

Chp 29: Electromagnetic Induction - (II)

Lenz's Law

Lenz's law is a convenient alternative method for determining the direction of an induced current or emf. Lenz's law is not an independent principle; it can be derived from Faraday's law.

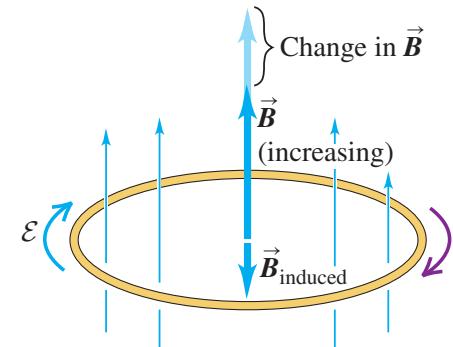
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Lenz's Law : The direction of any magnetic induction effect is such as to oppose the cause of the effect.

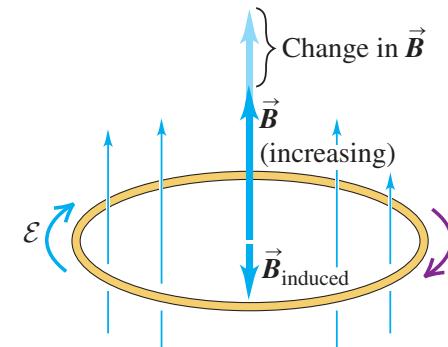


Lenz's Law

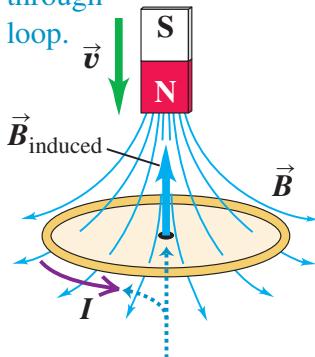
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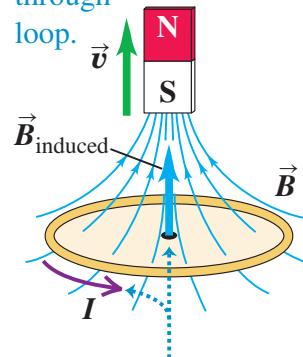
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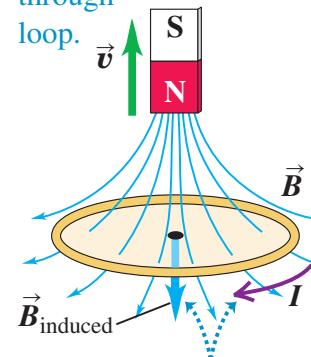
(a) Motion of magnet causes *increasing downward flux* through loop.



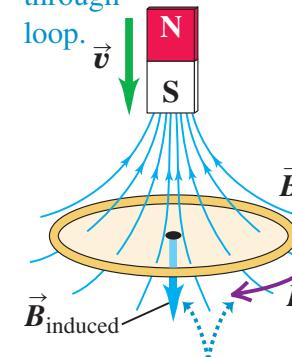
(b) Motion of magnet causes *decreasing upward flux* through loop.



(c) Motion of magnet causes *decreasing downward flux* through loop.



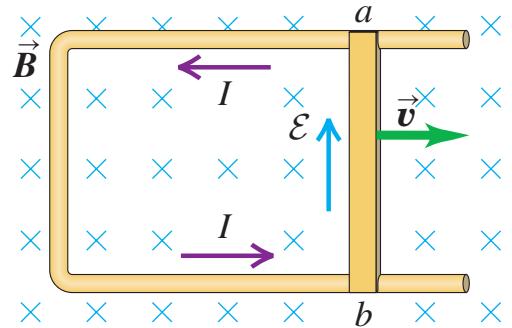
(d) Motion of magnet causes *increasing upward flux* through loop.



The induced magnetic field is *upward* to oppose the flux change. To produce this induced field, the induced current must be *counterclockwise* as seen from above the loop.

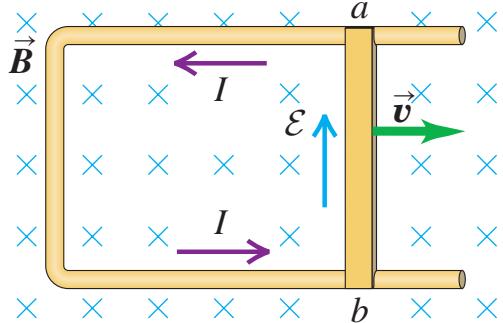
The induced magnetic field is *downward* to oppose the flux change. To produce this induced field, the induced current must be *clockwise* as seen from above the loop.

Motional electromotive force

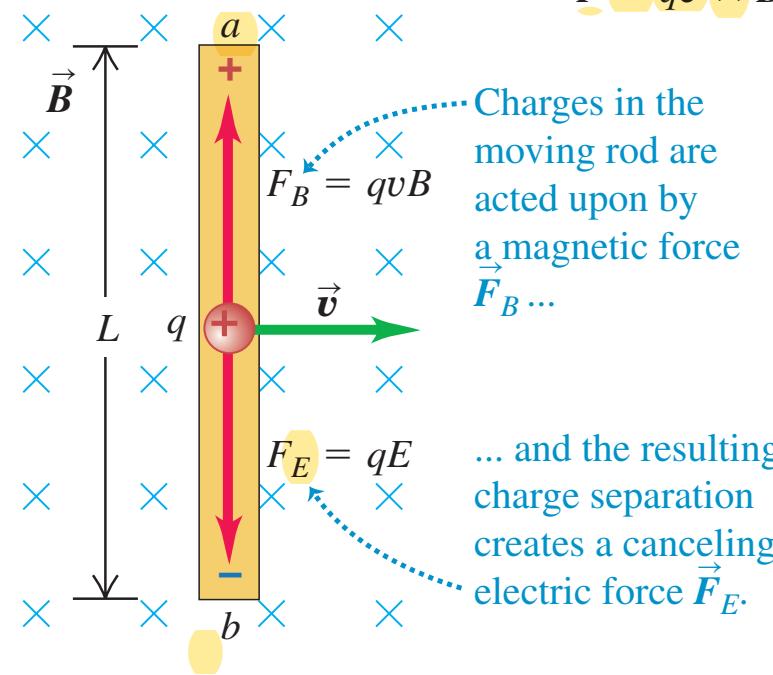


Motional electromotive force

$$\vec{F} = q\vec{v} \times \vec{B}$$



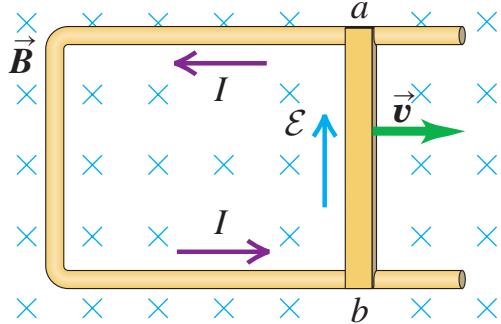
This magnetic force causes the free charges in the rod to move, creating an excess of positive charge at the upper end *a* and negative charge at the lower end *b*. This in turn creates an electric field *E* within the rod, in the direction from *a* toward *b* (opposite to the magnetic force).
 ➔ $qE = qvB$ and the charges are in equilibrium.



