

Corresponding Manuscript: Conditional Operation of Hole Spin Qubits above 1 K — [Notebook: F2-c-d_color_f1-f4_Two_Tone_J38.7MHZ_v040925_v9.nb]

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Project mirror: <https://github.com/carmig00/Publications-OpenAccess-Code-Conditional2Q>

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How to run: Evaluate top-to-bottom (Evaluation → Evaluate Notebook).

Outputs: HDF5 (.h5) written to ./exports. Human readable output available as PDF.

Figure mapping: This notebook reproduces the four coloured 2-tone traces in the lower two panels of Figure 2 c and d (SIM 2-Tone F2 c, d)

Comment/Note:

- The fast qubit (high f_{Rabi}), at lower Larmor frequency, is the one located physically on the right and is represented by pink/purple color tones, here qubit 1.
- The slow qubit (low f_{Rabi}), at higher Larmor frequency, is the one located physically on the left and is represented by orange/yellow color tones, here qubit 2.

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(* Two-tone spectroscopy, Parameters *)

(* f_Rabi,MHz *)

(* We assume this to be the bare Rabi frequency of the two ideally isolated
   spin qubits. Sometimes also referred to as bare Rabi frequency,for J=0 *)
fR1t = 24;
fR2t = 13;
(* Effective exchnage interaction J, MHz*)
Jval = 38.7;
(*Larmor frequencies, 2*Pi*MHz*)
 $\omega_{z1} = 2 \pi 2075$ ;
 $\omega_{z2} = 2 \pi 2270$ ;

(* Experimentally chosen pulse durations to induce a Pi/2 flip for qubit 1
   and qubit 2, see supplementary material for details on the calibration *)
(* units in \mus *)
Tp1 =  $20 \times 10^{-3}$ ; (* corresponding to 20 ns *)
Tp2 =  $30 \times 10^{-3}$ ; (* corresponding to 30 ns *)

(* Empirical frequency choices for f1-f4, see supplementary information. *)
(* 2*Pi*MHz *)
 $\omega_{\text{exp1}} = 2 \pi 2057$ ;
 $\omega_{\text{exp2}} = 2 \pi 2094$ ;
 $\omega_{\text{exp3}} = 2 \pi 2257$ ;
 $\omega_{\text{exp4}} = 2 \pi 2284$ ;

(*To simulate noise,we averaged over different Larmor
   frequencies sampled from a random,normal distributed set.*)
(*sigma of normal distributuion 2*Pi*MHz*)
 $\sigma_{\omega_{z1}} = 25 \times 2 \pi$ ;
 $\sigma_{\omega_{z2}} = 25 \times 2 \pi$ ;

(*Resolution of the array of random Larmor frequencies MHz*)
res $\omega_{z1}$  = 150;
res $\omega_{z2}$  = 150;
(*Arrays of random Larmor freqs. MHz*)
 $\omega_{z1\text{arr}} = \text{Abs}[\text{RandomVariate}[\text{NormalDistribution}[\omega_{z1}, \sigma_{\omega_{z1}}], \{\text{res}\omega_{z1}\}]]$ ;
 $\omega_{z2\text{arr}} = \text{Abs}[\text{RandomVariate}[\text{NormalDistribution}[\omega_{z2}, \sigma_{\omega_{z2}}], \{\text{res}\omega_{z2}\}]]$ ;

(*Resolution of frequency scanned (x-axis) MHz*)
res $\omega$  = 160;
(*Define the range of scanned frequencies (x-axis) MHz*)
(*initial freq MHz*)
 $\omega_{\text{in}} = 2 \pi 1800$ ;
(*final freq MHz*)
 $\omega_{\text{fin}} = 2 \pi 2600$ ;
(*array of freqs. scanned over, MHz*)
 $\omega_{\text{arr}} = \text{Table}[t, \{t, \omega_{\text{in}}, \omega_{\text{fin}}, (\omega_{\text{fin}} - \omega_{\text{in}}) / (\text{res}\omega - 1)\}];$ 

(* Main calculation core *)

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(* Theta is the SO-angle. In an accurate and complete microscopic description,
his accounts for the angle by which the spin is rotated due to SOI,
i.e. the angle in the rotation matrix R. To avoid over-
fitting this angle is set to 0 as we consider an effective J,
J_eff between the two qubits *)
 $\theta_{SO} = 0;$ 
(* Phases of the drives. *)
 $\varphi_1 = 0;$ 
 $\varphi_2 = 0;$ 

(* Populations:
These are the states that are blocked in PSB and which we initialize in. *)
 $\psi_{0a} = \{1, 0, 0, 0\};$ 
 $\psi_{0b} = \{0, 0, 0, 1\};$ 

(* not used here, only for time scans *)
resTg = 100; (*not used here, since no time scan*)
(*initial time*)
Tg0 = 0; (*not used here, since no time scan*)
(*final time, 66.6ns*)
Tgf = 3 * 1 / 30; (*not used here, since no time scan*)
Tgarr = Table[t, {t, Tg0, Tgf,  $\frac{Tgf - Tg0}{resTg - 1}$  }]; (*not used here, since no time scan*)

