

Gravitational Decoherence Rate Scales as $\cos^2 \theta$ with Measurement Basis

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Gravitational decoherence rate, we predict, depends on measurement basis: $\Gamma(\theta) = \Gamma_0 \cos^2 \theta$. For a 10^{-12} kg nanoparticle in $1 \mu\text{m}$ superposition, this yields: $\tau = 0.14$ s (position measurement), $\tau = 0.28$ s (45° quadrature), $\tau \rightarrow \infty$ (momentum measurement). No existing model predicts basis-dependence; all assume collapse occurs in a fixed basis regardless of how the system is observed. We derive this result from a projection principle that generalizes observer-dependence from classical mechanics to quantum-gravitational decoherence. The prediction is falsifiable using current optomechanical technology. A null result confirms standard Penrose-Diósi; a positive result implies gravity's preferred basis emerges only at measurement.

INTRODUCTION

Gravity imposes classicality—or so we have assumed. The Diósi-Penrose model predicts that spatial superpositions collapse under their own gravitational self-energy [1–3]. Larger masses decohere faster; sufficiently large objects never exhibit quantum behavior. This framework has guided experimental searches for gravitational decoherence for three decades [4–6].

We test it here—and predict it fails.

Not because the mechanism is wrong, but because it is incomplete. The standard treatment assumes decoherence rate Γ_0 depends only on mass and superposition geometry. It does not consider that the *observed* rate might depend on *how* the system is measured.

We propose that it does. The core prediction is:

$$\Gamma(\theta) = \Gamma_0 \cos^2 \theta \quad (1)$$

where θ is the angle between the measurement quadrature and the position basis. At $\theta = 0^\circ$ (position measurement), the full gravitational decoherence rate Γ_0 is observed. At $\theta = 90^\circ$ (momentum measurement), gravitational decoherence is invisible: $\Gamma(90^\circ) = 0$.

The physical reasoning is simple. Gravitational decoherence in the Diósi-Penrose model couples to position—the mass distribution of the superposed states. A measurement that interrogates position sees this decoherence. A measurement that interrogates momentum does not. Gravity cannot decohere what it does not see.

For a 10^{-12} kg silica nanosphere in a $1 \mu\text{m}$ superposition—parameters within reach of current optomechanical experiments [7, 8]—this yields coherence times of $\tau = 0.14$ s for position measurement, $\tau = 0.28$ s at 45° , and $\tau \rightarrow \infty$ for momentum measurement.

If confirmed, this result implies that gravitational decoherence is not a fixed property of the quantum state, but depends on the observer's choice of measurement basis. If refuted—if $\Gamma(\theta)$ is constant—it confirms that gravity decoheres position regardless of how we look.

The paper is structured as follows. Section introduces the projection principle. Section derives Eq. (1) from the Diósi master equation. Sections and present quantitative predictions and an experimental protocol. Section discusses implications.

THE PROJECTION PRINCIPLE

Measurement defines what can be observed. This is not philosophy—it is geometry.

Consider a pendulum on a rotating platform, viewed by cameras at different angles. A camera aligned with the rotation axis cannot see axial displacement; a camera perpendicular to it sees full displacement. The pendulum's motion does not change. What the camera records does—because the measurement axis determines which component of motion appears in the data.

We call this the projection principle. It applies wherever observation geometry filters physical information.

In the Diósi-Penrose framework, gravity couples to mass distribution: position. The collapse mechanism monitors where the particle is—not how fast it moves, not its energy, not its spin. A measurement reveals gravitational decoherence only if it interrogates position. This constraint has been invisible because every experiment has measured position.

A homodyne detector measures the generalized quadrature $\hat{X}_\theta = \hat{x} \cos \theta + \hat{p} \sin \theta$. At $\theta = 0^\circ$, it asks “where?”—maximum overlap with what gravity monitors. At $\theta = 90^\circ$, it asks “how fast?”—zero overlap. A momentum measurement cannot see position-based decoherence, exactly as an axially-aligned camera cannot see axial motion.

No existing collapse model accounts for this. Standard treatments assume decoherence rate Γ_0 is fixed by mass and separation alone, independent of measurement basis [1–4]. We propose instead that Γ_0 is the rate revealed when measuring position. Other measurements reveal a

projected rate:

$$\Gamma(\theta) = \Gamma_0 \cos^2 \theta \quad (2)$$

The form follows from quantum measurement theory. The quadrature \hat{X}_θ has position-amplitude $\cos \theta$. Decoherence contributes to measurement statistics incoherently, scaling with probability—the square of amplitude. The factor $\cos^2 \theta$ is not a model choice; it is a consequence of projection.

