# Discrete well-balanced WENO finite difference schemes: global-flux quadrature method with multi-step ODE integrator weights

Maria Kazolea\*, Carlos Parés<sup>+</sup>, Mario Ricchiuto\*

\*Team CARDAMOM, Inria research center at University of Bordeaux (France)

<sup>+</sup> EDANYA Group, Universidad de Málaga (Spain)

Converence name Conference adress Conference date



$$\partial_t U + \partial_x F(U) = S(U)\partial_x H$$

Typical examples

- Shallow water equations with topography/friction/Coriolis/etc
- Euler equations with gravity
- GPR model and hyperbolic reformulations of viscous/dispersive systems
- etc.

$$\partial_t U + \partial_x F(U) = S(U) \partial_x H$$

Shopping list

- Arbitrary high order
- Possibly monotonicity/positivity preserving
- Steady state preserving



We want to solve numerically hyperbolic systems of balance laws

$$\partial_t U + \partial_x F(U) = S(U)\partial_x H$$

Steady state preserving: on a given mesh we want to the scheme embed in the discretization a

*Definition* of  $\partial_x F_{
m h} = S_{
m h}$  such that data verifying it is exactly preserved

Projection operator providing such data, possibly different from the scheme

Characterization of the above projection: e.g enhanced consistency, well posedness

There are numberless approaches, and vast literature. Related work:

- Reconstruction/evolution of fluctuations wrt discrete equilibria (approximate full well balanced):
  - Castro & Parés J.Sci.Comp. 2020, Gómez-Bueno et al, Appl.Math.Comp. 2021, Gómez-Bueno et al Mathematics 2021, Guerrero Fernández et al Mathematics 2021, Gómez-Bueno et al, Appl.Num.Math. 2023, etc. etc.
- 2 Well balanced via integration of the source term and global fluxes
  - Gascón & Corberán JCP 2001, Donat & Martinez-Gavara J.Sci.Comp. 2011, Chertock et al JCP 2018, Cheng et al J.Sci.Comp. 2019, Mantri&Noelle JCP 2021, Ciallella et al. J.Sci.Comp. 2023, Xu and Shu J.Sci.Comp. to appear, etc. etc.
- 3 Other approaches
  - •



Approach 1.1

Equilibria Preservation, J.Sci.Comp. 96, 2023

 $\partial_t U + \partial_x F(U) = S(U)\partial_x H$ 

 $rac{d\overline{U}_i}{dt} + rac{1}{\Delta x} \left( \widehat{G}_{i+1/2} - \widehat{G}_{i-1/2} 
ight) = 0$ 

<sup>1</sup>M. Ciallella. D. Torlo, MR, Arbitrary High Order WENO Finite Volume Scheme with Flux Globalization for Moving

We seek solutions of the hyperbolic system of balance laws

Approach 1.1

$$rac{dU_i}{dt} + rac{1}{\Delta x}\left(\widehat{G}_{i+1/2} - \widehat{G}_{i-1/2}
ight) = 0.$$

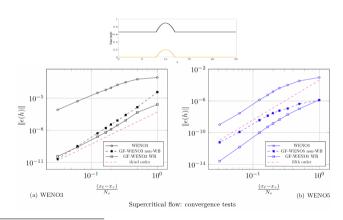
- 1 Reconstruct WENO polynomials  $U_i(x)$
- 2 Compute cell averaged source primitive  $\overline{R}_i = R_{i-1/2}^+ \sum_q \omega_q \int_{x_{i-1/2}}^{x_q} S(U_i(x), \varphi(x)) dx$
- 3 Compute cell averaged fluxes  $\overline{F}_i = \sum_q \omega_q F(U_i(x_q), \varphi(x_q)) dx$
- 4 Reconstruct WENO polynomials  $G_i(x) = (F + R)_i(x)$

<sup>1</sup>M. Ciallella. D. Torlo, MR, Arbitrary High Order WENO Finite Volume Scheme with Flux Globalization for Moving Equilibria Preservation, J.Sci.Comp. 96, 2023

<sup>5</sup> Compute upwind fluxes  $\hat{G}_{i+1/2} = (A^+A^{-1})_{i+1/2}G_i(x_{i+1/2}) + (A^-A^{-1})_{i+1/2}G_{i+1}(x_{i+1/2})$ 

$$\partial_t U + \partial_x F(U) = S(U) \partial_x H$$

Approach 1.1



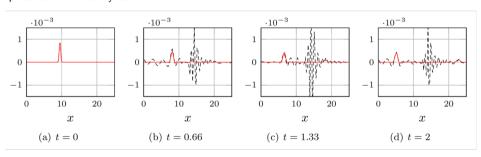
<sup>&</sup>lt;sup>1</sup>M. Ciallella, D. Torlo, MR, Arbitrary High Order WENO Finite Volume Scheme with Flux Globalization for Moving Equilibria Preservation, *J.Sci.Comp.* 96, 2023



