

**Corresponding Manuscript:** Conditional Operation of Hole Spin Qubits above 1 K — [Notebook: S3\_Two\_Tone\_vs\_Jeff\_v160925\_v1.nb]

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**Project mirror:** <https://github.com/carmig00/Publications-OpenAccess-Code-Conditional2Q>

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**How to run:** Evaluate top-to-bottom (Evaluation → Evaluate Notebook).

**Outputs:** HDF5 (.h5) written to ./exports.

**Figure mapping:** This notebook reproduces the simulation of how different values for  $J_{\text{eff}}$  (ranging from 0-80 MHz) manifest in the two tone spectroscopy experiment. We conclude that  $J_{\text{eff}}$  is best represented by the difference in minima of the two-tone traces.

#### Comment/Note:

-The fast qubit (high  $f_{\text{Rabi}}$ ), at lower Larmor frequency, is the one located physically on the right and is represented by pink/purple color tones, here qubit 1.

- The slow qubit (low  $f_{\text{Rabi}}$ ), at higher Larmor frequency, is the one located physically on the left and is represented by orange/yellow color tones, here qubit 2.

(\* Parameters \*)

(\*  $f_{\text{Rabi}}$ , MHz \*)

(\* We assume this to be the bare Rabi frequency of the two ideally isolated spin qubits. Sometimes also referred to as bare Rabi frequency, for  $J=0$  \*)

$fR1t = 24;$

$fR2t = 13;$

(\* NOT USED HERE (because burst time is swept in Chevron):

Experimentally chosen pulse durations to induce a  $\pi/2$  flip for qubit 1 and qubit 2, see supplementary material for details on the calibration \*)

(\* units in  $\mu\text{s}$  \*)

$Tp1 = 20 \times 10^{-3};$  (\* corresponding to 20 ns \*)

$Tp2 = 30 \times 10^{-3};$  (\* corresponding to 30 ns \*)

(\* Larmor frequencies,  $2\pi \times \text{MHz}$  \*)

(\*NOTE: these values for Larmor frequencies were chosen as a

coarse guess by orienting ourselves on the center position of the EDSR

peaks which are not the most accurate representation of the precise

Larmor frequency (due to convolution of two Chevrons). In addition,

the values were rounded down to the next 10 MHz. Only after,

in the simulation of the 2-tone experiment we could more precisely narrow-

down and adjust the larmor frequencies as fit parameters. For seeing the effect of  $J_{\text{eff}}$  on an ideal 2-tone experiment, this choice was sufficiently accurate. \*)

$\omega_{z1} = 2 \pi 2090$ ;

$\omega_{z2} = 2 \pi 2280$ ;

(\*To simulate noise, we averaged over different Larmor frequencies sampled from a random, normal distributed set. \*)

(\*sigma of normal distribuion  $2 \pi \text{MHz}$  \*)

$\sigma \omega_{z1} = 25 \times 2 \pi$ ;

$\sigma \omega_{z2} = 25 \times 2 \pi$ ;

(\* Resolution of the array of random Larmor frequencies MHz \*)

$\text{res} \omega_{z1} = 150$ ;

$\text{res} \omega_{z2} = 150$ ;

(\* Arrays of random Larmor freqs. MHz \*)

$\omega_{z1\text{arr}} = \text{Abs}[\text{RandomVariate}[\text{NormalDistribution}[\omega_{z1}, \sigma \omega_{z1}], \{\text{res} \omega_{z1}\}]]$ ;

$\omega_{z2\text{arr}} = \text{Abs}[\text{RandomVariate}[\text{NormalDistribution}[\omega_{z2}, \sigma \omega_{z2}], \{\text{res} \omega_{z2}\}]]$ ;

(\* Values of effective exchange interaction  $J_{\text{eff}}$  \*)

(\* largest value of  $J$  (we start from  $J_{\text{eff}} = 0$ ) \*)

$J_{\text{max}} = 80$ ;

(\* resolution of  $J$  \*)