(* SO-vector *)
{nx, ny, nz} = {1, 0, 0};

(*Effective exchange interaction*)
(*Relevant component parallel with respect to external magnetic field B.*)
Jpar[J_] = 2 Pi J (nz2 + (1 - nz2) Cos[ $\theta_{SO}$ ]);

(*2-Qubit Hamiltonian*)
HQ[ $\omega_{z1in}$ _,  $\omega_{z2in}$ _] = KroneckerProduct[PauliMatrix[3]  $\frac{\omega_{z1in}}{2}$ , PauliMatrix[0]] +
KroneckerProduct[PauliMatrix[0], PauliMatrix[3]  $\frac{\omega_{z2in}}{2}$ ];
Rso = {
{nx2 + (1 - nx2) Cos[ $\theta_{SO}$ ],
nx ny - nx ny Cos[ $\theta_{SO}$ ] - nz Sin[ $\theta_{SO}$ ], nx nz - nx nz Cos[ $\theta_{SO}$ ] + ny Sin[ $\theta_{SO}$ ]},
{nx ny - nx ny Cos[ $\theta_{SO}$ ] + nz Sin[ $\theta_{SO}$ ], ny2 + (1 - ny2) Cos[ $\theta_{SO}$ ],
ny nz - ny nz Cos[ $\theta_{SO}$ ] - nx Sin[ $\theta_{SO}$ ]}, {nx nz - nx nz Cos[ $\theta_{SO}$ ] - ny Sin[ $\theta_{SO}$ ],
2 Sin[ $\frac{\theta_{SO}}{2}$ ] (nx Cos[ $\frac{\theta_{SO}}{2}$ ] + ny nz Sin[ $\frac{\theta_{SO}}{2}$ ]), nz2 + (1 - nz2) Cos[ $\theta_{SO}$ ]}
};

(*Exchange Hamiltonian*)

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HJ[J_] =
  
$$\frac{J}{4} \text{Sum}[\text{KroneckerProduct}[\text{PauliMatrix}[i], (\text{Rso.Array}[\text{PauliMatrix}, 3])[[i]], \{i, 1, 3\}];$$


(*turn on time of pulses*)
dt0 = 0.001;
(*wait time between pulses*)
Tw = 0.001;

(*driving amplitudes, here normalised to 2 Pi and equal to each other. One could choose
to e.g. drive one qubit harder than the other one. (this could lead to over-fitting)*)
λx10 = 2 Pi; (*Drive amplitude Q1*)
λx20 = 2 Pi; (*Drive amplitude Q2*)

(*Generic case has two tones. Second tone is applied after
passing of the first duration Tflip1 and the waiting time.*)
(*First pulse for a time Tflip1. This amplitude
is an error function turning on/off the pulse.*)
λx1[t_] = λx10  $\left( \frac{1}{2} + \frac{1}{2} \text{Erf}[(\text{Tflip1} - t) / \text{dt0}] \right);$ 
(*second pulse after a time Tflip1+Tw*)
λx2[t_] = λx20  $\left( \frac{1}{2} + \frac{1}{2} \text{Erf}[(t - \text{Tflip1} - \text{Tw}) / \text{dt0}] \right);$ 

(* Drive Hamiltonian *)
Hλ[t_] = λx1[t] Cos[ωx1 t + φ1]
  (fR1 KroneckerProduct[{nx, ny, nz}.Array[PauliMatrix, 3], PauliMatrix[0]] +
   fR2 KroneckerProduct[PauliMatrix[0],
    {nx, ny, nz}.Array[PauliMatrix, 3]]) + λx2[t] Cos[ωx2 t + φ2]
  (fR1 KroneckerProduct[{nx, ny, nz}.Array[PauliMatrix, 3], PauliMatrix[0]] +
   fR2 KroneckerProduct[PauliMatrix[0], {nx, ny, nz}.Array[PauliMatrix, 3]]);

