

$$1. \frac{\partial u}{\partial t} = -c \frac{\partial u}{\partial x} + V \frac{\partial^2 u}{\partial x^2}$$

For forth-order centered FD

$$\frac{\partial u}{\partial t} = \frac{-u(x+2\Delta x) + 8u(x+\Delta x) - 8u(x-\Delta x) + u(x-2\Delta x)}{12\Delta x}$$

$$\frac{\partial^2 u}{\partial t^2} = \frac{-u(x+2\Delta x) + 16u(x+\Delta x) - 30u(x) + 16u(x-\Delta x) - u(x-2\Delta x)}{12\Delta x^2}$$

Then semi-discretized equation is:

$$\begin{aligned} \frac{\partial u}{\partial t} &= \frac{-c}{12\Delta x} [-u(x+2\Delta x) + 8u(x+\Delta x) - 8u(x-\Delta x) + u(x-2\Delta x)] \\ &\quad + \frac{V}{12\Delta x^2} [-u(x+2\Delta x) + 16u(x+\Delta x) - 30u(x) + 16u(x-\Delta x) - u(x-2\Delta x)] \end{aligned}$$

Write $u = \hat{u} e^{ikx}$

$$\text{Then } u(x+\Delta x) = \hat{u} e^{ik(x+\Delta x)} = \hat{u} e^{ikx} \cdot e^{ik\Delta x}$$

$$\text{Similarly: } u(x+2\Delta x) = \hat{u} e^{ikx} \cdot e^{i2k\Delta x}$$

$$u(x-\Delta x) = \hat{u} e^{ikx} \cdot e^{-ik\Delta x}$$

$$u(x-2\Delta x) = \hat{u} e^{ikx} \cdot e^{-i2k\Delta x}$$

Plug them into equation above.

$$\begin{aligned} \frac{\partial \hat{u}}{\partial t} e^{ikx} &= -\frac{c}{12\Delta x} [\hat{u} e^{ikx} \cdot e^{i2k\Delta x} - 8\hat{u} e^{ikx} \cdot e^{ik\Delta x} - 8\hat{u} e^{ikx} \cdot e^{-ik\Delta x} \\ &\quad + \hat{u} e^{ikx} \cdot e^{-i2k\Delta x}] \end{aligned}$$

$$\begin{aligned} &\quad + \frac{V}{12\Delta x^2} [\hat{u} e^{ikx} \cdot e^{i2k\Delta x} + 16\hat{u} e^{ikx} \cdot e^{ik\Delta x} - 30\hat{u} e^{ikx} \\ &\quad + 16\hat{u} e^{ikx} \cdot e^{-ik\Delta x} - \hat{u} e^{ikx} \cdot e^{-i2k\Delta x}] \end{aligned}$$

$$\Rightarrow \frac{\partial \hat{u}}{\partial t} = \hat{u} \left[\frac{-c}{12\Delta x} [e^{i2k\Delta x} + e^{-i2k\Delta x}] + 8e^{ik\Delta x} - 8e^{-ik\Delta x} \right] \\ - 2i \sin 2k\Delta x \frac{8(2i \sin k\Delta x)}{32 \cos k\Delta x}$$

$$\Rightarrow \hat{u} \left[\frac{V}{12\Delta x^2} [e^{i2k\Delta x} - e^{-i2k\Delta x}] + 16[e^{ik\Delta x} - e^{-ik\Delta x}] - 30 \right] \\ - 2 \cos 2k\Delta x \frac{32 \cos k\Delta x}{32 \cos k\Delta x}$$

$$\Rightarrow \frac{d\hat{A}}{dt} = \hat{U} \left[i \left(\frac{c}{6\pi} \sin 2k_0 x - i \frac{4}{3} \frac{c}{\Delta x} \sin k_0 x \right) \right. \\ \left. + \hat{U} L - \frac{1}{6} \frac{V}{\Delta x^2} \cos 2k_0 x + \frac{8}{3} \frac{V}{\Delta x^2} \cos k_0 x - \frac{5}{2} \frac{V}{\Delta x^2} \right] \\ = \lambda \hat{U}$$

Thus $\lambda = i \frac{c}{\Delta x} \left[L \left(\frac{1}{6} \sin 2k_0 x - \frac{4}{3} \sin k_0 x \right) \right] + \frac{V}{\Delta x^2} L - \frac{1}{6} \cos 2k_0 x + \frac{8}{3} \cos k_0 x - \frac{5}{2}$

$$\Rightarrow \lambda_{st} = i \frac{Cst}{\Delta x} \left[L \left(\frac{1}{6} \sin 2k_0 x - \frac{4}{3} \sin k_0 x \right) \right] + \frac{Vst}{\Delta x^2} L - \frac{1}{6} \cos 2k_0 x + \frac{8}{3} \cos k_0 x - \frac{5}{2}$$