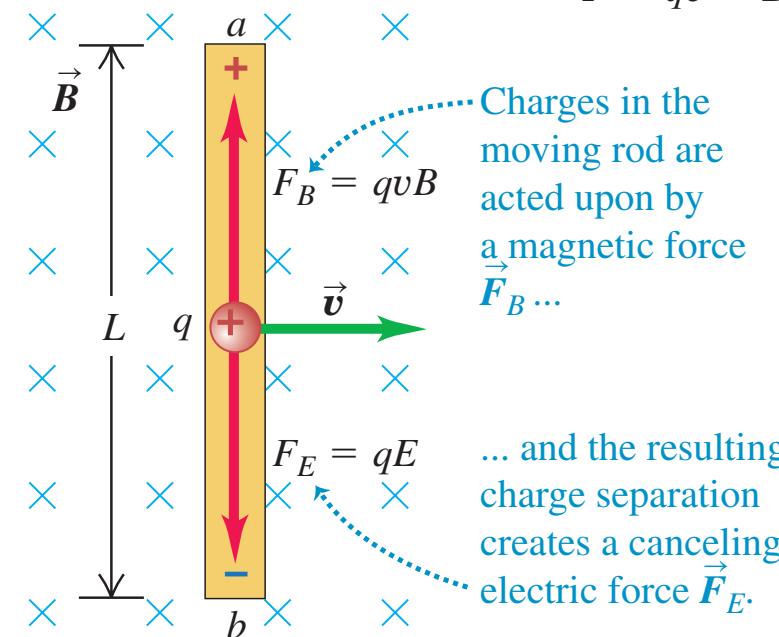
Charges in the moving rod are acted upon by a magnetic force \vec{F}_B ...

... and the resulting charge separation creates a canceling electric force \vec{F}_E .

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This magnetic force causes the free charges in the rod to move, creating an excess of positive charge at the upper end a and negative charge at the lower end b . This in turn creates an electric field E within the rod, in the direction from a toward b (opposite to the magnetic force).
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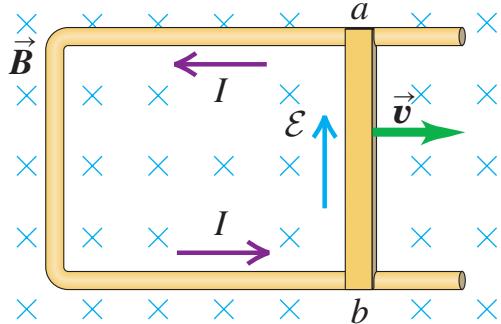
Charges in the moving rod are acted upon by a magnetic force $\vec{F}_B \dots$

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The magnitude of the potential difference $V_{ab} = V_a - V_b$ is equal to the electric-field magnitude E multiplied by the length L of the rod. From the above discussion, $E = vB$

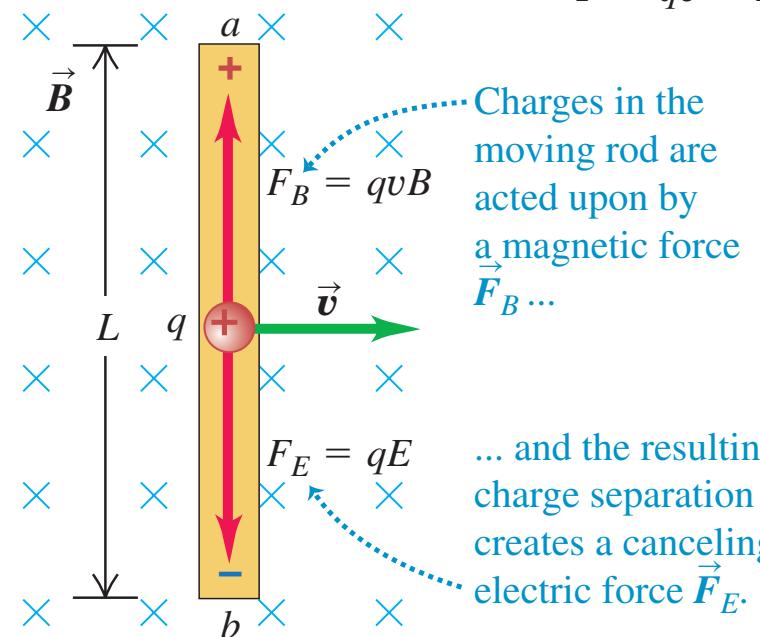
$$V_{ab} = EL = vBL$$

Motional electromotive force



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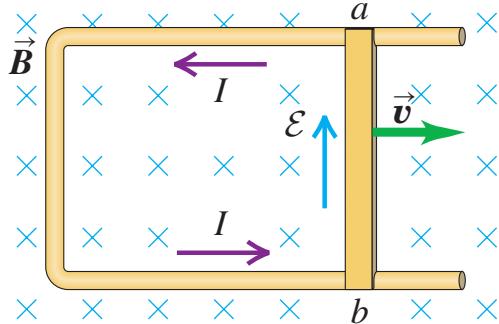
$$V_{ab} = EL = vBL$$

$$\mathcal{E} = vBL$$

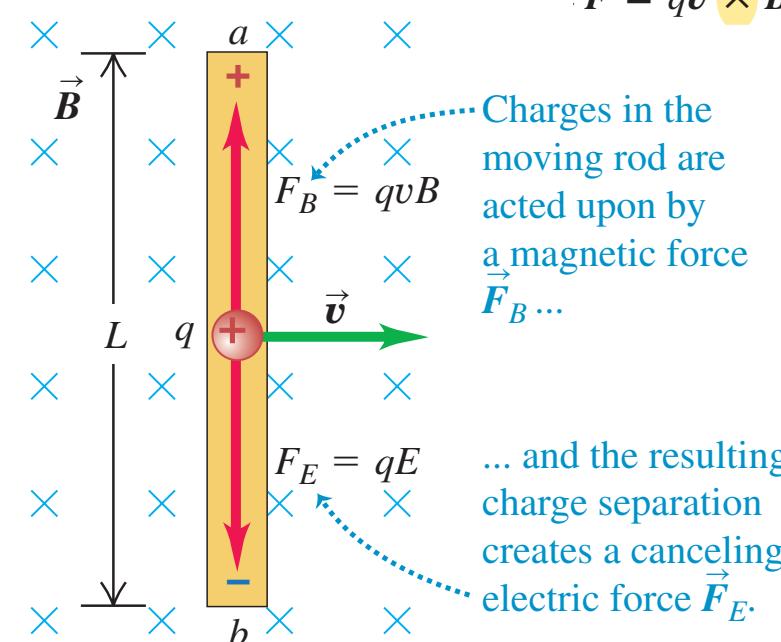
(motional emf; length and velocity perpendicular to uniform \vec{B})

$$1 \text{ V} = 1 \text{ J/C}$$

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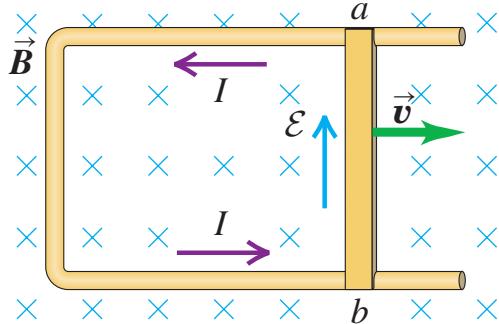
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$$d\mathcal{E} = (\vec{v} \times \vec{B}) \cdot d\vec{l}$$

$$\mathcal{E} = \oint (\vec{v} \times \vec{B}) \cdot d\vec{l} \quad (\text{motional emf; closed conducting loop})$$

