

On supersymmetric AdS_4 orientifold vacua

Joan Quirant



Based on: [2003.13578](#) with F. Marchesano, E. Palti and A. Tomasiello



SEMINAR SERIES ON STRING PHENO

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Mathematically: ✓

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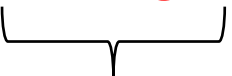
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$$m \sim \Lambda^{1/2}$$

Lüst, Palti, Vafa '19

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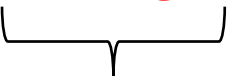
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But... These vacua are **not solving** the **10D** equations of motion. Difficulties arise because of the O-planes.

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Recent related work:
Junghans '20;
Buratti, Calderón,
Mininno, Uranga '20

We find an **approximate** solution to the **uplift** problem.
Approximate: at first order in an expansion parameter.

SUSY AdS_4 IIA flux vacua

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 - 1) We want a solution to the e.o.m with a metric: $ds^2 = e^{2A} ds_{AdS_4}^2 + ds^6$
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Graña , Minasian , Petrini , Tomasiello '04,...

 - 2.1) Restricts X_6 to be a manifold with either $SU(3)$ or $SU(3) \times SU(3)$ structure
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SUSY AdS_4 IIA flux vacua

- A key characteristic of this family of vacua is

$$\langle G_6 \rangle \propto \int_{X_6} G_6 = 0$$

($\rho_0 = 0$ in the language of Herráez, Ibáñez, Marchesano, Zoccarato, '18)

$SU(3)$ or $SU(3) \times SU(3)$

DGKT vacua?



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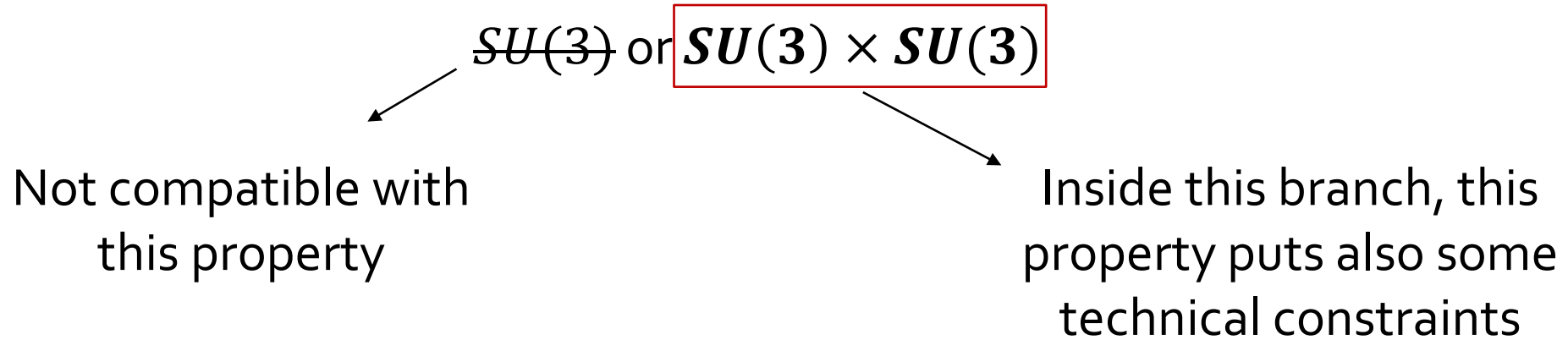
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$SU(3)$ or $SU(3) \times SU(3)$?

DGKT vacua?

