Astronomy C3101 Modern Stellar Astrophysics Homework 3

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Tuesday, November 18, 2014

Problem 1 - Luminosity variation from stellar pulsation

(a) Estimate the resulting change in luminosity δL .

$$\begin{split} L &= 4\pi R^2 \sigma T^4 \\ L_0 \left(1 + \frac{\delta L}{L_0}\right) &= 4\pi \left[R_0 \left(1 + \frac{\delta R}{R_0}\right)\right]^2 \sigma \left[T_0 \left(1 + \frac{\delta T}{T_0}\right)\right]^4 \\ L_0 \left(1 + \frac{\delta L}{L_0}\right) &= 4\pi \left[R_0^2 \left(\frac{\delta R}{R_0}^2 + 2\frac{\delta R}{R_0} + 1\right)\right] \sigma \left[T_0^4 \left(\frac{\delta T}{T_0}^4 + 4\frac{\delta T}{T_0}^3 + 6\frac{\delta T}{T_0}^2 + 4\frac{\delta T}{T_0} + 1\right)\right] \\ \text{Discarding higher-order terms (or using the Taylor expansion for brevity),} \\ L_0 \left(1 + \frac{\delta L}{L_0}\right) &\approx 4\pi \left[R_0^2 \left(1 + 2\frac{\delta R}{R_0}\right)\right] \sigma \left[T_0^4 \left(1 + 4\frac{\delta T}{T_0}\right)\right] \\ \frac{\delta L}{L_0} &\approx \left(1 + 2\frac{\delta R}{R_0}\right) \left(1 + 4\frac{\delta T}{T_0}\right) - 1 \end{split}$$

(b) Derive a relationship between δR and δT .

Gas is adiabatic, relating δP to $\delta \rho$:

Taylor expanding,
$$P = K_a \rho^{\gamma_a}$$

$$P_0 \left(1 + \frac{\delta P}{P_0} \right) = K_a \left[\rho_0 \left(1 + \frac{\delta \rho}{\rho_0} \right) \right]^{\gamma_a}$$
Taylor expanding,
$$\approx K_a \rho_0^{\gamma_a} \left(1 + \gamma_a \frac{\delta \rho}{\rho_0} \right)$$
But since

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$$\delta P = K_a \rho_0^{\gamma_a} \gamma_a \frac{\delta \rho}{\rho_0},$$

$$\frac{\delta P}{P_0} = \gamma_a \frac{\delta \rho}{\rho_0}$$

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Gas is ideal, relating δP to δT :

$$P = \frac{R}{\mu} \rho T$$

$$\begin{split} P &= \frac{\kappa}{\mu} \rho T \\ P_0 \left(1 + \frac{\delta P}{P_0} \right) &= \frac{R}{\mu} \rho_0 \left(1 + \frac{\delta \rho}{\rho_0} \right) T_0 \left(1 + \frac{\delta T}{T_0} \right) \\ P_0 &= \frac{R}{\mu} \rho_0 T_0 \\ 1 &+ \frac{\delta P}{P_0} &= \left(1 + \frac{\delta \rho}{\rho_0} \right) \left(1 + \frac{\delta T}{T_0} \right) \\ 1 &+ \frac{\delta P}{P_0} &= 1 + \frac{\delta \rho}{\rho_0} + \frac{\delta T}{T_0} + \frac{\delta \rho}{\rho_0} \frac{\delta T}{T_0} \\ \text{Dropping higher-order terms,} \\ \frac{\delta P}{P_0} &= \frac{\delta \rho}{\rho_0} + \frac{\delta T}{T_0} \\ \frac{\delta T}{T_0} &= \frac{\delta P}{P_0} - \frac{\delta \rho}{\rho_0} \end{split}$$

$$P_0 = \frac{R}{\mu} \rho_0 T_0$$

$$1 + \frac{\delta P}{P_0} = \left(1 + \frac{\delta \rho}{\rho_0}\right) \left(1 + \frac{\delta T}{T_0}\right)$$

$$1 + \frac{\delta P}{P_0} = 1 + \frac{\delta \rho}{\rho_0} + \frac{\delta T}{T_0} + \frac{\delta \rho}{\rho_0} \frac{\delta T}{T_0}$$

$$\frac{\delta P}{P_0} = \frac{\delta \rho}{\rho_0} + \frac{\delta T}{T_0}$$

$$\frac{\delta T}{T_0} = \frac{\delta P}{P_0} - \frac{\delta P}{\rho_0}$$

Use homologous expansion relation to relate $\delta \rho$ to δR (just like in Lecture 15): $dm = 4\pi R^2 \rho dR$

$$dm = 4\pi R^2 \rho dR$$

$$= 4\pi \left[R_0 (1 + \frac{\delta R}{R_0}) \right]^2 \rho_0 \left(1 + \frac{\delta \rho}{\rho_0} \right) dR_0 \left(1 + \frac{\delta R}{R_0} \right)$$