$\text{res} J = 40$ ;

(\*array of  $J_{\text{eff}}$  to be used in simulation\*)

$J_{\text{arr}} = \text{Table}[j_{\text{val}}, \{j_{\text{val}}, 0, J_{\text{max}}, J_{\text{max}} / (\text{res} J - 1)\}]]$ ;

(\*Resolution of frequency scanned (x-axis) MHz\*)

$\text{res} \omega = 160$ ;

(\*initial freq MHz\*)

$\omega_{\text{in}} = 2 \pi 1800$ ;

(\*final freq MHz\*)

$\omega_{\text{fin}} = 2 \pi 2600$ ;

(\*array of freqs. scanned over, MHz\*)

$\omega_{\text{arr}} = \text{Table}[t, \{t, \omega_{\text{in}}, \omega_{\text{fin}}, (\omega_{\text{fin}} - \omega_{\text{in}}) / (\text{res} \omega - 1)\}]]$ ;

(\* Main functions \*)

$\theta_{\text{so}} = 0$ ;

$\varphi_1 = 0$ ;

$\varphi_2 = 0$ ;

$\psi_{\theta a} = \{1, 0, 0, 0\}$ ;

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ψ0b = {0, 0, 0, 1};
resTg = 100;
(*initial time*)
Tg0 = 0;
(*final time, 66.6ns*)
Tgf = 3 * 1 / 30;
Tgarr = Table[t, {t, Tg0, Tgf,  $\frac{Tgf - Tg0}{resTg - 1}$ }}];

{nx, ny, nz} = {1, 0, 0};
Jpar[J_] = 2 Pi J (nz2 + (1 - nz2) Cos[θso]);
(*λx00=2 Pi 30; (*Russ/Burkard resonance condition for driving*)*)
HQ[ωz1in_, ωz2in_] = KroneckerProduct[PauliMatrix[3]  $\frac{\omega z1in}{2}$ , PauliMatrix[0]] +
    KroneckerProduct[PauliMatrix[0], PauliMatrix[3]  $\frac{\omega z2in}{2}$ ];
Rso = {
    {nx2 + (1 - nx2) Cos[θso],
    nx ny - nx ny Cos[θso] - nz Sin[θso], nx nz - nx nz Cos[θso] + ny Sin[θso]},
    {nx ny - nx ny Cos[θso] + nz Sin[θso], ny2 + (1 - ny2) Cos[θso],
    ny nz - ny nz Cos[θso] - nx Sin[θso]}, {nx nz - nx nz Cos[θso] - ny Sin[θso],
    2 Sin[ $\frac{\theta so}{2}$ ] (nx Cos[ $\frac{\theta so}{2}$ ] + ny nz Sin[ $\frac{\theta so}{2}$ ]), nz2 + (1 - nz2) Cos[θso]}
};
HJ[J_] =
 $\frac{J}{4}$  Sum[KroneckerProduct[PauliMatrix[i], (Rso.Array[PauliMatrix, 3])[[i]]], {i, 1, 3}];

(*turn on time of pulses*)
dt0 = 0.001;
(*wait time between pulses*)
Tw = 0.001;

(*driving amplitudes, here normalised to 2 Pi and equal to each other. One could choose to
e.g. drive one qubit harder than the other one. (this could lead to over-fitting)*)
λx10 = 2 Pi; (*Drive amplitude Q1*)
λx20 = 2 Pi; (*Drive amplitude Q2*)

(*first pulse for a time Tflip1*)

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$$\lambda x1[t_] = \lambda x10 \left( \frac{1}{2} + \frac{1}{2} \text{Erf}[(Tflip1 - t) / dt0] \right);$$

(*second pulse after a time Tflip1+Tw*)

$$\lambda x2[t_] = \lambda x20 \left( \frac{1}{2} + \frac{1}{2} \text{Erf}[(t - Tflip1 - Tw) / dt0] \right);$$