DERIVATION

We derive $\Gamma(\theta) = \Gamma_0 \cos^2 \theta$ from the Diósi master equation. No new physics is introduced. We ask only what the standard model predicts for measurements in bases other than position.

The Diósi master equation

For a particle of mass m in one dimension, gravitational decoherence is governed by [2, 3]:

$$\frac{d\hat{\rho}}{dt} = -\frac{i}{\hbar}[\hat{H}, \hat{\rho}] - \Gamma_0[\hat{x}, [\hat{x}, \hat{\rho}]] \quad (3)$$

with $\Gamma_0 = Gm^2/(\hbar d)$ for superposition separation d . The double commutator with \hat{x} destroys coherences between different position eigenstates at rate Γ_0 .

Rotated quadrature

Define:

$$\hat{X}_\theta = \hat{x} \cos \theta + \hat{p} \sin \theta \quad (4)$$

with conjugate $\hat{P}_\theta = -\hat{x} \sin \theta + \hat{p} \cos \theta$. Inverting gives:

$$\hat{x} = \hat{X}_\theta \cos \theta - \hat{P}_\theta \sin \theta \quad (5)$$

Transformation of the decoherence term

Substituting Eq. (5) into Eq. (3):

$$\begin{aligned} [\hat{x}, [\hat{x}, \hat{\rho}]] &= \cos^2 \theta [\hat{X}_\theta, [\hat{X}_\theta, \hat{\rho}]] \\ &+ \sin^2 \theta [\hat{P}_\theta, [\hat{P}_\theta, \hat{\rho}]] \\ &+ \text{cross terms} \end{aligned} \quad (6)$$

Three contributions emerge:

1. $\cos^2 \theta [\hat{X}_\theta, [\hat{X}_\theta, \hat{\rho}]]$ — decoheres \hat{X}_θ eigenstates
2. $\sin^2 \theta [\hat{P}_\theta, [\hat{P}_\theta, \hat{\rho}]]$ — decoheres \hat{P}_θ eigenstates, not \hat{X}_θ

3. Cross terms — produce phase-space shearing, not decoherence

Only the first term destroys coherence between \hat{X}_θ eigenstates. The effective decoherence rate in the measured basis is therefore:

$$\boxed{\Gamma(\theta) = \Gamma_0 \cos^2 \theta} \quad (7)$$

Figure 1 displays this result.

Limiting cases

Measurement	θ	$\Gamma(\theta)$
Position	0°	Γ_0
Symmetric quadrature	45°	$\Gamma_0/2$
Momentum	90°	0

These cases mark the endpoints and midpoint of the $\cos^2 \theta$ curve in Fig. 1.

Why $\cos^2 \theta$, not $\cos \theta$

Decoherence rates scale with the square of operator couplings—parallel to Fermi’s golden rule, where transition rates involve matrix elements squared. The projection of \hat{x} onto \hat{X}_θ has amplitude $\cos \theta$; the resulting decoherence rate scales as $\cos^2 \theta$.

What this derivation is—and is not

We have not modified the Diósi-Penrose model. We have applied it to a measurement basis other than position. The result—that observed decoherence depends on measurement basis—was always implicit in the formalism. It has simply never been extracted.

PREDICTIONS

The prediction $\Gamma(\theta) = \Gamma_0 \cos^2 \theta$ becomes testable once we specify parameters. We calculate coherence times for systems within reach of current optomechanical technology.

Diósi-Penrose decoherence rate

For a sphere of mass m and radius R in superposition of separation $d \gg R$ [2, 3]:

$$\Gamma_0 = \frac{Gm^2}{\hbar R} \quad (8)$$

Geometric corrections for finite d/R appear in Ref. [5].

