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Phenomenological Study of Heavy Neutrinos at the LHC Using the Vector Boson Fusion Technique

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Chapter 1

Introduction

The standard model (SM) gathers the entire understanding about fundamental particles and their interactions. Although the model has successfully explained various physical phenomena observed experimentally, there are still multiple unanswered questions concerning particle physics. For example, experiments [1] have shown that accelerator and reactor, solar, and atmospheric neutrinos have mass by proving the existence of neutrino oscillations. The fact that there are neutrino oscillations contradicts the SM, because this model predicts that neutrinos are massless. Some specific experiments for each neutrino category are: Super-Kamiokande [2] for solar and atmospheric neutrino oscillations, KamLAND [3] for reactor neutrinos, and K2K [4] for accelerator neutrino oscillations [5]. An additional open question about neutrinos is the fact that only neutrinos with left helicity have been observed. Helicity is defined as the projection of the particle's momentum vector over its spin direction. Therefore, only neutrinos with spin anti-parallel to its linear momentum have been detected.

In order to provide neutrinos with mass, several theories that extend the predictions of the SM have been proposed. One of the most known models is the "see-saw" or balance mechanism [6]. This mechanism postulates the existence of a yet undetected particle called the heavy neutrino. This heavy neutrino would have a mass inversely proportional to the one of the neutrino, thus making the mass of the heavy neutrino considerably large. Also, this neutrino would have a right he-

licity therefore restoring the right-left symmetry in the standard model. Other models that try to provide mass to the neutrinos by extending the SM postulate the conservation of the $(B - L)$ number, where B is the baryon number and L is the lepton number. Unlike these models, the see-saw mechanism proposes a breaking in the $(B - L)$ symmetry with consequences discussed in the next paragraph.

The see-saw mechanism includes three sub-models that provide mass to neutrinos. These three models come from the fact that is necessary to take into account the effects of breaking the $(B - L)$ symmetry. This effects can be parametrized using a Weinberg operator of the form $\lambda_{ll'} L_l L_{l'} \Phi \Phi / \Lambda$ where $\Phi = (\phi^+, \phi^0)^T$ is the doublet associated with the SM Higgs Boson and $L_l = (\nu_l, l)_L^T$ the representation of a doublet field associated with the lepton number $+1$ [6]. Since there are only three ways in which this Weinberg operator can be obtained at tree-level, there are also three types of the see-saw mechanism. In the type I see-saw mechanism the product between L_l and Φ results in a fermionic singlet state. In the type II see-saw mechanism, the product between L_l and $L_{l'}$ forms a scalar triplet. Finally, the product between L_l and Φ in the type III sub-model results in a fermionic triplet state. The state formed by the products described in each see-saw model would correspond to the definition of heavy neutrino in each sub-model. Taking into account the elements described above regarding the see-saw mechanism, if heavy neutrinos are observed, the left and right symmetry in the SM would be restored and the mechanism by which the neutrinos acquire mass would be explained.

Heavy neutrinos searches have been conducted in multiples experiments, but none has been able to prove that heavy neutrinos exist. Examples of these experiments can be found in collaborations such as LEP [7], CMS and ATLAS [8]. Searches at CMS and ATLAS have focused in final states with associated leptons and jets. However, these searches until now have not included the Vector Boson Fusion topology.

The main objective of this monograph is to perform a phenomenological study about the feasibility of conducting an experimental analysis for the detection of heavy neutrinos in the Large Hadron Collider (LHC) using a technique known as vector boson fusion (VBF). This technique has been recently used in the LHC [9] in searches for new physics. In high

energy physics, the bosons W^\pm , Z^0 and γ are known as vector bosons. The process of vector boson fusion occurs through an electroweak interaction of associated quarks with the LHC proton beams. The VBF topology consists in requiring two highly energetic jets in the longitudinal region of the detector and in opposite hemispheres thereof. It has been shown that by requiring this type of event, the noise level (background) is reduced considerably in regions of difficult study in searches of new physics.

In order to conduct the analysis, it is important to simulate signal and background processes and to perform a detailed physical study of the variables that allow to distinguish signal for experimental noise. It is necessary to use a quantitative estimator commonly known as figure of merit to determine optimal cuts in the mentioned variables. The latter with the objective of reducing the amount of experimental noise by finding the optimal cuts in the relevant variables. Furthermore, it is important to establish the expected experimental sensitivity using maximum likelihood limits or the calculation of the final significance for different hypothetical signal points. The procedure described would allow to conclude whether a study for the detection of heavy neutrinos at the LHC is feasible or not.

Chapter 2

Objectives

2.1 General Objectives

Conduct a phenomenological study to determine the possible experimental sensitivity of heavy neutrino searches in the LHC, using the VBF topology, in channels with high-mass resonance production.

2.2 Specific Objectives

- Develop the signal events and experimental noise simulations using MadGraph, Pythia, and Delphes software.
- Write an analysis code using ROOT software to analyze the simulated data.
- Conduct a physical study of the appropriate cinematic and topological variables that show strong separation between signal and background.
- Find the optimal cut points of the relevant physical variables using a significance figure.
- Conduct a statical analysis of the results.

Chapter 3

Definitions

3.1 Jet

A jet can be considered as a collimated spray of stable particles that arises from the fragmentation and hadronization of a quark or a gluon [10]. Since a jet is a shower of particles, algorithms of reconstruction have to be defined to gather all the information from the shower and define a jet as single element observed in the detector. The most common algorithms that define jets assume that the particle shower will arrive in a conical region in the detector. Hence, this algorithms establish a cluster in the $\eta - \phi$ space that defines where all the particles that make up the jet would arrive.

A schematic description of the production and definition of a jet is presented in Figure 3.1. The diagram shows how, as a result from the proton-proton collision, gluons and quarks are generated. Then, as described later in Section 7.1, the partons hadronize resulting in stable particles that are then observed in the detector. It can be seen in the figure the importance of defining correctly the cone of the jet. If the cone radius is too small, some particles that came from the initial parton may be left out, but if the radius is too large some particles that came from a different parton can be mistaken as part of the jet.

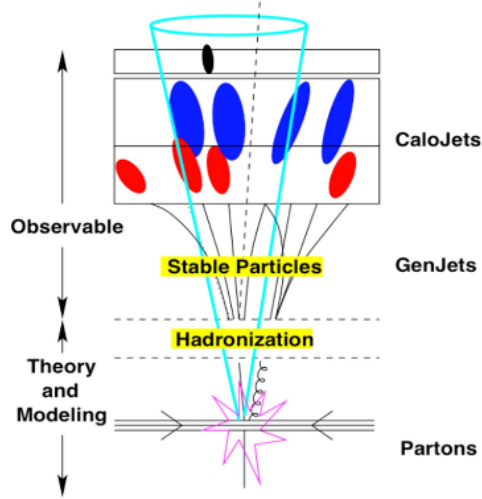


Figure 3.1: Schematic diagram of the creation and definition of a jet. (Taken from [10])

Variations of the cone algorithms are defined to gain more precision in the jet definition. Being able to reconstruct correctly all the particles making up the jet is important, because the jet provides a link between the observed particles in the detector and the physics at the partonic level [10]. In other words, since gluons and quarks decay rapidly, the only way to obtain information in detectors about processes related to partons are through jets. This is why more sophisticated algorithms to fully reconstruct jets in detectors and simulations are developed. For example, the iterative cone with progressive removal algorithm, or IC-PR algorithm, is used in CMS and the software Pythia [10]. The algorithm starts by finding the cell in the detector with largest transverse momentum and makes it a seed. Then, a cone is defined around the seed and a trial jet axis is calculated. If the trial jet axis is the same as the axis of the cone defined from the seed, then the cone is labelled as stable and all the particles that lie inside the cone are defined as the jet. If the trial jet axis does not equal the seed axis, the trial jet axis is defined as the new seed axis. The process described above is repeated until a

convergence of the axes occurs. This algorithm is repeated until there are no seeds above a threshold energy left [10].

3.2 Variable Definitions

The transverse momentum or p_T , is defined as the momentum component that a particle has in the plane perpendicular to the beam line. In the coordinate system of the LHC, the transverse plane corresponds to the $x - y$ plane.

