

A filtered Boris algorithm for charged-particle dynamics in a strong magnetic field

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Outline

- 1. Introduction
- Presentation of algorithms
 Standard Boris algorithm
 Filtered Boris algorithm
 Two-point filtered Boris algorithm
- 3. Results
- 4. Conclusion

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Introduction

Propose and analyse a numerical integrator for the equations of motion of a charged particle

$$\ddot{x}(t) = \dot{x}(t) \times B(x(t), t) + E(x(t), t), \quad x(t), B(x, t), E(x, t) \in \mathbb{R}^3$$
 (1)



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 (1)

in a strong inhomogeneous magnetic field

$$B(x,t) = \frac{1}{\varepsilon} B_0(\varepsilon x) + B_1(x,t) \quad \text{for } 0 < \varepsilon \ll 1$$
 (2)



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Standard Boris algorithm

De facto algorithm for charged particles in a magnetic field.

$$(x^n, v^{n-1/2}) \mapsto (x^{n+1}, v^{n+1/2})$$

$$v_{+}^{n-1/2} = v^{n-1/2} + \frac{h}{2}E(x^{n}, t^{n})$$
(3)

$$v_{-}^{n+1/2} - v_{+}^{n-1/2} = \frac{h}{2} \left(v_{-}^{n+1/2} + v_{+}^{n-1/2} \right) \times B(x^{n}, t^{n})$$
 (4)

$$v^{n+1/2} = v_{-}^{n+1/2} + \frac{h}{2}E(x^n, t^n)$$
 (5)

$$x^{n+1} = x^n + hv^{n+1/2} (6)$$

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Guiding center approximation

$$x^n_\odot = x^n + \frac{v^n \times B(x^n,t^n)}{|B(x^n,t^n)|^2} \quad \text{guiding center approximation}$$

$$\bar{x}^n = \theta^n x^n + (1-\theta^n) x^n_\odot \quad \text{point on the straight line connecting } x \text{ and } x_\odot$$

$$\theta^n = \theta(h|B(x^n,t^n)|)$$

$$\theta(\xi) = \frac{1}{\operatorname{sinc}(\xi/2)^2} \quad \text{with} \quad \operatorname{sinc}(\xi) = \frac{\sin(\xi)}{\xi}$$



Notation

$$B^{n} = B(x^{n}, t^{n})$$
$$\bar{B}^{n} = B(\bar{x}^{n}, t^{n})$$
$$E^{n} = E(x^{n}, t^{n})$$

We will also express $v \times B = -\hat{B}v$, where $B = (b_1, b_2, b_3)$, and so

$$\hat{B} = \begin{pmatrix} 0 & -b_3 & b_2 \\ b_3 & 0 & -b_1 \\ -b_2 & b_1 & 0 \end{pmatrix} \tag{7}$$

$$(x^n, v^{n-1/2}) \mapsto (x^{n+1}, v^{n+1/2})$$

$$v^{n+1/2} = v^{n-1/2} + \frac{h}{2}\Psi(h\hat{B^n})E^n$$
 (8)

$$v_{-}^{n+1/2} = \exp(-h\hat{B}^{n})v^{n-1/2}$$
(9)

$$v^{n+1/2} = v_{-}^{n+1/2} + \frac{h}{2}\Psi(h\hat{B^n})E^n$$
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where $\Psi(\zeta) = \tanh(\zeta/2)$ with $\tanh(\zeta) = \tanh(\zeta)/\zeta$.

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 $\Psi(h\hat{B})$ can therefore be expressed as:

$$\Psi(h\hat{B}) = I + \frac{1 - \operatorname{tanc}(h|B|/2)}{|B|^2}\hat{B}^2$$
 (12)



Since we have

$$v_{-}^{n+1/2} = \exp(-h\hat{B}^{\hat{n}})v^{n-1/2},\tag{9}$$

and

$$\bar{x}^n = \theta^n x^n + (1 - \theta^n) x_{\odot}^n,$$

$$x_{\odot}^n = x^n + \frac{v^n \times B^n}{|B^n|^2},$$

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$$x_{\odot}^n = x^n + \frac{v^n \times B^n}{|B^n|^2},$$

we need to compute v^n with the following approximation,

$$v^{n} = \Phi_{1}(hB^{n}) \frac{v_{-}^{n+1} - v_{+}^{n-1}}{2} - h\Gamma(hB^{n})E^{n},$$
(13)

where $\Phi_1(\zeta) = \frac{1}{\mathrm{sinch}(\zeta)}$ with $\mathrm{sinch}(\zeta) = \frac{\mathrm{sinh}(\zeta)}{\zeta}$, and $\Gamma(\zeta) = \Phi_1(\zeta) - 1$.

