We want to solve numerically hyperbolic systems of balance laws

$$\partial_t U + \partial_x F(U) = S(U) \partial_x H$$

Typical examples that we will use here

- Burger's equations
- Shallow water equations with topography

Intro

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We want to solve numerically hyperbolic systems of balance laws

$$\partial_t U + \partial_x F(U) = S(U) \partial_x H$$

Shopping list

- Arbitrary high order
- Possibly monotonicity/positivity preserving
- Steady state preserving

We want to solve numerically hyperbolic systems of balance laws

$$\partial_t U + \partial_x F(U) = S(U) \partial_x H$$

The objective of well balanced schemes is to preserve exactly or with enhanced accuracy some of the steady state solutions.

We focus here on the development of well-balanced high order FD methods for general systems of balance laws that preserve all the stationary solutions.

End

# Some approaches in the Bibliography for 1D

There are numberless approaches, and vast literature. Related work:

- Reconstruction/evolution of fluctuations wrt discrete equilibria (approximate full well balanced):
  - Castro & Parés J.Sci.Comp. 2020. Gómez-Bueno et al. Appl.Math.Comp. 2021. Gómez-Bueno et al athematics 2021. Guerrero Fernández et al Mathematics 2021, Gómez-Bueno et al. Appl.Num.Math. 2023, etc. etc.
- Well balanced via integration of the source term and global fluxes
  - Gascón & Corberán JCP 2001. Donat & Martinez-Gavara J.Sci.Comp. 2011. Chertock et al JCP 2018. Cheng et al J.Sci.Comp. 2019. Mantri&Noelle JCP 2021, Ciallella et al. J.Sci.Comp. 2023, Xu and Shu J.Sci.Comp. to appear, etc. etc.
- Other approaches
  - Well balanced via reconstruction/evolution of fluctuations wrt to a given equilibrium
  - Preserving  $V(u,\phi(x))=V_0$  via WB differencing or generalized polynomial approximations
  - Fully well balanced Riemann solver with 0-wave to enforce integral steady balance.



$$\partial_t U + \partial_x F(U) = S(U) \partial_x H$$

$$rac{dU_i}{dt} + rac{1}{\Delta x} \left( \widehat{\mathcal{F}}_{i+1/2} - \widehat{\mathcal{F}}_{i-1/2} 
ight) = 0$$

(1)

We seek solutions of the hyperbolic system of balance laws

$$\partial_t U + \partial_x F(U) = S(U) \partial_x H$$

$$\frac{dU_i}{dt} + \frac{1}{\Delta x} \left( \widehat{\mathcal{F}}_{i+1/2} - \widehat{\mathcal{F}}_{i-1/2} \right) = 0$$

① Look for the solution of the Cauchy problem  $U_i^*(x)$ 

$$egin{array}{lcl} F(U_i^*)_x & = & S(U_i^*(x)) H_x \ U_i^*(x_i) & = & U_i. \end{array}$$

- **2** Define  $\mathcal{F}_i = F(U_i) F(U_i^*(x_i)), j \in \{i r 1, \dots, i s 1\}$
- 3 Compute  $\hat{\mathcal{F}}_{i,i+1/2}$ ,  $\hat{\mathcal{F}}_{i,i-1/2}$  through a flux-reconstruction operator, eg WENO reconstruction on the stencil  $\mathcal{S}_i = i r, \dots, i + s$
- 4 Find the numerical fluxes using (for eg.) an upwind scheme.

### **Exactly well balanced**

We seek solutions of the hyperbolic system of balance laws

$$\partial_t U + \partial_x F(U) = S(U) \partial_x H$$

Some options:

- ① The solution  $U_i^*(x)$  of (1) if known or easy to find.
- ② Solve (1) numerically w.r.t U i.e  $U_x = J(U)^{-1}S(U)H_x$
- 3 or..

We seek solutions of the hyperbolic system of balance laws

$$\partial_t U + \partial_x F(U) = S(U) \partial_x H$$

$$rac{dU_i}{dt} + rac{1}{\Delta x} \left( \widehat{G}_{i+1/2} - \widehat{G}_{i-1/2} 
ight) = 0$$

$$\partial_t U + \partial_x F(U) = S(U) \partial_x H$$

$$rac{dU_i}{dt} + rac{1}{\Delta x} \left( \widehat{G}_{i+1/2} - \widehat{G}_{i-1/2} 
ight) = 0$$

In this approach we seek to embed directly in the discrete equation a consistency condition with a discrete approximation of the flux form of the Cauchy problem. So we integrate directly the latter with some approach.

$$\partial_t U + \partial_x F(U) = S(U) \partial_x H$$

#### Main result

**Proposition** (Discrete steady state). The WENO-FD scheme with global flux quadrature preserves exactly <u>continuous</u> discrete steady states  $U_i^* = U(F_i)$  with F obtained by integrating the ODE

$$F' = S(U(F))\partial_x H$$

using the multi-step ODE integrator with weights  $\{\omega_q\}_{q\geq 0}$  on spatial slabs of size h. If U(F) is a one to one mapping,  $U^*(x)$  verifies the consistency estimates of the multi-step scheme.