The variable related with the polar angle in the LHC is called pseudorapidity, or η , defined as in Equation 3.1. The use of this variable is justified for mainly two reasons. The first one is that $\Delta\eta$, contrary to $\Delta\theta$, is a Lorentz invariant. This makes $\Delta\eta$ a more natural variable than $\Delta\theta$ for relativistic calculations. The second reason is that the distribution of the values of η in the barrel region, where the multiplicity of particle is less than in the end-caps, is wider allowing the η particle distribution to be approximately constant.

$$\eta = -\ln \left[\tan \left(\frac{\theta}{2} \right) \right] \quad (3.1)$$

Sometimes it is necessary to determine the total angular separation between two particle in the detector, that is, determining the separation in both ϕ and η . That is why the variable ΔR , shown in Equation 3.2, is defined. As stated earlier, it is generally used to determine how close in the detector two particle have been observed. Taking this into account, ΔR is useful to establish whether two particles may have overlapped detection points in the detector.

$$\Delta R = \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2} \quad (3.2)$$

Another useful variable defined in particle physics is the missing transverse energy or \cancel{E}_T . Given that the protons in the LHC collide head to head in the z axis, the initial momentum in the x and y components is zero. Since momentum in all directions must be conserved during the collision, the momentum after the collision must also be zero. However, there are particles resulting from the proton-proton collision,

like the neutrino, that cannot be observed by the components in the detector. That is why the energy corresponding to these unobserved particles is going to be missing from the energy observed in the detector. Considering what has been said in this paragraph, the momentum equation in the $x - y$ plane must be

$$\sum_{i=1}^m p_T(i)^{observed} + \sum_{j=1}^n p_T(j)^{missing} = 0$$

Taking the last equation into account, it is natural to define the missing transverse energy as in Equation 3.3, i.e. the negative of the sum of transverse momenta of the observed particles in the detector.

$$\cancel{E}_T = - \sum_{i=1}^m p_T(i)^{observed} \quad (3.3)$$

With the idea of exploiting the possible difference between signal and background in the p_T for jets and τ 's, two new variables shown in Equations 3.4 and 3.5 were defined to check for possible further separation between signal and background. As shown in equation 3.4, the H_T variable is defined as the scalar sum of the jets in the event that are not B-jets. S_T is defined as the scalar sum of jets that fulfill the same conditions of H_T , added to the p_T of the τ 's in the event.

$$H_T = \sum_{i=1}^n p_T(jet_i) \quad (3.4)$$

$$S_T = \sum_{i=1}^n p_T(jet_i) + \sum_{j=1}^m p_T(\tau_j) \quad (3.5)$$

Throughout the analysis it is important to establish a quantitative estimator to define the amount of signal compared to the amount of background. For that purpose a quantitative estimator commonly known as figure of merit to determine optimal cuts in the relevant variables is defined. This figure of merit allows to find the optimal cuts in the relevant variables by giving a quantitative result of the amount of signal events compared to the number of background events. For this

particular analysis, the significance formula that was used is the one shown in Equation 3.6, where S is the significance, $N(s)$ is the number of signal events, and $N(B)$ is the number of background events.

$$S = \frac{N(s)}{\sqrt{N(s) + N(B)}} \quad (3.6)$$

Chapter 4

The Neutrino Mass

The Standard Model predicts that the neutrino is a massless particle. However, experiments like the ones mentioned in Chapter 1 have proven that neutrino oscillations exist. These oscillations show that, contrary to the predictions of the Standard Model, neutrinos are particles with mass. In this chapter, the mechanisms by which neutrinos can acquire mass in the Standard Model are given followed by the reasons that do not allow to add mass to the neutrinos into the SM. Also, the See Saw mechanism is described in greater detail.

4.1 The L and B symmetries

The baryon number was first suggested by Ernst Stueckelberg in 1938 to explain why the proton does not decay into a positron and a gamma ray [11]. Stueckelberg assigned a baryon number of $B = 1$ to the proton and the neutron and $B = -1$ to the antiproton and antineutron. Leptons, photons, and mesons are assigned $B = 0$. The definition of the baryon number was later generalized to be as the one shown in Equation 4.1, where n_q is the number of quarks and $n_{\bar{q}}$ is the number of antiquarks.

$$B = \frac{1}{3} (n_q - n_{\bar{q}}) \quad (4.1)$$

The baryon number conservation states that the baryon number in a process must be conserved. That is $\sum B = \text{const.}$ This conservation

law can arise more formally from the global gauge transformation shown in Equation 4.2 [11].

$$\psi' = \psi e^{i\epsilon B} \quad (4.2)$$

In 1953, E. J. Konopinski and H.M. Mahmoud introduced the lepton number L . They assigned $L = 1$ to e^- , μ^- , ν_e and ν_μ , $L = -1$ to the antileptons, and $L = 0$ to all the other particles. This lepton number was later generalized to the τ and the ν_τ . The lepton number conservation, as the baryon number conservation, can also arise from a global gauge transformation like the one shown in Equation 4.3 [11].

$$\psi' = \psi e^{i\epsilon L} \quad (4.3)$$

4.2 The Dirac Mass

A standar Dirac free fermion can be described using the Langrangian shown is Equation 4.4 [12], where $\not{\partial} = \gamma^\mu \partial_\mu$ and m is the mass of the particle. Since the mass term in this Langragian is $-m\bar{\psi}\psi$, a mass term of this form is always called a Dirac mass term. The best known description for a neutrino is the a neutral standard Dirac particle. That is, a neutrino can be described by the Langrangian showed in Equation 4.4, where ψ would be four component Dirac field ν describing the neutrino.

$$\mathcal{L} = \bar{\psi} (i\not{\partial} - m) \psi \quad (4.4)$$

If the field ψ is decomposed into the corresponding left and right chiral states, the Dirac mass term can be written as

$$-m(\overline{\psi_L + \psi_R})(\psi_L + \psi_R)$$

As shown in Appendix A, the terms $\bar{\psi}_R\psi_R$ and $\bar{\psi}_L\psi_L$ are zero. Therefore, the Dirac mass term can be written as shown in Equation 4.5.

$$-m(\bar{\psi}_L\psi_R + \bar{\psi}_R\psi_L) \quad (4.5)$$

For the particular case of the neutrinos, and taking into account the three neutrino flavors or generations, ψ_L and ψ_R would correspond to the ν_L and ν_R shown in Equation 4.6 [12].

$$\nu_L = \begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix}_L, \quad \nu_R = \begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix}_R \quad (4.6)$$

In the generalization for the three neutrino generations, the Dirac mass Lagrangian, or Dirac mass term, is like the one shown in Equation 4.7. The matrix \mathcal{M}_D is in general a 3×3 complex mass matrix [12].

$$\mathcal{L}_{mass}^D = -(\bar{\nu}_R \mathcal{M}_D \nu_L + \bar{\nu}_L \mathcal{M}_D \nu_R) \quad (4.7)$$

4.3 The Majorana Mass

The Majorana neutrino (ν_M) is a neutrino proposed by Majorana in 1937 that satisfies the property shown in Equation 4.8, where \mathbf{C} is the charge conjugation operator [12]. This charge conjugation operator transforms a free neutrino state into an anti-neutrino state. Taking the latter into account, Equation 4.8 implies that the Majorana neutrino is its own anti-particle.

$$\nu_M = \nu_M^C \equiv \mathbf{C} \nu_M \mathbf{C}^{-1} \quad (4.8)$$

As in the case of a Dirac fermion, a Lagrangian like the shown in Equation 4.9 can be defined. The factor $1/2$ is added to take into account double counting when an interaction term is added to the Lagrangian.

$$\mathcal{L} = \frac{1}{2} \bar{\nu}_M (i \not{\partial} - m) \nu_M \quad (4.9)$$

Once a neutrino that is its own anti-particle is defined, additional terms of the form $\overline{\nu^C} \nu^C$, $\overline{\nu^C} \nu$, and $\bar{\nu} \nu^C$ can be added to the mass term shown in Equation 4.5 [12]. The addition of these terms does not violate charge conservation, because neutrinos are neutral particles. Hence, the Lorentz invariance in the Lagrangian is being conserved. The term $\overline{\nu^C} \nu^C$

is identical to $\bar{\nu}\nu$ except for an irrelevant surface term, so it would not be necessary to include it in the mass term in Equation 4.5 [12]. However, the other two terms mentioned above are not included in this mass term, so they should be added to the general mass lagrangian. These two new mass terms, taking into account an hypothetical neutrino with right helicity, would be like the ones shown in Equations 4.10 and 4.11.

$$\mathcal{L} \sim m_L \left(\overline{\nu_L^C} \nu_L + \bar{\nu}_L \nu_L^C \right) \quad (4.10)$$

$$\mathcal{L} \sim m_R \left(\overline{\nu_R^C} \nu_R + \bar{\nu}_R \nu_R^C \right) \quad (4.11)$$

If a Majorana neutrino is constructed as $\nu_M \equiv \nu_L + \nu_L^C$, Equation 4.10 can be written like in Equation 4.12. A similar definition can be made with an hypothetical neutrino with right helicity, so Equation 4.11 can also be written as 4.12. As in the Dirac case, Equation 4.12 is defined as the Majorana mass term. However, contrary to the Dirac mass term, the Majorana mass term can be constructed using either only ν_L or ν_R [12].

$$\mathcal{L} \sim m \bar{\nu}_M \nu_M \quad (4.12)$$

The Lagrangians shown in Equations 4.10 and 4.11 are valid for only one neutrino flavor. If the mass Lagrangian is generalized to the three neutrino generations as in the Dirac case, the Majorana mass Lagrangian would be the one shown in Equation 4.13. In this equation, h.c. stands for the hermitian conjugate of the other term in the Lagrangian and ν_L is the same as the one defined in Equation 4.6. As in the Dirac case, \mathcal{M}_M is, in general, a complex 3×3 matrix with the additional property that this matrix is always symmetric, i.e. $\mathcal{M}_M = \mathcal{M}_M^T$ [12].

$$\mathcal{L}_{mass}^M = -\frac{1}{2} \overline{\nu_L^C} \mathcal{M}_M \nu_L + \text{h.c.} \quad (4.13)$$

4.4 Neutrino Mass in the Standard Model

As mentioned in Chapter 1, the Standard Model only includes a neutrino with left chirality. As shown in Equation 4.5, the absence of a

neutrino with right chirality makes that all the terms in the Dirac mass Lagrangian vanish. Therefore, it is not possible to add mass to the neutrinos in the Standard Model using a Dirac mass term.