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We therefore need to solve the equation with a fixed-point iteration to get v^n .

Variants

When
$$\theta(x,t)=1$$
 \rightarrow $\bar{x^n}=x^n$, and so the algorithm is explicit, we will call it **Exp-A**



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When
$$\theta(x,t)=rac{1}{\sin c(h|B(x,t)|/2)^2}$$
 o algorithm is implicit, we will call it **Imp-A**



Two-Point Filtered Boris Algorithm

$$(x^n, v^{n-1/2}) \mapsto (x^{n+1}, v^{n+1/2})$$

$$v_{+}^{n-1/2} = v^{n-1/2} + \frac{h}{2}\Psi(h\hat{B^n})E^n$$
(14)

$$\Phi_2(h\hat{B}^n_{\odot})(v_-^{n+1/2} - v_+^{n-1/2}) = \frac{h}{2}\Phi_1(h\hat{B}^n)(v_-^{n+1/2} + v_+^{n-1/2}) \times B^n$$
(15)

$$v^{n+1/2} = v_{-}^{n+1/2} + \frac{h}{2}\Psi(h\hat{B^n})E^n$$
(16)

$$x^{n+1} = x^n + hv^{n+1/2} (17)$$

Two-Point Filtered Boris Algorithm

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$$x^{n+1} = x^n + hv^{n+1/2} (17)$$

where $\Psi(\zeta)=\mathrm{tanch}(\zeta/2)$ and $\Phi_1(\zeta)=\frac{1}{\mathrm{sinch}(\zeta)}$ and $\Phi_2(\zeta)=\frac{1}{\mathrm{sinch}(\zeta/2)^2}.$

Two-Point Filtered Boris Algorithm

This method is also implicit.

The fixed-point iteration for x_{\odot}^n requires:

- the evaluation of matrix functions by the Rodriguez formula
- and the solution of a linear system with the 3x3 matrix $\Phi_2(h\hat{B}^n_{\odot}) + \frac{1}{2}h\hat{B}^n\Phi_1(h\hat{B}^n)$.

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Numerical experiment

$$B(x,t) = \nabla \times \frac{1}{\epsilon} (0 \ x_1 \ 0) + \nabla \times (0 \ x_1 x_3 \ 0)$$

$$= \frac{1}{\epsilon} (0 \ 0 \ 1) + (-x_1 \ 0 \ x_3)$$

$$E(x,t) = -\nabla_x U(x)$$

$$U(x) = \frac{1}{\sqrt{x_1^2 + x_2^2}}$$

Initial conditions:

$$x(0) = \left(\frac{1}{3}, \frac{1}{4}, \frac{1}{2}\right)^{\top}$$
$$v(0) = \left(\frac{2}{5}, \frac{2}{3}, 1\right)^{\top}$$
$$0 \le t \le 1$$

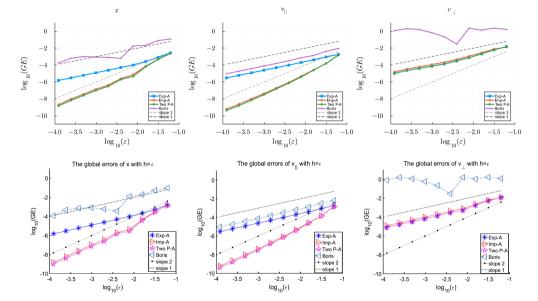


Figure: The logarithm of the global error at t=1 against the logarithm of ϵ for $h=1\epsilon$

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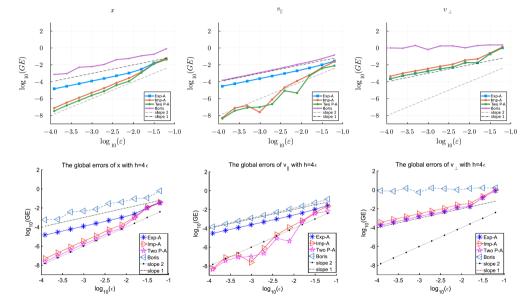


Figure: The logarithm of the global error at t=1 against the logarithm of ϵ for $h=4\epsilon$