Multi-step schemes so far: Adams methods ABn and AMn

$$\partial_t U + \partial_x F(U) = S(U) \partial_x H$$

① Compute approximations  $R_j = R_{j-1} + \mathcal{I}_{j-1/2}(U), \ \forall i$  where

$$\mathcal{I}_{j-1/2} pprox \int_{x_{j-1}}^{x_j} S(U) H_x dx$$

- 2 Define  $\mathcal{F}_i = F(U_i) R_i$ ,  $j \in \{i r 1, \dots, i s 1\}$
- 3 Reconstruct WENO polynomials  $G_{i+1/2}^{\pm}(x)$  using  $\mathcal{F}_j$
- $oldsymbol{G}$  Compute upwind fluxes  $\widehat{G}_{i+1/2} = P_{i+1/2}^+ G_{i+1/2}^-(x_{i+1/2}) + P_{i+1/2}^- G_{i+1/2}^+(x_{i+1/2})$

Difference with exactly well balanced: The way in which the reference flux subtracted is obtained-Integration of the source

$$\partial_t U + \partial_x F(U) = S(U) \partial_x H$$

Here we compute

$$\mathcal{I}_{j-1/2} \approx \int_{x_{j-1}}^{x_j} S(U) H_x dx$$

using Adams ODE Integrators ie:

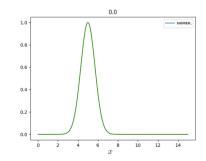
$$\mathcal{I}_{j-1/2} = \sum_{k=-1}^{m2} eta_k S(U_{j+k}) H_x(x_{j+k})$$

Since we will work with schemes of formal accuracy  $2k+1 \ge 3$  we consider ODE solvers whose solutions have accuracy at least  $2k+2 \ge 4$ .

- $\bullet$  Adams-Bashforth :  $m_1=4,6,8$  and  $m_2=-1$
- Adams-Moulton:  $m_1=3,5,7$  and  $m_2=0$

$$\partial_t U + \partial_x rac{U^2}{2} = S(U) \partial_x H$$

Moving solution (safety check): S(U)=U-C and  $H(x,t)=e^{-(x-x_0-Ct)^2}$ ,  $U_{\sf ex}=H(x,t)$ 



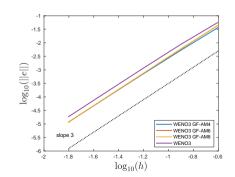
For WENOn-ABm or WENOn-AMn we expect convergence rates  $\min(m,v)$ 

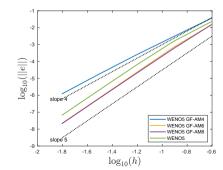
$$\epsilon = \Delta x^{\min(m,n)}$$

We seek solutions of the hyperbolic system of balance laws

$$\partial_t U + \partial_x rac{U^2}{2} = S(U) \partial_x H$$

Moving solution (safety check): S(U)=U-C and  $H(x,t)=e^{-(x-x_0-Ct)^2}$  ,  $U_{\sf ex}=H(x,t)$ 



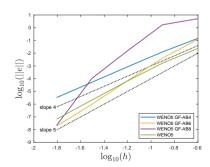


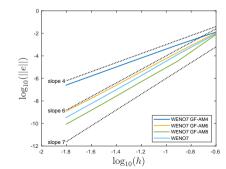


We seek solutions of the hyperbolic system of balance laws

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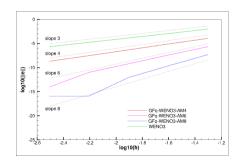