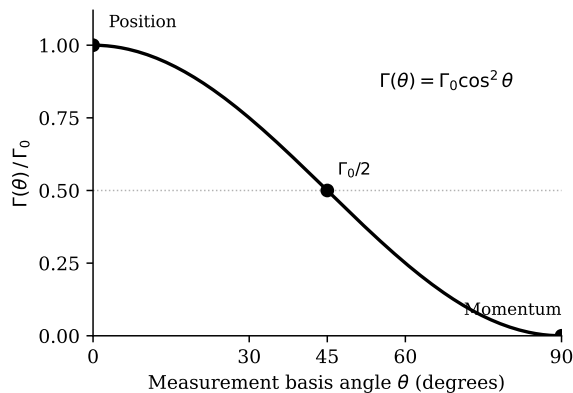


FIG. 1. Predicted decoherence rate versus measurement basis angle. Position measurement ($\theta = 0^\circ$) reveals full gravitational decoherence; momentum measurement ($\theta = 90^\circ$) reveals none. The ratio $\Gamma(45^\circ)/\Gamma(0^\circ) = 1/2$ (equivalently $\tau(45^\circ)/\tau(0^\circ) = 2$) is independent of system parameters.

Reference system

We adopt parameters consistent with recent ground-state cooling experiments [7]:

Parameter	Value
Material	Silica nanosphere
Mass	10^{-12} kg
Diameter	~ 500 nm
Superposition separation	$1 \mu\text{m}$
Mechanical frequency	$\omega_m/2\pi \sim 100$ kHz

Applying the Diósi formula: $\Gamma_0 \approx 7 \text{ s}^{-1}$, giving $\tau_0 = 0.14 \text{ s}$.

Basis-dependent predictions

From Eq. (7):

Measurement	θ	$\Gamma(\theta) [\text{s}^{-1}]$	$\tau(\theta) [\text{s}]$
Position	0°	7.0	0.14
45° quadrature	45°	3.5	0.28
60° quadrature	60°	1.75	0.57
Momentum	90°	0	∞

The critical ratio

The ratio $\tau(45^\circ)/\tau(0^\circ) = 2$ holds for any mass, separation, or Γ_0 . It depends only on $\cos^2(45^\circ) = 1/2$ —the midpoint of Fig. 1. A single measurement at $\theta = 45^\circ$ distinguishes our prediction from any constant-rate model.

Required precision

Distinguishing $\tau(0^\circ)$ from $\tau(45^\circ)$ at 3σ requires coherence time uncertainty $< 20\%$. Current experiments achieve $\sim 10\%$ precision in decoherence measurements [10]. The challenge is not measurement precision but isolation: suppressing environmental decoherence until gravitational effects dominate.

Current experimental status

Ground-state cooling of levitated nanoparticles has been achieved [7]. Sustained macroscopic superpositions at the 0.1 s timescale remain beyond current reach; environmental sources (blackbody radiation, residual gas, photon recoil) presently limit coherence to milliseconds [11]. Progress in cryogenic levitation and coherent scattering protocols suggests gravitationally-limited regimes may become accessible within this decade [8, 12].

Confirmation and refutation

- *Confirmation:* $\tau(\theta)$ varies with basis. A measured ratio $\tau(45^\circ)/\tau(0^\circ) = 2.0 \pm 0.4$ at 3σ would strongly support basis-dependence.
- *Refutation:* $\tau(\theta)$ constant across bases. This confirms standard Diósi-Penrose and rules out our prediction.

Both outcomes constrain fundamental physics. Neither is excluded by current theory.

EXPERIMENTAL PROTOCOL

We propose a protocol to test basis-dependent gravitational decoherence using established optomechanical techniques. The key measurement—correlating coherence decay with homodyne phase—has not previously been performed.

Platform

Levitated dielectric nanospheres provide the required mass ($\sim 10^{-12}$ kg), thermal isolation, and optical read-out [7, 8]. Alternative platforms (trapped ions, superconducting circuits) access smaller masses where gravitational decoherence rates are negligible.

State preparation

Cool the center-of-mass motion to the quantum ground state via coherent scattering [7] or cavity sideband cooling [13]. Create spatial superposition through pulsed optical potentials [14]. Target separation $d \sim 1 \mu\text{m}$.

Quadrature measurement

Optical homodyne detection accesses arbitrary quadratures [15, 16]:

$$\hat{X}_\theta = \hat{x} \cos \theta + \hat{p} \sin \theta \quad (9)$$

The local oscillator phase ϕ_{LO} sets θ . Phase stability $< 1^\circ$ is routine in continuous-variable quantum optics.