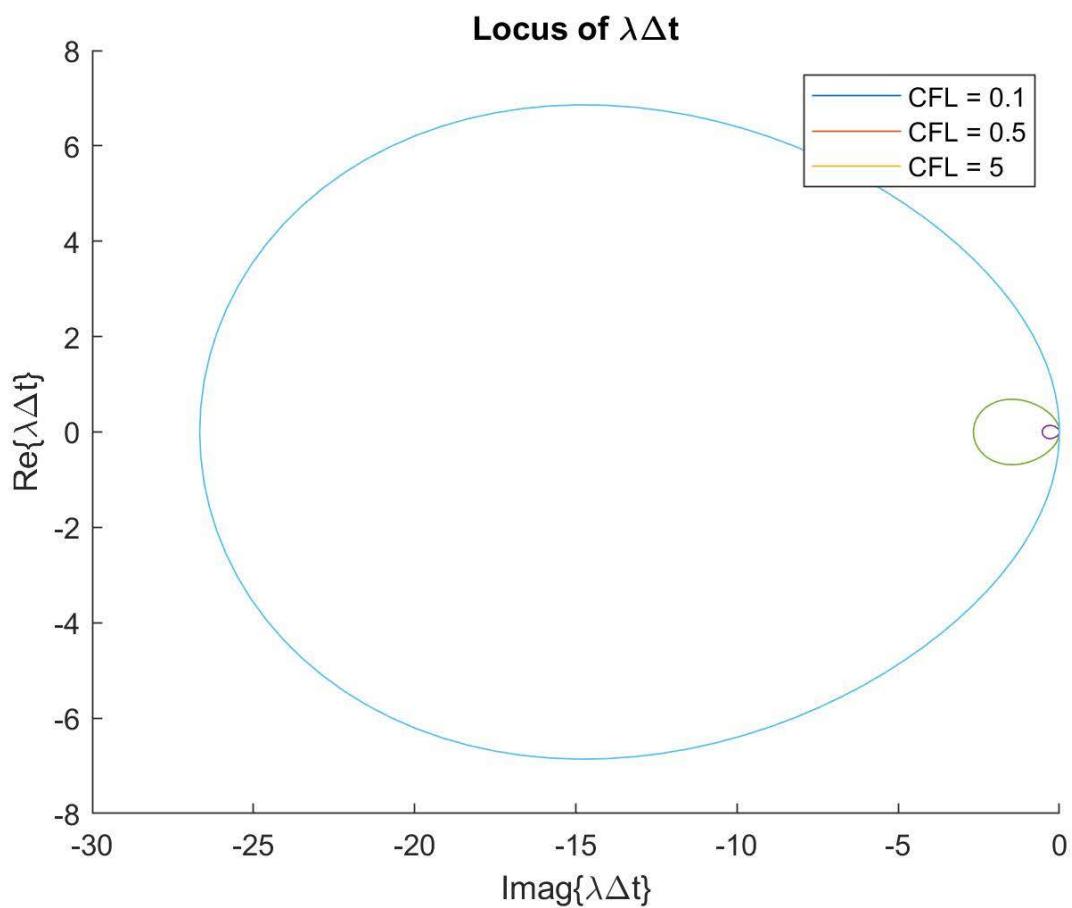
Since $\frac{Vst}{\Delta x^2} = \frac{Cst}{\Delta x} \frac{V}{\Delta x C} = \frac{CFL}{Re}$

$$\Rightarrow \lambda_{st} = i CFL \underbrace{\left[L \left(\frac{1}{6} \sin 2k_0 x - \frac{4}{3} \sin k_0 x \right) \right]}_{\lambda_{st}} + \underbrace{\frac{CFL}{Re} L \left[-\frac{1}{6} \cos 2k_0 x + \frac{8}{3} \cos k_0 x - \frac{5}{2} \right]}_{\lambda_{Rst}}$$

where $k_0 x = \frac{2\pi n}{L} \frac{L}{\Delta x} = \frac{2\pi n}{\Delta x}$. $n = -\frac{N}{2}, -1, \frac{N}{2}-1$

$$\Rightarrow k_0 x = -\pi, \dots, \pi - \frac{2\pi}{\Delta x}$$

Problem1. Plot



$$\frac{dh}{dt} = H(u^n) + Lu^n$$

Problem 2

2 (c)

Exact Taylor Expansion of u^{n+1} with respect to u^n is

$$u^{n+1} = u^n + \Delta t \left. \frac{du}{dt} \right|_{u=u^n} + \frac{\Delta t^2}{2} \left. \frac{d^2u}{dt^2} \right|_{u=u^n} + \dots$$

$$\text{By } \left. \frac{du}{dt} \right|_{u=u^n} = H(u^n) + Lu^n$$

$$\left. \frac{d^2u}{dt^2} \right|_{u=u^n} = \frac{d}{dt} \left(\left. \frac{du}{dt} \right|_{u=u^n} \right) = \frac{d}{dt} (Lu^n + H(u^n))$$

$$= L \left. \frac{du}{dt} \right|_{u=u^n} + \left. \frac{dH(u)}{du} \right|_{u=u^n} \left. \frac{du}{dt} \right|_{u=u^n}$$

$$= L \left. \frac{du}{dt} \right|_{u=u^n} + \left. \frac{dH(u)}{du} \right|_{u=u^n} (Lu^n + H(u^n))$$

$$= L(Lu^n + H(u^n)) + \left. \frac{dH(u)}{du} \right|_{u=u^n} (Lu^n + H(u^n))$$

$$= \left[L + \left. \frac{dH(u)}{du} \right|_{u=u^n} \right] [Lu^n + H(u^n)]$$

$$\Rightarrow u^{n+1} = u^n + \Delta t [Lu^n + H(u^n)] + \frac{\Delta t^2}{2} \left[L^2 u^n + LH(u^n) + \left. \frac{dH(u)}{du} \right|_{u=u^n} Lu^n + \left. \frac{dH(u)}{du} \right|_{u=u^n} H(u^n) \right]$$

For two step RK-2 we have :

$$u^* = u^n + \Delta t + L u^n + \underbrace{\Delta t \beta_1 L u^*}_{\rightarrow \Delta t f, H(u^n)}$$

$$\Rightarrow (1 - \Delta t \beta_1 L) (u^* = u^n + \Delta t + L u^n + \Delta t f, H(u^n))$$

$$\Rightarrow u^* = (1 - \Delta t \beta_1 L)^{-1} (u^n + \Delta t + L u^n + \Delta t f, H(u^n))$$

Expand $(1 - \Delta t \beta_1 L)^{-1}$ by Taylor series and plug into Eq :

$$u^* = (1 + \Delta t \beta_1 L + \Delta t^2 \beta_1^2 L^2 + \dots) (u^n + \Delta t + L u^n + \Delta t f, H(u^n))$$

$$\Rightarrow \tilde{U}^* = U^n + \Delta t \mathcal{J}_2 L U^n + \Delta t \mathcal{J}_1 H(U^n)$$

$$\quad \quad \quad + \Delta t \beta_1 L U^n + \Delta t^2 \beta_1 \mathcal{J}_2 L^2 U^n + \Delta t^2 \beta_1 \mathcal{J}_1 H(U^n)$$

$$\quad \quad \quad + \Delta t^2 \beta_1^2 L^2 U^n + \Delta t^3 \beta_1^2 L^3 U^n + \Delta t^3 \beta_1^2 \mathcal{J}_1 L^2 H(U^n) + \dots$$

