## Notes for Analytical Mechanics

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#### 1 Mathematics

**Theorem 1** (Chain rule). If  $F(q_1, \dots, q_n)$  is a function of n variables, then

$$dF = \sum_{i} \frac{\partial F}{\partial q_i} dq_i.$$

In particular, if  $F(q_1, \dots, q_n, t)$  is a function of coordinates and time, then

$$\dot{F} = \frac{dF}{dt} = \sum_i \frac{\partial F}{\partial q_i} \dot{q}_i + \frac{\partial F}{\partial t}.$$

**Definition 1** (Levi-Cevita symbol). The Levi-Civita symbol  $\epsilon_{i_1 \cdots i_n}$  is the sign of the permutation  $(i_1 \cdots i_n)$ . It has the following properties:

- 1. It is always 1, -1 or 0.
- 2. Swapping any two indexes inverts the Levi-Civita symbol.
- 3.  $\epsilon_{i_1\cdots i_n} = 0$  if  $i_{\alpha} = i_{\beta}$  for any  $\alpha \neq \beta$ .
- 4. A cyclic permutation does not change the Levi-Civita symbol.

#### 2 Newtonian mechanics

**Theorem 2.** The work done by a force equals the change in kinetic energy:

$$W_1 2 = \int_1^2 \mathbf{F} d\mathbf{s} = T_2 - T_1$$

**Definition 2** (Conservative force). A conservative force is the gradient of a scalar potential:

$$\mathbf{F} = -\nabla V(\mathbf{r})$$

**Theorem 3.** A conservative force conserves energy:

$$T_2 - T_1 = W_{12} = V_1 - V_2$$

**Definition 3** (Lagrange multipliers). content

#### 3 Constraints, coordinates and displacements

**Definition 4** (Holonomic constraint). Consider a system of n particles with (actual, physical) coordinates  $\mathbf{r_1}, \cdots \mathbf{r_n}$  and the time t. A *holonomic* constraint on the system can be expressed as a set of equations

$$f(\mathbf{r_1}, \cdots \mathbf{r_n}, t) = 0.$$

Remark 1. We generally limit ourselves to holonomic constraints, and at a microscopic scale, these are generally the only ones involved.

**Definition 5** (Rheonomous and scleronomous constraints). A constraint is called *rheonomous* if it contains the time as an explicit variable. If it does *not*, it is called *scleronomous*.

**Definition 6** (Virtual displacement). content

**Postulate 1** (d'Alembert's principle). In the physical systems that we are interested in, reaction forces do no work under a virtual displacement.

**Definition 7** (Generalized force).

$$Q_{\sigma} = \sum_{i=1}^{3N} F_i^{(a)} \frac{\partial x_i}{\partial q_{\sigma}}$$

### 4 Lagrangian mechanics

**Definition 8** (Lagrangian). The Lagrangian is the kinetic energy T minus the potential energy V:

$$L = T - V$$

Theorem 4 (Euler-Lagrange equation, without constraints).

$$\frac{\partial L}{\partial q} = \frac{d}{dt} \left( \frac{\partial L}{\partial \dot{q}} \right)$$

**Definition 9** (Action). The *action* S of a system evolving from  $t_1$  to  $t_2$  is the integral of the Lagrangian:

$$S = \int_{t_1}^{t_2} L(t) \ dt$$

**Definition 10** (Variation). Variation:

$$\delta S =$$

what is this really?

**Postulate 2** (Hamilton's principle). A system will evolve in such a way that the action has a stationary value at the actual path of motion:

$$\delta S = 0$$

#### 5 Hamiltonian mechanics

**Definition 11** (Canonical momentum from the Lagrangian).

$$p_i = \frac{\partial L}{\partial \dot{q}_i}$$

**Definition 12** (Hamiltonian). The Hamiltonian (*i* running over all variables) is calculated from a Lagrangian  $L(\{q_i, \dot{q}_i\}, t)$  and the canonical coordinates  $\{q_i\}$  and momenta  $\{p_i\}$  as

$$H = \sum_{i} p_i \dot{q}_i - L$$

**Theorem 5** (Hamiltonian equations of motion). Let  $H(\{q_i, p_i\}, t)$  be a Hamiltonian (derived from a Lagrangian L, or from another Hamiltonian by a canonical transformation). Then, for all i,

$$\dot{q}_i = \frac{\partial H}{\partial p_i},$$
 
$$\dot{p}_i = -\frac{\partial H}{\partial q_i}.$$

Furthermore, if H is derived from a Lagrangian L, then

$$\frac{dH}{dt} = \frac{\partial H}{\partial t} = -\frac{\partial L}{\partial t}.$$

**Definition 13** (Canonical transformation). A transformation  $(p_i, q_i) \to (P_i, Q_i)$  is *canonical* if it preserves the form of the Hamiltonian equations of motion, i.e.

$$\dot{Q} = \frac{\partial H}{\partial P}$$

$$\dot{P} = -\frac{\partial H}{\partial Q}$$

**Definition 14** (Poisson bracket). The Poisson bracket of two functions f and g, with respect to coordinates  $\{q_i, p_i\}$ , is defined as

$$\{f,g\}_{\{q_i,p_i\}} = \sum_i \left(\frac{\partial f}{\partial q_i}\frac{\partial g}{\partial p_i} - \frac{\partial f}{\partial p_i}\frac{\partial g}{\partial q_i}\right)$$

**Theorem 6.** A transformation is canonical iff it preserves Poisson brackets:

$$\{f,g\}_{\{q_i,p_i\}} = \{f,g\}_{\{Q_i,P_i\}} \quad \textit{for all functions } f,g$$

**Theorem 7** (Hamilton-Jacobi equation). Let H be a Hamiltonian function and S a generating function such that  $H + \frac{\partial S}{\partial t} = 0$ . Then

$$H\bigg(q_1,\cdots,q_n,\frac{\partial S}{\partial q_1},\cdots,\frac{\partial S}{\partial q_1},t\bigg)+\frac{\partial S}{\partial t}=0$$

Not much of a theorem, really just by construction. But useful.

### 6 Central force problem

Theorem 8 (Virial theorem).

$$\bar{T} = -\frac{1}{2} \overline{\sum_{i} \vec{F}_{i} \cdot \vec{r}_{i}}$$

# 7 Rotating coordinate systems

**Theorem 9** (Effective force). If a particle is applied a force of  ${\bf F}$  in a reference system, an observer in a coordinate system rotating with angular velocity  ${\boldsymbol \omega}$  will see it as affected by a force

$$\mathbf{F}_{eff} = \mathbf{F} - 2m(\boldsymbol{\omega} \times \mathbf{v_r}) - m\boldsymbol{\omega} \times (\boldsymbol{\omega} \times \mathbf{r})$$

where  $\mathbf{v_r}$  is the velocity of the particle in the rotating coordinate system.