Weyl to go

Weyl spinor definitions

$$\begin{split} \psi_A &= \begin{pmatrix} \psi_1 \\ \psi_2 \end{pmatrix}, \quad \bar{\psi}_A &= \begin{pmatrix} \bar{\psi}_1 \\ \bar{\psi}_2 \end{pmatrix} \\ \psi^A &= \epsilon^{AB} \psi_B = \begin{pmatrix} \psi_2 \\ -\psi_1 \end{pmatrix}^T \\ \bar{\psi}^{\dot{A}} &= \epsilon^{\dot{A}\dot{B}} \bar{\psi}_B = \begin{pmatrix} \bar{\psi}_2 \\ -\bar{\psi}_1 \end{pmatrix}^T \\ \epsilon_{AB} &= \epsilon_{\dot{A}\dot{B}} = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} \\ \epsilon^{AB} &= \epsilon^{\dot{A}\dot{B}} = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \\ (\psi_A)^\dagger &= \bar{\psi}_{\dot{A}} \\ \psi\chi &\equiv \psi^A \chi_A = \psi_2 \chi_1 - \psi_1 \chi_2 = \psi^2 \chi^1 - \psi^1 \chi^2 \\ \bar{\psi}\bar{\chi} &\equiv \bar{\psi}_{\dot{A}} \bar{\chi}^{\dot{A}} = -\bar{\psi}_2 \bar{\chi}_1 + \bar{\psi}_1 \bar{\chi}_2 = -\bar{\psi}^2 \bar{\chi}^1 + \bar{\psi}^1 \bar{\chi}^2 \\ \psi^2 &= -2\psi_1 \psi_2 \\ \bar{\psi}^2 &= 2\bar{\psi}_1 \bar{\psi}_2 \\ \sigma^\mu &= (1, \vec{\sigma}) \\ \bar{\sigma}^\mu &= (1, -\vec{\sigma}) \\ \sigma^{\mu\nu} &= \frac{i}{4} (\sigma^\mu \bar{\sigma}^\nu - \sigma^\nu \bar{\sigma}^\mu) \\ \sigma_1 &= \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_2 &= \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \sigma_3 &= \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \end{split}$$

Weyl spinor contractions

$$\begin{split} \eta\psi &= \psi\eta \\ \bar{\eta}\bar{\psi} &= \bar{\psi}\bar{\eta} \\ (\eta\psi)^\dagger &= \bar{\psi}\bar{\eta} \\ (\eta\psi) (\eta\phi) &= -\frac{1}{2} (\eta\eta) (\psi\phi) \\ \eta\sigma^\mu\bar{\psi} &= -\bar{\psi}\bar{\sigma}^\mu\eta \\ (\sigma^\mu\bar{\eta})_A(\eta\sigma^\nu\bar{\eta}) &= \frac{1}{2} g^{\mu\nu}\eta_A(\bar{\eta}\bar{\eta}) \\ (\eta\sigma^\mu\bar{\eta}) (\eta\sigma^\nu\bar{\eta}) &= \frac{1}{2} (\eta\eta) (\bar{\eta}\bar{\eta}) \\ (\eta\sigma^\mu\partial_\mu\bar{\psi}) (\eta\psi) &= -\frac{1}{2} (\psi\sigma^\mu\partial_\mu\bar{\psi}) (\eta\eta) \\ (\partial_\mu\sigma^\mu\bar{\eta}) (\bar{\eta}\bar{\psi}) &= -\frac{1}{2} (\partial_\mu\psi\sigma^\mu\bar{\psi}) (\bar{\eta}\bar{\eta}) \\ (\bar{\eta}\bar{\psi}) (\eta\sigma^\mu\bar{\eta}) (\eta\psi) &= \frac{1}{4} (\eta\eta) (\bar{\eta}\bar{\eta}) (\psi\sigma^\mu\bar{\psi}) \\ \eta\sigma^{\mu\nu}\psi &= -\psi\sigma^{\mu\nu}\eta \end{split}$$

Charge conjugation

$$\Gamma^{r} = \left\{ \mathbb{1}, \gamma^{5}, \gamma^{\mu}, \gamma^{5} \gamma^{\mu}, \gamma^{\mu\nu} \right\}$$

$$C^{\dagger} = C^{-1}$$

$$C^{T} = -C$$

$$C^{-1}\Gamma^rC = \eta^r\Gamma^{rT}$$

$$\eta^r = \begin{cases} 1 & 1 \leq r \leq 6 \\ -1 & 7 \leq r \leq 16 \end{cases}$$

