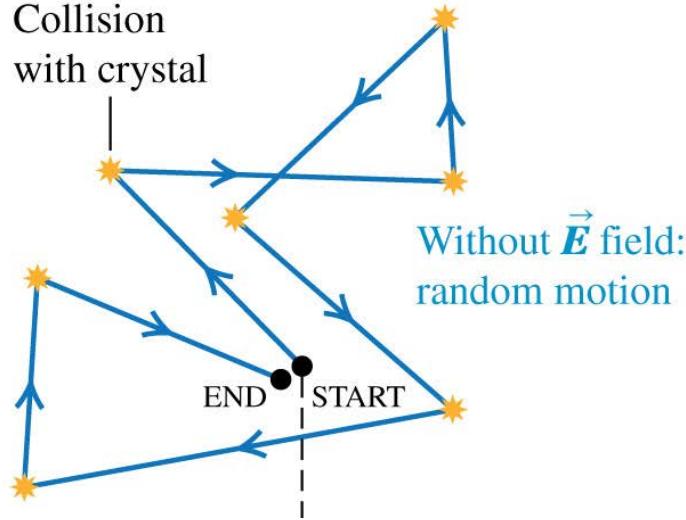
Average speed of random motion is  
 $10^6$  m/s



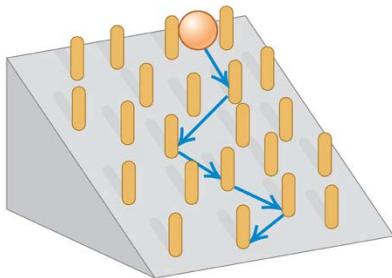
If there is no electric field, the electrons move in straight lines between collisions, the directions of their velocities are random, and on average net displacement is **ZERO**.

# Theory of metallic conduction

(a) Collision with crystal

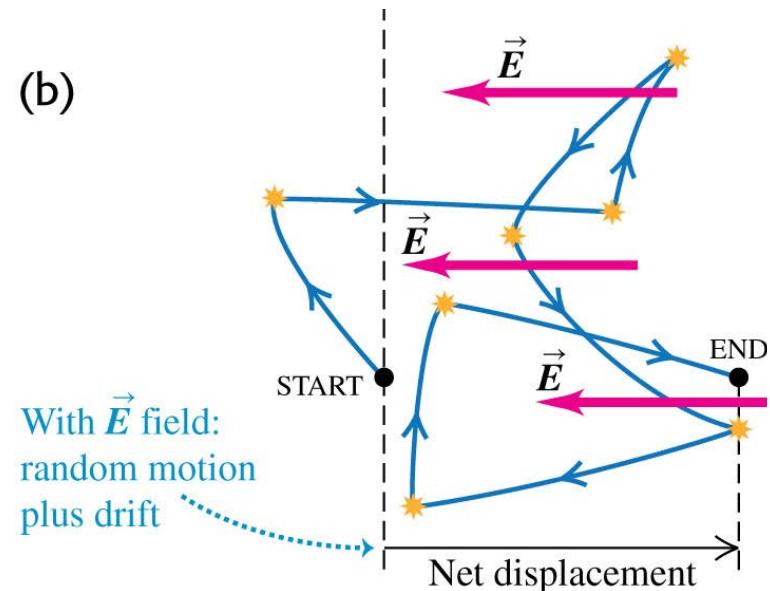


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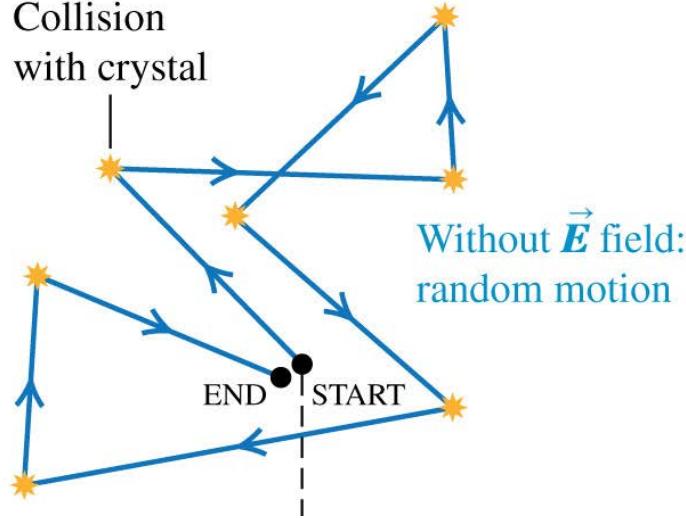
(b)