- This is a 4d property... But it restricts the possible 10d uplift:



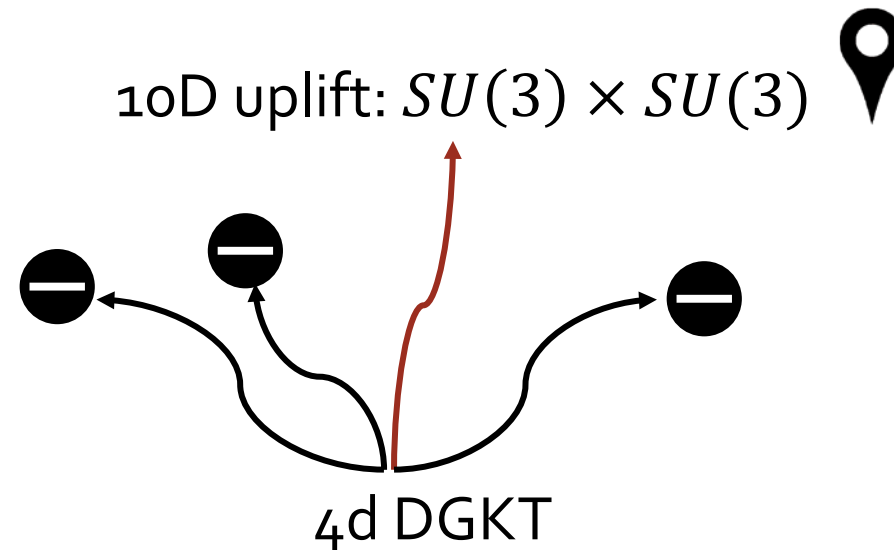
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Having set the
framework we
are ready to
deal with the
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Large volume approximation

- Inspired by Saracco, Tomasiello '12 we will solve the equations in the limit $\hat{\mu} = l_s \mu = l_s \sqrt{-\Lambda/3}$ small or, equivalently, g_s small, R_{internal} big:

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- If the source is **smeared** $\delta_{O6} \rightarrow \# \text{Re}\Omega$ one has the following scaling:

$$F_4 \sim N, \quad R_{X_6} \sim N^{1/4}, \quad g_s \sim \mu \sim R_{X_6}^{-3} \sim N^{-3/4}$$

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- We will expand the solution to the vanishing SUSY variations in terms of g_s in the limit $g_s \sim 0$

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- 4) For example, at leading order the dilaton ϕ and the warp factor A are no longer constant

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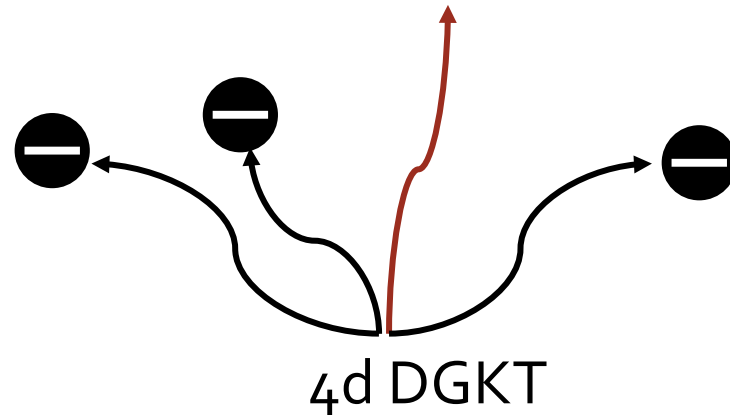
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Vanishing SUSY variations