Another mechanism that could provide mass to the neutrino in the Standard Model is the Majorana mass term. However this term would not conserve the Lepton number, i.e. would violate the L symmetry, that is conserved throughout the Standard Model[13]. This happens because the Majorana fermions, in this case neutrinos, are their own antiparticles. A lepton number of $L = +1$ is assigned to a fermion and $L = -1$ to its antiparticle. Since the Majorana neutrinos are their own antiparticles, they do not have a well defined lepton number (it could be either $+1$ or -1). That is why the lepton number for both Majorana neutrinos would be either $L = 1$ or $L = -1$, so the Majorana mass term would violate the lepton conservation number by $\Delta L = \pm 2$. Also, some non-perturbative effects in the Standard Model can violate L symmetry but conserve the current $B - L$, where B is the baryon number. To summarize, since the Majorana terms violate both $B - L$ and L symmetries, these terms can not be used to introduce mass in the Standard Model using perturbation theory or non-perturbative effects [13].

Since these two mechanisms are not useful to explain the origin of the neutrino mass inside the Standard Model, a minimal extension to the SE is proposed in order to provide the neutrinos with mass. This extension consists in inserting right handed neutrinos to the model to explain the origin of the neutrino mass. In principle, the addition of this right handed neutrino would make the Dirac mass term different from zero, and the neutrino mass would be of the form of the charged leptons. Since the mass terms of the leptons is of the form $m = Yv/\sqrt{2}$, where Y is the Yukawa coupling constant, this constant is required to be $Y \simeq 10^{-11}$ or less to explain the small mass of the neutrino. Since the introduction of such a small coupling would require a symmetry reason for the small value, only introducing the right handed neutrino would not solve entirely the neutrino mass problem [14]. That is why an introduction of a model that explains both the neutrino mass and its small value must be introduced.

4.5 The See Saw Mechanism

The See-Saw mechanism assumes the existence of a neutrino with right helicity. With the existence of this neutrino with right helicity and combining the Dirac and Majorana neutrinos mass terms discussed in the previous sections, the resulting mass Lagrangian would be the shown in Equation 4.14, where m_D is a Dirac mass, m_L and m_R are Majorana masses, and h.c is the hermitian conjugate.

$$\mathcal{L}_{mass} = -m_D \bar{\nu}_L \nu_R - \frac{1}{2} \left(m_L \bar{\nu}^C_L \nu_L + m_R \bar{\nu}^C_R \nu_R \right) + \text{h.c} \quad (4.14)$$

If the vector ν is defined as

$$\nu \equiv \begin{pmatrix} \nu_L \\ \nu_R^C \end{pmatrix}$$

then the Lagrangian in Equation 4.14, can be written as in Equation 4.15, where \mathcal{M} is defined in Equation 4.16.

$$\mathcal{L}_{mass} = -\frac{1}{2} \bar{\nu}^C \mathcal{M} \nu + \text{h.c} \quad (4.15)$$

$$\mathcal{M} = \begin{pmatrix} m_L & m_D \\ m_D & m_R \end{pmatrix} \quad (4.16)$$

If matrix \mathcal{M} is diagonalized, the two left handed neutrino mass eigenstate Majorana neutrinos, ν_1 and ν_2 , can be expressed as in Equation 4.17 where θ is the mixing angle. A more detailed calculation of the mass eigenstate Majorana neutrinos, can be found in Appendix B.

$$\begin{aligned} \nu_L &= \nu_1 \cos \theta + \nu_2 \sin \theta \\ \nu_R^C &= -\nu_1 \sin \theta + \nu_2 \cos \theta \end{aligned}$$

or

$$\begin{aligned} \nu_1 &= \nu_L \cos \theta - \nu_R^C \sin \theta \\ \nu_2 &= \nu_L \sin \theta + \nu_R^C \cos \theta \end{aligned} \quad (4.17)$$

To guarantee that the mass eigenvalues are positive, and insertion of i , the imaginary number, in the equation for ν_1 in Equation 4.17 is necessary. With this insertion, ν_1 would now be defined as

$$\nu_1 = i\nu_L \cos \theta - i\nu_R^C \sin \theta$$

Using the calculations performed in Appendix B, the mixing angle value would be the one shown in Equation 4.18, and the two mass eigenvalues would be the ones shown in Equation 4.19 [12].

$$\tan 2\theta = \frac{2m_D}{m_R - m_L} \quad (4.18)$$

$$\begin{aligned} m_1 &= \frac{1}{2} \sqrt{4m_D^2 + (m_R - m_L)^2} - \frac{m_R + m_L}{2} \\ m_2 &= \frac{1}{2} \sqrt{4m_D^2 + (m_R - m_L)^2} + \frac{m_R + m_L}{2} \end{aligned} \quad (4.19)$$

The values of m_D , m_R , and m_L for the See Saw mechanism are values in which m_R is much larger than m_D and m_L . That is, $m_R \gg m_D, m_L$. Furthermore, the value of m_L is chosen to be zero. The latter with the objective of leaving the standard weak interaction theory unchanged, since $m_L \neq 0$ would imply the addition of a left Majorana neutrino. Setting $m_L = 0$ and making the approximation $m_R \gg m_D$ in Equation 4.19, the values for m_1 and m_2 are the ones shown in Equation 4.20. On one hand, it can be seen from Equation 4.20 that if m_R has a large value, m_2 would also have a large value. On the other hand, if m_R has a large value, m_1 has a small value. Because of this inverse proportionality between m_1 and m_2 this mechanism is called the See Saw mechanism. Also, and more important, since these eigenvalues are the ones that are actually observed in measurements, the See Saw mechanism would successfully explain the origin of the neutrino mass as well as the small values for the neutrino mass observed in experiments. In this mechanism, m_D is the mass scale associated with the Standard Model, whereas m_R is the scale provided by models beyond the SM [12]. It is also noteworthy that for the values of m_R and m_D considered, the mixing angle approaches to zero. This fact indicates that the eigenstates ν_1 and ν_2 would be completely decoupled.

$$m_1 \simeq \frac{m_D^2}{m_R}, \quad m_2 \simeq m_R \quad (4.20)$$

To end the analysis of the See Saw mechanism, it is worth mentioning that the analysis followed in this section was for only one neutrino flavor. Firstly, Equation 4.20 would turn into the expression in Equation 4.21, where each of the elements is a 3×3 matrix. Also, the matrix in Equation 4.16, in the case of three neutrino generations would be the one shown in Equation 4.22 where each of the elements in the matrix is a 3×3 matrix [12].

$$m_1 = m_D \frac{1}{m_R} m_D^\top \quad (4.21)$$

$$\mathcal{M}_6(6 \times 6) = \begin{pmatrix} 0 & m_D \\ m_D^\top & \mathcal{M}_3 \end{pmatrix} \quad (4.22)$$

If the matrix in 4.22 is diagonalized by blocks and if the mixings between flavors are neglected, a term for the neutrino masses in each generation like the one in Equation 4.23 would be obtained with $m_{f,i} = (m_e, m_\mu, m_\tau)$. In the process of obtaining Equation 4.23, \mathcal{M}_3 is assumed to be of the form shown in Equation 4.24.

$$m(\nu_i) \simeq \frac{m_{f,i}^2}{M_R} \quad (4.23)$$

$$\mathcal{M}_3 \simeq M_R \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \quad (4.24)$$

However it is also possible that the eigenvalues of are not all M_R as in Equation 4.23, but M_1 , M_2 , and M_3 . This difference in the eigenvalues would lead to a similar hierarchy to the one observed in the charged lepton masses [12].

As a final note, due to the inclusion of Majorana mass terms in the mechanism and as stated earlier in this chapter, a breaking of the L symmetry is induced in the See Saw mechanism. Nevertheless, this does not rises a problem for the validity of the See Saw mechanism since

the energy scales of the heavy neutrino, scales at which the L symmetry would break, would be very large. These large scales have not been yet reached by the particle physics experiments, so that would explain the lack of experimental evidence of the breaking of the L symmetry.

Chapter 5

Heavy Neutrino Searches

5.1 Neutrino Searches in the LHC

Heavy neutrino searches in the LHC have mainly focused in production via Drell-Yan mechanism. The Feynman diagram for this process is shown in Figure 5.2. As shown in the diagram, the heavy neutrino results from a W boson that decays into a lepton and the heavy neutrino. In this model, the heavy neutrino decays into a lepton and a W boson. The latter decays into two quarks that result in two jets observed by the detector. Advances in the collision energies provided by the LHC have allowed probing for masses of heavy neutrino at larger scales. This fact is shown in Figure 5.1 where a comparison in the cross section of the Drell-Yan heavy neutrino process in the Tevatron and the LHC is shown. It can be seen clearly that the LHC, due to its higher center of mass energy in the collision, has a greater cross section than the one present at the Tevatron for the same process [15].