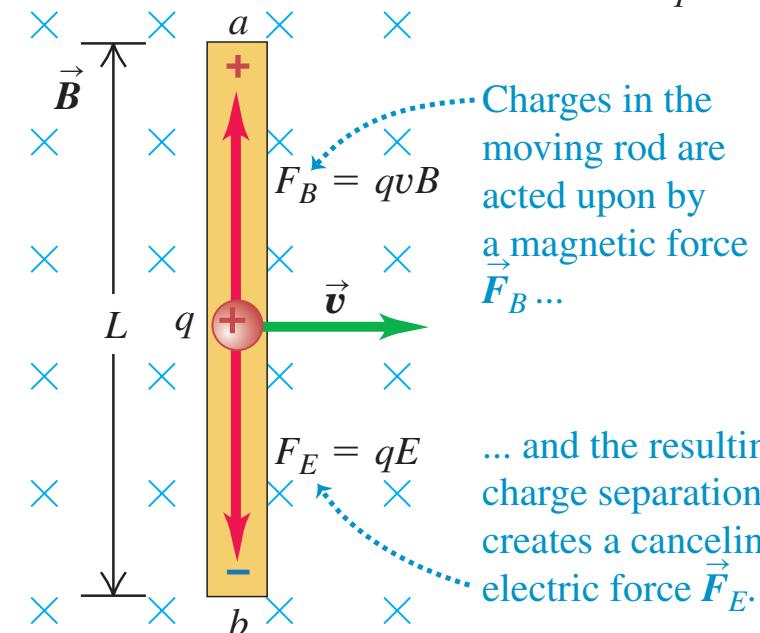
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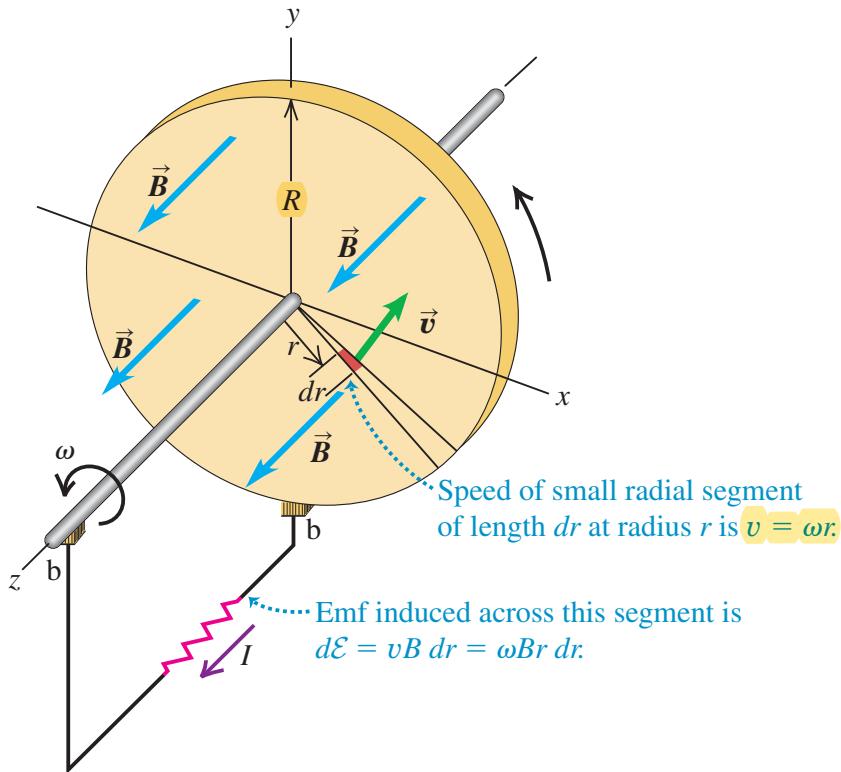
$$\mathcal{E} = \oint (\vec{v} \times \vec{B}) \cdot d\vec{l} \quad (\text{motional emf; closed conducting loop})$$

Only for moving conductors!

$$\mathcal{E} = -d\Phi_B/dt$$

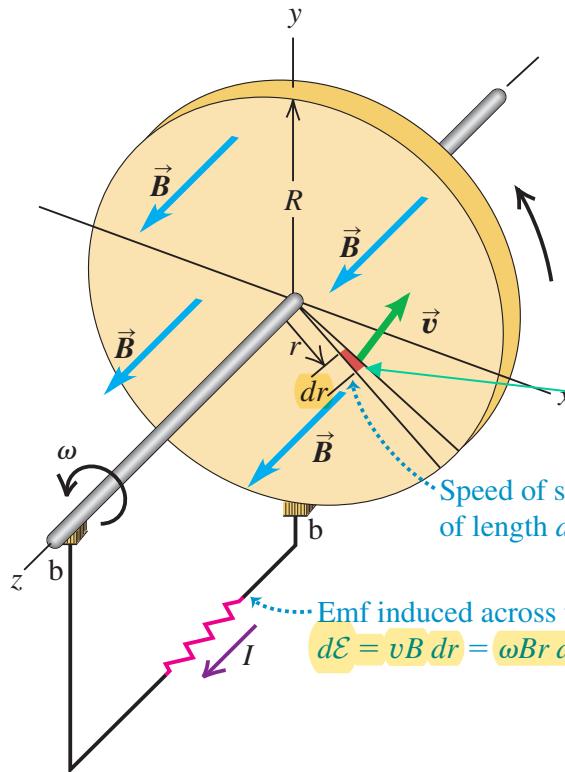
(stationary conductors:
Faraday's law)

The Faraday disk dynamo



a conducting disk with radius R that lies in the xy -plane and rotates with constant angular velocity v about the z -axis. The disk is in a uniform, constant B field in the z -direction. Find the induced emf between the center and the rim of the disk.

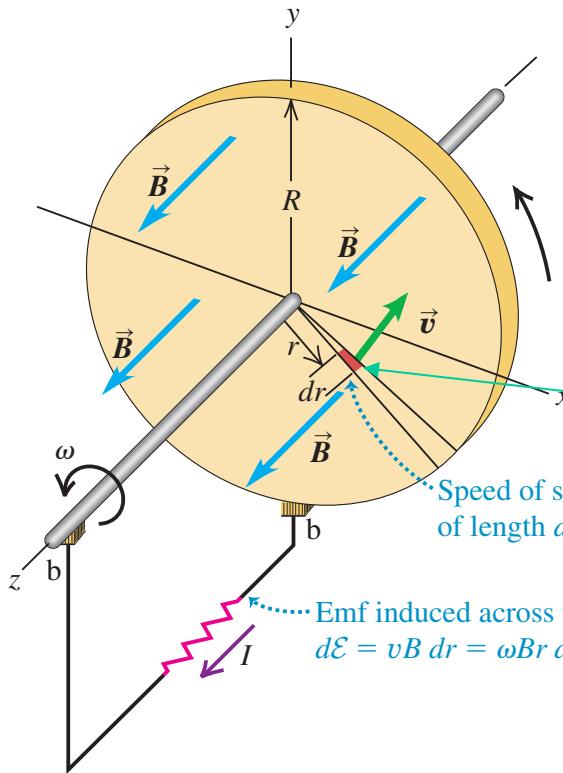
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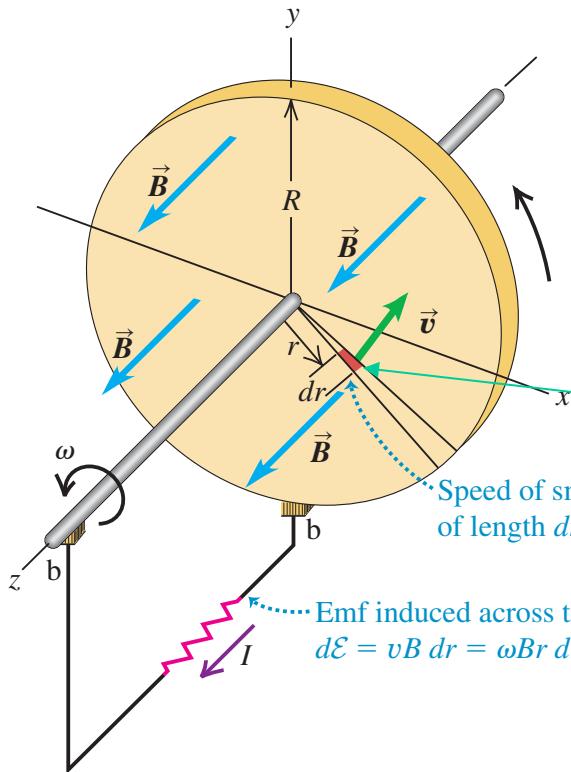