(*Total Hamiltonian*)
Htot[t_, J_, ωz1in_, ωz2in_] = HQ[ωz1in, ωz2in] + HJ[J] + Hλ[t];
ρSf =
  Function[{ψ0, ψro, Tflip1in, Tflip2in, ωz1in, ωz2in, ωx1in, ωx2in, fR1in, fR2in, Jin},
    sub = {Tflip1 → Tflip1in, Tflip2 → Tflip2in, ωx1 → ωx1in,
      ωx2 → ωx2in, fR1 → fR1in, fR2 → fR2in};
    Tfin = Tflip1 + Tw + Tflip2 /. sub;
    ψfin[t_] = {a[t], b[t], c[t], d[t]} /. NDSolve[
      {Htot[t, Jin, ωz1in, ωz2in].{a[t], b[t], c[t], d[t]} ==
         $\mathbb{I} D[\{a[t], b[t], c[t], d[t]\}, t] /. sub,$ 
      a[0] == ψ0[[1]], b[0] == ψ0[[2]], c[0] == ψ0[[3]], d[0] == ψ0[[4]],
      {a[t], b[t], c[t], d[t]}, {t, 0, Tfin},
      Method → "PDEDiscretization" → {"MethodOfLines", "SpatialDiscretization" →
        "FiniteElement", MaxCellMeasure → 0.0005 / fR1}][[1]];
    (*readout signal*)

```

$$\sum_{i=1}^{\text{Dimensions}[\psi_{\text{ro}}][1]} \text{Abs}[(\psi_{\text{ro}}[i] \cdot \psi_{\text{fin}}[T_{\text{fin}}])]^2];$$

(* Two-tone experiment with empirical frequencies*)

(* NOTE: In an ideal experiment performed on a perfectly characterized system for which all parameters are known, the four frequencies f1 to f4 would be chosen to be $\omega_{z1}(2) \pm \frac{J_{\text{par}}[J_{\text{val}}]}{2}$.*)

(* We performed the experiment with the fixed frequency guesses introduced in the supplementary information, denoted as ω_{xexp1} to ω_{xexp4} above. We further use ω_{z1} and ω_{z2} as fit parameters and adjusted them such that the simulation run best matched our data. The values for the Larmor frequencies ω_{z1} and ω_{z2} found herein are the ones we settled with.*)

(* We read out in the 0,1,0,0 and *)

(* 0,0,1,0 states (anti-parallel) which are not blocked in PSB *)

(*NOTE: Below, we apply the Pi pulse Tp1 for f1 and f2 and Tp2 for f3 and f4. *)

```
resF1 = ParallelTable[{j, warr[[i]] / (2 Pi 10^3),
  (ρSf[ψ0a, {{0, 1, 0, 0}, {0, 0, 1, 0}}, Tp1, Tp1, ωz1arr[[j]], ωz2arr[[j]], ωxexp1,
    warr[[i]], fR1t, fR2t, Jpar[Jval]] + ρSf[ψ0b, {{0, 1, 0, 0}, {0, 0, 1, 0}}, Tp1,
    Tp1, ωz1arr[[j]], ωz2arr[[j]], ωxexp1, warr[[i]], fR1t, fR2t, Jpar[Jval]]) / 2},
  {i, 1, resω}, {j, 1, resωz1}]; // AbsoluteTiming
resF2 = ParallelTable[{j, warr[[i]] / (2 Pi 10^3),
  (ρSf[ψ0a, {{0, 1, 0, 0}, {0, 0, 1, 0}}, Tp1, Tp1, ωz1arr[[j]], ωz2arr[[j]], ωxexp2,
    warr[[i]], fR1t, fR2t, Jpar[Jval]] + ρSf[ψ0b, {{0, 1, 0, 0}, {0, 0, 1, 0}}, Tp1,
    Tp1, ωz1arr[[j]], ωz2arr[[j]], ωxexp2, warr[[i]], fR1t, fR2t, Jpar[Jval]]) / 2},
  {i, 1, resω}, {j, 1, resωz1}]; // AbsoluteTiming
resF3 = ParallelTable[{j, warr[[i]] / (2 Pi 10^3),
  (ρSf[ψ0a, {{0, 1, 0, 0}, {0, 0, 1, 0}}, Tp2, Tp2, ωz1arr[[j]], ωz2arr[[j]], ωxexp3,
    warr[[i]], fR1t, fR2t, Jpar[Jval]] + ρSf[ψ0b, {{0, 1, 0, 0}, {0, 0, 1, 0}}, Tp2,
    Tp2, ωz1arr[[j]], ωz2arr[[j]], ωxexp3, warr[[i]], fR1t, fR2t, Jpar[Jval]]) / 2},
  {i, 1, resω}, {j, 1, resωz1}]; // AbsoluteTiming
resF4 = ParallelTable[{j, warr[[i]] / (2 Pi 10^3),
  (ρSf[ψ0a, {{0, 1, 0, 0}, {0, 0, 1, 0}}, Tp2, Tp2, ωz1arr[[j]], ωz2arr[[j]], ωxexp4,
    warr[[i]], fR1t, fR2t, Jpar[Jval]] + ρSf[ψ0b, {{0, 1, 0, 0}, {0, 0, 1, 0}}, Tp2,
    Tp2, ωz1arr[[j]], ωz2arr[[j]], ωxexp4, warr[[i]], fR1t, fR2t, Jpar[Jval]]) / 2},
  {i, 1, resω}, {j, 1, resωz1}]; // AbsoluteTiming
Export["C:\\Users\\exports\\two_tone_average_Larmor_sigma" <>
  TextString[σωz1] <> "MHz_J" <> TextString[Jval] <> "MHz_vDDMMYY_vX.h5",
  {resF1, resF2, resF3, resF4,  $\frac{\omega_{z1arr}}{2 \text{ Pi}}$ ,  $\frac{\omega_{z2arr}}{2 \text{ Pi}}$ ,  $\frac{warr}{2 \text{ Pi}}$ },
  {"Datasets", {"resF1", "resF2", "resF3", "resF4", "fR1arr", "fR2arr", "farr"}}];
```

Out[8]=

{104.53, Null}

Out[9]=

{98.8663, Null}

Out[*n*]=

{144.783, Null}

Out[*n*]=

{142.309, Null}

(* Note MJC v9 is 2075 and 2270 and J 38.7 exact*)