Step 2: $U^{n+1} = U^* + \Delta t \mathcal{J}_2 L \tilde{U}^* + \Delta t \beta_2 L U^{n+1} + \Delta t \mathcal{J}_2 H(U^*) - \Delta t^3 H(U^n)$

Since $H(U^*) = H(U^n) + \frac{dH}{dU}(U^* - U^n) + \dots$

$$\rightarrow \tilde{U}^{n+1} = U^* + \underbrace{\Delta t \mathcal{J}_2 L \tilde{U}^*}_{\Delta t \mathcal{J}_2 H(U^n)} + \underbrace{\Delta t \beta_2 L U^{n+1}}_{\Delta t \mathcal{J}_2 \frac{dH}{dU}(U^* - U^n)}$$

$$\quad \quad \quad + \Delta t^3 H(U^n) \quad \quad \quad (\Delta t \mathcal{J}_2 L + \Delta t^2 \beta_1 \mathcal{J}_2 L^2 + \Delta t^3 \beta_1^2 \mathcal{J}_2 L^3) \tilde{U}^*$$

$$\rightarrow U^{n+1} = (1 - \Delta t \beta_2 L)^{-1} U^* + (1 - \Delta t \beta_2 L)^{-1} \Delta t \mathcal{J}_2 L \tilde{U}^*$$

$$\quad \quad \quad + (1 - \Delta t \beta_2 L)^{-1} [\Delta t \mathcal{J}_2 H(U^n) + \Delta t \mathcal{J}_2 \frac{dH}{dU}(U^* - U^n)]$$

$$\quad \quad \quad + (1 - \Delta t \beta_2 L)^{-1} \Delta t^3 H(U^n)$$

Again expand $(1 - \Delta t \beta_2 L)^{-1}$ by Taylor series and plug in \tilde{U}^* :

$$U^{n+1} = (1 + \Delta t \beta_2 L + \Delta t^2 \beta_2^2 L^2 + \Delta t \mathcal{J}_2 L + \Delta t^2 \beta_2 \mathcal{J}_2 L^2 + \Delta t^3 \beta_2^2 \mathcal{J}_2 L^3) [U^n + \Delta t \mathcal{J}_1 L U^n]$$

$$\quad \quad \quad + \Delta t \mathcal{J}_1 H(U^n) + \Delta t \beta_1 L U^n + \Delta t^2 \beta_1 \mathcal{J}_2 L^2 U^n + \Delta t^2 \beta_1 \mathcal{J}_1 H(U^n) + \Delta t^2 \beta_1^2 L^2 U^n$$

$$\quad \quad \quad + \Delta t^3 \beta_1^2 L^3 U^n + \Delta t^3 \beta_1^2 \mathcal{J}_1 L^2 H(U^n)]$$

$$\quad \quad \quad + (1 + \Delta t \beta_2 L + \Delta t^2 \beta_2^2 L^2) [\Delta t \mathcal{J}_2 H(U^n) + \Delta t \mathcal{J}_2 \frac{dH}{dU} (\Delta t \mathcal{J}_1 L U^n + \Delta t \mathcal{J}_1 H(U^n))]$$

$$\quad \quad \quad + \Delta t \beta_1 L U^n + \Delta t^2 \beta_1 \mathcal{J}_2 L^2 U^n + \Delta t^2 \beta_1 \mathcal{J}_1 H(U^n) + \Delta t^2 \beta_1^2 L^2 U^n + \Delta t^3 \beta_1^2 L^3 U^n$$

$$\quad \quad \quad + \Delta t^3 \beta_1^2 \mathcal{J}_1 L^2 H(U^n)]$$

$$\quad \quad \quad + (1 + \Delta t \beta_2 L + \Delta t^2 \beta_2^2 L^2) (\Delta t^3 H(U^n))$$

Ignore all of these terms, and ignore terms not in exact Taylor expansion:

$$U^n, H^{un}, L^{un}, L^2 U^n, L H^{un}, \frac{\partial H}{\partial u} L^{un}, \frac{\partial H}{\partial u} H^{un} \quad \text{page 5}$$

U^{un}

$$\begin{aligned}
 &= \left[U^n + \Delta t \dot{f}_1 L^{un} + \Delta t \dot{f}_1 H^{un} + \Delta t \beta_1 L^{un} + \Delta t^2 \beta_1 \dot{f}_1 L^2 U^n \right. \\
 &\quad \left. + \Delta t^2 \beta_1 \dot{f}_1 L H^{un} + \Delta t^2 \beta_1^2 L^2 U^n \right] \\
 &\quad + \left[\Delta t \beta_2 L^{un} + \Delta t^2 \beta_2 \dot{f}_1 L^2 U^n + \Delta t^2 \beta_2 \dot{f}_1 L H^{un} + \Delta t^2 \beta_2 \beta_1 L^2 U^n \right] \\
 &\quad + \left[\Delta t \dot{f}_2 L^{un} + \Delta t^2 \dot{f}_2 \dot{f}_1 L^2 U^n + \Delta t^2 \dot{f}_2 \dot{f}_1 L H^{un} + \Delta t^2 \dot{f}_2 \beta_1 L^2 U^n \right] \\
 &\quad + \left[\Delta t^2 \beta_2^2 L^2 U^n \right] \\
 &\quad + \left[\Delta t^2 \beta_2 \dot{f}_2 L^2 U^n \right] \\
 &\quad + \left[\Delta t \dot{f}_1 H^{un} + \Delta t^2 \dot{f}_1 \dot{f}_1 \frac{\partial H}{\partial u} L^{un} + \Delta t^2 \dot{f}_1 \dot{f}_1 \frac{\partial H}{\partial u} H^{un} + \Delta t^2 \dot{f}_1 \beta_1 \frac{\partial H}{\partial u} L^{un} \right] \\
 &\quad + \left[\Delta t^2 \beta_2 \dot{f}_2 L H^{un} \right] \\
 &\quad + \left[\Delta t \beta_1 H^{un} + \Delta t^2 \beta_1 \beta_1 L H^{un} \right]
 \end{aligned}$$