Protocol

1. Prepare superposition state $|\psi_0\rangle$
2. Free evolution for time t
3. Homodyne measurement at phase θ
4. Repeat $N \sim 10^3$ times
5. Extract coherence $C(t, \theta)$ from quadrature statistics
6. Scan t ; fit decay curve $C(t) = C_0 e^{-\Gamma(\theta)t}$
7. Scan θ from 0° to 90° in $\sim 15^\circ$ increments

Data analysis

Fit measured rates to:

$$\Gamma(\theta) = \Gamma_{\text{grav}} \cos^2 \theta + \Gamma_{\text{env}} \quad (10)$$

The parameter Γ_{grav} quantifies basis-dependent gravitational decoherence; Γ_{env} captures basis-independent environmental sources.

Discrimination criteria

- *Confirmation*: $\Gamma_{\text{grav}} > 0$ at 3σ , with residuals consistent with $\cos^2 \theta$ functional form.
- *Refutation*: Γ_{grav} consistent with zero; $\Gamma(\theta)$ flat within uncertainty.

Control experiments

1. *Mass scaling*: Repeat with varied particle mass. Gravitational: $\Gamma_{\text{grav}} \propto m^2$. Environmental: weaker mass dependence.
2. *Temperature*: Γ_{env} increases with temperature; Γ_{grav} should not.
3. *Pressure*: Γ_{env} from residual gas scales with pressure; Γ_{grav} remains constant.

Feasibility

Every component exists independently: ground-state cooling [7], spatial superposition creation [14], phase-resolved homodyne detection [15, 16]. The protocol combines established techniques to perform a measurement that tests a prediction not previously considered. No new apparatus is required—only a new analysis correlating θ with Γ .

DISCUSSION

The prediction $\Gamma(\theta) = \Gamma_0 \cos^2 \theta$ has not appeared in prior literature. We address why, and consider the implications of confirmation or refutation.

Why has this not been noticed?

The Diósi-Penrose model was built to explain classicality—the emergence of definite positions from quantum superpositions. Its natural application is position measurement. No theory predicted basis-dependent decoherence, so no experiment tested for it. The gap is conceptual: the question was not asked.

Relation to other collapse models

Continuous spontaneous localization (CSL) and GRW [4, 17] also couple to position. Any collapse model with position-localization will exhibit $\cos^2 \theta$ scaling when measured in rotated quadratures. The prediction is generic to position-coupled dynamics, not specific to gravitational origin.

If confirmed

A positive result would establish that gravitational decoherence depends on measurement basis—the rate Γ is not intrinsic to the quantum state but emerges in the act of observation.

This echoes Wheeler’s delayed-choice experiment [18], now realized in the laboratory [19]: a photon’s wave-or-particle character becomes definite only at detection. Basis-dependent gravitational decoherence would imply an analogous result: gravity’s decohering action becomes definite only when—and in the basis where—we measure. The preferred basis emerges from the conjunction of dynamics and observation, not from gravity alone.

The full implications extend beyond collapse phenomenology. We develop them in subsequent work; here we note only that they exist.

If refuted

A null result— $\Gamma(\theta)$ constant across bases—confirms standard Diósi-Penrose. Gravitational decoherence proceeds at Γ_0 regardless of measurement basis. This rules out observer-dependence and constrains interpretations assigning measurement a constitutive role. Null results are not failures; they sharpen theoretical boundaries.

Limitations

Our derivation assumes the standard Diósi master equation with position coupling. Models with different operator structure would yield different angular dependence. The prediction tests position-coupled collapse specifically.

Realistic homodyne detection introduces inefficiency and noise, contributing to Γ_{env} . These do not alter the $\cos^2 \theta$ signature; basis-dependent gravitational decoherence remains extractable from the fit.

Future directions

Extension to two-dimensional phase space yields $\Gamma(\theta, \varphi)$. Delayed-choice protocols—selecting measurement basis after superposition evolution—test whether basis-dependence is fixed at preparation or at detection. These directions are reserved for subsequent work.

CONCLUSION

Gravitational decoherence rate, we predict, depends on measurement basis: $\Gamma(\theta) = \Gamma_0 \cos^2 \theta$. This follows from the Diósi-Penrose master equation without modification. The question—what does an observer measuring a rotated quadrature see?—was not previously asked.