Average drift speed is  $10^{-4}$  m/s

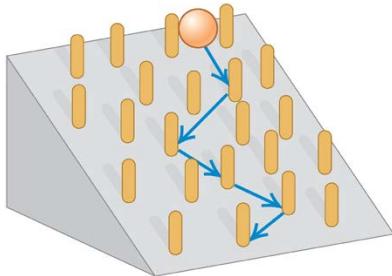
# Theory of metallic conduction

(a) Collision with crystal



Without  $\vec{E}$  field:  
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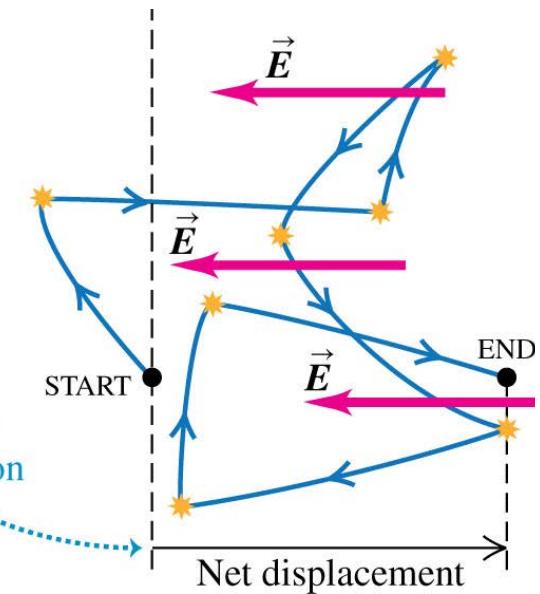
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With  $\vec{E}$  field:  
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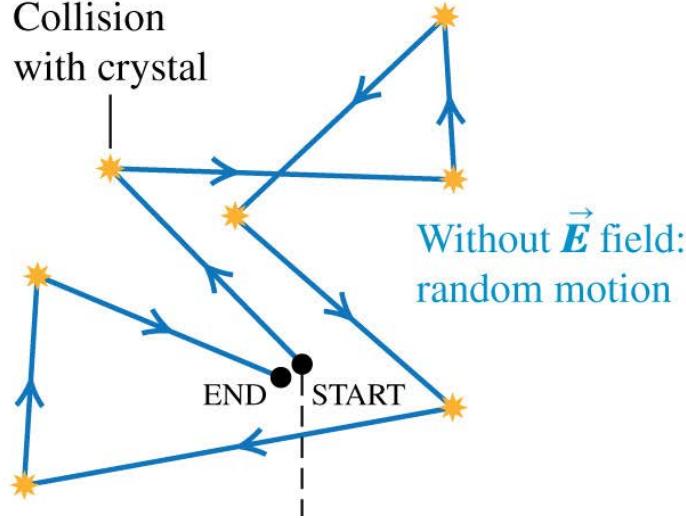


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The average time between collisions is called the mean free time, denoted by  $\tau$

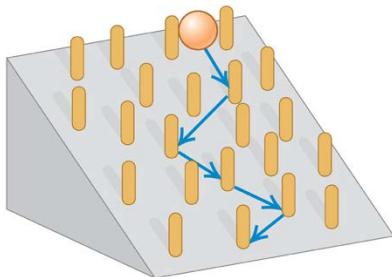
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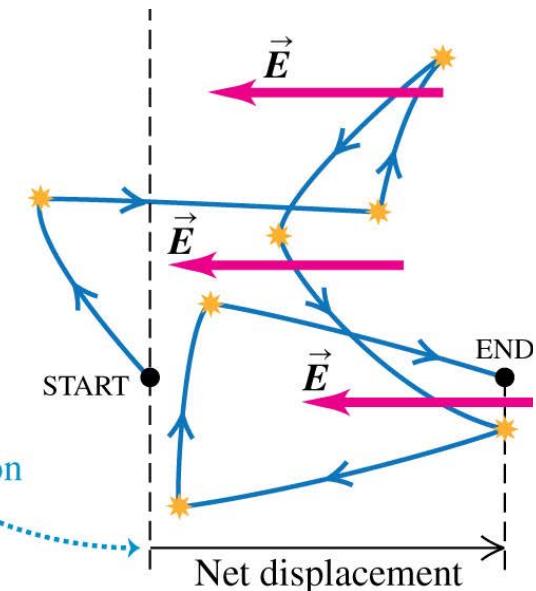
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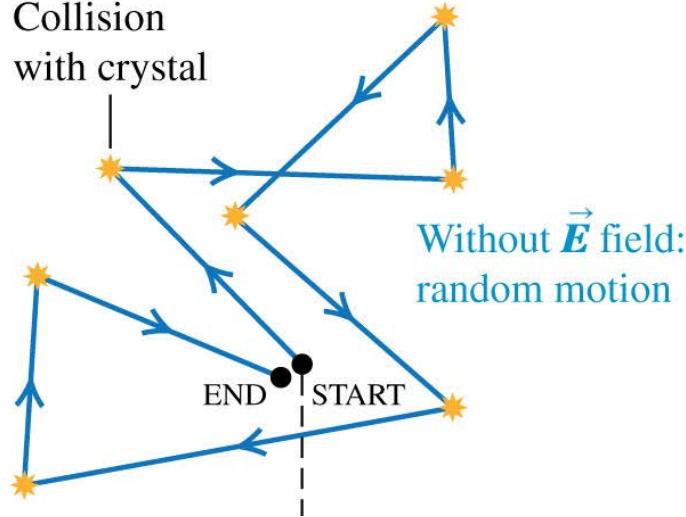
Resistivity