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$$\Rightarrow d\mathcal{E} = \omega Br dr$$

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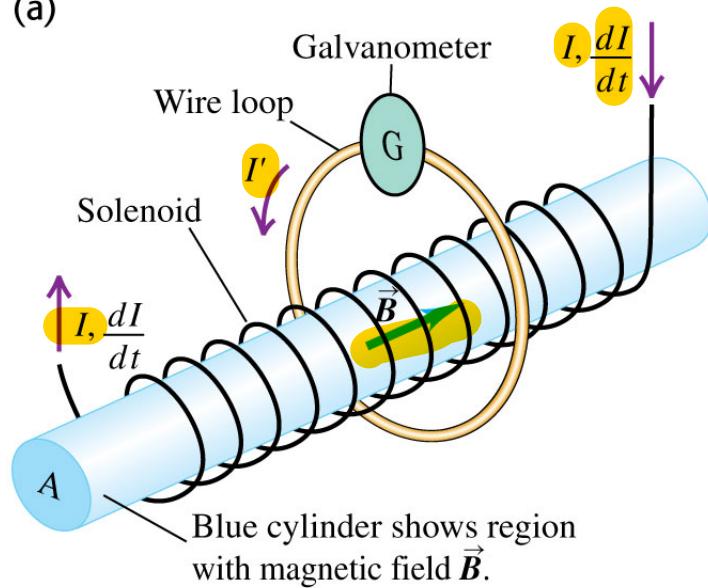
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$$d\mathcal{E} = (\vec{v} \times \vec{B}) \cdot d\vec{l}$$

$$\Rightarrow d\mathcal{E} = \omega Br dr \quad \mathcal{E} = \int_0^R \omega Br dr = \frac{1}{2} \omega BR^2$$

Induced electric fields

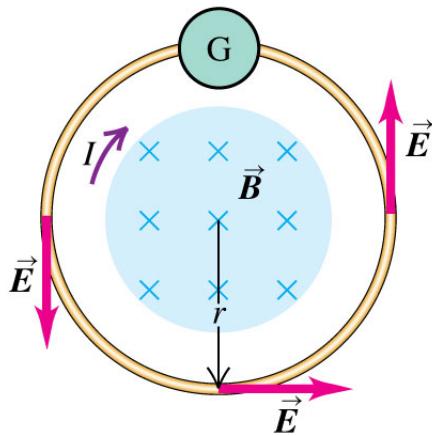
(a)



When a conductor moves in a magnetic field, we can understand the induced emf on the basis of magnetic forces on charges in the conductor.

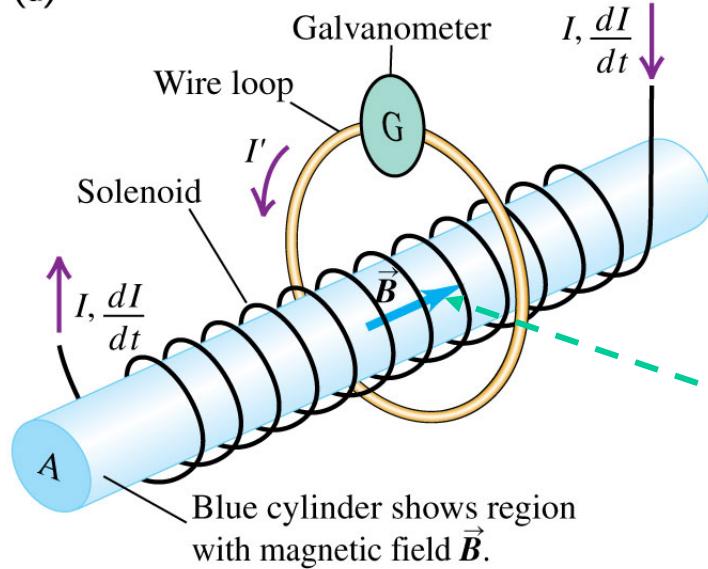
But an induced emf also occurs when there is a **changing flux through a stationary conductor**. What is it that pushes the charges around the circuit in this type of situation?

(b)



Induced electric fields

(a)



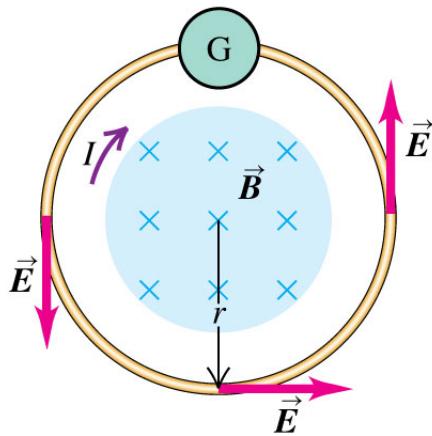
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Magnetic field due to the current I : $B = \mu_0 n I$

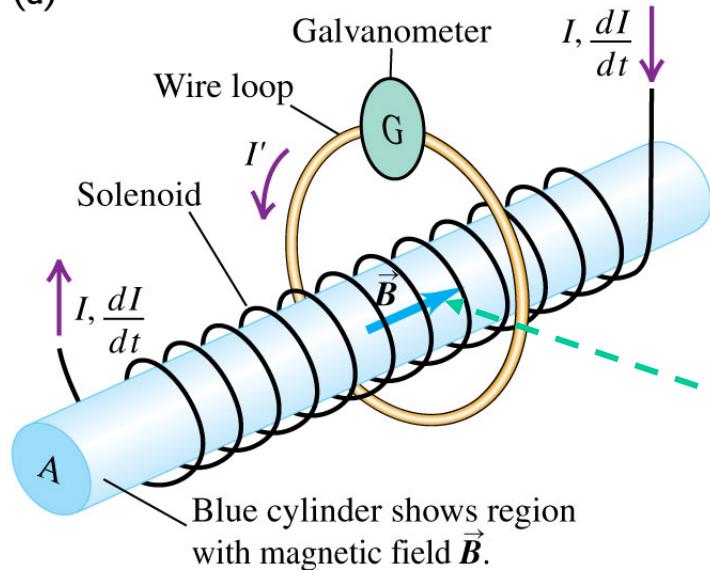
$$\Phi_B = BA = \mu_0 n I A$$

(b)



Induced electric fields

(a)



Blue cylinder shows region with magnetic field \vec{B} .

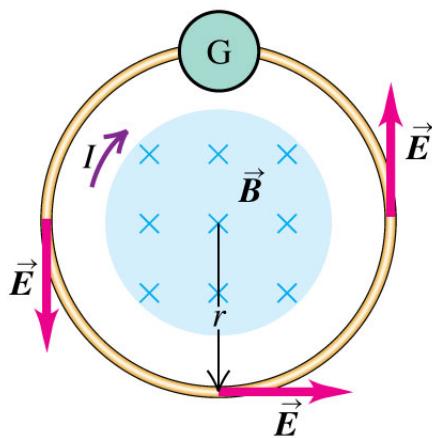
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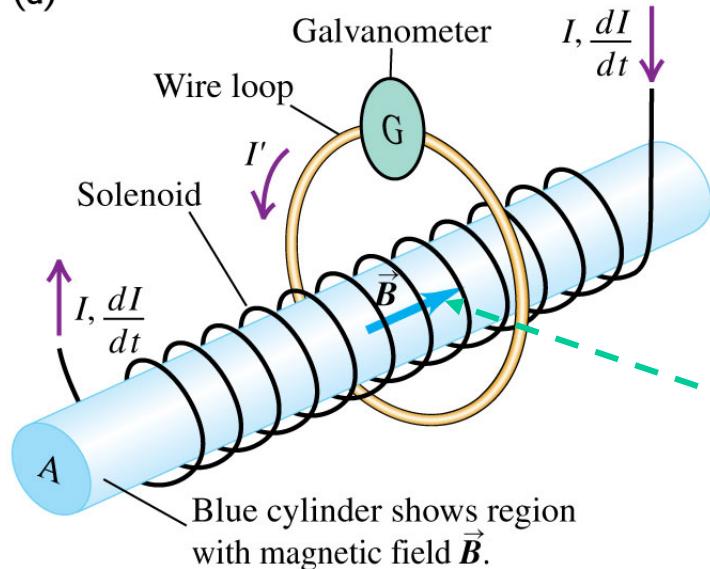