Bianchi identities

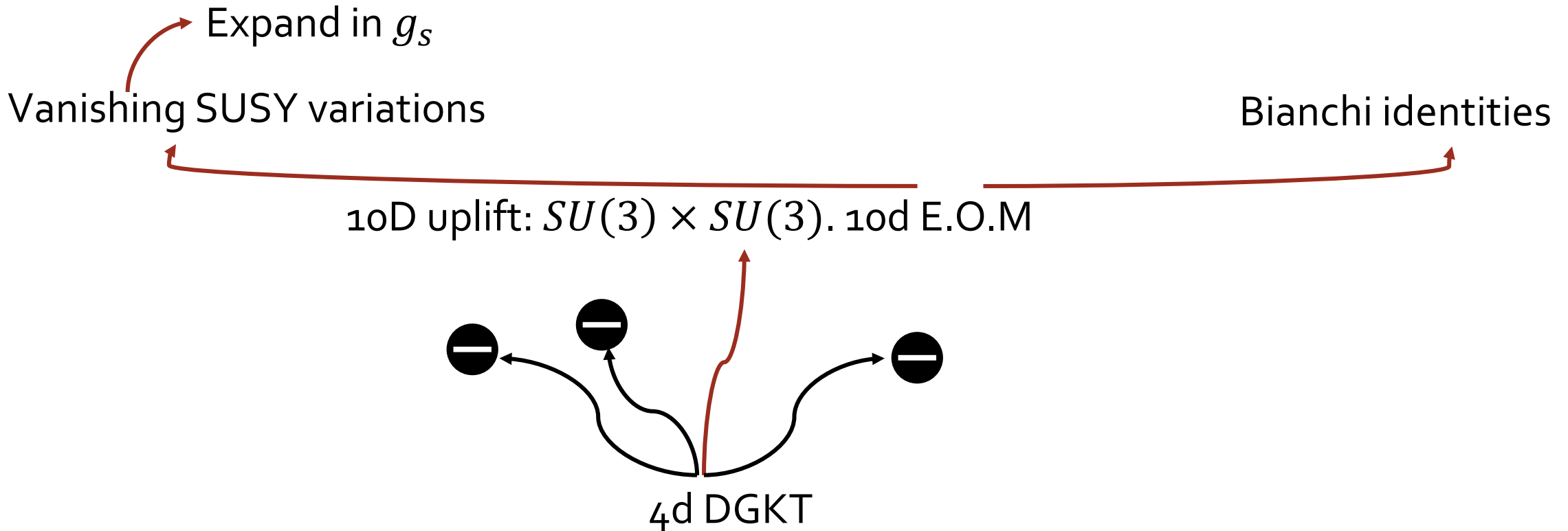
10D uplift: $SU(3) \times SU(3)$. 10d E.O.M