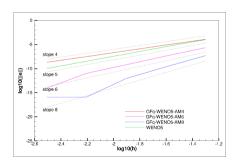


We seek solutions of the hyperbolic system of balance laws

$$\partial_t U + \partial_x rac{U^2}{2} = S(U) \partial_x H$$

Steady state:  $S(U)=U^2$  and H(x)=x, exact solution  $U_{\mathsf{ex}}(x)=e^x$ 

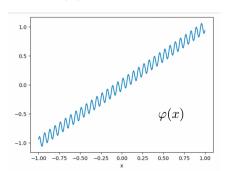


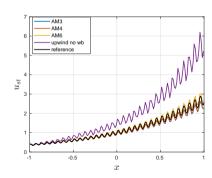




$$\partial_t U + \partial_x rac{U^2}{2} = S(U) \partial_x H$$

Steady state:  $S(U) = U^2$  and  $\varphi = x + 100 \sin x$ 

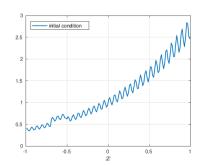


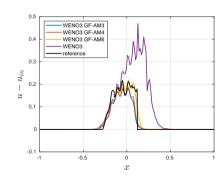




$$\partial_t U + \partial_x rac{U^2}{2} = S(U) \partial_x H$$

Perturbed steady state:  $S(U) = U^2$  and  $\varphi = x + 100 \sin x + \text{top hat } \delta u = 0.2\chi_{[-0.7, -0.5]}$ 



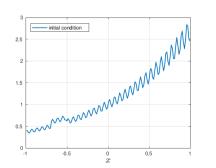


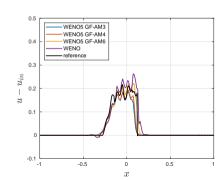
End



$$\partial_t U + \partial_x rac{U^2}{2} = S(U) \partial_x H$$

Perturbed steady state:  $S(U)=U^2$  and  $\varphi=x+100\sin x+$  top hat  $\delta u=0.2\chi_{[-0.7,-0.5]}$ 

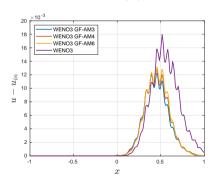


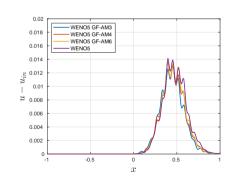


We seek solutions of the hyperbolic system of balance laws

$$\partial_t U + \partial_x rac{U^2}{2} = S(U) \partial_x H$$

Perturbed steady state:  $S(U)=U^2$  and  $\varphi=x+100\sin x+$  Gaussian with amplitude  $\alpha=0.005$ 



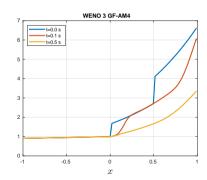


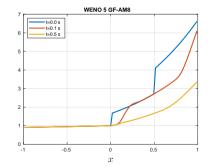


We seek solutions of the hyperbolic system of balance laws

$$\partial_t U + \partial_x rac{U^2}{2} = S(U) \partial_x H$$

**Discontinuous H**: 
$$S(U) = U^2$$
,  $H(x) = 0.1x|_{(x \le 0)}$ ,  $0.9 + x|_{(x > 0.5)}$ ,  $(0.5 + x)|_{(0 < x \le 0.5)}$ ,  $U_{\text{ex}}(x) = e^H$ 







$$\partial_t U + \partial_x rac{U^2}{2} = S(U) \partial_x H$$

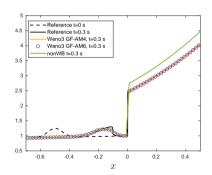
**Discontinuous H**:  $S(U) = U^2$ ,  $H(x) = 0.1x|_{(x < 0)}$ ,  $0.9 + x|_{(x > 0.5)}$ ,  $U_{\text{ex}}(x) = e^H$ 

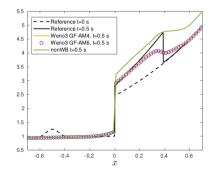
Lets assume that the discontinuity is at the middle of the cell  $[x_{i-1},\ x_i]$ 

- $\, \bullet \,$  Far from the discontinuity we use our standard  $\it n\text{-}step$  selected method.
- In the cell  $[x_{i-1}, x_i]$  we set  $I_{i-1/2} = \tilde{S}_{i-1/2}[[H]]_{i-1/2}$  where  $\tilde{S}_{i-1/2}$  linearization, such as  $[[F]]_{i-1/2} = \tilde{S}_{i-1/2}[[H]]_{i-1/2}$ .
- ullet For n cells after the discontinuity we use AM2.