When the solenoid current I changes with time, the magnetic flux Φ_B also changes, and according to Faraday's law the *induced emf* in the loop is given by

$$\mathcal{E} = -\frac{d\Phi_B}{dt} = -\mu_0 nA \frac{dI}{dt}$$

Induced electric fields

(a)



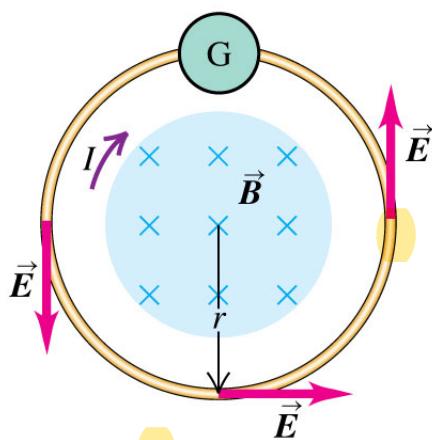
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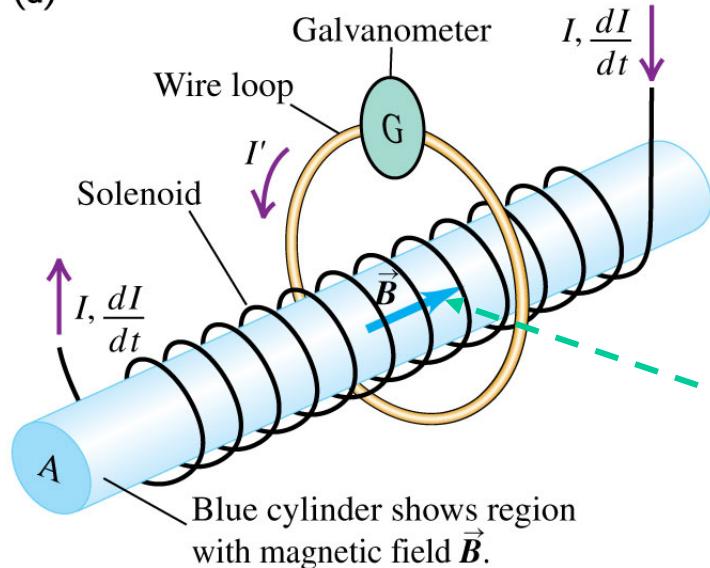
$$\mathcal{E} = -\frac{d\Phi_B}{dt} = -\mu_0 nA \frac{dI}{dt} \quad I' = \mathcal{E}/R$$

But what force makes the charges move around the wire loop? It can't be a magnetic force because the loop isn't even in a magnetic field → induced electric field in the conductor caused by the changing magnetic flux.

$$\oint \vec{E} \cdot d\vec{l} = \mathcal{E}$$

Induced electric fields

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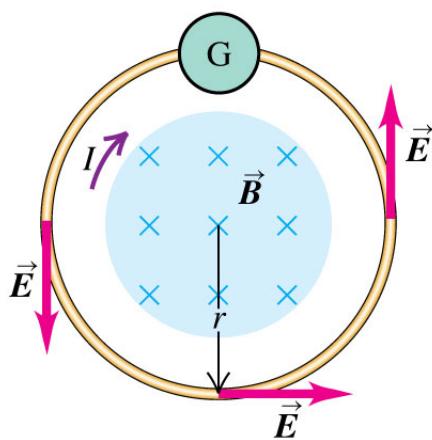
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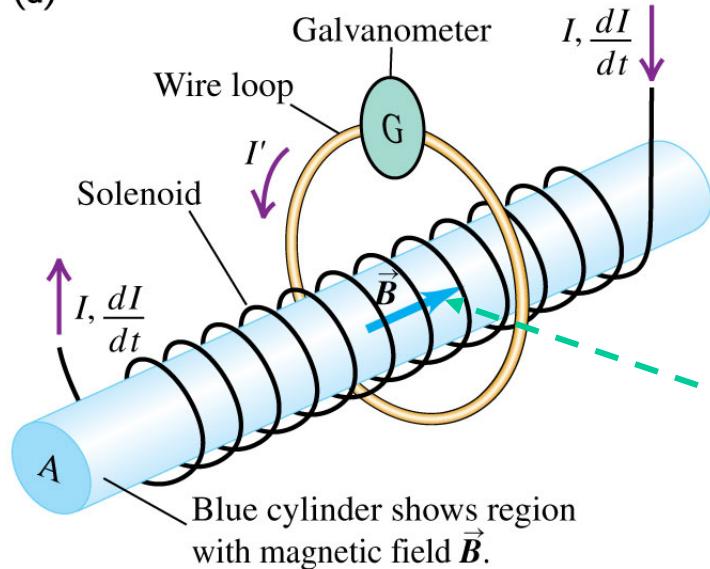
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$$\oint \vec{E} \cdot d\vec{l} = -\frac{d\Phi_B}{dt} \quad (\text{stationary integration path})$$

Induced electric fields

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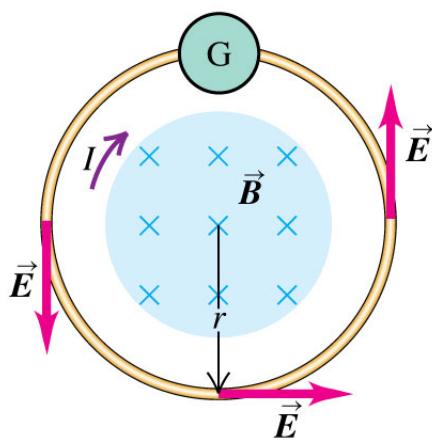
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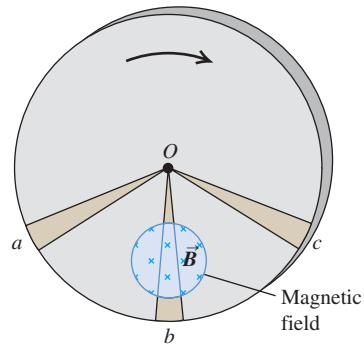
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$$E = \frac{1}{2\pi r} \left| \frac{d\Phi_B}{dt} \right|$$

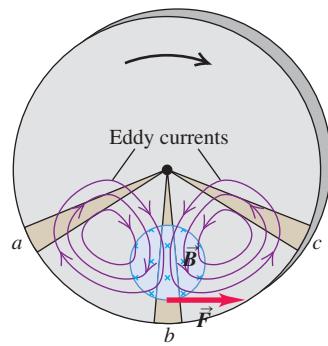
Eddy Currents

(a) Metal disk rotating through a magnetic field



many pieces of electrical equipment contain masses of metal moving in magnetic fields or located in changing magnetic fields. In situations like these we can have induced currents that circulate throughout the volume of a material.