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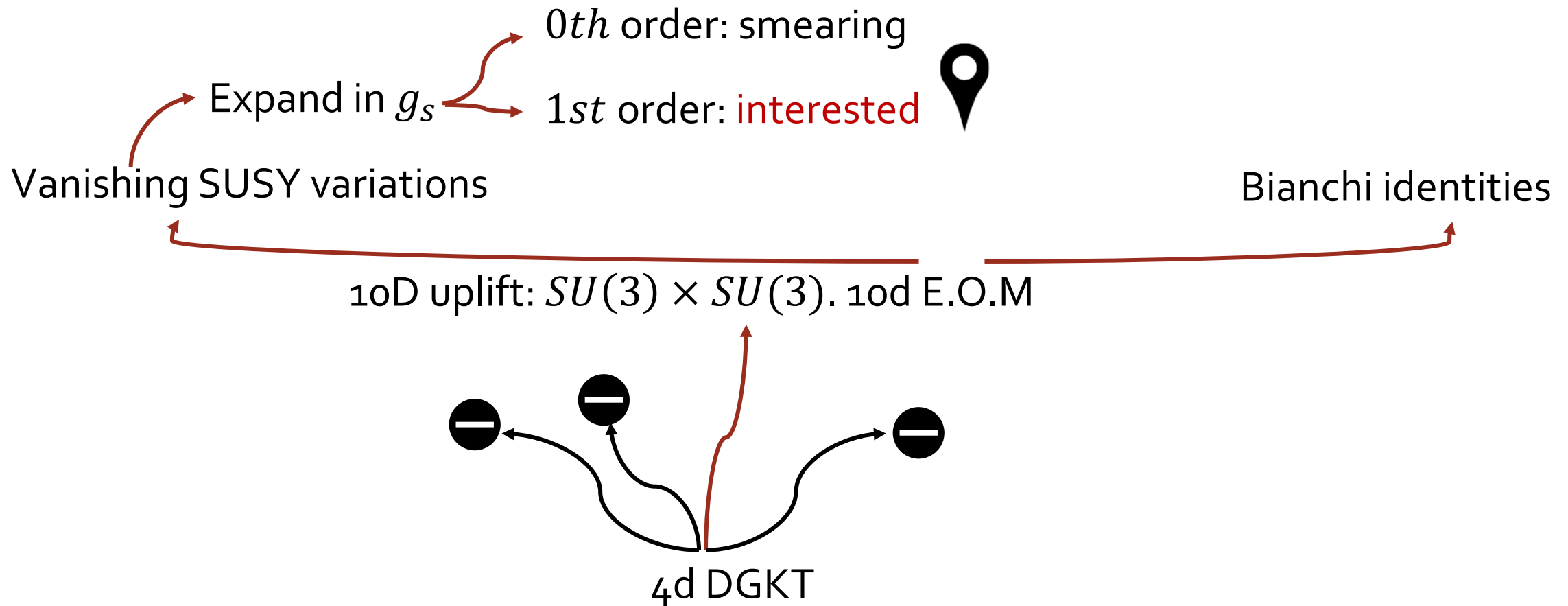
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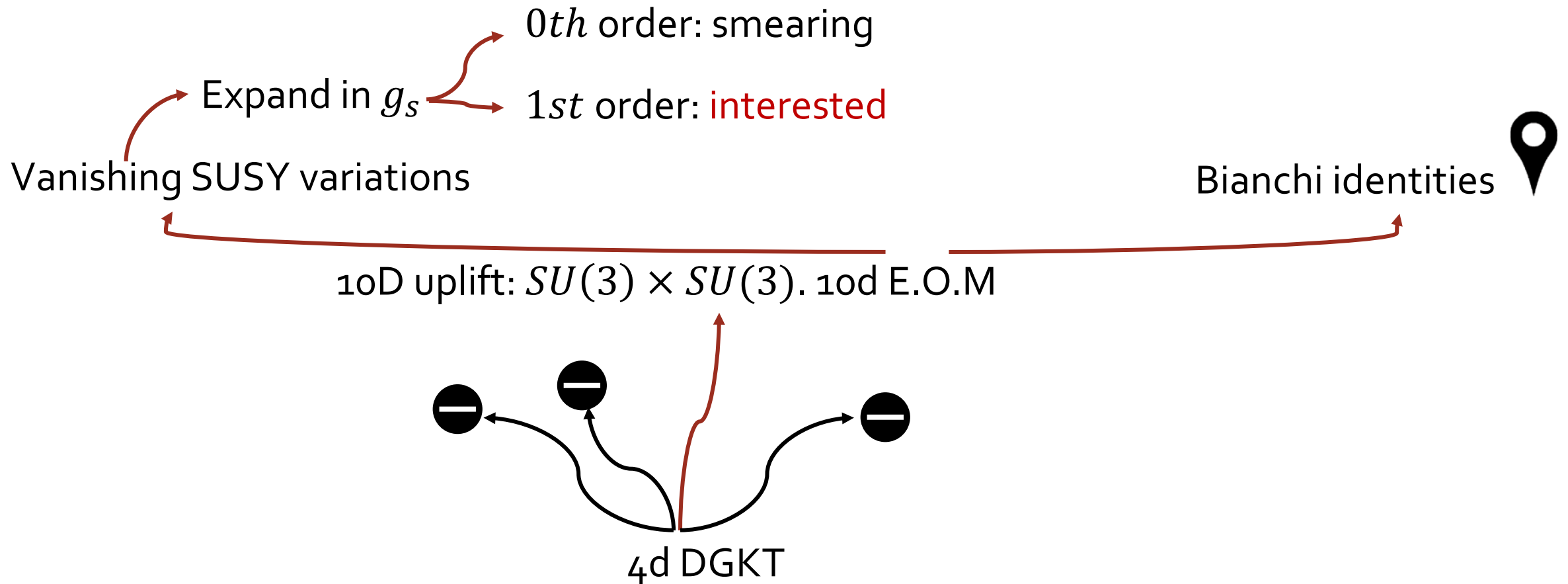