$$= 4\pi R_0^2 \rho_0 dR_0 \left(1 + 3\frac{\delta R}{R_0} + \frac{\delta \rho}{\rho_0} \right)$$
But since $dm = 4\pi R_0^2 \rho_0 dR_0 = 0$,
$$3\frac{\delta R}{R_0} + \frac{\delta \rho}{\rho_0} = 0$$

$$=4\pi R_0^2 \rho_0 dR_0 (1+3\frac{\delta R}{R_0}+\frac{\delta \rho}{20})$$

But since
$$dm = 4\pi R_0^2 \rho_0 dR_0 = 0$$

$$\frac{\delta\rho}{\rho_0} = -3\frac{\delta R}{R_0}$$

Taking those three relations and substituting/combining.

(c) Arrive at an estimate for the relationship between δL and δR . Does the peak luminosity of a pulsating star with $\gamma_a = 5/3$ occur when its radius is at its maximum value or its

minimum value?
$$\frac{\delta L}{L_0} \approx \left(1 + 2\frac{\delta R}{R_0}\right) \left(1 + 4\frac{\delta T}{T_0}\right) - 1$$

$$\frac{\delta T}{T_0} = -3\frac{\delta R}{R_0} (\gamma_a - 1)$$
 Substituting,
$$\frac{\delta L}{L_0} \approx \left(1 + 2\frac{\delta R}{R_0}\right) \left(1 + 4(-3\frac{\delta R}{R_0}(\gamma_a - 1))\right) - 1$$

$$\frac{\delta L}{L_0} \approx \left(1 + 2\frac{\delta R}{R_0}\right) \left(1 - 12\frac{\delta R}{R_0}(\gamma_a - 1)\right) - 1$$
 When $\gamma_a = 5/3$,
$$\frac{\delta L}{L_0} \approx \left(1 + 2\frac{\delta R}{R_0}\right) \left(1 - 12\frac{\delta R}{R_0}(2/3)\right) - 1$$

$$\frac{\delta L}{L_0} \approx \left(1 + 2\frac{\delta R}{R_0}\right) \left(1 - 8\frac{\delta R}{R_0}\right) - 1$$

$$\frac{\delta L}{L_0} \approx \left(1 - 6\frac{\delta R}{R_0} - 16(\frac{\delta R}{R_0})^2\right) - 1$$
 Removing higher-order terms,
$$\frac{\delta L}{L_0} \approx -6\frac{\delta R}{R_0}$$
 There is an inverse relationship, so the peak pulsating luminosity occurs when the radius is at its minimum value.

is at its minimum value.

Problem 2 - Generating the T vs. ρ diagram

(a) Calculate μ_I, μ_e , and μ .

$$\begin{array}{l} X = 0.75, Y = 0.25 \\ \frac{1}{\mu_I} = X + \frac{1}{4}Y = 0.8125 \\ \mu_I = 1.2308 \\ \frac{1}{\mu_e} = \frac{1}{2}(1+X) = 0.875 \\ \mu_e = 1.1429 \\ \frac{1}{\mu} = \frac{1}{\mu_I} + \frac{1}{\mu_e} = 1.6875 \\ \mu = 0.5926 \end{array}$$

(b) Solve for the boundary between the ideal gas zone and the non-relativistic degenerate

$$\begin{split} &K_0 \rho T = K_1 \frac{1}{\mu_e^{5/3}} \rho^{5/3} \\ &T = \frac{K_1}{K_0} \frac{1}{\mu_e^{5/3}} \rho^{2/3} \\ &log(T) = log\left(\frac{K_1}{K_0} \frac{1}{\mu_e^{5/3}}\right) + log(\rho^{2/3}) \\ &\frac{2}{3}log(\rho) = log(T) - log\left(\frac{K_1}{K_0} \frac{1}{\mu_e^{5/3}}\right) \\ &log(\rho) = \frac{3}{2}log(T) - \frac{3}{2}log\left(\frac{K_1}{K_0} \frac{1}{\mu_e^{5/3}}\right) \\ &R = 8.3145 \times 10^7 erg/K \cdot mol \\ &K_0 = \frac{R}{\mu} = 1.4031 \times 10^8 erg/K \cdot mol \\ &K_1 = 1.0 \times 10^7 m^4 kg^{-2/3} s^{-2} = 1 \times 10^{13} cm^4 g^{-2/3} s^{-2} \\ &\mu_e = 1.1429 \\ &log(\rho) = \frac{3}{3}log(T) - 7.1344 \end{split}$$

(c) Solve for the boundary between the non-relativistic and relativistic degenerate gas zones.