Compare with Exact Taylor Expansion

① U^n term: $1 = 1$

② H^{un} term: $\Delta t \dot{f}_1 + \Delta t \dot{f}_2 + \Delta t \dot{f}_2 = \Delta t$

③ L^{un} term: $\Delta t \dot{f}_1 + \Delta t \beta_1 + \Delta t \beta_2 + \Delta t \dot{f}_2 = \Delta t$

④ $L^2 U^n$ term: $\Delta t^2 \beta_1 \dot{f}_1 + \Delta t^2 \beta_1^2 + \Delta t^2 \beta_2 \dot{f}_1 + \Delta t^2 \beta_2 \beta_1 + \Delta t^2 \dot{f}_2 \dot{f}_1 + \Delta t^2 \dot{f}_2 \beta_1 = \frac{\Delta t^2}{2}$

⑤ $L H^{un}$ term: $\Delta t^2 \beta_1 \dot{f}_1 - \underbrace{\Delta t^2 \beta_2 \dot{f}_1}_{\cancel{\Delta t^2 \dot{f}_2 \dot{f}_1}} + \cancel{\Delta t^2 \dot{f}_2 \dot{f}_1} + \Delta t^2 \beta_2 \dot{f}_2 + \Delta t^2 \beta_2 \beta_1 = \frac{\Delta t^2}{2}$

⑥ $\frac{\partial H}{\partial u} L^{un}$ term: $\Delta t^2 \dot{f}_2 \dot{f}_1 - \Delta t^2 \dot{f}_2 \beta_1 = \frac{\Delta t^2}{2}$

⑦: $\frac{\partial H}{\partial u} H^{un}$ term: $\Delta t^2 \dot{f}_2 \dot{f}_1 = \frac{\Delta t^2}{2}$

Thus in total we have 6 equations with 7 unknowns

$$\text{Hun} \quad ① \quad \begin{cases} \gamma_1 + \gamma_2 + \beta_1 = 1 \end{cases}$$

$$\text{Lvn} \quad ② \quad \alpha_1 + \alpha_2 + \beta_1 + \beta_2 = 1$$

$$\text{L}^2\text{vn} \quad ③ \quad \beta_1 \alpha_1 + \beta_1^2 + \alpha_1 \beta_2 + \beta_1 \beta_2 + \alpha_1 \alpha_2 + \alpha_2 \beta_1 + \beta_2 \alpha_2 + \beta_2^2 = \frac{1}{2}$$

$$\text{LHvn} \quad ④ \quad \beta_1 \gamma_1 + \beta_2 \gamma_1 + \alpha_2 \gamma_1 + \beta_2 \gamma_2 + \beta_2 \beta_1 = \frac{1}{2}$$

$$\text{Lvn} \quad ⑤ \quad \alpha_1 \gamma_2 + \beta_1 \gamma_2 = \frac{1}{2}$$

$$\frac{\partial H}{\partial n} \text{ Hun} \quad ⑥ \quad \gamma_1 \gamma_2 = \frac{1}{2}$$

Unknowns: $\alpha_1, \alpha_2, \beta_1, \beta_2, \gamma_1, \gamma_2, \beta_1$

Problem 2

(b) (ii) Since we have 7 undetermined variables, but only have 6 equations. This problem is undetermined.

two step RK-2 gives:

$$\begin{aligned}
 U^{n+1} = & U^n + \Delta t \left[(\alpha_1 + \alpha_2 + \beta_1 + \beta_2) L U^n + (\gamma_1 + \gamma_2 + \gamma_3) H(U^n) \right] \\
 & + \frac{\Delta t^2}{2} \left[(\alpha_1 \alpha_2 + \alpha_1 \beta_1 + \alpha_2 \beta_2 + \beta_1 \beta_2 + \alpha_1 \gamma_1 + \alpha_2 \gamma_2 + \beta_1 \gamma_1 + \beta_2 \gamma_2 + \beta_1^2 + \beta_2^2 + \beta_1 \beta_2) L^2 U^n \right. \\
 & \quad \left. + (\beta_1 \gamma_1 + \beta_2 \gamma_1 + \beta_1 \gamma_2 + \beta_2 \gamma_2 + \gamma_2 \gamma_3) L H(U^n) \right. \\
 & \quad \left. + (\gamma_1 \gamma_2 + \beta_1 \gamma_2) \frac{\partial H}{\partial u} L U^n + \gamma_1 \gamma_2 \frac{\partial H}{\partial u} H(U^n) \right]
 \end{aligned}$$

To see what will happen when $\Delta t \rightarrow \infty$ we can only look at viscous term, since it solely contributed to the rear part of λ , which is treated with viscous term.

Problem 2

(b) Then Apply this two-step RK-2 Scheme to Model

$$\text{Equation } \frac{du}{dt} = \lambda u$$

$$\text{Thus } u^* = u^n + \Delta t \beta_1 \lambda u^n + \Delta t \beta_1 \lambda u^*$$

$$\Rightarrow u^* = (1 - \Delta t \beta_1 \lambda)^{-1} (1 + \beta_1 \lambda \Delta t) u^n$$

$$\text{Plug this into } u^{n+1} = u^* + \Delta t \beta_2 \lambda u^* + \Delta t \beta_2 \lambda u^{n+1}$$

$$\Rightarrow u^{n+1} = (1 - \beta_2 \lambda \Delta t)^{-1} (1 + \beta_2 \lambda \Delta t) (1 - \Delta t \beta_1 \lambda)^{-1} (1 + \beta_1 \lambda \Delta t) u^n$$

The Amplification factor gives:

$$\sigma = \frac{u^{n+1}}{u^n} = \frac{(1 + \beta_2 \lambda \Delta t)(1 + \beta_1 \lambda \Delta t)}{(1 - \beta_2 \lambda \Delta t)(1 - \beta_1 \lambda \Delta t)}$$