We seek solutions of the hyperbolic system of balance laws

$$\partial_t U + \partial_x F(U) = S(U) \partial_x H$$

# **Approach 1.**<sup>1</sup> Small perturbation of steady state



<sup>&</sup>lt;sup>1</sup>M. Ciallella, D. Torlo, MR, Arbitrary High Order WENO Finite Volume Scheme with Flux Globalization for Moving Equilibria Preservation, *J.Sci.Comp.* 96, 2023

$$\partial_t U + \partial_x F(U) = S(U) \partial_x H$$

# Approach 1.1

- ✓ No a-priori knowledge of equilibrium, all steady states!
- No need of compute the solution of the Cauchy problem .. (maybe for initialization)
- Considerable accuracy enhancements at steady state (3 to 4 orders of magnitude)
  - No super convergence
- Very hard to charcterize the steady state and generate one

$$\overline{R}_i = R_{i-1/2}^+ - \sum_q \omega_q \int_{x_{i-1/2}}^{x_q} S( \overline{U_i(x)}, arphi(x)) dx$$

<sup>&</sup>lt;sup>1</sup>M. Ciallella. D. Torlo, MR, Arbitrary High Order WENO Finite Volume Scheme with Flux Globalization for Moving Equilibria Preservation, J.Sci.Comp. 96, 2023

$$\partial_t U + \partial_x F(U) = S(U) \partial_x H$$

$$rac{dU_i}{dt} + rac{1}{\Delta x} \left( \widehat{G}_{i+1/2} - \widehat{G}_{i-1/2} 
ight) = 0$$

We seek solutions of the hyperbolic system of balance laws

$$\partial_t U + \partial_x F(U) = S(U) \partial_x H$$

$$rac{dU_i}{dt} + rac{1}{\Delta x} \left( \widehat{G}_{i+1/2} - \widehat{G}_{i-1/2} 
ight) = 0$$

- 1 Reconstruct WENO polynomials  $U_i(x)$
- 2 Compute nodal source primitive  $R_i = R_{i-1} \Delta x \sum_{\sigma} \omega_q S(U_{i-q}, \varphi(x_{i-\sigma}))$
- 3 Compute cell averaged fluxes  $\overline{F}_i = \sum_q \omega_q F(U_i(x_q), \varphi(x_q)) dx$
- 4 Reconstruct WENO polynomials  $G^{\pm}_{i+1/2}(x) = (F+R)^{\pm}_{i+1/2}(x)$
- 5 Compute upwind fluxes  $\hat{G}_{i+1/2} = P_{i+1/2}^+ G_{i+1/2}^- (x_{i+1/2}) + P_{i+1/2}^- G_{i+1/2}^+ (x_{i+1/2})$

$$\partial_t U + \partial_x F(U) = S(U) \partial_x H$$

New approach

#### Main result

Proposition (Discrete steady state). The WENO-FD scheme with global flux quadrature preserves exactly continuous discrete steady states  $U_i^* = U(F_i)$  with F obtained by integrating the ODE

$$F' = S(U(F))\partial_x H$$

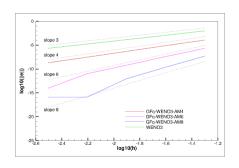
using the multi-step ODE integrator with weights  $\{\omega_q\}_{q\geq 0}$  on spatial slabs of size h. If U(F) is a one to one mapping,  $U^*(x)$  verifies the consistency estimates of the multi-step scheme.

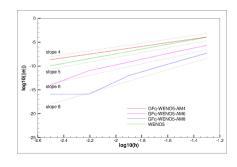
Multi-step schemes so far: Adams methods ABn and AMn

We seek solutions of the hyperbolic system of balance laws

$$\partial_t U + \partial_x rac{U^2}{2} = S(U) \partial_x H$$

For 
$$S(U) = U^2$$
 and  $H(x) = x$  exact steady state  $u = e^x$ 





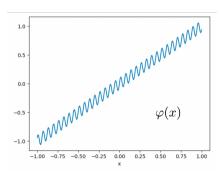


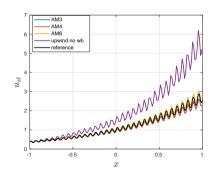
We seek solutions of the hyperbolic system of balance laws

$$\partial_t U + \partial_x rac{U^2}{2} = S(U) \partial_x H$$

# New approach

For  $s(U)=U^2$  and  $arphi=x+100\sin x$ 



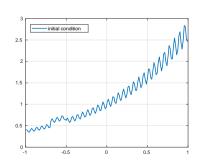


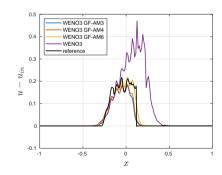


$$\partial_t U + \partial_x rac{U^2}{2} = S(U) \partial_x H$$

# New approach

For  $S(U)=u^2$  and  $H=x+100\sin x+$  top hat of perturbation  $\delta u=0.2\chi_{[-0.7,-0.5]}$ 