and current density:

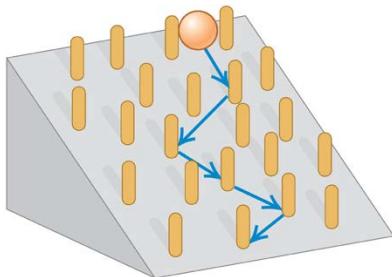
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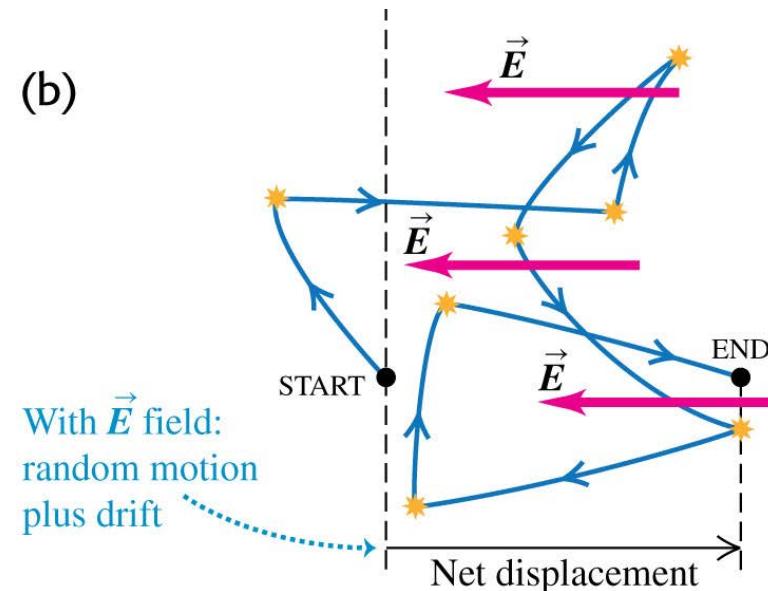


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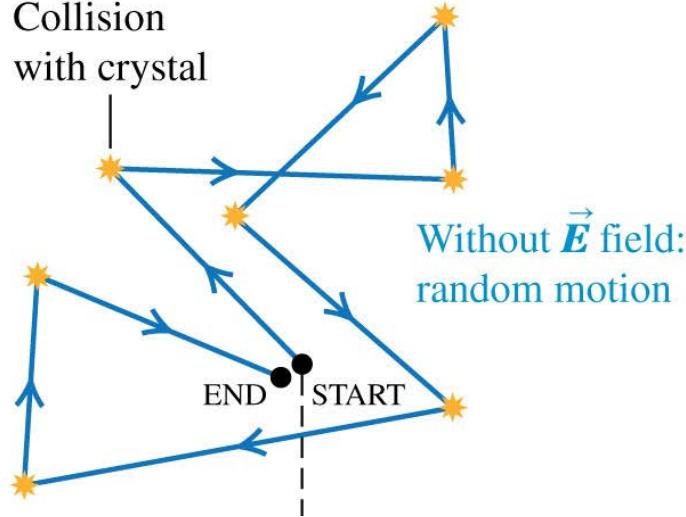
$$\vec{v}_d = \vec{a}\tau = \frac{q\tau}{m}\vec{E}$$