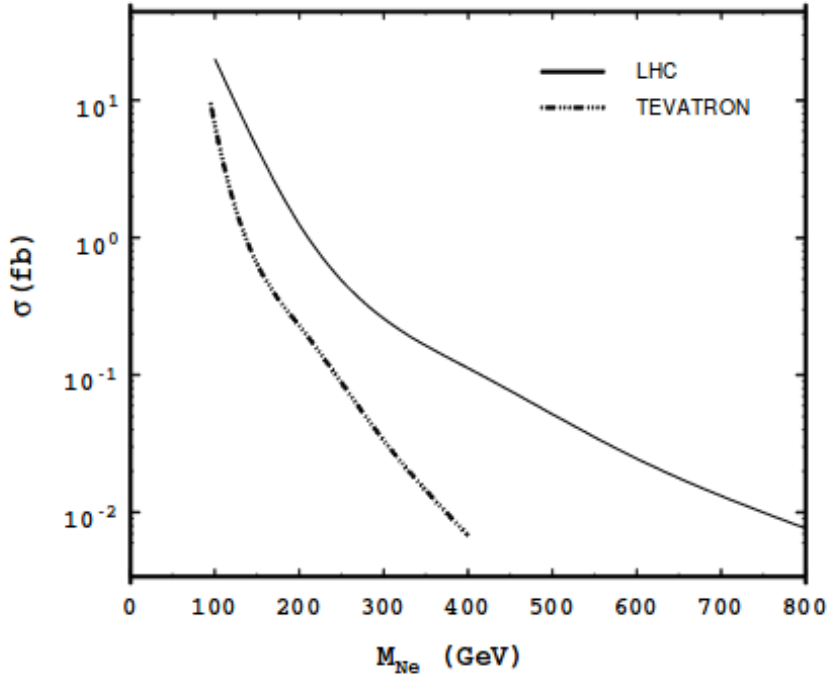


Figure 5.1: Cross sections for heavy neutrino production at the LHC and the Tevatron (Taken from [15])

As mentioned in [16], the process mentioned above is the dominant signal for a heavy neutrino with a mass between the 100 GeV and 500 GeV. The main problem with this type of process is that, given there is a W boson in the heavy neutrino production, the cross section for large values of the heavy neutrino mass decreases rapidly. The fast decrease in the signal cross section can be observed in Figure 5.1. This small cross section rises a problem for an observable signal with a heavy neutrino mass greater than 400 GeV [16].

Although the production of heavy neutrino through the Drell-Yan process has a small cross section, it has the advantage of having a clean

signal if the heavy neutrino has a decay of the form $N \rightarrow l^\pm W^\mp \rightarrow l^\pm jj$ [16]. This type of process allows the detection of all the products of the final state of the process. That is, signal from the two leptons and the two jets that result at the end of the process should be observed in the detector.

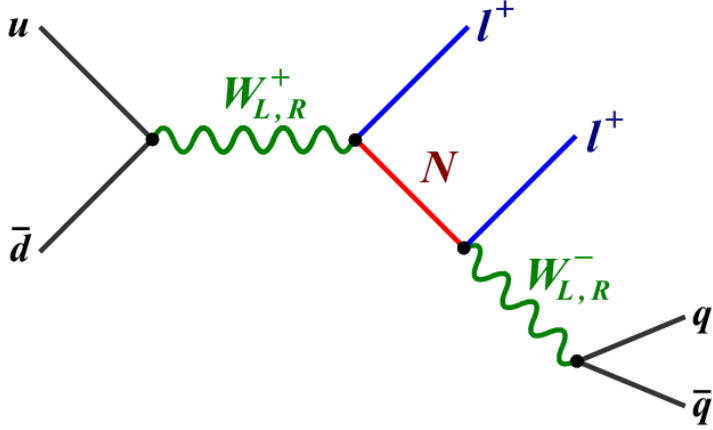


Figure 5.2: Feynman diagram of heavy neutrino production by Drell-Yan process

As noted in [6], all the heavy neutrino searches in the LHC have focused in the Drell-Yan process. It is also mentioned that studies involving final states with taus have not been explored because of the challenges of reconstructing and tagging correctly a τ . However it is also mentioned that this can be solved with the improvement of τ -tagging algorithms. The latter makes a study like the one presented in this monograph novel, because of the inclusion of taus instead of muons and electrons in the final state.

5.2 Heavy Neutrino Search with Vector Boson Fusion

The first difference between most of the studies mentioned in the previous section and the one done in this study is the use of τ instead of μ for the lepton in the process. The second and main difference is the use of the Vector Boson Fusion technique. As stated earlier, neutrino searches have been mainly performed using the Drell-Yan mechanism shown in Figure 5.2. However, a search using the Vector Boson Fusion technique has not been conducted. The main difference between the Drell-Yan process and the one including the VBF topology is the addition of two jets resulting from the fusion of two bosons. These two jets would have the characteristic of traveling in opposite directions and in the longitudinal region of the detector when they are created. As a result, the VBF topology includes two jets with a big angular separation in the detector, because they should be detected in opposite hemispheres (high $\Delta\eta$ separation) and in the end cap region (large η values). Also, the jets coming from the Vector Boson Fusion process would typically have greater transverse momentum p_T than jets coming, for example, from the quarks resulting from the decay of bosons. These characteristics mentioned above would allow to identify these jets coming from the VBF process, hence increasing the possibility of identifying the signal from experimental noise in the detector.

The addition of Vector Boson Fusion to the Drell-Yan process of the heavy neutrino would result in Feynman diagrams like the ones shown in Figures 5.3 and 5.4. Figure 5.3 shows the production of a heavy neutrino including the fusion of boson Z with boson W, and Figure 5.4 shows the fusion of boson gamma, or photon, with boson W. These two examples show that the bosons that fuse in the VBF processes can be multiple combinations of vector bosons. Also, it is noteworthy that the Feynman diagram portion regarding the production of the heavy neutrino remains unchanged compared to the one of the Drell-Yan production process shown in Figure 5.2. The latter shows how the Vector Boson Fusion is an addition to the process that would help with the identification of the signal, but does not affect the original process that is being studied.

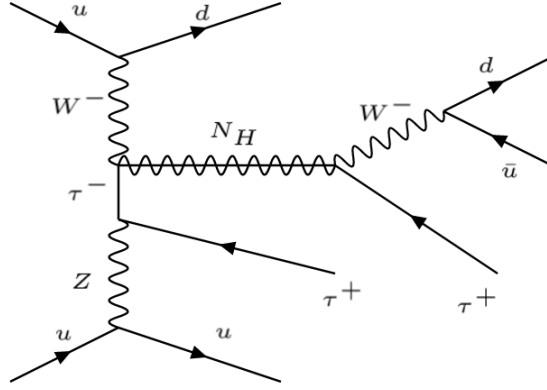


Figure 5.3: Feynman diagram of heavy neutrino production including VBF and involving Z boson

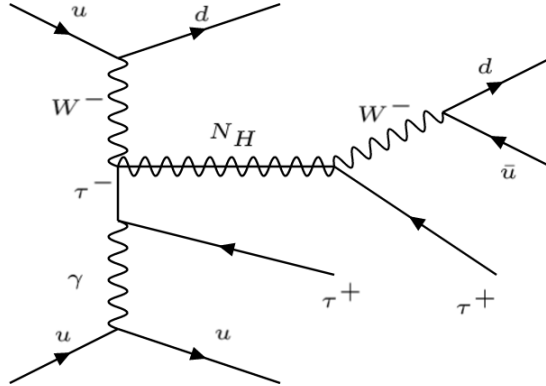


Figure 5.4: Feynman diagram of heavy neutrino production including VBF and involving photon

A search using the VBF topology would be promising, because there could be a potential enhancement of the signal cross section compared

with the heavy neutrino production via Drell-Yan process. A comparison of the cross section of both processes is shown in Figure 5.5. Although in both cases the cross section decreases when the mass of the heavy neutrino increases, the cross section of the VBF process decay is slower than the one of the Drell-Yan process. This causes that from a value for the heavy neutrino mass of approximately 1.4 TeV the VBF process has a greater cross section compared to the cross section of the heavy neutrino produced by the Drell-Yan process. This plot shows clearly why it is relevant to perform the analysis of a potential search of heavy neutrinos at the LHC using the Vector Boson Fusion technique.

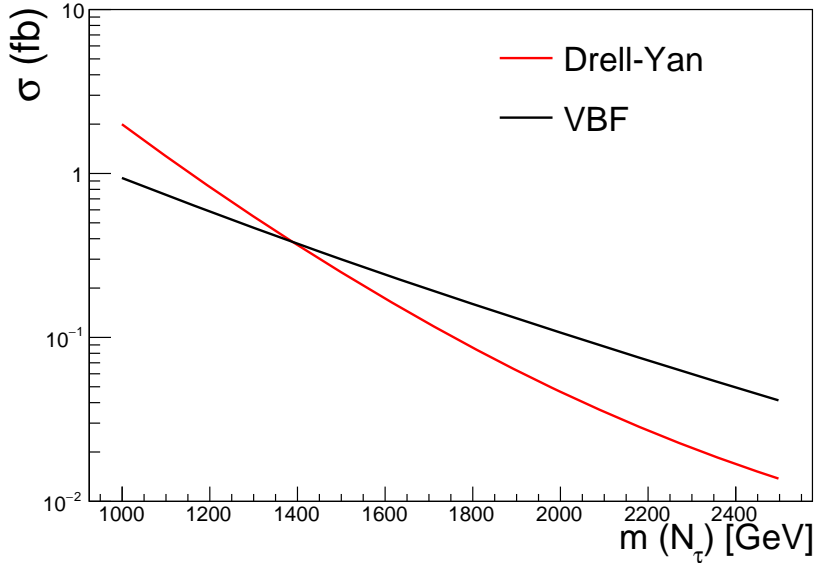


Figure 5.5: Feynman diagram of heavy neutrino production by Drell-Yan process

Chapter 6

Backgrounds

Backgrounds in particle physics experimental studies are processes that have similar final states to the process that is being studied. Because of their similar final states with the process of interest, background signal can be mistaken as the signal of the process being studied. That is why backgrounds are experimental noise that is not desirable in a study and that has to be eliminated as much as possible. In this particular study three backgrounds were considered because of their similarity in the final states with the one of the heavy neutrino production.