$$\partial_t U + \partial_x rac{U^2}{2} = S(U) \partial_x H$$

**Discontinuous H**:  $S(U) = U^2$ ,  $H(x) = 0.1x|_{(x < 0)}$ ,  $0.9 + x|_{(x > 0.5)}$ ,  $U_{\text{ex}}(x) = e^H$ 



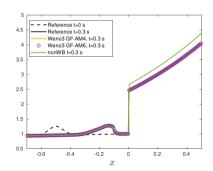


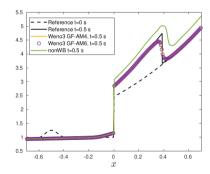


We seek solutions of the hyperbolic system of balance laws

$$\partial_t U + \partial_x rac{U^2}{2} = S(U) \partial_x H$$

**Discontinuous H**:  $S(U) = U^2$ ,  $H(x) = 0.1x|_{(x < 0)}$ ,  $0.9 + x|_{(x > 0.5)}$ ,  $U_{\text{ex}}(x) = e^H$ 



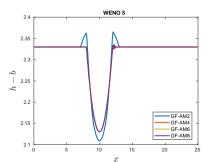


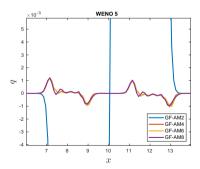


We seek solutions of the hyperbolic system of balance laws

$$\partial_t \left[ egin{array}{c} h \ hu \end{array} 
ight] + \partial_x \left[ egin{array}{c} hu \ hu^2 + gh^2/2 \end{array} 
ight] = - \left[ egin{array}{c} 0 \ h \end{array} 
ight] b'(x)$$

Lake at rest on a parabolic hump.





We seek solutions of the hyperbolic system of balance laws

$$\partial_t \left[ egin{array}{c} h \ hu \end{array} 
ight] + \partial_x \left[ egin{array}{c} hu \ hu^2 + gh^2/2 \end{array} 
ight] = - \left[ egin{array}{c} 0 \ h \end{array} 
ight] b'(x)$$

Method modification (with both AB and AM):

$$\hat{b}(x) = \sum_m L_m(x) b(x_m)$$

with  $\{x_m\}$  the points generating the AMm method. We proceed as follows

$$\int_{x_i}^{x_{i+1}} ghb'(x) dx \approx \int_{x_i}^{x_{i+1}} g\eta \hat{b}'(x) dx - \int_{x_i}^{x_{i+1}} g\left(\frac{\hat{b}^2}{2}\right)' dx = g \sum_m \omega_m \eta_m \hat{b}'(x_m) - g\left(\frac{b_{i+1}^2}{2} - \frac{b_{i}^2}{2}\right)$$

We seek solutions of the hyperbolic system of balance laws

$$\partial_t \left[ egin{array}{c} h \ hu \end{array} 
ight] + \partial_x \left[ egin{array}{c} hu \ hu^2 + gh^2/2 \end{array} 
ight] = - \left[ egin{array}{c} 0 \ h \end{array} 
ight] b'(x)$$

For lake at rest:

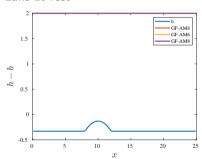
Then it easy to show that

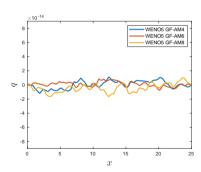
$$F_{i+1} - F_i + g \sum_m \omega_m \eta_m \hat{b}'(x_m) - g \left( \frac{b_{i+1}^2}{2} - \frac{b_i^2}{2} \right) = 0$$

We seek solutions of the hyperbolic system of balance laws

$$\partial_t \left[ egin{array}{c} h \ hu \end{array} 
ight] + \partial_x \left[ egin{array}{c} hu \ hu^2 + gh^2/2 \end{array} 
ight] = - \left[ egin{array}{c} 0 \ h \end{array} 
ight] b'(x)$$

#### Lake at rest

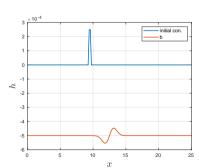


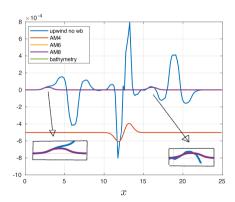


We seek solutions of the hyperbolic system of balance laws

$$\partial_t \left[ egin{array}{c} h \ hu \end{array} 
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ight] = - \left[ egin{array}{c} 0 \ h \end{array} 
ight] b'(x)$$

#### Lake at rest with perturbation



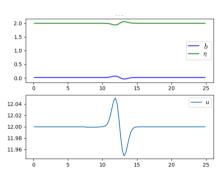


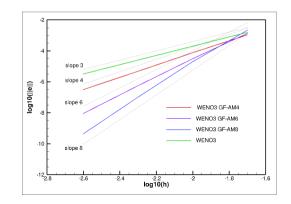


We seek solutions of the hyperbolic system of balance laws

$$\partial_t \left[ egin{array}{c} h \ hu \end{array} 
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ight] = - \left[ egin{array}{c} 0 \ h \end{array} 
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### Supercritical flow on a smooth hump