For $\operatorname{Re}(\lambda \Delta t) = \lambda \Delta t \rightarrow \infty$. If $\sigma \rightarrow 0$, it need denominator

of σ should be much larger than numerator, or numerator

goes to 0; Since the former one is not valid, to let the order of magnitude of denominator much larger than numerator,

it requires either $\beta_2 \lambda \Delta t$ or $\beta_1 \lambda \Delta t$, or both of them vanish,

This imply $\beta_2 = 0$ or $\beta_1 = 0$ or Both $\beta_1 = \beta_2 = 0$

Unknowns $\alpha_1, \beta_1, \gamma_1, \gamma_2, \beta_1,$

Problem 2. (b) (ii)

Set $\alpha_2 = 0$, and $\beta_2 = \beta_1$. The equation becomes:

$$\textcircled{1} \quad \gamma_1 + \gamma_2 + \beta_1 = 1$$

$$\textcircled{2} \quad \alpha_1 + 2\beta_1 = 1$$

$$\textcircled{3} \quad \beta_1 \alpha_1 + \beta_1^2 + \alpha_1 \beta_1 + \beta_1^2 = \frac{1}{2}$$

$$\Rightarrow \beta_1(3\beta_1 + 2\alpha_1) = \frac{1}{2}$$

$$\textcircled{4} \quad \beta_1 \gamma_1 + \beta_1 \gamma_1 + \beta_1 \gamma_2 + \beta_1 \beta_1 = \frac{1}{2}$$

$$\Rightarrow \beta_1(2\gamma_1 + \gamma_2 + \beta_1) = \frac{1}{2}$$

$$\textcircled{5} \quad \alpha_1 \gamma_2 + \beta_1 \gamma_2 = \frac{1}{2}$$

$$\textcircled{6} \quad \gamma_1 \gamma_2 = \frac{1}{2}$$

With the aid of Matlab, we find

$$\begin{bmatrix} \alpha_1 \\ \beta_1 \\ \gamma_1 \\ \gamma_2 \\ \beta_1 \end{bmatrix} = \begin{bmatrix} \sqrt{2}-1 \\ 1-\frac{\sqrt{2}}{2} \\ \frac{\sqrt{2}}{2} \\ \frac{\sqrt{2}}{2} \\ 1-\sqrt{2} \end{bmatrix}$$

$$\text{And } \alpha_2 = 0, \quad \beta_2 = \beta_1 = 1 - \frac{\sqrt{2}}{2}$$

Problem 3

3. (a) Apply Modern Eqn $\frac{\partial u}{\partial t} = \lambda u$ to Analyse the stability of two step RK-O with any H term:

$$\begin{aligned} U^* &= U^n + \gamma_1 H u^n \Delta t \\ &= U^n + \gamma_1 \lambda U^n \Delta t \end{aligned}$$

$$\begin{aligned} U^{n+1} &= U^* + [\gamma_2 H(U^*) + \gamma_3 H(U^n)] \Delta t \\ &= U^* + [\gamma_2 \lambda U^* + \gamma_3 \lambda U^n] \Delta t \\ &= U^n + \gamma_1 \lambda \Delta t U^n + \gamma_2 \lambda \Delta t (U^n + \gamma_1 \lambda \Delta t U^n) \\ &\quad + \gamma_3 \lambda \Delta t U^n \end{aligned}$$

$$\begin{aligned} \rightarrow |\zeta| &= \frac{|U^{n+1}|}{|U^n|} = 1 + \gamma_1 \lambda \Delta t + \gamma_2 \lambda \Delta t + \gamma_1 \gamma_2 (\lambda \Delta t)^2 + \gamma_3 \lambda \Delta t \\ &= 1 + (\gamma_1 + \gamma_2 + \gamma_3) \lambda \Delta t + \gamma_1 \gamma_2 (\lambda \Delta t)^2 \end{aligned}$$

From 2(b) (ii), we have

$$\gamma_1 + \gamma_2 + \gamma_3 = \frac{\sqrt{2}}{2} + \frac{\sqrt{2}}{2} + 1 - \sqrt{2} = 1$$

$$\gamma_1 \gamma_2 = \frac{\sqrt{2}}{2} \cdot \frac{\sqrt{2}}{2} = \frac{1}{2}$$

$$\rightarrow |\zeta| = 1 + \lambda \Delta t + \frac{1}{2}(\lambda \Delta t)^2$$

for stability At Requre $| \zeta | \leq 1$

Since H only contribute to imaginary part, we can write

$$\lambda \Delta t = i \lambda_{\text{R}} \Delta t, \quad |\zeta| = 1 + i \lambda_{\text{R}} \Delta t + \frac{1}{2} (i \lambda_{\text{R}} \Delta t)^2$$

$$\rightarrow |\zeta| = \left| \left[1 - \frac{1}{2} (\lambda_{\text{R}} \Delta t)^2 \right] + i \lambda_{\text{R}} \Delta t \right| \leq 1$$