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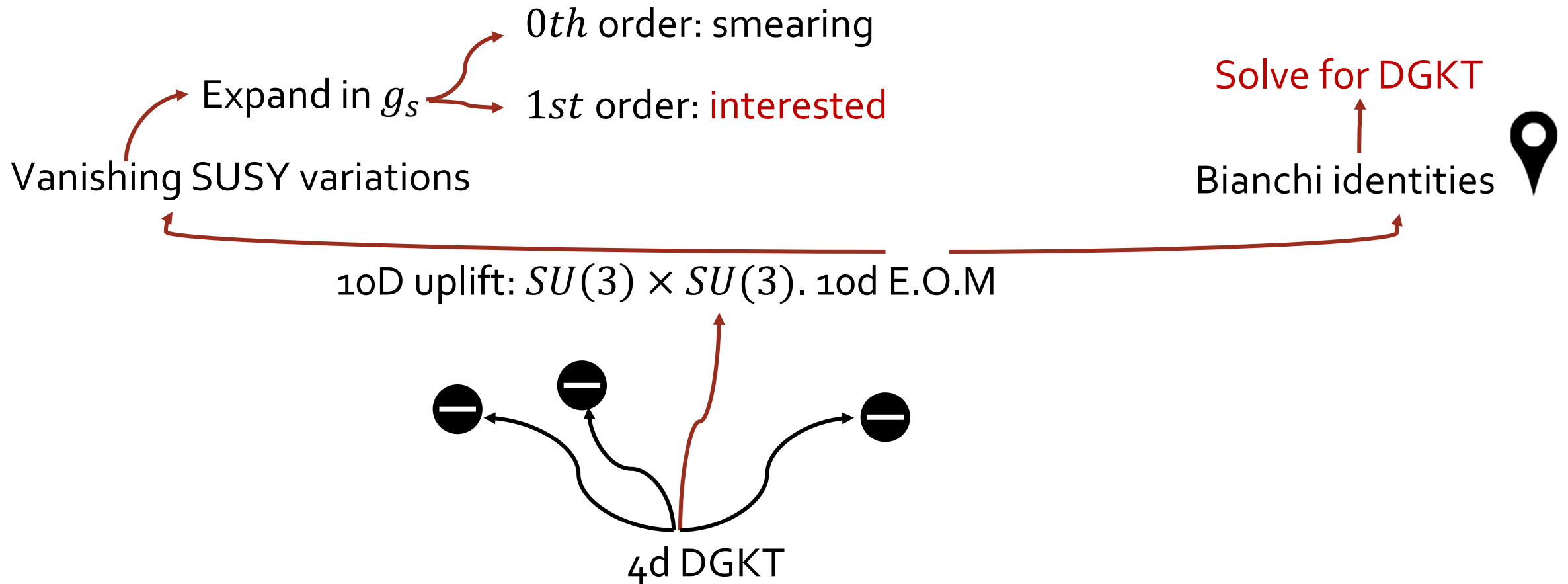
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Solving the equations: B.I.

