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Lecture Notes Advanced Quantum Mechanics

Transforming $\psi^{\dagger}(x)$

$$\psi'(x') = \left(1 - \frac{i}{2} \omega_{\mu\nu} S^{\mu\nu}\right) \psi(x)$$

$$\psi'^{\dagger}(x') = \psi^{\dagger}(x) \left(1 + \frac{i}{2} \omega_{\mu\nu} S^{\mu\nu\dagger} \right)$$

 $S^{\mu\nu}$: Hermitian for rotations $(S^{ij};\ i,j>0)$ and anti-Hermitian for boosts $(S^{0j} \text{ for } j>0)$

Property: $S^{\mu\nu\dagger} \rightarrow \gamma^0 S^{\mu\nu} \gamma^0$ (Why?)

So,

$$\begin{split} \psi'^\dagger(x')\gamma^0 &= \psi^\dagger(x) \left(1 + \frac{i}{2}\omega_{\mu\nu}\gamma^0 S^{\mu\nu}\gamma^0\right)\gamma^0 \\ \bar{\psi}'(x') &= \psi^\dagger(x) \left(\gamma^0 + \frac{i}{2}\omega_{\mu\nu}\gamma^0 S^{\mu\nu}\right) \\ \bar{\psi}'(x') &= \bar{\psi}(x) \left(\mathbb{1} + \frac{i}{2}\omega_{\mu\nu}S^{\mu\nu}\right) \end{split}$$

The product

$$\bar{\psi}'(x')\psi'(x') = \bar{\psi}(x)\left(\mathbb{1} + \frac{i}{2}\omega_{\mu\nu}S^{\mu\nu}\right)\left(\mathbb{1} - \frac{i}{2}\omega_{\mu\nu}S^{\mu\nu}\right)$$
$$\sim \bar{\psi}(x)\psi(x)$$
$$(\bar{\psi}\psi)'(x') = (\bar{\psi}\psi)(x)$$

Therefore $\bar{\psi}\psi$ transforms like a scalar field.

Transformation of $\bar{\psi}(x)\gamma^{\mu}\psi(x)$ under Lorentz transform

Transformed: $\bar{\psi}'(x')\gamma^{\mu}\psi'(x')$ use the transformations if ... (I didn't understand this part) The commutator for $S^{\mu\nu}$ and γ^{μ} pops up (figure it out) Takes up a form,

$$\bar{\psi}(x')\gamma^{\mu}\psi(x') = \left(\delta^{\mu}_{\sigma} - \frac{i}{2}\omega_{\lambda\nu}(J^{\lambda\nu})^{\mu}_{\sigma}\right)\bar{\psi}(x)\gamma^{\mu}\psi(x)$$

same transformation is used for vector field

 $\bar{\psi}(x)\gamma^{\mu}\partial_{\mu}\psi(x)$ will transform like a scalar field.

Covariant vectors:
$$\begin{split} x'_{\mu} &= (\Lambda^{-1})^{\nu}_{\mu} x_{\nu} \\ &\text{Contra: } x'^{\mu} = \Lambda^{\mu}_{\nu} x^{\nu} \\ &x'_{\mu} x'^{\mu} = x_{\mu} x^{\mu}, \quad \left[\because (\Lambda^{-1})^{\nu}_{\mu} (\Lambda)^{\mu}_{\lambda} = \delta^{\nu}_{\lambda} \right] \\ &\frac{\partial}{\partial x'^{\mu}} = \left((\Lambda^{-1})^{\nu}_{\mu} \frac{\partial}{\partial x^{\nu}} \right) \end{split}$$

 \hookrightarrow this transformation was missing

Constructing a Lagrangian

For a scalar field we used $\phi\phi^*$ (i.e. $\partial^{\mu}\phi^*\partial_{\mu}\phi$). For the fields under a Lorentz transformation

$$\begin{split} \mathcal{L} = & i \bar{\psi} \gamma^{\mu} \partial_{\mu} \psi - k \bar{\psi} \psi \\ \hookrightarrow & \bar{\psi} \gamma^{0} \frac{1}{c} \frac{d\psi}{dt} + \bar{\psi} \vec{\gamma} \cdot \nabla \psi \\ \text{with } \vec{\gamma} = \begin{pmatrix} \gamma^{1} \\ \gamma^{2} \\ \gamma^{3} \end{pmatrix} \end{split}$$

 $\gamma^0 \to {
m Hermitian}$ (time-like part) $\vec{\gamma} \to {
m anti-Hermitian}$ (spatial part)

$$\gamma^0 = \gamma_0$$

$$\gamma_\mu = -\gamma^\mu$$

$$\gamma_\mu = g_{\mu\nu}\gamma^\nu$$

the metric tensor transforms like the Dirac matrices

Equation of motion for the constructed lagrangian of ψ

$$\partial_{\mu} \left(\frac{\partial \mathcal{L}}{\partial (\partial_{\mu} \phi_{l})} \right) = \frac{\partial \mathcal{L}}{\partial \phi_{l}}, \quad l = 1, ..N$$

We are treating ψ and $\bar{\psi}$ as independent variables So, for the Lagrangian,

$$(i\gamma^\mu\partial_\mu-\kappa)\psi(x)=0\to {\rm Dirac}$$
 equation

The equation for $\bar{\psi}$ turns out to be

$$\bar{\psi}(x)(i\gamma^{\mu}\partial_{\mu}+\kappa)=0$$

 \hookrightarrow Derive this

(Just multiply Dirac equation by γ^0 and pull out a γ^0 after taking conjugate). So for QED, there would be extra interaction terms in the Lagrangian

$$\mathcal{L} = \bar{\psi} i \gamma^{\mu} (\partial_{\mu} - i e A_{\mu}) \psi - m \bar{\psi} \psi - \frac{1}{4} F_{\mu\nu} F^{\mu\nu}$$

Relevance of κ in the Dirac equation

$$\begin{split} (i\gamma^{\mu}\partial_{\mu}-\kappa)\psi(x) &= 0 &\qquad \times \left(i\gamma^{\mu}\partial_{\mu}+\kappa\right) \text{ on the left} \\ &\qquad (i\gamma^{\mu}\partial_{\mu}+\kappa)(i\gamma^{\mu}\partial_{\mu}-\kappa)\psi(x) = 0 \\ &\qquad (-\underbrace{\gamma^{\nu}\gamma^{\mu}}_{\text{becomes}}\partial_{\nu}\partial_{\mu}-\kappa^{2})\psi(x) = 0 \\ &\hookrightarrow \frac{1}{2}(\gamma^{\nu}\gamma^{\mu}+\gamma^{\mu}\gamma^{\nu}) = \frac{1}{2}\{\gamma^{\nu},\gamma^{\mu}\} = g^{\mu\nu}\mathbb{1} \\ &\qquad \Rightarrow (-g^{\mu\nu}\partial_{\nu}\partial_{\mu}-\kappa^{2})\psi(x) = 0 \\ &\qquad (\partial^{\mu}\partial_{\mu}-\kappa^{2})\psi(x) = 0 \\ &\hookrightarrow \text{Klein-Gordon equation} \end{split}$$

Therefore, the Dirac equation implies the Klein-Gordon equation. $\mathcal L$