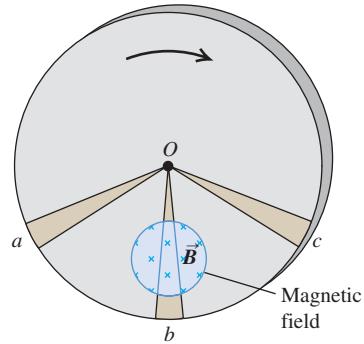
(b) Resulting eddy currents and braking force



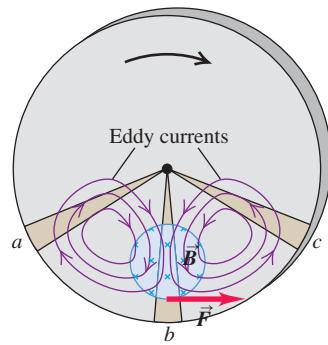
Eddy Currents

many pieces of electrical equipment contain masses of metal moving in magnetic fields or located in changing magnetic fields. In situations like these we can have induced currents that circulate throughout the volume of a material. Because their flow patterns resemble swirling eddies in a river, we call these **eddy currents**.

(a) Metal disk rotating through a magnetic field



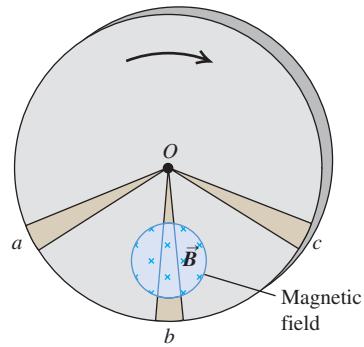
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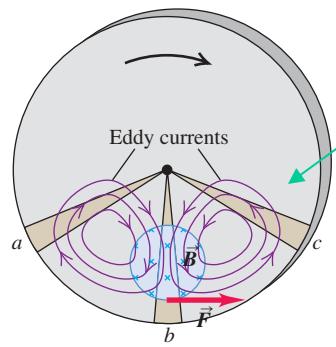
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(b) Resulting eddy currents and braking force

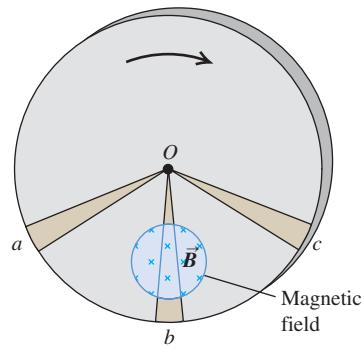


The return currents lie outside the field, so they do not experience magnetic forces. The interaction between the eddy currents and the field causes a braking action on the disk. Such effects can be used to stop the rotation of a circular saw quickly when the power is turned off.

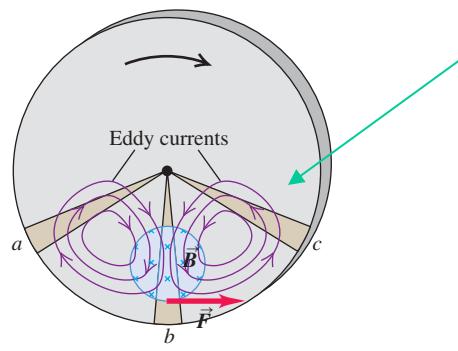
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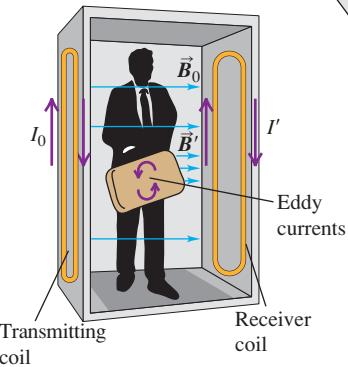


(b) Resulting eddy currents and braking force

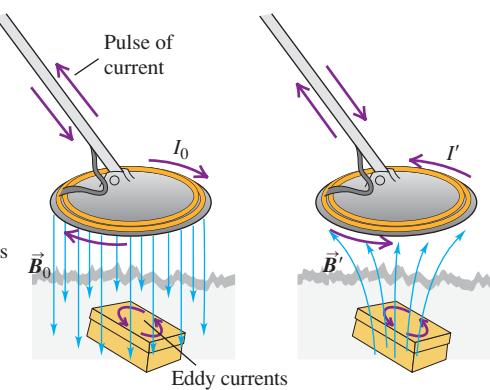


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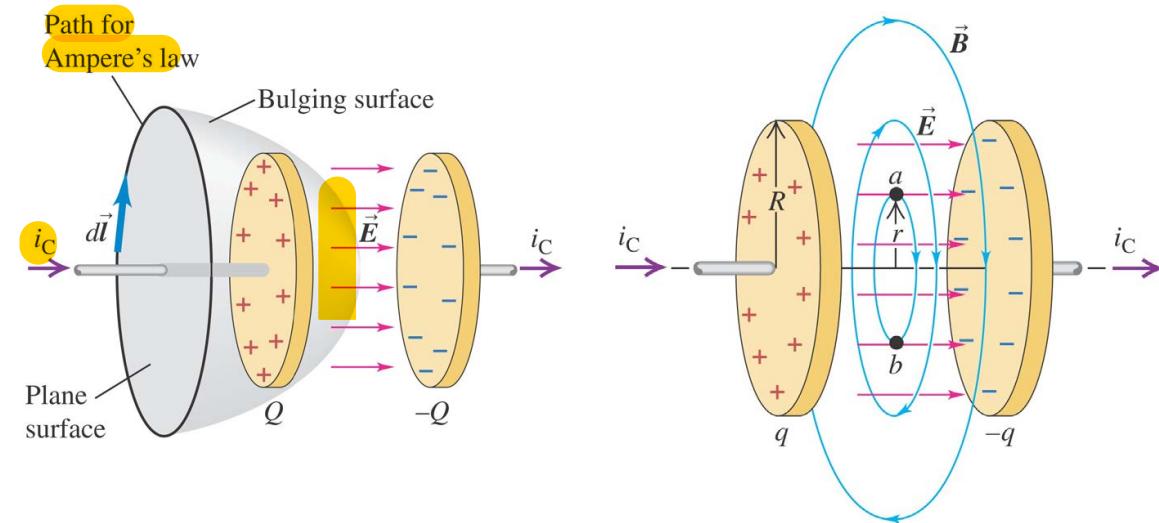


(b)



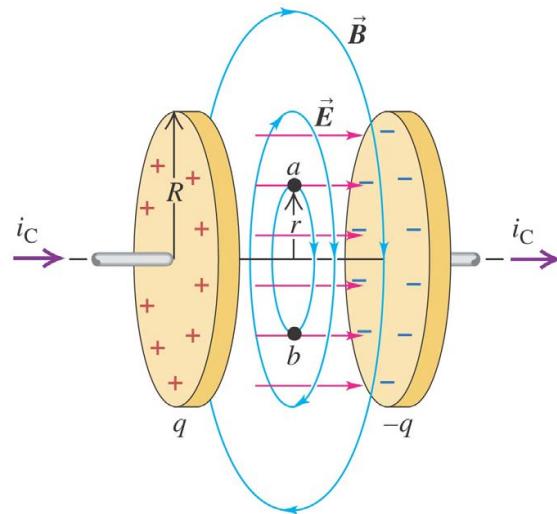
A metal detector at an airport security checkpoint generates an alternating magnetic field \vec{B} . This induces eddy currents in a conducting object carried through the detector. The eddy currents in turn produce an alternating magnetic field \vec{B}' , and this field induces a current in the detector's receiver coil. (b) Portable metal detectors work on the same principle.