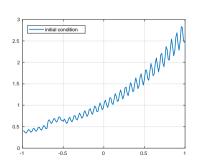


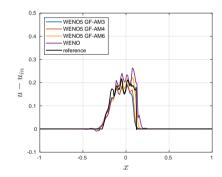


$$\partial_t \left[ egin{array}{c} h \ hu \end{array} 
ight] + \partial_x \left[ egin{array}{c} hu \ hu^2 + gh^2/2 \end{array} 
ight] = - \left[ egin{array}{c} 0 \ h \end{array} 
ight] b'(x)$$

New approach

For  $S(U)=u^2$  and  $H=x+100\sin x+$  top hat of perturbation  $\delta u=0.2\chi_{[-0.7,-0.5]}$ 



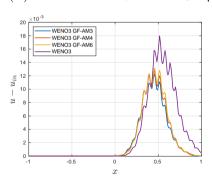


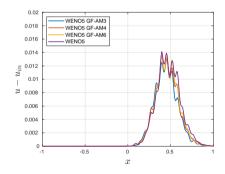
We seek solutions of the hyperbolic system of balance laws

$$\partial_t \left[ egin{array}{c} h \ hu \end{array} 
ight] + \partial_x \left[ egin{array}{c} hu \ hu^2 + gh^2/2 \end{array} 
ight] = - \left[ egin{array}{c} 0 \ h \end{array} 
ight] b'(x)$$

# New approach

For  $S(U)=u^2$  and  $H=x+100\sin x+$  top hat of Gaussian perturbation  $\alpha=0.005$ 



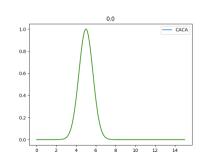


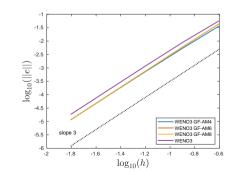


$$\partial_t \left[ egin{array}{c} h \ hu \end{array} 
ight] + \partial_x \left[ egin{array}{c} hu \ hu^2 + gh^2/2 \end{array} 
ight] = - \left[ egin{array}{c} 0 \ h \end{array} 
ight] b'(x)$$

# New approach

For S(U) = u - 1 and  $H(x) = e^{-(x - x_0 - ct)^2}$ ,  $u(x) = e^{-(x - x_0 - ct)^2}$ . The wave is centered at  $x_0 = 0.5$ .

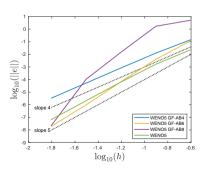


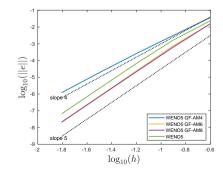


$$\partial_t \left[ egin{array}{c} h \ hu \end{array} 
ight] + \partial_x \left[ egin{array}{c} hu \ hu^2 + gh^2/2 \end{array} 
ight] = - \left[ egin{array}{c} 0 \ h \end{array} 
ight] b'(x)$$

### New approach

For S(U)=u-1 and  $H(x)=e^{-(x-x_0-ct)^2}$ ,  $u(x)=e^{-(x-x_0-ct)^2}$ . The wave is centered at  $x_0=0.5$ .



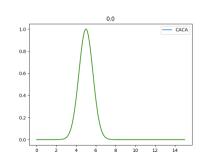


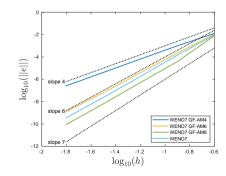


$$\partial_t \left[ egin{array}{c} h \ hu \end{array} 
ight] + \partial_x \left[ egin{array}{c} hu \ hu^2 + gh^2/2 \end{array} 
ight] = - \left[ egin{array}{c} 0 \ h \end{array} 
ight] b'(x)$$

## New approach

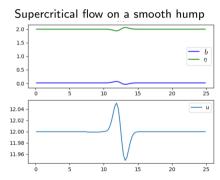
For S(U) = u - 1 and  $H(x) = e^{-(x - x_0 - ct)^2}$ ,  $u(x) = e^{-(x - x_0 - ct)^2}$ . The wave is centered at  $x_0 = 0.5$ .

