## Turbulence HW3

## Yu Cang 018370210001

November 6, 2018

## 1 Exercise1

When the gravity is considered, the N-S equation is written as

$$\rho \frac{\partial u_j}{\partial t} + \rho u_k \frac{\partial u_j}{\partial x_k} = -\frac{\partial p}{\partial x_j} + \mu \frac{\partial^2 u_j}{\partial x_k \partial x_k} + \rho g_j$$
 (1.1)

Decompose the density into mean and fluctuation as  $\rho = \rho_0 + \tilde{\rho}$ , and substitute into the equation above yields

$$\rho_0 \frac{\partial u_j}{\partial t} + \tilde{\rho} \frac{\partial u_j}{\partial t} + \rho_0 u_k \frac{\partial u_j}{\partial x_k} + \tilde{\rho} \frac{\partial u_j}{\partial x_k} = -\frac{\partial p}{\partial x_j} + \mu \frac{\partial^2 u_j}{\partial x_k \partial x_k} + \rho_0 g_j + \tilde{\rho} g_j$$
(1.2)

The 2nd and 4th term in the LHS of (1.2) can be neglected as  $\tilde{\rho} << \rho_0$ . Assuming  $\frac{\mu}{\rho_0} \approx \frac{\mu}{\rho} = v$  yields

$$\frac{\partial u_j}{\partial t} + u_k \frac{\partial u_j}{\partial x_k} = -\frac{1}{\rho_0} \frac{\partial p}{\partial x_i} + v \frac{\partial^2 u_j}{\partial x_k \partial x_k} + g_j + g_j \frac{\tilde{\rho}}{\rho_0}$$
(1.3)

When the EOS for ideal gas( $p = \rho RT$ ) is adopted,  $\rho T$  can be assumed to be constant when the change of p is negligible and velocity is small. Hence

$$(\rho_0 + \tilde{\rho})(\bar{T} + \tilde{T}) = \rho_0 \bar{T} \tag{1.4}$$

Thus

$$\frac{\tilde{\rho}}{\rho_0} + \frac{\tilde{T}}{\bar{T}} = 0 \tag{1.5}$$

This is achieved by neglecting the 2nd-order small quantities.

Substitute into (1.3) yields

$$\frac{\partial u_j}{\partial t} + u_k \frac{\partial u_j}{\partial x_k} = -\frac{1}{\rho_0} \frac{\partial p}{\partial x_j} + v \frac{\partial^2 u_j}{\partial x_k \partial x_k} + g_j - g_j \frac{\tilde{T}}{\tilde{T}}$$
(1.6)

Decompose the velocity and pressure into mean and fluctuation as

$$u_j = U_j + u'_j$$

$$p = P + p'$$
(1.7)

Substitute into (1.6)

$$\frac{\partial (U_j + u_j')}{\partial t} + (U_k + u_k') \frac{\partial (U_j + u_j')}{\partial x_k} = -\frac{1}{\rho_0} \frac{\partial (P + p_j')}{\partial x_j} + v \frac{\partial^2 (U_j + u_j')}{\partial x_k \partial x_k} + g_j - g_j \frac{\tilde{T}}{\tilde{T}}$$
(1.8)

Multiply  $u_i'$  and take ensemble average yields

$$LHS \triangleq \overline{u'_{j}\frac{\partial u'_{j}}{\partial t}} + U_{k}\overline{u'_{j}\frac{\partial u'_{j}}{\partial x_{k}}} + \overline{u'_{j}u'_{k}\frac{\partial U_{j}}{\partial x_{k}}} + \overline{u'_{j}u'_{k}\frac{\partial u'_{j}}{\partial x_{k}}} = -\frac{1}{\rho_{0}}\overline{u'_{j}\frac{\partial p'}{\partial x_{j}}} + v\overline{u'_{j}\frac{\partial^{2}u'_{j}}{\partial x_{k}\partial x_{k}}} + \frac{g_{j}}{\bar{T}}\overline{u'_{j}}\tilde{T} \triangleq RHS$$
(1.9)

Denote  $k_T = \frac{1}{2} \overline{u'_i u'_j}$  and follow from the continuity equation

$$LHS = \frac{\bar{D}k_T}{\bar{D}t} + \frac{\bar{u}_j'u_k'}{\partial x_k} \frac{\partial U_j}{\partial x_k} + \frac{1}{2} \frac{\partial}{\partial x_k} \frac{\bar{u}_k'u_j'u_j'}{\bar{u}_j'}$$
(1.10)

Since

$$\overline{u'_j \frac{\partial^2 u'_j}{\partial x_k \partial x_k}} = \overline{\frac{\partial}{\partial x_k} \left( u'_j \frac{\partial u'_j}{\partial x_k} \right)} - \overline{\frac{\partial u'_j}{\partial x_k} \frac{\partial u'_j}{\partial x_k}} = \overline{\frac{\partial^2 k_T}{\partial x_k \partial x_k}} - \overline{\frac{\partial u'_j}{\partial x_k} \frac{\partial u'_j}{\partial x_k}} \tag{1.11}$$

then, also follows from the continuity equation

$$RHS = -\frac{1}{\rho_0} \frac{\partial \overline{u_j' p'}}{\partial x_j} + \nu \left( \frac{\partial^2 k_T}{\partial x_k \partial x_k} - \frac{\overline{\partial u_j'}}{\partial x_k} \frac{\partial u_j'}{\partial x_k} \right) + \frac{g_j}{\overline{T}} \overline{u_j' \widetilde{T}}$$
(1.12)

Combining (1.10) and (1.12) yields the turbulent kinetic energy equation for the buoyancy case

$$\frac{\bar{D}k_T}{\bar{D}t} = -\overline{u'_j u'_k} \frac{\partial U_j}{\partial x_k} - \frac{1}{2} \frac{\partial}{\partial x_k} \overline{u'_k u'_j u'_j} - \frac{1}{\rho_0} \frac{\partial \overline{u'_j p'}}{\partial x_j} + \nu \left( \frac{\partial^2 k_T}{\partial x_k \partial x_k} - \frac{\overline{\partial u'_j}}{\partial x_k} \frac{\partial u'_j}{\partial x_k} \right) + \frac{g_j}{\bar{T}} \overline{u'_j \tilde{T}}$$
(1.13)

2 EXERCISE 2