Displacement current

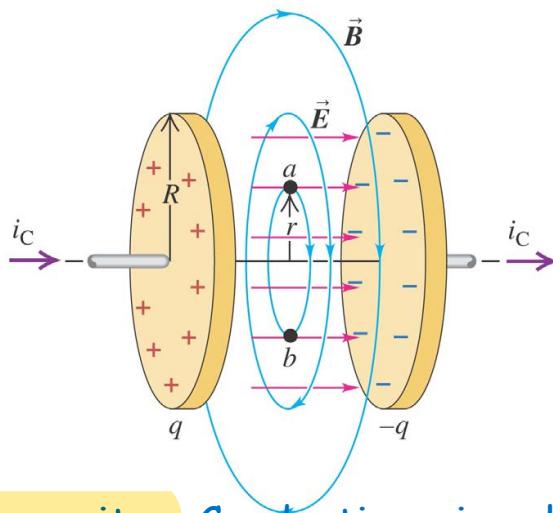
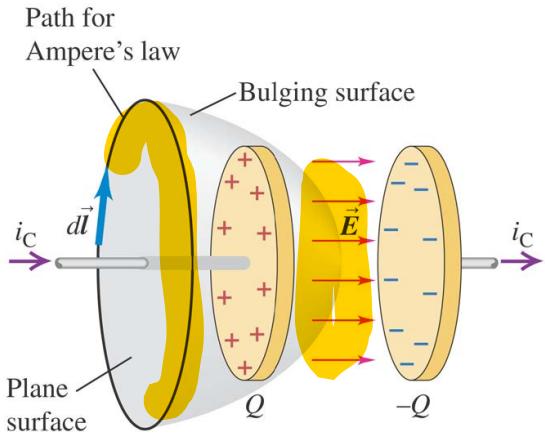


Ampere's Law (incomplete):

$$\oint \vec{B} \cdot d\vec{l} = \mu_0 I_{\text{encl}}$$



Displacement current



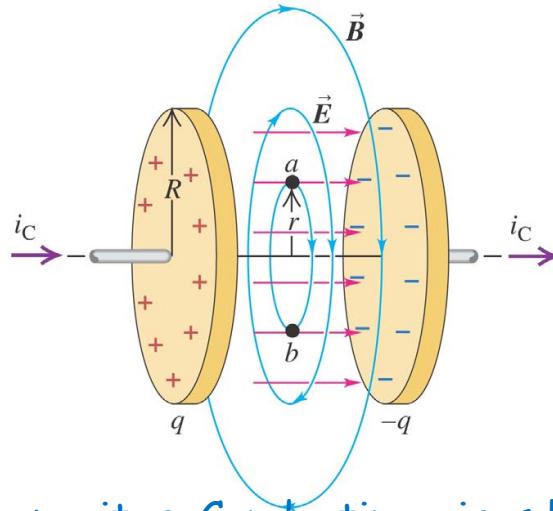
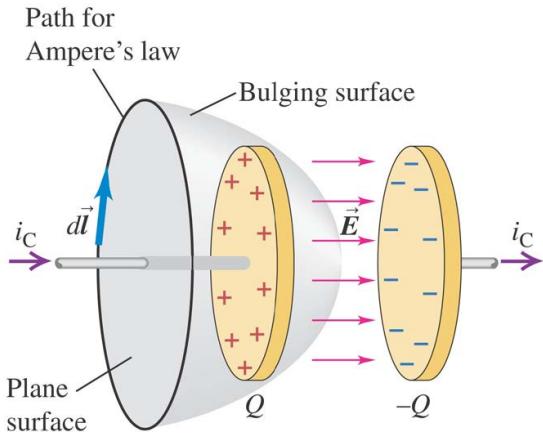
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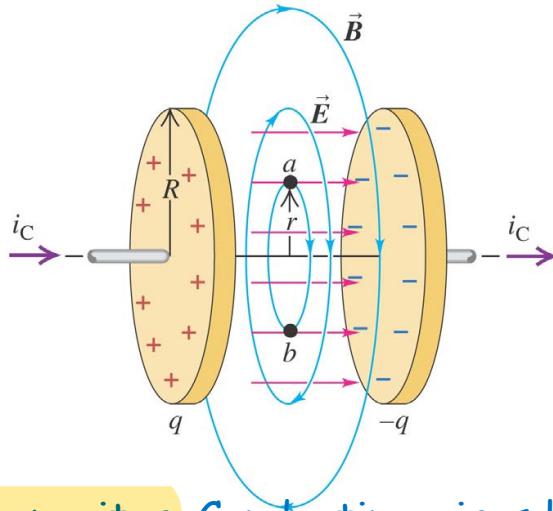
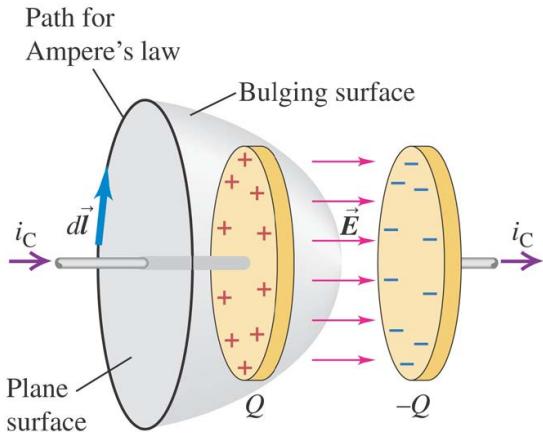
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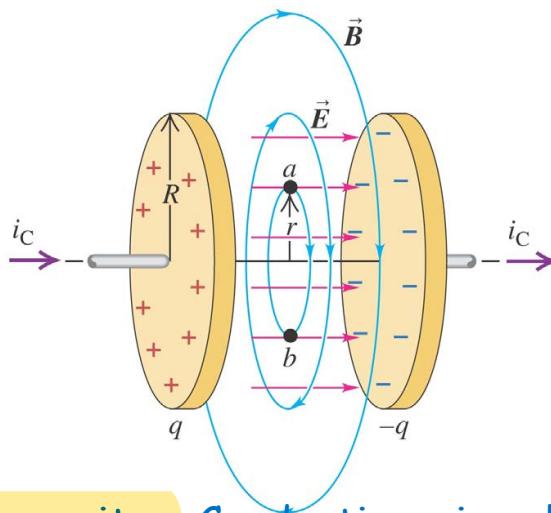
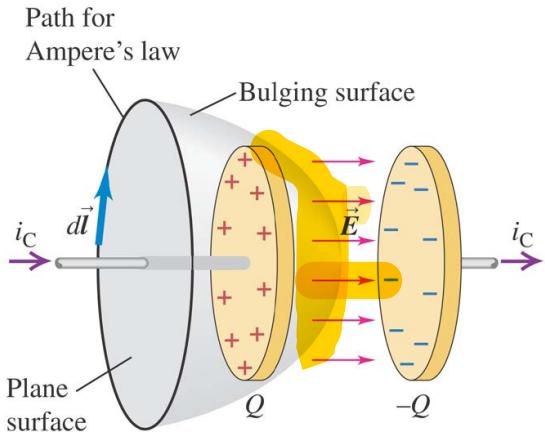
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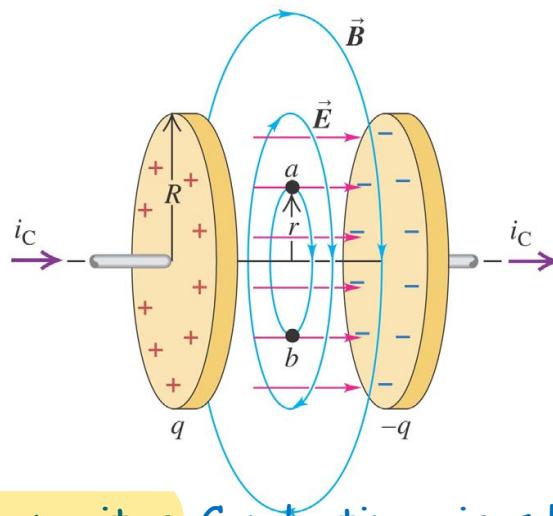
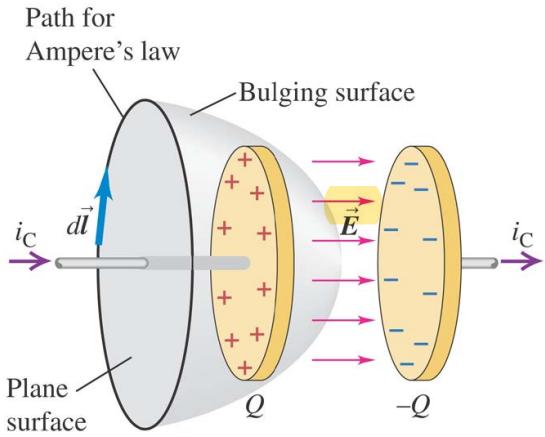
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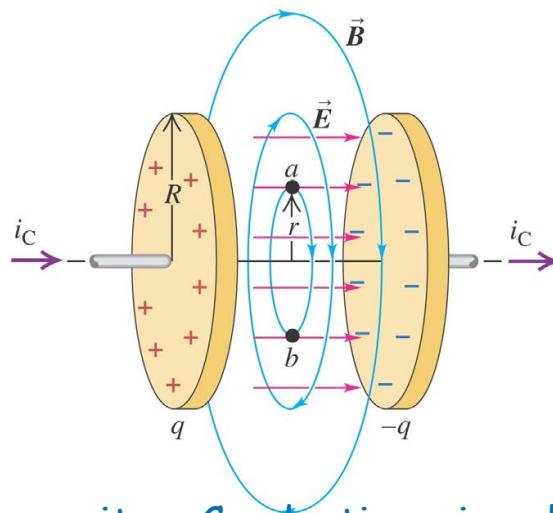
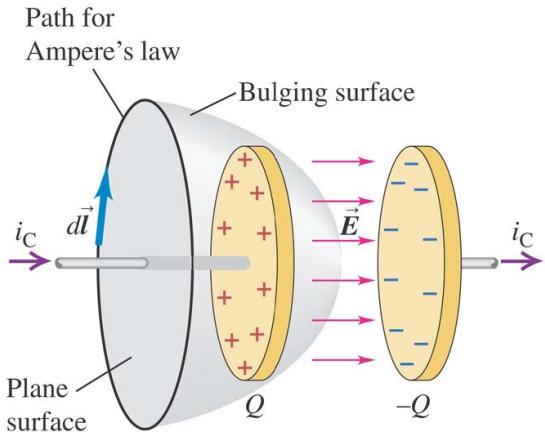
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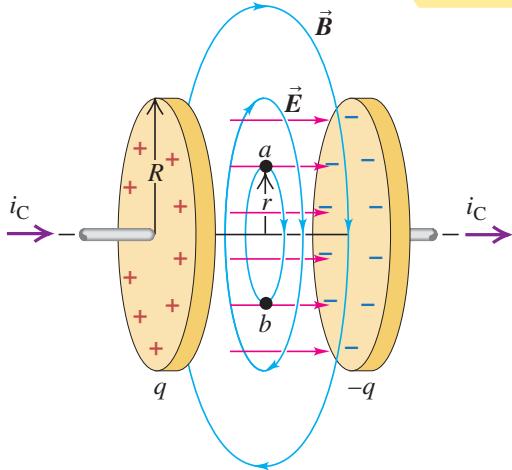
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$$i_D = \epsilon \frac{d\Phi_E}{dt} \rightarrow \oint \vec{B} \cdot d\vec{l} = \mu_0(i_C + i_D)_{\text{encl}}$$