One of the backgrounds considered for this analysis was the $W + \text{jets}$ process shown in Figure 6.1. This process is considered as a background for the heavy neutrino production because of the presence of one lepton coming from the decay of the W boson in the process. It can be noted that in principle this process would not have the second lepton present in the heavy neutrino production process. However, it has to be remembered that taus are observed as jets in the detector. There is a possibility of tagging a jet as a tau incorrectly. Since the gluon, the quark b and the other quark present in the process are also observed in the detector as jets, these could be identified incorrectly as a tau. If one of these jets is tagged incorrectly, then the final state of the $W + \text{jets}$ process would be the same as the one of the heavy neutrino process because it would be identified as a process with two leptons and two jets.

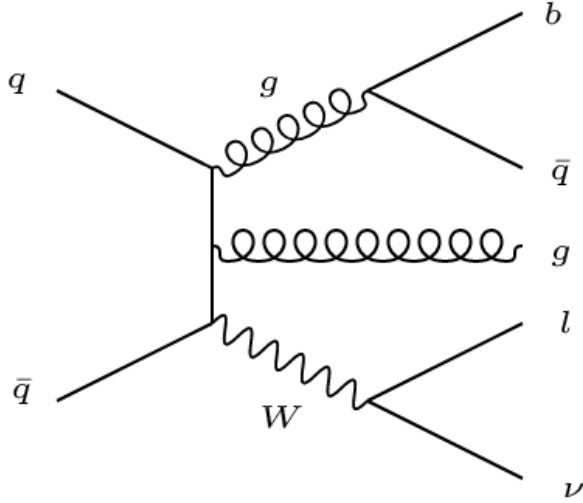
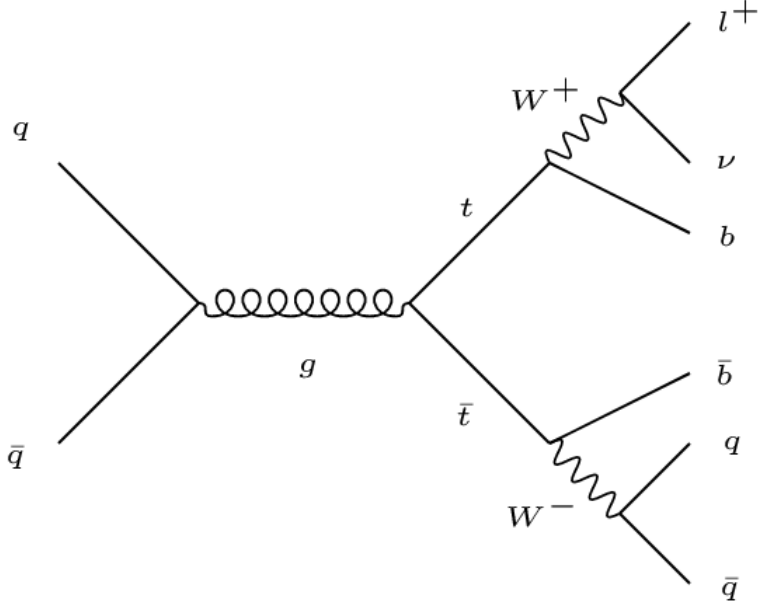


Figure 6.1: Feynman diagram of W+jets process

Another background considered was the $t\bar{t}$ semileptonic process that contains the decay of a gluon into the quarks top and anti-top as shown in Figure 6.2. As in the case of the W + jets process, the $t\bar{t}$ semileptonic background has only one lepton in its final state. However, as mentioned earlier, this process would be observed in the detector as two jets coming from the b quark, or B-jets, and two jets coming from the other two quarks. If one of the four is tagged incorrectly as a τ , the final state would be identical to the one of the heavy neutrino process studied in the analysis.

Figure 6.2: Feynman diagram of $t\bar{t}$ process

The other background considered for the analysis is the Drell-Yan process shown in Figure 6.3. The reason for this process to be considered as a background is the two leptons in its final state as well as the two jets coming from the initial state radiation jets. It can be seen that from the three backgrounds, Drell-Yan is the one that is the most similar to the heavy neutrino production, because the heavy neutrino process considered for this analysis also has two jets and two leptons in the final state. This is why the Drell-Yan background was expected to be the most difficult to separate from the signal of interest in the analysis.

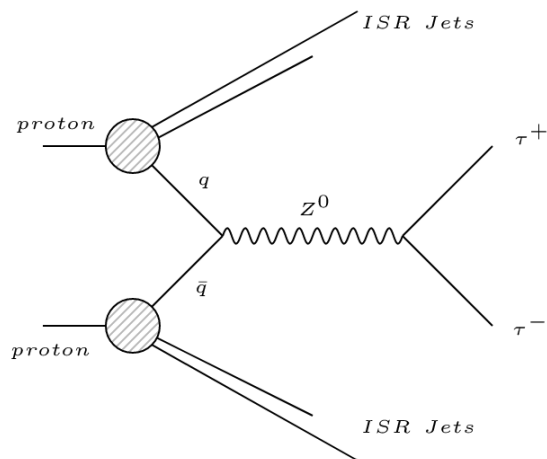


Figure 6.3: Feynman diagram of Drell-Yan process

Chapter 7

Signal Simulation and Analysis

7.1 Simulation Computational Resources

The project required computational work, because simulations of events from the different processes were needed. Also, an analysis of the samples using the analysis code was required. The background and signal samples were simulated using the software MadGraph [17], Pythia [18] and Delphes [19]. The data analysis and all the subsequent kinematic, topological, and optimal cuts analyses were performed using ROOT software [20].

MadGraph is an event generator software that allows the simulation of collision between two particle beams. For this analysis in particular, the simulations will consist in proton collision at 13 TeV in order to reproduce the actual conditions of the LHC. MadGraph includes the physical parameters that determine the production probability of a given process, as well as the possible decays of the simulated particles. Besides providing the necessary matrices to calculate the cross sections of the processes, MadGraph also creates the pictorial representations of the Feynman Diagrams from the generated processes. To this end, the software uses perturbation theory in the calculations of production and generation of physical processes.

Pythia is a software that allows the simulation of various strong processes models that evolve from a few bodies to final states with high particle multiplicity. Particularly, in this case Pythia will be used for the simulation of quark and gluon fragmentation processes. This fragmentation process occurs when, due to an intrinsic characteristic of the strong interaction, there is an energy gain caused by the increase of the distance of two bound quarks. If the separation is enough to reach a critical energy, a pair quark-antiquark is created. The Pythia simulation is necessary, because processes like the ones mentioned above occur during a proton collision at the LHC.

Delphes is a software used to add the effects that a multipurpose detector, like ATLAS or CMS, may have on the particles to the Monte Carlo simulations performed for different processes. In this particular case, Delphes is necessary to simulate the interaction of the particles coming from the generated processes in MadGraph and Pythia with the CMS components. Namely, reproducing the conditions of the detector and the uncertainties coming from the measuring process is achieved by using Delphes. The changes in the cinematic variables due to their interaction with matter, errors caused by the electronics of the detector, and the additional particles generated because of the interaction between the particles and the detector components can be accounted for using Delphes. Other functionalities included in Delphes are: simulation of the detector geometry, the effect of the magnetic field over the particles, and the particle identification and reconstruction efficiencies, among others.

7.2 Simulation Process

The MadGraph signal simulation was performed for four values of the heavy neutrino masses: 1.0 TeV, 1.25 TeV, 1.5 TeV, and 2.0 TeV. Also, taking into account that the analysis was going to be performed using Vector Boson Fusion, the parameter of minimum pseudorapidity separation ($\Delta\eta$) between two jets was set to 3.5. This parameter allowed to have more events with the VBF topology in the simulation.

The commands used to generate the desired signal were the following:

- import model SM_HeavyN_NLO
- generate p p > n3 ta+ jj, QCD= 0, n3 > ta+ jj
- add process p p > n3 ta+ jj, QCD= 0, n3 > ta- jj
- add process p p > n3 ta- jj, QCD= 0, n3 > ta- jj
- add process p p > n3 ta- jj, QCD= 0, n3 > ta+ jj

The first command imports the theoretical model that includes the interactions related with the heavy neutrino formation and decay. The next command specifies the processes that are going to be simulated. $pp > n3 \tau^+ jj$ stands for the proton-proton collision that decays into a heavy neutrino, a τ with positive charge, and two jets. The flag $QCD=0$ is used to exclude all strong interactions that can be involved in the process. Finally, $n3 > \tau^+ jj$ is used to force the decay of the heavy neutrino into a τ charged positively and two jets. The subsequent commands are used to take into account all the possible combinations of the electrical charge that the τ may have. The additional parameter of a separation greater than 3.5, i.e. $\Delta\eta > 3.5$, between VBF jets was included in run card of the simulation.