(* Drive Hamiltonian *)
H $\lambda$ [t_] =  $\lambda x1[t]$  Cos[ $\omega x1$  t +  $\phi 1$ ]
  (fR1 KroneckerProduct[{nx, ny, nz}.Array[PauliMatrix, 3], PauliMatrix[0]] + fR2
    KroneckerProduct[PauliMatrix[0], {nx, ny, nz}.Array[PauliMatrix, 3]]) +  $\lambda x2[t]$  Cos[
     $\omega x2$  t +  $\phi 2$ ] (fR1 KroneckerProduct[{nx, ny, nz}.Array[PauliMatrix, 3], PauliMatrix[0]] +
    fR2 KroneckerProduct[PauliMatrix[0], {nx, ny, nz}.Array[PauliMatrix, 3]]);

(*Total Hamiltonian*)
Htot[t_, J_,  $\omega z1in$ _,  $\omega z2in$ _] = HQ[ $\omega z1in$ ,  $\omega z2in$ ] + HJ[J] + H $\lambda$ [t];
 $\rho Sf$  =
Function[{ $\psi 0$ ,  $\psi ro$ , Tflip1in, Tflip2in,  $\omega z1in$ ,  $\omega z2in$ ,  $\omega x1in$ ,  $\omega x2in$ , fR1in, fR2in, Jin}, sub =
  {Tflip1 → Tflip1in, Tflip2 → Tflip2in,  $\omega x1$  →  $\omega x1in$ ,  $\omega x2$  →  $\omega x2in$ , fR1 → fR1in, fR2 → fR2in};
  Tfin = Tflip1 + Tw + Tflip2 /. sub;
   $\psi fin[t_] = \{a[t], b[t], c[t], d[t]\} /. \text{NDSolve}[$ 
    {Htot[t, Jin,  $\omega z1in$ ,  $\omega z2in$ ].{a[t], b[t], c[t], d[t]} ==
       $i \hbar D[\{a[t], b[t], c[t], d[t]\}, t] /. sub,$ 
    a[0] ==  $\psi 0[[1]]$ , b[0] ==  $\psi 0[[2]]$ , c[0] ==  $\psi 0[[3]]$ , d[0] ==  $\psi 0[[4]]$ },
    {a[t], b[t], c[t], d[t]}, {t, 0, Tfin},
    Method → "PDEDiscretization" → {"MethodOfLines",
      "SpatialDiscretization" → "FiniteElement", MaxCellMeasure → 0.0005 / fR1}][[1]];
  (*readout signal*)
   $\sum_{i=1}^{\text{Dimensions}[\psi ro][[1]]} \text{Abs}[(\psi ro[[i]].\psi fin[Tfin])]^2];$ 

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(*Two-tone spectroscopy vs J, averaged over different Larmor frequencies*)
(* The four frequencies f1 to f4 are chosen to represent
the experimentally "ideal" case, where:  $\omega z1(2) \pm \frac{Jpar[Jval]}{2}$  *)
(* We read out in the 0,1,0,0 and *)
(* 0,0,1,0 states (anti-parallel) which are not blocked in PSB *)
(*NOTE: Below, we apply the Pi pulse Tp1 for f1 and f2 and Tp2 for f3 and f4. *)