Problem 3

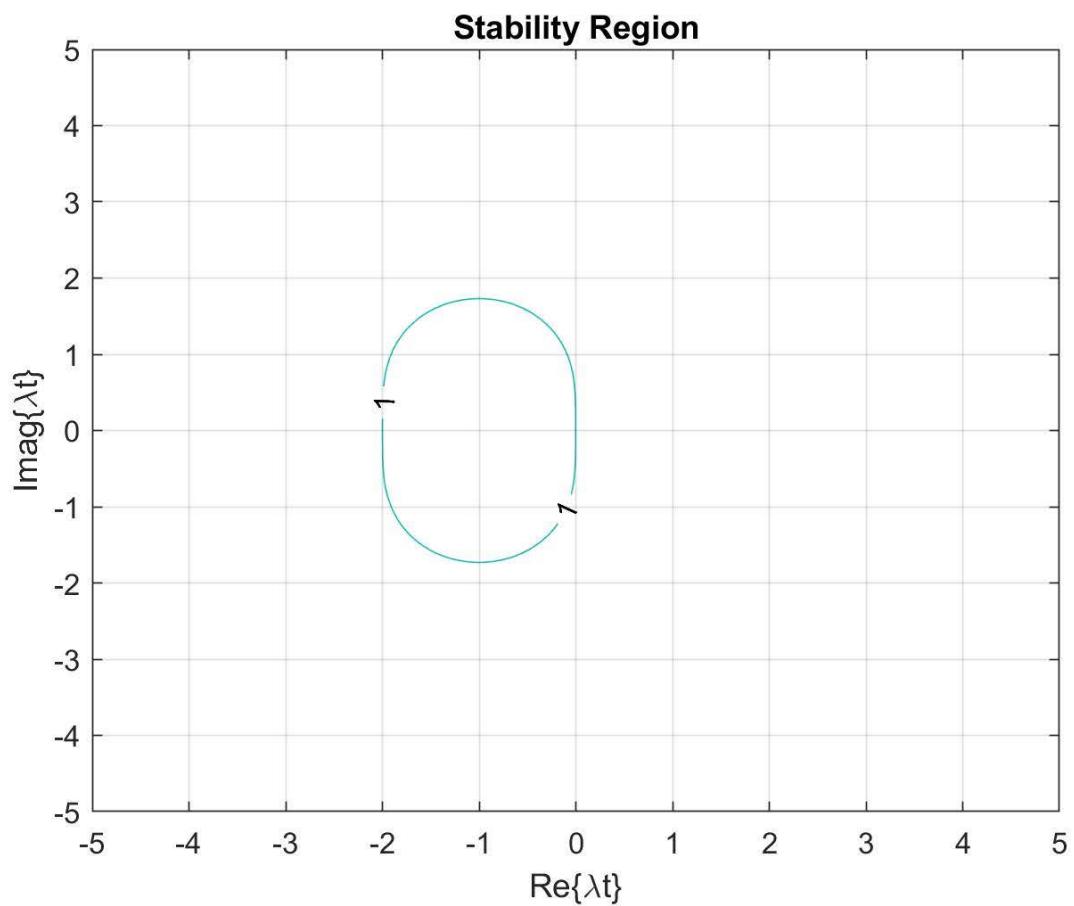
3. (a) (Ansatz) $|z| = \sqrt{[1 - \frac{1}{2}(\lambda_i \Delta t)^2]^2 + (\lambda_i \Delta t)^2} \leq 1$

$$1 + \frac{1}{4}(\lambda_i \Delta t)^4 - (\lambda_i \Delta t)^2 + (\lambda_i \Delta t)^2 \leq 1$$

$$\rightarrow \frac{1}{4}(\lambda_i \Delta t)^4 \leq 0$$

Then the step $Rk=0$ is unconditionally unstable for any terms.

\Rightarrow No stable Region



Problem 3

3. (b) Applying stability analysis to model (9) as what has been done in 2(b), we find

$$\sigma = \frac{v^{n+1}}{v^n} = \frac{(1 + \lambda_2 \Delta t)(1 - \lambda_1 \Delta t)}{(1 - \beta_2 \lambda \Delta t)(1 - \beta_1 \lambda \Delta t)}$$

$$= \frac{1 + (\sqrt{2} - 1)\lambda \Delta t}{[1 - (1 - \frac{\sqrt{2}}{2})\lambda \Delta t]^2}$$

Since L term only contribute to Real part of $\lambda \Delta t$
we can write $\lambda \Delta t = \lambda_R \Delta t$

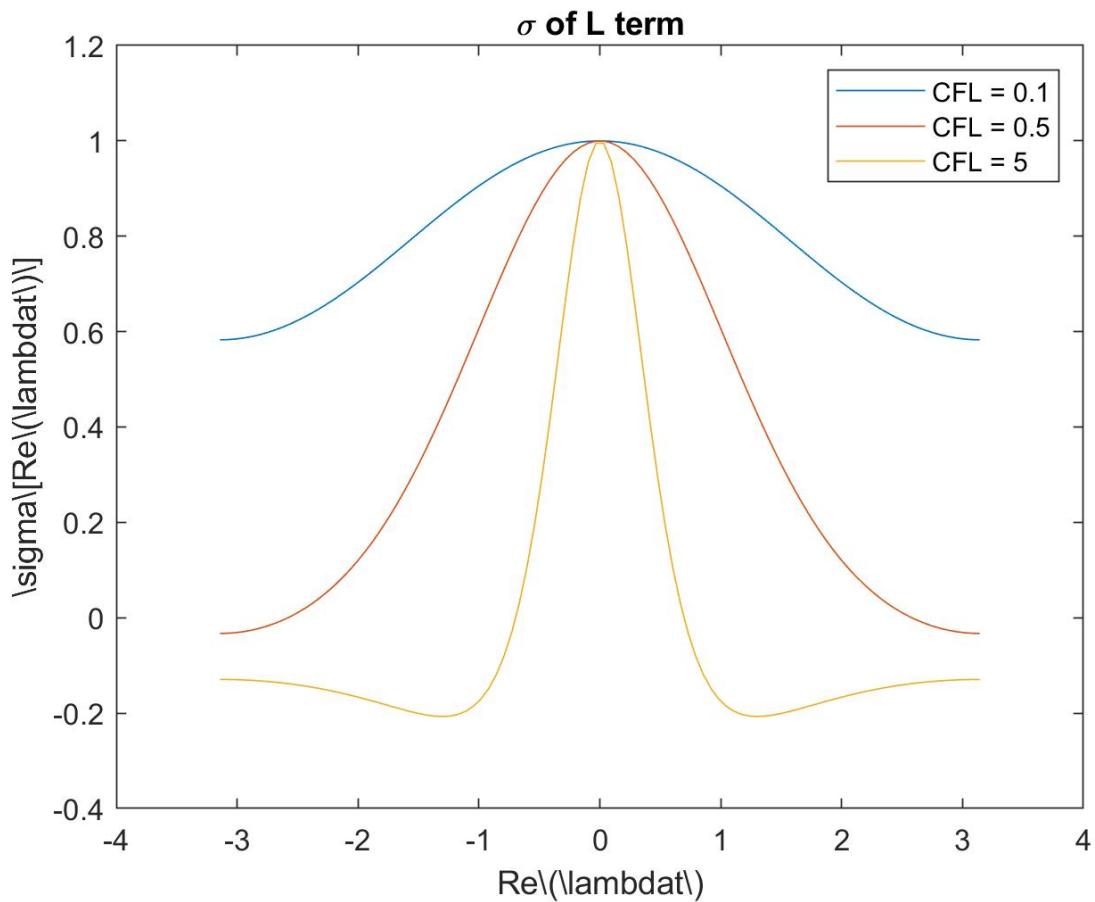
$$\text{Then } \sigma = \frac{1 + (\sqrt{2} - 1)\lambda_R \Delta t}{[1 - (1 - \frac{\sqrt{2}}{2})\lambda_R \Delta t]^2}$$