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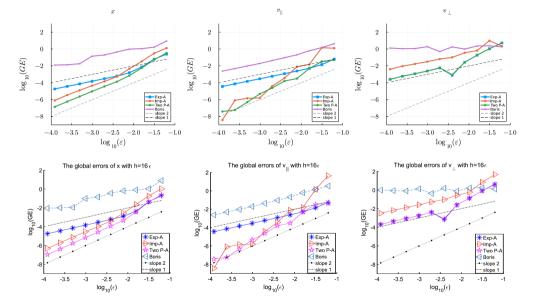


Figure: The logarithm of the global error at t=1 against the logarithm of ϵ for $h=16\epsilon$ Seminar in Physics for CSE

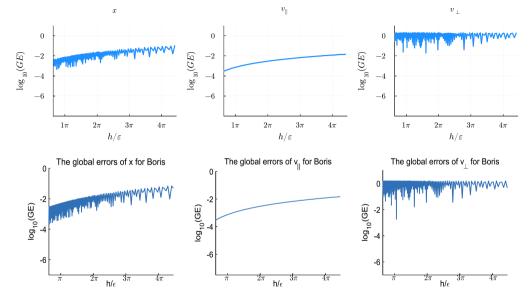


Figure: The logarithm of the global error at t=1 against h/ϵ for $\epsilon=(1/2)^{10}$ and h=1/k, k=60,61,...,600**ETH** zürich

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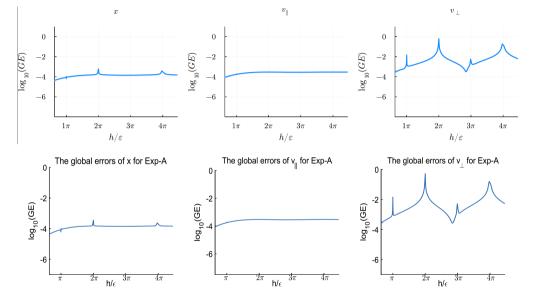


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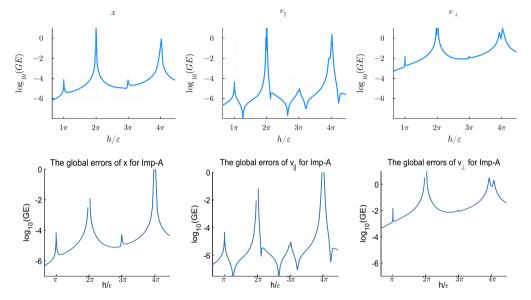


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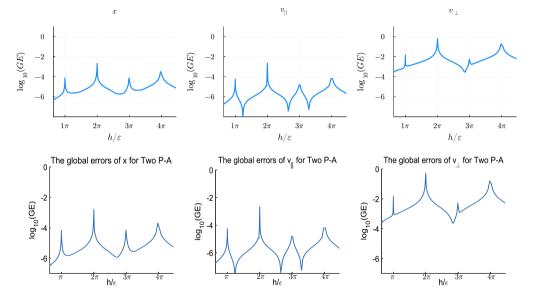


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Challenges

- Reference solution of the ODE was not stated, solved it using RK ode45 equivalent in Julia with appropriate tolerance.
- Most standard libraries/ programming languages (Julia, Python, Matlab but not Mathematica!) implement $\operatorname{sinc}(\zeta)$ as $\operatorname{sin}(\pi\zeta)/(\pi\zeta)$ which gave incorrect results since the authors used $\operatorname{sin}(\zeta)/\zeta$.
- Derivation of the efficient computation of filter functions using a Rodriguez type formula was included in the Appendix of the paper but does not include intermediate steps, and not much literature is found regarding this.
- This and the miscalculation of sinc lead me to believe that the derivation of the matrix filter functions using Rodriguez like formula was incorrect.
- Which in turn lead me to try other computation methods for the matrix functions using matrix identities which yielded correct results.
- However, with one concrete example, the derivation of the rest of the filter functions using the Rodriguez like formula is straightforward and I was able to verify the results and derive the not explicitly stated filter functions.

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Conclusion

- It is observed that all three filtered Boris methods improve considerably over the standard Boris method, and the optimally filtered methods Imp-A and Two P-A show second order, whereas method Exp-A only shows first order.
- Methods Imp-A and Two P-A behave very similar away from stepsize resonances, but method Two P-A appears more robust near stepsize resonances.

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Thank you for your attention!

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