resF1vsJ =
ParallelTable[( $\rho Sf[\psi 0a, \{\{0, 1, 0, 0\}, \{0, 0, 1, 0\}\}, Tp1, Tp1, \omega z1arr[[j]], \omega z2arr[[j]],$ 
 $\omega z1 - Jpar[Jarr[[k]] / 2, \omega arr[[i]], fR1t, fR2t, Jpar[Jarr[[k]]] +$ 
 $\rho Sf[\psi 0b, \{\{0, 1, 0, 0\}, \{0, 0, 1, 0\}\}, Tp1, Tp1, \omega z1arr[[j]], \omega z2arr[[j]],$ 
 $\omega z1 - Jpar[Jarr[[k]] / 2, \omega arr[[i]], fR1t, fR2t, Jpar[Jarr[[k]]] ) / 2,$ 
{j, 1, res $\omega z1$ }, {i, 1, res $\omega$ }, {k, 1, resJ}]; // AbsoluteTiming
resF2vsJ =
ParallelTable[( $\rho Sf[\psi 0a, \{\{0, 1, 0, 0\}, \{0, 0, 1, 0\}\}, Tp1, Tp1, \omega z1arr[[j]], \omega z2arr[[j]],$ 
 $\omega z1 + Jpar[Jarr[[k]] / 2, \omega arr[[i]], fR1t, fR2t, Jpar[Jarr[[k]]] +$ 
 $\rho Sf[\psi 0b, \{\{0, 1, 0, 0\}, \{0, 0, 1, 0\}\}, Tp1, Tp1, \omega z1arr[[j]], \omega z2arr[[j]],$ 
 $\omega z1 + Jpar[Jarr[[k]] / 2, \omega arr[[i]], fR1t, fR2t, Jpar[Jarr[[k]]] ) / 2,$ 
{j, 1, res $\omega z1$ }, {i, 1, res $\omega$ }, {k, 1, resJ}]; // AbsoluteTiming
resF3vsJ =
ParallelTable[( $\rho Sf[\psi 0a, \{\{0, 1, 0, 0\}, \{0, 0, 1, 0\}\}, Tp2, Tp2, \omega z1arr[[j]], \omega z2arr[[j]],$ 
 $\omega z2 - Jpar[Jarr[[k]] / 2, \omega arr[[i]], fR1t, fR2t, Jpar[Jarr[[k]]] +$ 
 $\rho Sf[\psi 0b, \{\{0, 1, 0, 0\}, \{0, 0, 1, 0\}\}, Tp2, Tp2, \omega z1arr[[j]], \omega z2arr[[j]],$ 
 $\omega z2 - Jpar[Jarr[[k]] / 2, \omega arr[[i]], fR1t, fR2t, Jpar[Jarr[[k]]] ) / 2,$ 
{j, 1, res $\omega z1$ }, {i, 1, res $\omega$ }, {k, 1, resJ}]; // AbsoluteTiming
resF4vsJ =
ParallelTable[( $\rho Sf[\psi 0a, \{\{0, 1, 0, 0\}, \{0, 0, 1, 0\}\}, Tp2, Tp2, \omega z1arr[[j]], \omega z2arr[[j]],$ 
 $\omega z2 + Jpar[Jarr[[k]] / 2, \omega arr[[i]], fR1t, fR2t, Jpar[Jarr[[k]]] +$ 
 $\rho Sf[\psi 0b, \{\{0, 1, 0, 0\}, \{0, 0, 1, 0\}\}, Tp2, Tp2, \omega z1arr[[j]], \omega z2arr[[j]],$ 
 $\omega z2 + Jpar[Jarr[[k]] / 2, \omega arr[[i]], fR1t, fR2t, Jpar[Jarr[[k]]] ) / 2,$ 
{j, 1, res $\omega z1$ }, {i, 1, res $\omega$ }, {k, 1, resJ}]; // AbsoluteTiming

Export["C:\\Users\\export\\two_tone_vs_J_average_Larmor_sigma" <> TextString[ $\sigma \omega z1$ ] <>
"MHZ_J" <> TextString[Jval] <> "MHZ_vDDMMYY_vX.h5", {resF1vsJ, resF2vsJ, resF3vsJ,
resF4vsJ,  $\omega z1arr / (2 \text{ Pi})$ ,  $\omega z2arr / (2 \text{ Pi})$ ,  $\omega arr / (2 \text{ Pi})$ , Jarr}, {"Datasets",
{"resF1vsJ", "resF2vsJ", "resF3vsJ", "resF4vsJ", "fR1arr", "fR2arr", "farr", "Jarr"}}];

Out[ ]=
{5491.59, Null}

Out[ ]=
{5476.5, Null}

Out[ ]=
{8083.16, Null}

Out[ ]=
{7948.77, Null}

```