From problem 1, we have

$$\lambda_R \Delta t = \frac{CFL}{Re} \left[1 - \frac{1}{6} \cos 2k \omega x + \frac{3}{2} \cos k \omega x - \frac{5}{2} \right]$$

Then σ can be found for this discretization

Problem3 (b) Plot



Problem 3

3.(c) For stability analysis we can apply Rk=0 method to model

$$\text{Eqn } \frac{du}{dt} = \lambda u = i\lambda I u + \lambda_R u = H u \rightarrow Lu$$

$$\text{Then } U^* = u + [d_1 L u^n + \beta_1 L u^{n+1}] \Delta t + [f_1 H u] \Delta t$$

$$= u + d_1 \lambda_R \Delta t u^n + \beta_1 \lambda_R \Delta t U^* + i f_1 \lambda_I \Delta t u^n$$

$$\rightarrow (1 - \beta_1 \lambda_R \Delta t) U^* = (1 + d_1 \lambda_R \Delta t + i f_1 \lambda_I \Delta t) u^n$$

$$\rightarrow U^* = (1 - \beta_1 \lambda_R \Delta t)^{-1} (1 + d_1 \lambda_R \Delta t + i f_1 \lambda_I \Delta t) u^n$$

Plug this into next step.

$$U^{n+1} = U^* + (d_2 L U^* + \beta_2 L U^{n+1}) \Delta t + [f_2 H U^* + \beta_3 H u^n] \Delta t$$

$$= U^* + d_2 \lambda_R \Delta t U^* + \beta_2 \lambda_R \Delta t U^{n+1}$$

$$+ i f_2 \lambda_I \Delta t U^* + i \beta_3 \lambda_I \Delta t u^n$$

$$\Rightarrow (1 - \beta_2 \lambda_R \Delta t) U^{n+1} = (1 + d_2 \lambda_R \Delta t + i f_2 \lambda_I \Delta t) U^* + i \beta_3 \lambda_I \Delta t u^n$$

$$\Rightarrow U^{n+1} = \frac{1 + d_2 \lambda_R \Delta t + i f_2 \lambda_I \Delta t}{1 - \beta_2 \lambda_R \Delta t} \frac{1 + f_1 \lambda_R \Delta t + i f_1 \lambda_I \Delta t}{1 - \beta_1 \lambda_R \Delta t} u^n$$

$$\rightarrow \frac{i \beta_3 \lambda_I \Delta t}{1 - \beta_2 \lambda_R \Delta t} u^n$$

$$\Rightarrow \sigma = \frac{U^{n+1}}{u^n} = \frac{(1 + d_2 \lambda_R \Delta t + i f_2 \lambda_I \Delta t)(1 + f_1 \lambda_R \Delta t + i f_1 \lambda_I \Delta t)}{(1 - \beta_2 \lambda_R \Delta t)(1 - \beta_1 \lambda_R \Delta t)}$$

$$\rightarrow \frac{i \beta_3 \lambda_I \Delta t}{1 - \beta_2 \lambda_R \Delta t}$$

Problem 3

3.(c) (contd.) Apply the condition $\beta_2 = \alpha$, $\beta_1 = \beta$.

$$\sigma = \frac{(1 + i\gamma_1 \lambda_{I\text{at}})(1 + \gamma_1 \lambda_{R\text{at}} + i\gamma_1 \lambda_{I\text{at}}) - i\gamma_1 \lambda_{I\text{at}}(1 - \beta_1 \lambda_{R\text{at}})}{(1 - \beta_1 \lambda_{R\text{at}})^2}$$

$$= \frac{1}{(1 - \beta_1 \lambda_{R\text{at}})^2} [1 + \gamma_1 \lambda_{R\text{at}} + i\gamma_1 \lambda_{I\text{at}} + i\gamma_2 \gamma_1 \lambda_{I\text{at}} \lambda_{R\text{at}} \\ + (-\gamma_2 \gamma_1 \lambda_{I\text{at}} \lambda_{R\text{at}}) + i\gamma_1 \lambda_{I\text{at}} - i\gamma_1 \beta_1 \lambda_{I\text{at}} \lambda_{R\text{at}}]$$

$$= \frac{1 + \gamma_1 \lambda_{R\text{at}} - \gamma_1 \gamma_2 (\lambda_{I\text{at}})^2}{(1 - \beta_1 \lambda_{R\text{at}})^2} + i \frac{\lambda_{I\text{at}}(\gamma_1 + \gamma_1 \gamma_2 \lambda_{R\text{at}} + \gamma_1 - \gamma_1 \beta_1 \lambda_{R\text{at}})}{(1 - \beta_1 \lambda_{R\text{at}})^2}$$

When $\gamma_1 = \gamma_2 = \frac{\sqrt{2}}{2}$, $\beta_1 = 1 - \frac{\sqrt{2}}{2}$, $\gamma_1 = 1 - \sqrt{2}$, $\gamma_2 = \sqrt{2}$

$$\lambda_I = \text{CFL} \left(\frac{1}{6} \sin 2k_0 x - \frac{4}{3} \sin k_0 x \right)$$

$$\lambda_R = \frac{\text{CFL}}{\text{Re}} \left(-\frac{1}{6} \cos 2k_0 x + \frac{8}{3} \cos k_0 x - \frac{5}{2} \right)$$