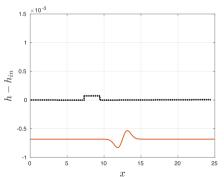


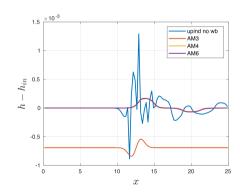


We seek solutions of the hyperbolic system of balance laws

$$\partial_t \left[ egin{array}{c} h \ hu \end{array} 
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### Supercritical flow with perturbation



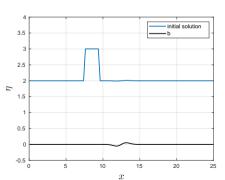


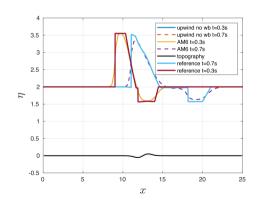


We seek solutions of the hyperbolic system of balance laws

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### Supercritical flow with perturbation

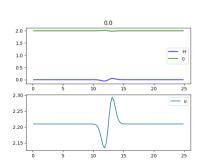


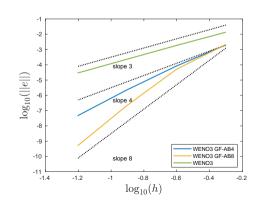




$$\partial_t \left[ egin{array}{c} h \ hu \end{array} 
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#### Subcritical flow on a smooth hump



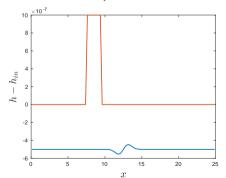


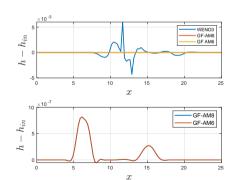


We seek solutions of the hyperbolic system of balance laws

$$\partial_t \left[ egin{array}{c} h \ hu \end{array} 
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### Subcritical flow with perturbation



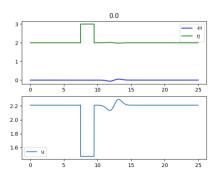


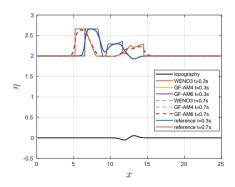


We seek solutions of the hyperbolic system of balance laws

$$\partial_t \left[ egin{array}{c} h \ hu \end{array} 
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### Subcritical flow with perturbation



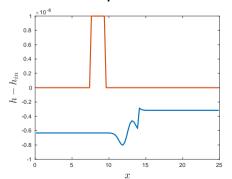


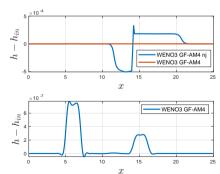


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$$\partial_t \left[egin{array}{c} h \ hu \end{array}
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ight] = - \left[egin{array}{c} 0 \ h \end{array}
ight] b'(x)$$

### Subcritical flow with perturbation and Discontinuous H



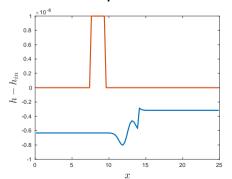


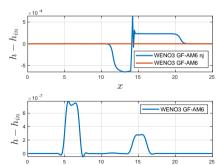


We seek solutions of the hyperbolic system of balance laws

$$\partial_t \left[ egin{array}{c} h \ hu \end{array} 
ight] + \partial_x \left[ egin{array}{c} hu \ hu^2 + gh^2/2 \end{array} 
ight] = - \left[ egin{array}{c} 0 \ h \end{array} 
ight] b'(x)$$

#### Subcritical flow with perturbation and Discontinuous H





x



- √ No a-priori knowledge of equilibrium, all steady states!
- √ No need of compute the solution of the Cauchy problem .. (maybe for initialization)
- √ Convergence at steady state arbitrarily accurate by changing ODE weights
- ✓ Discrete initial state can be generated
- × Non-compact quadrature

#### Take home message

- GFq in 1D: clever quadrature of the source based on high order ODE integrators
- Tremendous error reduction at stady-state (super-covergence)
- In one dimension discrete equilibria can be generated a-priori if necessary with the ODE solver

#### **Future**

- Sonic points ?
- adaptive ODE weights
- multiD

Thank you!