Figures 5.3 and 5.4 show two of the main possible diagrams generated by MadGraph for the processes simulated. These diagrams show the vector boson fusion process, occurring in Figures 5.3 and 5.4 in the fusion of the W boson with the Z boson and the photon (γ) respectively. In these last two diagrams, the decay of the W boson coincides with the one shown in Figure 5.2 for the decay of the W boson resulting in a heavy neutrino and a lepton, which in this case is a τ . The final states for the processes shown in these two diagrams would include two jets coming from the VBF process, two τ 's, and two jets coming from the quarks resulting of the W^- decay

Each simulation was for 100,000 events. The simulation for the 1.5 TeV mass value was performed in 10 different simulation batches each one containing 10,000 events. Every batch was generated with a different random seed to guarantee the independence of the events between each one of the generated batches. This independance was necessary because the 10,000 event files were merged to form a single file with 100,000

events. The other three simulations were performed in a single run of 100,000 events each.

As explained earlier, after the events were simulated in MadGraph they were passed to Pythia and then to Delphes so the signal resembled one that could be found at CMS. The Pythia process using Pythia 6 and MadGraph 5 was done seamlessly, because Pythia 6 is intergrated directly into MadGraph 5. This functionality makes that to perform the Pythia simulation process it is enough to turn on Pythia at the beginning of the MadGraph 5 simulation commands. To perform the Delphes simulation part, the output files from Pythia and MadGraph were used to execute a Delphes simulation using a Delphes card where all the parameters of the CMS detector were introduced.

7.3 Analysis Computational Resources

As mentioned earlier, the ROOT framework was used to perform all the analyses regarding kinematic and topological variables, and optimal cuts. ROOT is a software library developed by CERN to perform data analyses related with particle physics. One of the main characteristics of this library is the possibility of handling large volumes of data efficiently. The latter is achieved by using a tree structure in which the information related with the particles is stored and can be accessed easily using ROOT functionalities. Other features included in the library are the creation of histograms from data trees, multivariate analysis, four-vector calculations, among others. By using ROOT functionalities, it is also possible to estimate optimal cuts in variables to reduce experimental noise to its minimum. This is why the entire final analysis involved using tools provided by ROOT. A more detailed description of the analysis performed using ROOT is presented in the following chapters.

Chapter 8

Event Selection Criteria

8.1 Preselection Criteria

When an event is analyzed, several preselection criteria are established to count a signal observed in the detector as a valid particle. This is done to be certain that the signal received is indeed from the type of the particle that is expected. In this section, the preselection criteria used in this analysis are given. The summary of the criteria mentioned in this section is presented in Table 8.1.

In each event, a maximum of six jets that had a p_T greater than 15 GeV and its absolute value of η less than 5.0 were stored to be analyzed later. Among this list of jets, the two jets whose summed masses resulted in the greatest mass combination were stored and defined as the Di-Jet Pair. The jet with greater momentum in the Di-Jet Pair is the leading jet and the other one in the pair is the sub-leading jet. These two jets are the ones related with the VBF process. Another variable defined regarding the Di-Jet Pair was the Di-Jet mass and corresponds to the sum of the masses from the jets in the Di-Jet Pair.

Since the τ selection is important for this analysis, it is relevant to provide a further description of the preselection criteria for the τ 's in the simulated events. For starters, a jet identified as a tau is considered a valid τ if it has a transverse momentum greater than 20 GeV. Also, it was required that a valid τ should not overlap with an electron, a muon, or a jet. That is, the ΔR , defined as in Equation 3.2, should not

Variable	Criteria
$p_T(j)$	$> 40 \text{ GeV}$
$ \eta(j) $	< 5
$p_T(\tau)$	$> 20 \text{ GeV}$
$p_T(j_B)$	$> 30 \text{ GeV}$
ΔR	> 0.3

Table 8.1: Preselection criteria table

be less than 0.3. This condition guarantees that the jet identified as a τ does not overlap with other leptons or jets. Since the final state for this analysis includes two τ 's, the two taus with greater p_T are selected among a maximum of three taus stored for each event. The leading τ is the one with highest p_T and the sub-leading τ is the one with second highest p_T .

8.2 Variable Cut Optimization

To determine the optimal values for some of the variables, a study of the significance through multiple cuts was performed using the relation showed in Equation 3.6 for each cut value. The histograms analyzed were normalized to the luminosity of each process. Taking this analyses into account, the optimal cut values for the variables $\Delta\eta$ from the Di-Jet Pair and the mass from the Di-Jet Pair, $m(jj)$, was determined.

The result for the significance study for the $\Delta\eta$ from the Di-Jet Pair is shown in Figure 8.1. The first cut value considered was 3.8, because, as mentioned earlier, it is expected that the jets from the VBF process have a big separation in η . It is clear from the plot in Figure 8.1 that the best value for the cut is 3.8. The significance for $\Delta\eta$ values greater than 3.8 tend to decrease. This is why a selection value of 3.8 for the $\Delta\eta$ of the diJet pair was established.

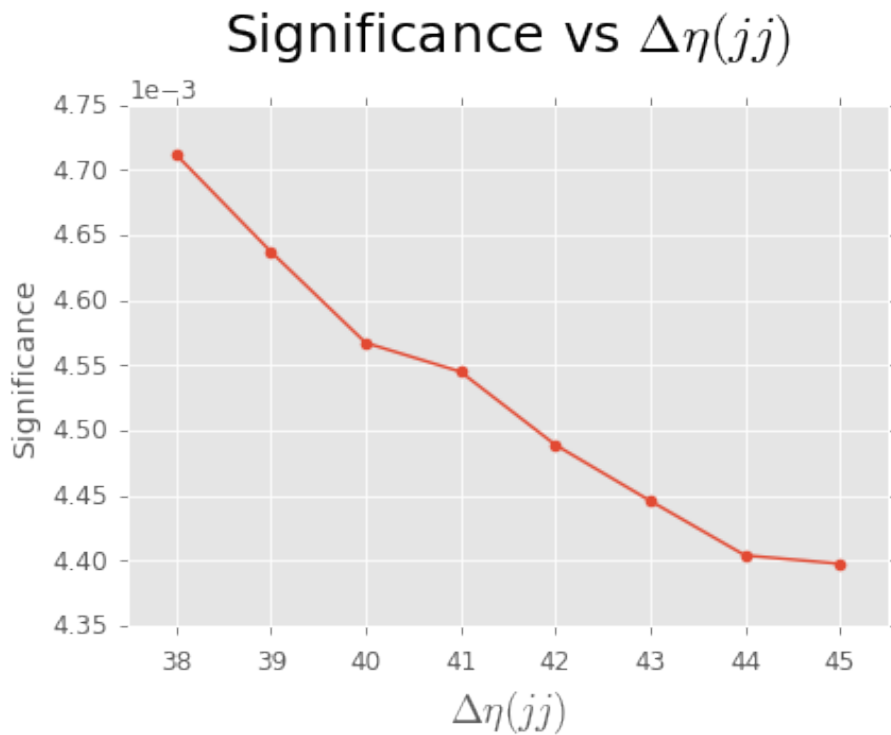
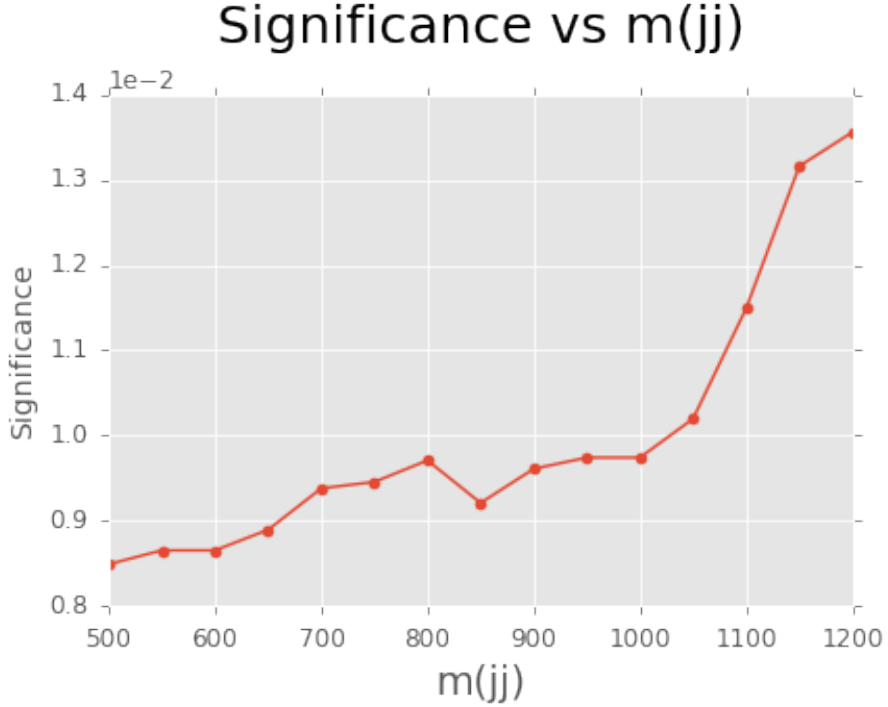


Figure 8.1: Significance of multiple cuts in $\Delta\eta$ diJet variable

Figure 8.2: Significance of multiple cuts in $m(jj)$

8.3 Event Selection Criteria

In order to achieve a separation between background and signal, several successive requirements for the variables of the particles in the event were made. These requirements are defined as cuts, and for each cut the events that do not comply with the established condition are not taken into account to fill the histograms. Eight cuts were made to the histograms, storing in each cut the resulting distributions to analyze them later. The first four cuts were related with jets and τ 's in the event, and the subsequent four were related with the VBF topology. In the next paragraphs of this section a description of each one of the cuts is given as well as the order in which they were performed. A summary of the selection criteria can be found in Table 8.2.