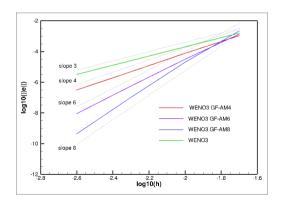






$$\partial_t \left[ egin{array}{c} h \ hu \end{array} 
ight] + \partial_x \left[ egin{array}{c} hu \ hu^2 + gh^2/2 \end{array} 
ight] = - \left[ egin{array}{c} 0 \ h \end{array} 
ight] b'(x)$$

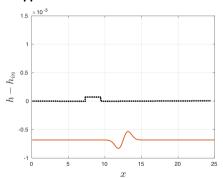


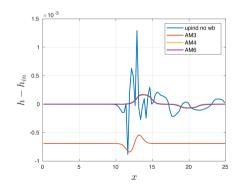




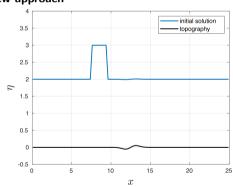
We seek solutions of the hyperbolic system of balance laws

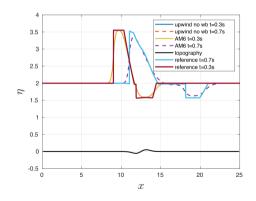
$$\partial_t \left[ egin{array}{c} h \ hu \end{array} 
ight] + \partial_x \left[ egin{array}{c} hu \ hu^2 + gh^2/2 \end{array} 
ight] = - \left[ egin{array}{c} 0 \ h \end{array} 
ight] b'(x)$$





$$\partial_t \left[ egin{array}{c} h \ hu \end{array} 
ight] + \partial_x \left[ egin{array}{c} hu \ hu^2 + gh^2/2 \end{array} 
ight] = - \left[ egin{array}{c} 0 \ h \end{array} 
ight] b'(x)$$

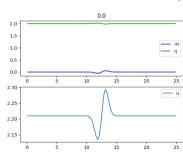


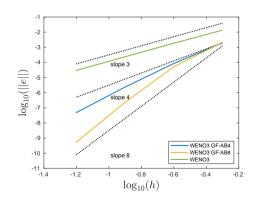


$$\partial_t \left[ egin{array}{c} h \ hu \end{array} 
ight] + \partial_x \left[ egin{array}{c} hu \ hu^2 + gh^2/2 \end{array} 
ight] = - \left[ egin{array}{c} 0 \ h \end{array} 
ight] b'(x)$$

# New approach

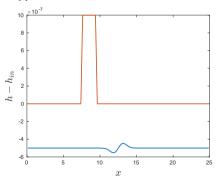
Subcritical flow on a smooth hump

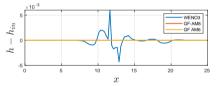


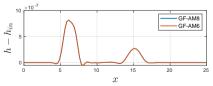


We seek solutions of the hyperbolic system of balance laws

$$\partial_t \left[ egin{array}{c} h \ hu \end{array} 
ight] + \partial_x \left[ egin{array}{c} hu \ hu^2 + gh^2/2 \end{array} 
ight] = - \left[ egin{array}{c} 0 \ h \end{array} 
ight] b'(x)$$



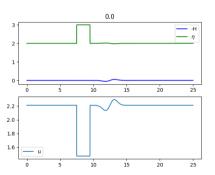


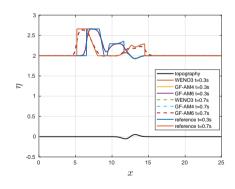




We seek solutions of the hyperbolic system of balance laws

$$\partial_t \left[ egin{array}{c} h \ hu \end{array} 
ight] + \partial_x \left[ egin{array}{c} hu \ hu^2 + gh^2/2 \end{array} 
ight] = - \left[ egin{array}{c} 0 \ h \end{array} 
ight] b'(x)$$





We seek solutions of the hyperbolic system of balance laws

$$\partial_t \left[egin{array}{c} h \ hu \end{array}
ight] + \partial_x \left[egin{array}{c} hu \ hu^2 + gh^2/2 \end{array}
ight] = - \left[egin{array}{c} 0 \ h \end{array}
ight] b'(x)$$

- ✓ No a-priori knowledge of equilibrium, all steady states!
- √ No need of compute the solution of the Cauchy problem .. (maybe for initialization)
- ✓ Convergence at steady state arbitrarily accurate by changing ODE weights
- ✓ Discrete initial state can be generated
- × Non-compact quadrature

Maria Kazolea\* Carlos Parés<sup>+</sup>, Mario Ricchiuto<sup>\*</sup>

Intro WENG GFq

### Take home message

- GFq in 1D: clever quadrature of the source based on high order ODE integrators
- Tremendous error reduction at stady-state (super-covergence)
- In one dimension discrete equilibria can be generated a-priori if necessary with the ODE solver

#### **Future**

- Sonic points ?
- discontinuous data ?
- adaptive ODE weights
- multiD