$$\vec{J} = nq\vec{v}_d = \frac{nq^2\tau}{m}\vec{E}$$

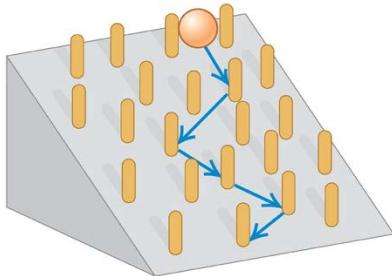
$$\rho = \frac{m}{ne^2\tau}$$

# Theory of metallic conduction

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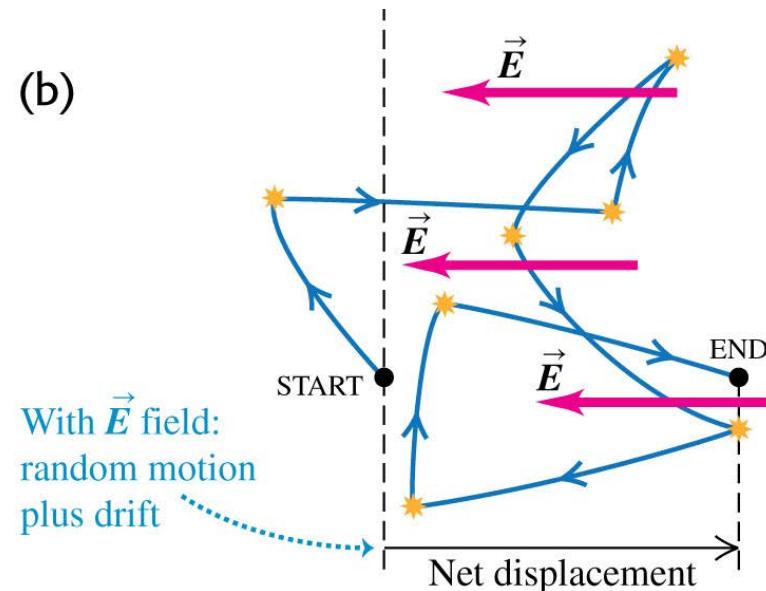


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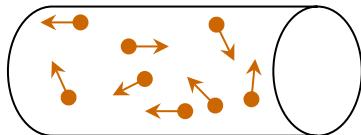
$$\vec{J} = nq\vec{v}_d = \frac{nq^2\tau}{m}\vec{E}$$

$$\rho = \frac{m}{ne^2\tau}$$

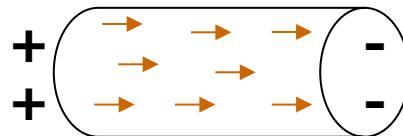
If  $n$  and  $\tau$  are independent of  $E$ , then the resistivity is independent of  $E$  and the conducting material obeys Ohm's law.

# What makes current flow?

conductivity  $\sigma = 1/\rho$ .



APPLIED FIELD = 0  
Random motion  
flow left = flow right

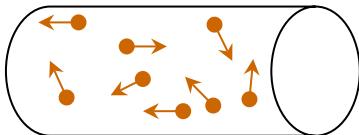


APPLIED FIELD NOT ZERO  
Moving charges collide with fixed ions  
and flow with **drift velocity**

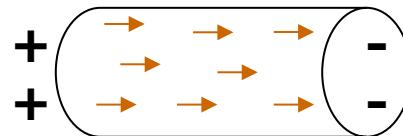
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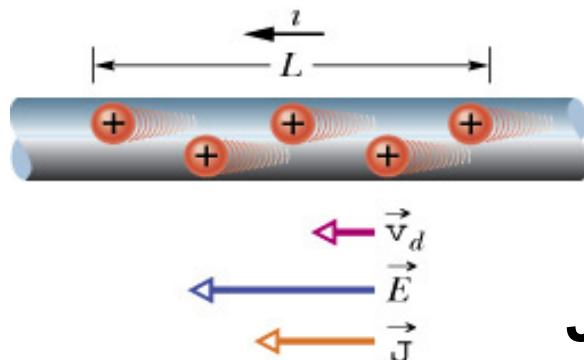


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APPLIED FIELD NOT ZERO  
Moving charges collide with fixed ions  
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$$\mathbf{J} = \sigma \mathbf{E}$$



- $n \equiv \text{density of charge carriers }$  Units : # /volume
- $nv_D = \# \text{ of charge carriers crossing unit area per unit time}$

$$\mathbf{J} = qn\mathbf{v}_D = \text{net charge crossing area A per unit time}$$