The first cuts that were made to the histograms were to require that the leading and sub-leading τ 's should have a minimum transverse momentum of 20 GeV and a maximum of 2.1 for the absolute value of η . The latter guarantees that the τ 's left are detected by the barrel and not the end-caps of the detector. That is an important condition because the detection components, such as the tracker, are in the barrel section and are more accurate than the ones in the end-caps. As a result, a signal detected in the barrel is most certain to be accurate than one detected in an end-cap.

The next cut requires that the event does not have any B-jet. This cut is justified by the fact that one of the main backgrounds for the signal is the top anti-top ($t\bar{t}$) process. The interaction between the top and anti-top quark is related with the production of jets associated with the b quark or B-jets. That is why, much of the $t\bar{t}$ should be eliminated by requiring no B-jets in the event. This fact will be later analyzed further in chapter 9. The cut that follows the one regarding the B-jets selects the events that have a minimum of two jets with transverse momentum greater than 30 GeV. These two jets must be different from the ones used in the Di-Jet Pair and are the ones that should result from the fragmentation of the quarks resulting from the W decay shown in Figure 5.2.

The last three cuts made to the histograms are related to the VBF topology. The first of the three selects events in which the product of η from the leading and sub-leading jets is negative. This condition guarantees that the jets in the Di-Jet Pair are in opposite hemispheres. The next cut requires that the leading and sub-leading jets of the event have a $\Delta\eta$ greater than 3.8. Finally, the last cut requires that the Di-Jet mass of the event is greater than 500 GeV.

Variable	Criteria
$p_T(\tau)$	$> 20 \text{ GeV}$
$ \eta(\tau) $	< 2.1
\cancel{E}_T	$> 30 \text{ GeV}$
Number of B-jets	$= 0$
Number of jets	≥ 2
$\eta(j_l) * \eta(j_s)$	< 0
$\Delta\eta(jj)$	> 3.8
$m(jj)$	$> 500 \text{ GeV}$

Table 8.2: Selection criteria table

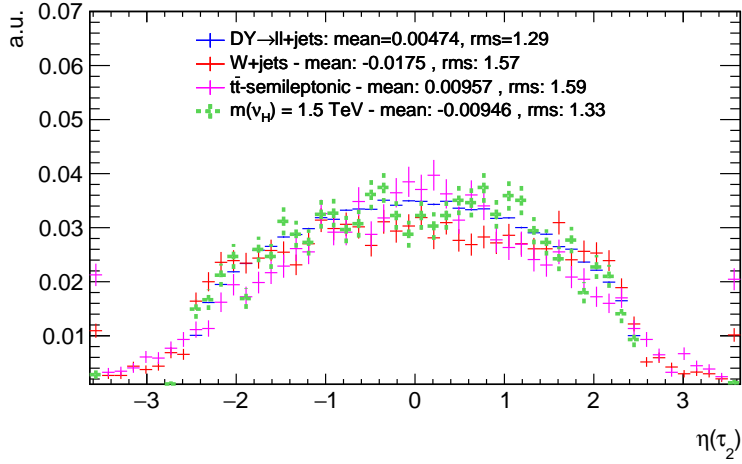
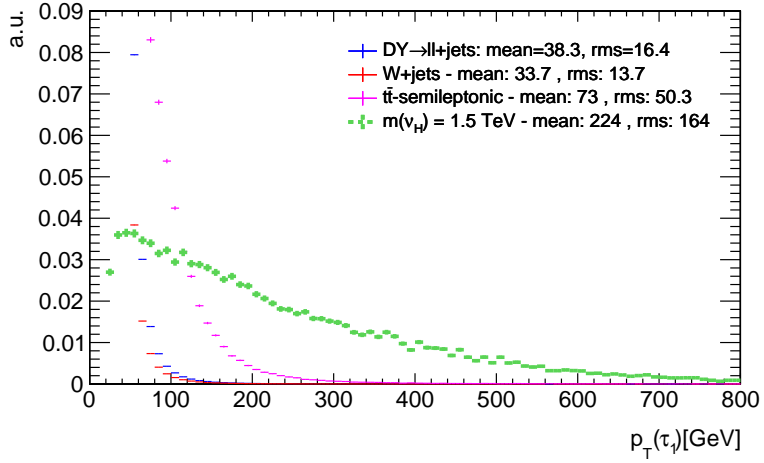
Chapter 9

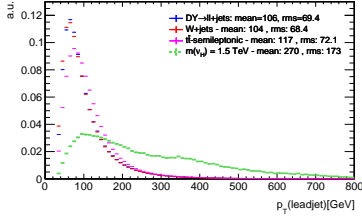
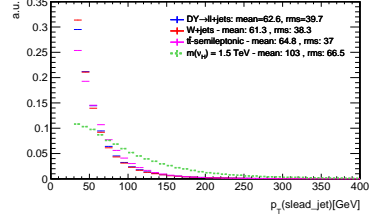
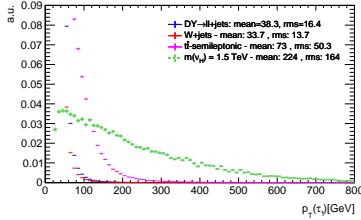
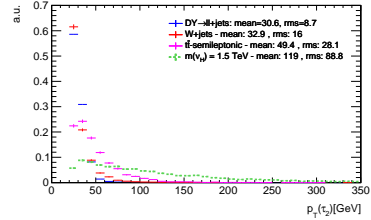
Distribution Analysis

9.1 Normalized Distributions

The normalized plots are useful to check the shape of backgrounds and signal distributions. These shapes help to identify which variables are not useful for the analysis because the backgrounds overlap with the signal and which variables have to be studied with greater detail because the signal separates from the backgrounds from certain points. An example for both cases is provided in Figures 9.1 and 9.2. In figure 9.1 it can be seen that the signal overlaps for all values with the background distributions. That is why this plot can be used to conclude that the η variable from the sub-leading τ is not useful to isolate the signal from the background. In contrast, the plot in Figure 9.2 show that from around 150 GeV the signal separates from the background distributions. This separation for the p_T of the leading τ suggests that this variable should be examined more closely through the subsequent cuts.

To understand the definition of H_T and S_T mentioned in chapter 3, the plots in Figure 9.3 are shown. The four plots show a separation, in some cases a smaller than others, between the signal and background distributions. A tendency of the signal jets to have greater transverse momentum than the ones in the backgrounds is shown. This tendency is also displayed for both τ 's. Hence, the distributions of H_T and S_T should show a similar behaviour because this variables are the result of

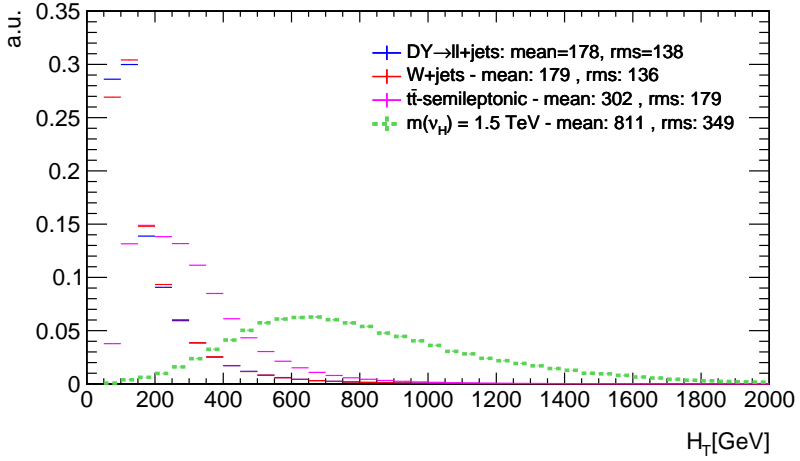
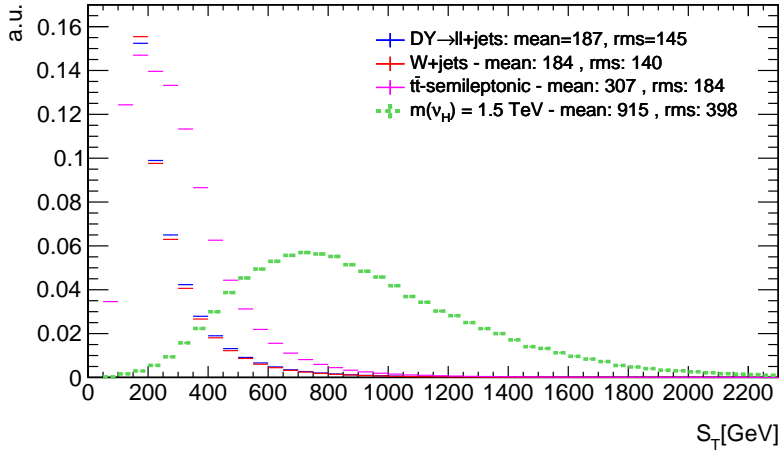
Figure 9.1: Unit plot of η from the sub-leading τ with no cutsFigure 9.2: Unit plot of p_T of leading τ with no cuts