- Forget for the moment the SUSY equations. Let us solve the B.I. for DGKT. Then we will use the expansion of the SUSY solution to accommodate this result ☺. Consider:

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with dC_1 exact, \widetilde{F}_2^h CY harmonic and $\Delta K = 2m\hat{\mu} \text{Re } \Omega_{CY} + \delta(\Pi_{06})$

- The rest of BI follows easily from this solution.

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This H is an approximation of the *actual* H. In the expansion it will correspond to the leading term

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+ expansion $F_2 = -\frac{1}{g_s} J \cdot d(e^{-3A_0} \text{Im} \Omega) + \mathcal{O}(g_s), \quad (\Omega, J) \text{ defining an } SU(3) \text{ structure,}$

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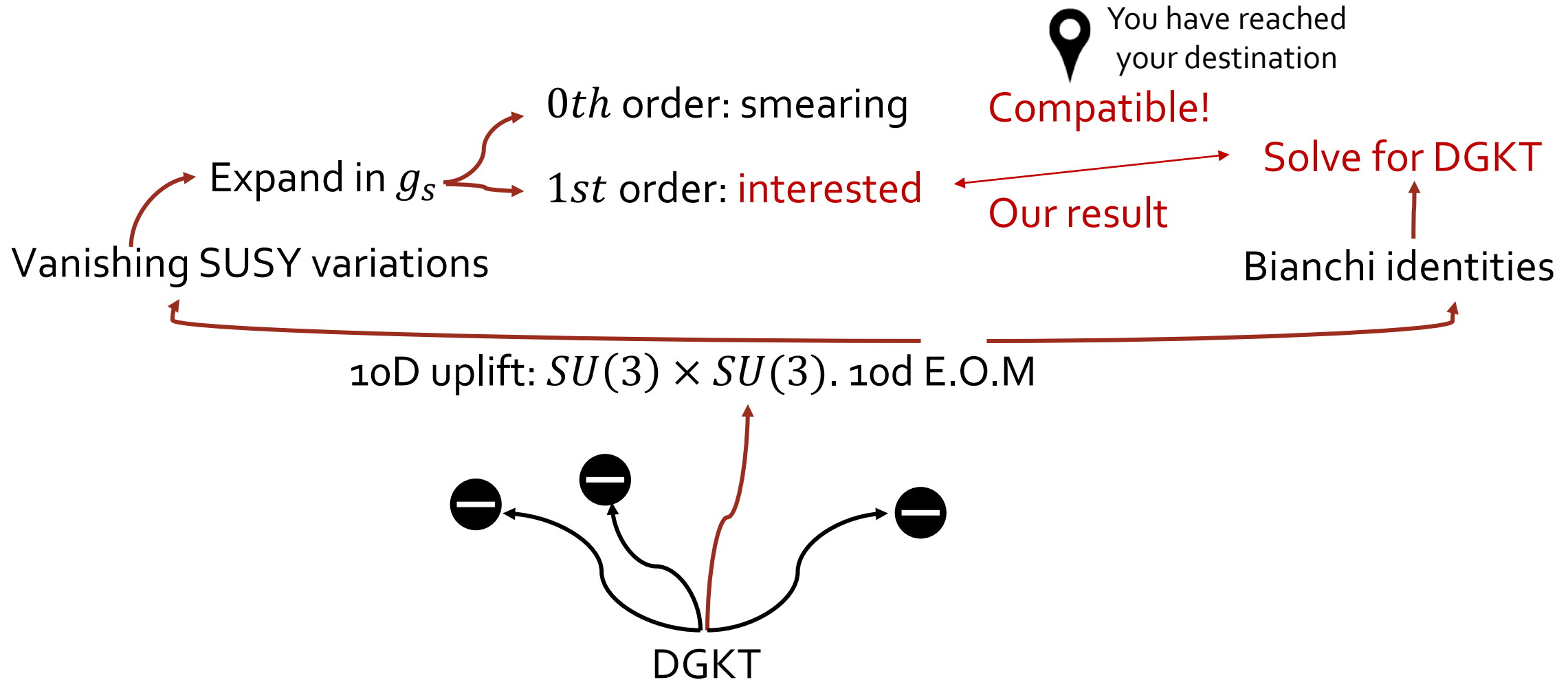
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- Same procedure for the rest of functions...
- Explicit toroidal example in section 6.2

Solving the equations: B.I.+SUSY

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Solving the equations: B.I.+SUSY

- Final result:

$$H = g_s \frac{2}{5} G_0 \text{Re } \Omega_{\text{CY}} \left(\mathbf{1} + O\left(g_s^{4/3}\right) \right), \quad e^{-A} = \mathbf{1} + O\left(g_s^{4/3}\right),$$

$$G_2 = O\left(g_s^{2/3}\right), \quad e^{\phi} = g_s \left(\mathbf{1} + O\left(g_s^{4/3}\right) \right),$$

$$G_4 = \frac{3}{10} G_0 J_{\text{CY}} \left(\mathbf{1} + O\left(g_s^{4/3}\right) \right), \quad J = J_{\text{CY}} \left(\mathbf{1} + O\left(g_s^{4/3}\right) \right),$$

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- At 0^{th} order we recover the usual **smeared** solution of Acharya, Benini, Valandro '06
- At **first order corrections** appear. The B.I. is solved with the *smeared* H
- The **obstruction** present for the $SU(3)$ case is **avoided** in this solution (at the level of approximation we are working) 

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Thank you for your attention! 