The prediction yields a clean signature. Sweeping homodyne phase from position ($\theta = 0^\circ$) to momentum ($\theta = 90^\circ$) should reveal coherence times that double at 45° and diverge at 90° . The ratio $\tau(45^\circ)/\tau(0^\circ) = 2$

holds independent of mass, separation, or absolute rate—requiring only relative measurement.

All components exist: ground-state cooling, superposition protocols, phase-resolved detection. The experiment combines established techniques to test a prediction not previously considered. No new apparatus is required.

Confirmation implies gravity’s preferred basis emerges at measurement. Refutation confirms standard Diósi-Penrose. Either outcome advances the quantum-gravitational frontier. The question is now empirical.

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- [1] R. Penrose, On gravity’s role in quantum state reduction, *General Relativity and Gravitation* **28**, 581 (1996).
- [2] L. Diósi, A universal master equation for the gravitational violation of quantum mechanics, *Physics Letters A* **120**, 377 (1987).
- [3] L. Diósi, Models for universal reduction of macroscopic quantum fluctuations, *Physical Review A* **40**, 1165 (1989).
- [4] G. C. Ghirardi, P. Pearle, and A. Rimini, Markov processes in Hilbert space and continuous spontaneous localization of systems of identical particles, *Physical Review A* **42**, 78 (1990).
- [5] S. Donadi *et al.*, Underground test of gravity-related wave function collapse, *Nature Physics* **17**, 74 (2021).
- [6] K. Hornberger, S. Uttenthaler, B. Brezger, L. Hackermüller, M. Arndt, and A. Zeilinger, Collisional decoherence observed in matter wave interferometry, *Physical Review Letters* **90**, 160401 (2003).
- [7] U. Delić, M. Reisenbauer, K. Dare, D. Grass, V. Vuletić, N. Kiesel, and M. Aspelmeyer, Cooling of a levitated nanoparticle to the motional quantum ground state, *Science* **367**, 892 (2020).
- [8] M. Aspelmeyer, T. J. Kippenberg, and F. Marquardt, Cavity optomechanics, *Reviews of Modern Physics* **86**, 1391 (2014).
- [9] [Author], Observer geometry in rotating systems: The Caprazli projection principle, *American Journal of Physics* (2025), submitted.
- [10] A. Vinante *et al.*, Narrowing the parameter space of collapse models with ultracold layered force sensors, *Physical Review Letters* **125**, 100404 (2020).
- [11] C. Gonzalez-Ballester, M. Aspelmeyer, L. Novotny, R. Quidant, and O. Romero-Isart, Levitodynamics: Levitation and control of microscopic objects in vacuum, *Science* **374**, eabg3027 (2021).
- [12] A. Bassi, K. Lochan, S. Satin, T. P. Singh, and H. Ul-

- bricht, Models of wave-function collapse, underlying theories, and experimental tests, *Reviews of Modern Physics* **85**, 471 (2013).
- [13] F. Tebbenjohanns, M. L. Mattana, M. Rossi, M. Frimmer, and L. Novotny, Quantum control of a nanoparticle optically levitated in cryogenic free space, *Nature* **595**, 378 (2021).
- [14] O. Romero-Isart, M. L. Juan, R. Quidant, and J. I. Cirac, Toward quantum superposition of living organisms, *New Journal of Physics* **12**, 033015 (2010).
- [15] M. R. Vanner, I. Pikovski, G. D. Cole, M. S. Kim, v. Brukner, K. Hammerer, G. J. Milburn, and M. Aspelmeyer, Pulsed quantum optomechanics, *Proceedings of the National Academy of Sciences* **108**, 16182 (2011).
- [16] H. M. Wiseman and G. J. Milburn, *Quantum Measurement and Control* (Cambridge University Press, 2010).
- [17] A. Bassi and G. C. Ghirardi, Dynamical reduction models, *Physics Reports* **379**, 257 (2003).
- [18] J. A. Wheeler, The “Past” and the “Delayed-choice” double-slit experiment, in *Mathematical Foundations of Quantum Theory*, edited by A. R. Marlow (Academic Press, 1978) pp. 9–48.
- [19] J.-S. Tang, Y.-L. Li, X.-Y. Xu, G.-Y. Xiang, C.-F. Li, and G.-C. Guo, Realization of quantum Wheeler’s delayed-choice experiment, *Nature Photonics* **6**, 600 (2012).