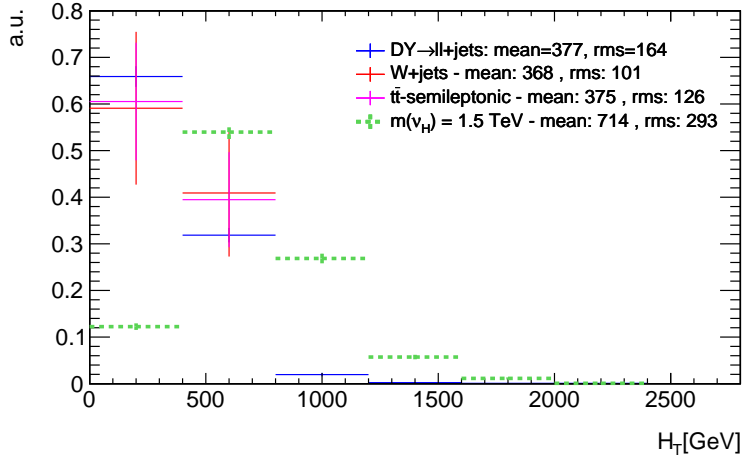
(a) Leading jet p_T unit plot(b) Sub-leading jet p_T unit plot(c) Leading τ p_T unit plot(d) Sub-leading τ p_T unit plotFigure 9.3: p_T unit plots for different bodies in the event

adding the transverse momentum of jets and τ 's in the event.

In Figures 9.4 and 9.5, the normalized plots with no cuts of H_T and S_T are shown. It can be seen that indeed a greater separation between signal and background was achieved. Unlike the distributions of the τ 's and jets transverse momentum, the maxima of H_T and S_T lie outside the backgrounds distributions. Furthermore, the background that overlaps at a greater energy with the signal corresponds to $t\bar{t}$. Taking into account what was mentioned in section ??, the overlap between the signal and this background could be reduced with the cut related with the number of B-jets in the event. This is why this two variables need to be studied closer in the analysis.

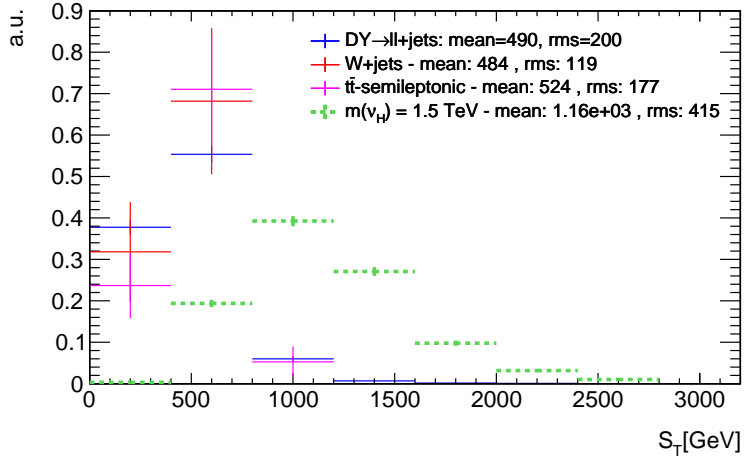
In Figures 9.6 and 9.7, the plots including cuts of H_T and S_T respectively are shown. Both plots include all the cuts mentioned in section 8.3. A rebining of the histograms was necessary, because due to the cuts, the number of background events decreased significantly. This made the error bars for these histograms larger making the analysis of the

Figure 9.4: Unit plot of H_T with no cutsFigure 9.5: Unit plot of S_T with no cuts

Figure 9.6: Unit plot of H_T with VBF cuts

plot more difficult. For that reason, a bin size of 400 GeV was defined for both of these plots. Comparing the two signal histograms, it can be seen that the maximum of the H_T distribution still overlaps with the background distributions. Although the signal in this case appears on top of the background, it has to be noted that this normalized plots are only useful to get an idea of the distribution shapes. That is why the fact the the signal appears on top of the background does not mean that the signal can be detected over the background for this value. On the contrary, the maximum of the S_T signal distribution has its maximum when the background distributions start to decay. This distribution shape difference between H_T and S_T is related with the fact that the signal has τ 's with higher p_T than the ones found in the backgrounds.

The mean values for S_T shown in Figures 9.5 and 9.7 are 915 GeV and 1160 GeV respectively. From the Feynman diagram in Figure 5.4, it can be seen that the jets in the event and one of the taus come from the heavy neutrino decay. Taking into account that the defined mass for the heavy neutrino is 1.5 TeV, it is expected that the maximum of the S_T distribution is around this value. This analysis leads to the

Figure 9.7: Unit plot of S_T with VBF cuts

conclusion that the cuts made to the histograms are helping to select the significant events, because the mean S_T is closer to the expected value after performing all the cuts.

Appendix A

Same Chirality Fields Terms

A field ψ can be decomposed in its chiral components

$$\psi = \psi_R + \psi_L$$

where

$$\psi_R = \frac{1}{2} (1 + \gamma^5) \psi \tag{A.1}$$

$$\psi_L = \frac{1}{2} (1 - \gamma^5) \psi \tag{A.2}$$

where γ^5 is the fifth Dirac matrix. Taking into account the Equations A.1 and A.2, the projection operators are defined as

$$P_R = \frac{1}{2} (1 + \gamma^5)$$

$$P_L = \frac{1}{2} (1 - \gamma^5)$$

As expected and as shown in Equations A.1 and A.2, P_R projects a field in its right chiral state and P_L projects the field in its left chiral state.

In a similar way the fields $\bar{\psi}_R$ and $\bar{\psi}_L$ are defined as

$$\bar{\psi}_R = \frac{1}{2} \bar{\psi} (1 - \gamma^5)$$

$$\bar{\psi}_L = \frac{1}{2} \bar{\psi} (1 + \gamma^5)$$

Taking into account the previous definitions, the term $\bar{\psi}_R \psi_R$ would be

$$\frac{1}{4} \bar{\psi} (1 - \gamma^5) (1 + \gamma^5) \psi$$

Remembering the property $(\gamma^5)^2 = 1$, the term $(1 - \gamma^5) (1 + \gamma^5)$ would be zero.

A similar analysis can be performed of the the field with left chirality to obtain the same conclusion. These analysis leads to the conclude that $\bar{\psi}_R \psi_R = \bar{\psi}_L \psi_L = 0$.

Appendix B

Mass Eigenstate Majorana Neutrinos Calculation

The eigenvalues of the matrix in Equation 4.16 can be obtained by using the standard method for calculating eigenvalues. That is, the eigenvalues if \mathcal{M} , are λ values that satisfy

$$\det(\mathcal{M} - \lambda) = 0$$

After some algebra, this determinant is equal to

$$\lambda^2 - \lambda(m_L + m_R) + m_L m_R - m_D^2 = 0$$

Using the quadratic equation to solve for λ values, the result is like the one shown in Equation B.1. In this Appendix, the λ with the negative sign will be defined as m'_1 and the value with the plus sign as m'_2 .

$$\lambda = \frac{m_R + m_L}{2} \pm \frac{\sqrt{4m_D^2 + (m_R - m_L)^2}}{2} \quad (\text{B.1})$$

Since \mathcal{M} is real and symmetric, a real orthogonal matrix \mathcal{O} that diagonalizes \mathcal{M} exists. In other words, there must be an \mathcal{O} such that the expression in Equation B.2

$$\mathcal{O}^\top \mathcal{M} \mathcal{O} = m_D \quad (\text{B.2})$$

where

$$m_D = \begin{pmatrix} m'_1 & 0 \\ 0 & m'_2 \end{pmatrix}$$

and

$$\mathcal{O} = \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix}$$

Before continuing with the derivation, it is worth noting that the value for m'_1 as defined earlier can take negative values. Since the physical masses have to be positive a new matrix

$$m_D = \begin{pmatrix} m_1 \eta_1 & 0 \\ 0 & m_2 \eta_2 \end{pmatrix}$$

where $\eta_1 = -1$, $\eta_2 = 1$, and $m'_i = m_i \eta_i$. That way it is assured that the values of the m_1 and m_2 are positive. This result can also be achieved by multiplying the first column of \mathcal{O} by i . The latter is the procedure followed in Chapter 4 to make the mass values positive.

Since \mathcal{O} is the matrix that diagonalizes \mathcal{M} , then the field ν_L is related to N_L by the expression shown Equation B.3 where N_L is the vector with definite mass fields, i.e. the fields with mass m_1 and m_2 defined as in Equation B.4.

$$\nu_L = \mathcal{O} N_L \tag{B.3}$$

$$N_L = \begin{pmatrix} \nu_1 \\ \nu_2 \end{pmatrix} \tag{B.4}$$

With the definitions presented in Equations B.3 and B.4, the equations shown in 4.17 become clear.

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