For w = u, v and w' = v, u

$$w = C\bar{w}'^T$$

$$\bar{w}_1 \Gamma^r w_2 = -\eta^r \bar{w}_2' \Gamma^r w_1'$$

Weyl-Dirac translation

$$\begin{split} \Psi_D &= \begin{pmatrix} \psi_L \\ \bar{\psi}_R \end{pmatrix}, \quad \chi_M = \begin{pmatrix} \chi \\ \bar{\chi} \end{pmatrix} \\ \gamma^\mu &= \begin{bmatrix} 0 & \sigma^\mu \\ \bar{\sigma}^\mu & 0 \end{bmatrix} \\ \gamma^5 &= i \gamma^0 \gamma^1 \gamma^2 \gamma^3 = \frac{i}{4!} \epsilon^{\mu\nu\rho\sigma} \gamma_\mu \gamma_\nu \gamma_\rho \gamma_\sigma \\ \bar{\Psi} &\equiv \Psi^\dagger \gamma^0 = \begin{pmatrix} \psi_R \\ \bar{\psi}_L \end{pmatrix}^T. \\ \bar{\Psi}_M \gamma^\mu P_{L/R} \Phi_M &= -\bar{\Phi}_M \gamma^\mu P_{R/L} \Psi_M \end{split}$$

Chiral identities

$$\Psi_{L/R} = \frac{1}{2} \left(\mathbb{1} \mp \gamma^5 \right)$$
$$(\psi_R \phi_L) = \bar{\Psi} P_L \Phi$$
$$(\bar{\psi}_L \bar{\phi}_R) = \bar{\Psi} P_R \Phi$$
$$(\bar{\psi}_L \bar{\sigma}^\mu \phi_L) = \bar{\Psi} \gamma^\mu P_L \Phi$$
$$(\psi_R \sigma^\mu \bar{\phi}_R) = \bar{\Psi} \gamma^\mu P_R \Phi$$

Dirac algebra

(Both in 4 and
$$d=4-\epsilon$$
 dimensions)
$$\{\gamma^{\mu}, \gamma^{\nu}\} = 2g^{\mu\nu}\mathbb{1}_{4\times4}$$

$$\{\gamma^{\mu}, \gamma^{5}\} = 0$$

$$(\gamma^{\mu})^{\dagger} = \gamma^{0}\gamma^{\mu}\gamma^{0}$$

$$(\gamma^{5})^{\dagger} = \gamma^{5}$$

$$\gamma^{\mu\nu} = \frac{i}{2}\left(\gamma^{\mu}\gamma^{\nu} - \gamma^{\nu}\gamma^{\mu}\right)$$

$$M_{n}^{\mu_{1}\dots\mu_{n}} \equiv \gamma^{\nu}\gamma^{\mu_{1}}\dots\gamma^{\mu_{n}}\gamma_{\nu}$$

$$M_{n+1}^{\mu_{1}\dots\mu_{n+1}} = 2\gamma^{\mu_{n+1}}\gamma^{\mu_{1}}\dots\gamma^{\mu_{n}} - M_{n}^{\mu_{1}\dots\mu_{n}}\gamma^{\mu_{n+1}}$$

$$M_{0} = d$$

$$M_{1}^{\mu} = (2-d)\gamma^{\mu}$$

$$M_{2}^{\mu\nu} = 4g^{\mu\nu}\mathbb{1}_{4\times4} + (d-4)\gamma^{\mu}\gamma^{\nu}$$

$$M_{2}^{\mu\nu\rho} = (4-d)\gamma^{\mu}\gamma^{\nu}\gamma^{\rho} - 2\gamma^{\rho}\gamma^{\nu}\gamma^{\mu}$$

Traces

(Both in 4 and $d = 4 - \epsilon$ dimensions.)

$$T_n^{\mu_1 \dots \mu_n} = \text{Tr} \left\{ \gamma^{\mu_1} \dots \gamma^{\mu_n} \right\}$$

$$T_{2n+1}^{\mu_1 \dots \mu_{2n+1}} = 0$$

$$T_{2n}^{\mu_1 \dots \mu_{2n}} = \sum_{i=2}^{2n} (-1)^i g^{\mu_1 \mu_i} T_{2n-2}^{\mu_2 \dots \mu_{i-1} \mu_{i+1} \dots \mu_{2n}} \qquad n \ge 1$$

$$T_0 = 4$$

$$T_2^{\mu\nu} = 4g^{\mu\nu}$$

$$T_4^{\mu\nu\rho\sigma} = 4 \left(g^{\mu\nu} g^{\rho\sigma} - g^{\mu\rho} g^{\nu\sigma} + g^{\mu\sigma} g^{\nu\rho} \right)$$

$$\tilde{T}_{n}^{\mu_{1}\dots\mu_{n}} = \operatorname{Tr}\left\{\gamma^{5}\gamma^{\mu_{1}}\dots\gamma^{\mu_{n}}\right\}$$

$$\tilde{T}_{2n+1}^{\mu_{1}\dots\mu_{2n+1}} = 0$$

$$\tilde{T}_{2n}^{\mu_{1}\dots\mu_{2n}} = \sum_{i=2}^{2n} (-1)^{i} g^{\mu_{1}\mu_{i}} \tilde{T}_{2n-2}^{\mu_{2}\dots\mu_{i-1}\mu_{i+1}\dots\mu_{2n}} \qquad n \geq 3$$

$$\tilde{T}_{0} = 0$$

$$\tilde{T}_{2}^{\mu\nu} = 0$$

$$\tilde{T}_{2}^{\mu\nu\rho\sigma} = -i\epsilon^{\mu\nu\rho\sigma}$$