$$\begin{split} K_1 \frac{1}{\mu_e^{5/3}} \rho^{5/3} &= K_2 \frac{1}{\mu_e^{4/3}} \rho^{4/3} \\ \rho^{1/3} &= \frac{K_2}{K_1} \mu_e^{1/3} \\ log(\rho^{1/3}) &= log\left(\frac{K_2}{K_1} \mu_e^{1/3}\right) \\ log(\rho) &= 3log\left(\frac{K_2}{K_1} \mu_e^{1/3}\right) \\ K_1 &= 1.0 \times 10^7 m^4 kg^{-2/3} s^{-2} = 1 \times 10^{13} cm^4 g^{-2/3} s^{-2} \\ K_2 &= 1.24 \times 10^{10} m^3 kg^{-1/3} s^{-1} = 1.24 \times 10^{15} cm^3 g^{-1/3} s^{-1} \\ \mu_e &= 1.1429 \\ log(\rho) &= 6.3383 \end{split}$$

(d) Solve for the boundary between the ideal gas zone and the relativistic degenerate gas zone.

$$K_{0}\rho T = K_{2} \frac{1}{\mu_{e}^{4/3}} \rho^{4/3}$$

$$T = \frac{K_{2}}{K_{0}} \frac{1}{\mu_{e}^{4/3}} \rho^{1/3}$$

$$log(T) = log(\frac{K_{2}}{K_{0}} \frac{1}{\mu_{e}^{4/3}} \rho^{1/3})$$

$$\frac{1}{3} log(\rho) = log(T) - log(\frac{K_{2}}{K_{0}} \frac{1}{\mu_{e}^{4/3}})$$

$$log(\rho) = 3log(T) - 3log(\frac{K_{2}}{K_{0}} \frac{1}{\mu_{e}^{4/3}})$$

$$K_{0} = \frac{R}{\mu} = 1.4031 \times 10^{8} erg/K \cdot mol$$

$$K_{2} = 1.24 \times 10^{10} m^{3} kg^{-1/3} s^{-1} = 1.24 \times 10^{15} cm^{3} g^{-1/3} s^{-1}$$

$$\mu_{e} = 1.1429$$

$$log(\rho) = 3log(T) - 20.607$$

- (e) Solve for the value of (T, ρ) where the ideal gas pressure, non-relativistic degenerate gas pressure, and relativistic gas pressure are all equal.
- (f) Solve for the boundary between the ideal gas zone and the radiation pressure taking the boundary as $P_{rad} = 10P_{gas}$.

$$10(K_0\rho T) = \frac{1}{3}aT^4$$

$$K_0\rho = \frac{1}{30}aT^3$$

$$log(\rho) = log(\frac{a}{30K_0}T^3)$$

$$log(\rho) = 3log(T) + log(\frac{a}{30K_0})$$

$$R = 8.3145 \times 10^7 erg/K \cdot mol$$

$$K_0 = \frac{R}{\mu} = 1.4031 \times 10^8 erg/K \cdot mol$$

$$a = 7.6 \times 10^{-15} erg/cm^3/K^4$$

$$log(\rho) = 3log(T) - 23.7434$$

(g) Plot the boundaries between the different zones for log T(K) = 6–10 and $log \rho(g \, cm^{-3}) = 0$ –10.

Problem 3 - Convection

(a) Show that the envelope of a star that has a Kramers Law opacity (with a = 1, b = -3.5) is stable against convection if the equation of state is that of an ideal gas with $\gamma = 5/3$.

Kramer's law:
$$\kappa = c \rho^a T^b$$

 $\kappa = c \rho T^{-3.5}$
Gas is adiabatic
Gas is ideal: $P = \frac{R}{\mu} \rho T$
 $\rho = \frac{\mu}{R} \frac{P}{T}, T = \frac{\mu}{R} \frac{P}{\rho}$
As related in lecture 16,
 $dP = \frac{R}{\mu} \left(\rho \frac{dT}{dr} + T \frac{d\rho}{dr} \right) dr$
 $= \left(\frac{P}{T} \frac{dT}{dr} + \frac{P}{\rho} \frac{d\rho}{dr} \right) dr$

$$P = K_a \rho^{\gamma}$$

$$dP = K_a \gamma \rho^{\gamma - 1} \frac{d\rho}{dr} dr$$

$$= \gamma \frac{P}{\rho} \frac{d\rho}{dr} dr$$

$$\gamma \frac{P}{\rho} \frac{d\rho}{dr} = \frac{P}{T} \frac{dT}{dr} + \frac{P}{\rho} \frac{d\rho}{dr}$$

$$\frac{dT}{dr} = (\gamma - 1) \frac{T}{P} \frac{P}{\rho} \frac{d\rho}{dr}$$

$$= (\frac{\gamma - 1}{\gamma}) \frac{T}{P} \frac{dP}{dr}$$

$$\frac{d\ln(T)}{d\ln(P)} = \frac{PdT}{TdP} = \frac{\gamma - 1}{\gamma} = \frac{5/3 - 1}{5/3} = \frac{2/3}{5/3} = \frac{2}{5}$$
Use this result to predict which main see

(b) Use this result to predict which main-sequence stars should have convective envelopes.

${\bf Problem~4~-~(see~attached)}$