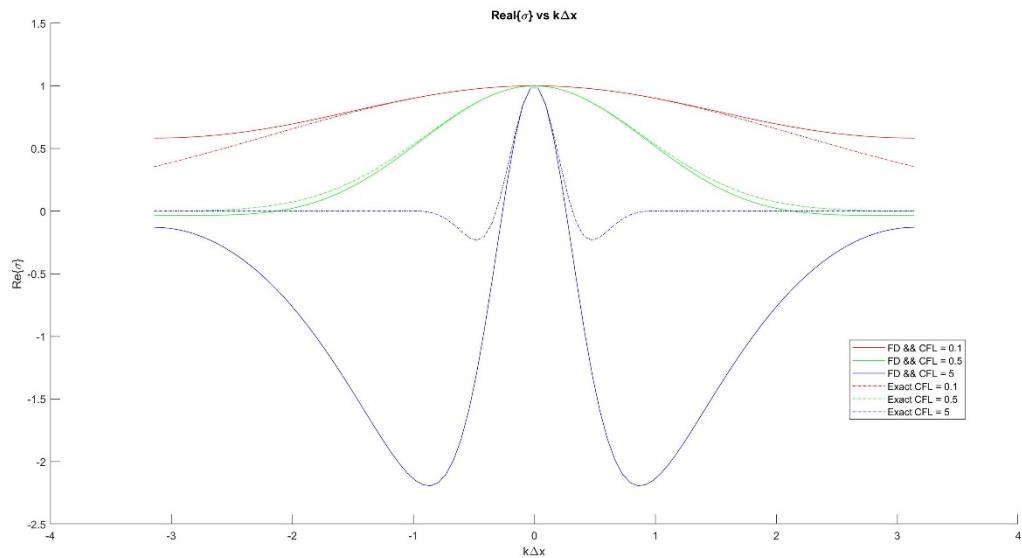
$$k_0 x = -\pi, \dots, \pi - \frac{2\pi}{N}$$

Thus, for $m=1$, $V^{m,1} = \sigma = \frac{1 + \gamma_1 \lambda_{R\text{at}} - \gamma_1 \gamma_2 (\lambda_{I\text{at}})^2}{(1 - \beta_1 \lambda_{R\text{at}})^2}$

$$+ i \frac{\lambda_{I\text{at}}(\gamma_1 + \gamma_1 \gamma_2 \lambda_{R\text{at}} + \gamma_1 - \gamma_1 \beta_1 \lambda_{R\text{at}})}{(1 - \beta_1 \lambda_{R\text{at}})^2}$$

Then we can use Matlab to plot σ against $k_0 x$

Problem3 (c) Plot



Problem 3

- 1d) Compare the σ given by FD and Exact Fourier.
- for $CFL = 0.1$. FD scheme works well for low wavenumber, but undershoot for high wavenumber
 - for $CFL = 0.5$. FD scheme also works well for low wavenumber, but overshoot for high wavenumber
 - for $CFL = 5$. FD scheme only converge for limited region of wavenumber. \times

The reason for this. When CFL is large. the range of dot will getting larger and larger and go further and further beyond the stable Region of the scheme

```

close all;clear;clc;
%Problem1
CFL = [0.1 0.5 5];
Re = 1;
kx = linspace(-pi,pi,100);
figure;hold on;
for i = 1:3
Real_lambdat(:,i) = (CFL(i)/Re).*((-1/6).*cos(2*kx)+(8/3).*cos(kx)-5/2);
Imag_lambdat(:,i) = CFL(i).*((1/6).*sin(2*kx)-(4/3).*sin(kx));
plot(Real_lambdat,Imag_lambdat);
end
legend('CFL = 0.1','CFL = 0.5','CFL = 5'); hold off;
title('Locus of
\lambda\Delta t'); xlabel('Imag\{\lambda\Delta t\}'); ylabel('R
e\{\lambda\Delta t\}');

%Problem3 (a)
x = linspace(-5,5,100); y = x; [X,Y]=meshgrid(x,y);
Z = abs(1+(X+i*Y)+(1/2)*(X+i*Y).^2);
figure; contour(X,Y,Z,[1,1], 'ShowText', 'on');
title('Stability Region');
xlabel('Re\{\lambda\Delta t\}'); ylabel('Imag\{\lambda\Delta t\}'); grid on;

%Problem3 (b)
sigma_b = (1+(sqrt(2)-1).*Real_lambdat)./(1-(1-sqrt(2)/2).*Real_lambdat).^2;
figure; plot(kx,sigma_b); title('\sigma of L
term'); xlabel('Re\{\lambda\Delta t\}');
ylabel('\sigma\{Re\(\lambda\Delta t\)\}'); legend('CFL = 0.1','CFL
= 0.5','CFL = 5');
ylim([-0.4 1.2])

%Problem3 (c)
alpha = sqrt(2)-1; beta1 = 1-sqrt(2)/2; gamma1 = sqrt(2)/2;
gamma2 = gamma1;
zeta = sqrt(2)-1;
%sigma = Amplification factor
sigma = (1+alpha.*Real_lambdat-
gamma1.*gamma2.* (Imag_lambdat).^2)./(1-
beta1.*Real_lambdat).^2 ...
+1i*Imag_lambdat.* (gamma1+gamma1.*gamma2.*Real_lambdat+zeta-
zeta.*beta1.*Real_lambdat)./(1-beta1.*Real_lambdat).^2;
for i = 1 : 3
sigma_exact(:,i) = exp((1i.*kx-kx.^2)*CFL(i));
end

```

```
figure; hold on; p1 = plot(kx,sigma); p2 =
plot(kx,sigma_exact,'-.');
set(p1,{'color'},{{[1 0 0];[0 1 0];[0 0 1]}});
set(p2,{'color'},{{[1 0 0];[0 1 0];[0 0 1]}});
xlabel('k\Delta x'); ylabel('Re{\sigma}');
title('Real{\sigma} vs k\Delta x');
legend('FD && CFL = 0.1','FD && CFL = 0.5','FD && CFL =
5','Exact CFL = 0.1','Exact CFL = 0.5','Exact CFL = 5');
ylim([-2.5 1.5])
```