Generalized Ampere's Law

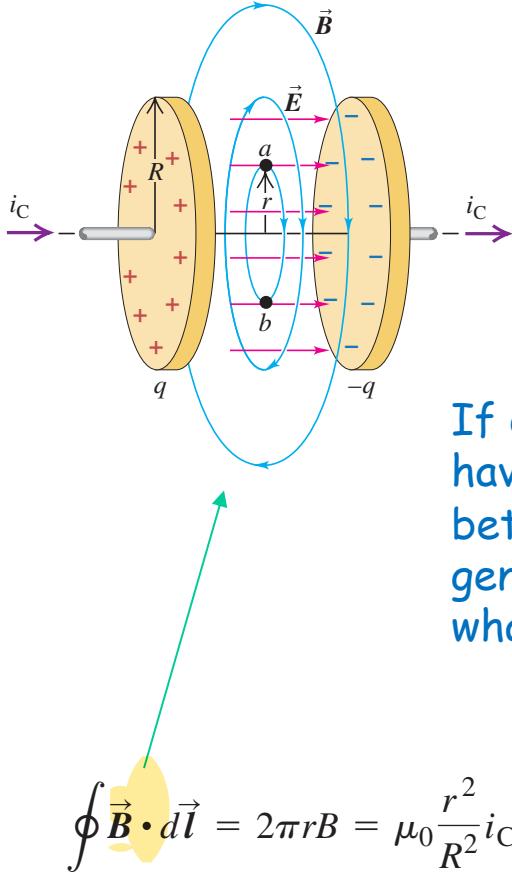
The Reality of Displacement Current



The fictitious current i_D was invented in 1865 by the Scottish physicist James Clerk Maxwell (1831-1879), who called it displacement current. There is a corresponding displacement current density $j_D = i_D / A$, using $\Phi_E = EA$

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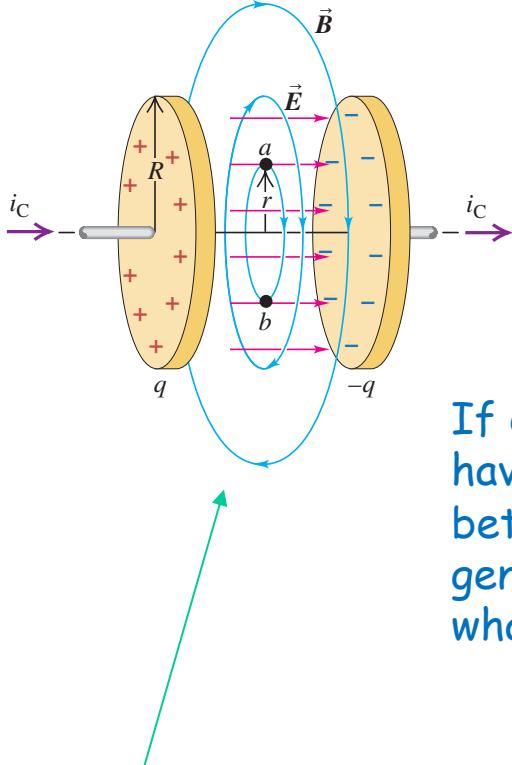


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$$\oint \vec{B} \cdot d\vec{l} = 2\pi r B = \mu_0 \frac{r^2}{R^2} i_C \quad \rightarrow \quad B = \frac{\mu_0}{2\pi} \frac{r}{R^2} i_C$$

When we *measure* the magnetic field in this region, we find that it really is there and that it behaves just as generalized Ampere law predicts. This confirms directly the role of displacement current as a source of magnetic field. It is now established beyond reasonable doubt that displacement current, far from being just an artifice, is a fundamental fact of nature. Maxwell's discovery was the bold step of an extraordinary genius.

Maxwell's Equations of Electromagnetism

$$\oint \vec{E} \cdot d\vec{A} = \frac{Q_{\text{encl}}}{\epsilon_0} \quad (\text{Gauss's law for } \vec{E})$$

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$$\Phi_E = \int \vec{E} \cdot d\vec{A} \quad \Phi_B = \int \vec{B} \cdot d\vec{A}$$

$$\vec{F} = q(\vec{E} + \vec{v} \times \vec{B})$$

In empty space $i_C = 0$; $Q_{\text{encl}} = 0$

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