# Dynamical Tadpoles, Stringy Cobordism and the SM from Spontaneous Compactification



Ginevra Buratti

GB, M. Delgado, A. Uranga [arXiv:2104.02091]

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### Plan of the talk

Introduction and motivation

Dynamical tadpoles and the swampland cobordism conjecture

■ Two tadpole lessons from string theory examples

Conifold

3-form flux models

Magnetized branes

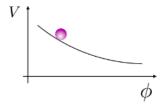
non susy 10d USp(32) theory

SM from spontaneous compactification

Conclusions



# Dynamical tadpoles



- theories sitting on the **slope** of some scalar potential
  - ⇒ dynamical tadpoles (as opposed to topological tadpoles)
- properties of the resulting spacetime-dependent solutions?

## Cobordism defects

■ n-dim manifold can/cannot be the boundary of (n+1)-dim manifold

 $\mathbf{S}^1$  is trivial



Point is non-trivial

cobordism charge can be removed by localized sources

 $S^2$  with a U(I)  $F_2$  flux can't just shrink



Flux removed by monopole





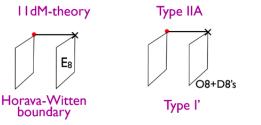
## Cobordism defects

swampland cobordism conjecture

[McNamara Vafa]

 $\Omega_{QG}=0$  related to no global symmetries

examples



Type IIB on S<sup>1</sup>
With -I SL(2,Z) WL

Etc...
Half T<sup>2</sup>/Z<sub>2</sub>

not all defects are known

## Two lessons

#### Finite Distance:

The running solution extends at most a distance  $\Delta$  scaling as

$$\Delta^{-n} \sim \mathcal{T}$$

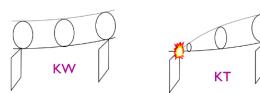
with the strength of the tadpole  $\mathcal{T}$ .

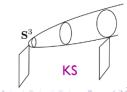
#### **Dynamical Cobordism:**

Spacetime is cut off at this distance by the cobordism defect of the swampland cobordism conjecture.

# $AdS_5 \times T^{1,1}$

- AdS<sub>5</sub> × T<sup>1,1</sup> with **N** units of RR 5-form flux near horizon limit of N D3's at conifold singularity  $T^{1,1} = S^2 \times S^3 \text{ is 5d base of 6d cone}$
- lacktriangle Add lacktriangle units of RR 3-form flux on  $S^3$  [Klebanov Tseytlin] running  $T^{1,1}$  geometry and singularity at finite distance  $r_0$
- $lue{S}$  Smooth out the singularity by finite size  $S^3$





[Klebanov Strassler]

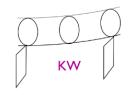
[Klebanov Witten]

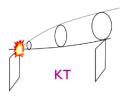
# $AdS_5 \times T^{1,1}$

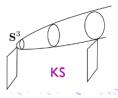
- $\blacksquare$  dilaton tadpole  $V(\phi)\sim M^2e^\phi$   $\nabla^2\phi=-e^{-6q-\phi}(\partial\Phi)^2+e^{-14q+\phi}M^2$
- solved by running NSNS axion  $\Phi = 3g_s M \log(r/r_0)$

# Finite Distance $\Delta^{-1} \sim M^2 e^\phi \sim \mathcal{T}$

## **Dynamical Cobordism**







## 3-form flux models

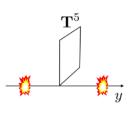
- type IIB on T<sub>5</sub> with RR 3-form flux  $F_3 = Ndx^1dx^2dx^3$
- dilaton tadpole  $\nabla^2\phi\sim e^\phi(F_3)^2-e^{-\phi}(H_3)^2$  solved by  $H_3=Ndy^1dy^2dy$   $\Rightarrow$   $\Phi\sim Ny$

#### **Finite Distance**

$$\begin{split} ds_{10}^2 &= Z^{-\frac{1}{2}}ds_4^2 + Z^{\frac{1}{2}}R^2 \left( dz^1 d\overline{z}^1 + dz^2 d\overline{z}^2 + dz^3 d\overline{z}^3 \right) \\ \text{with } &-\tilde{\nabla}^2 Z = \frac{g_s}{6} \left( F_3 \right)^2 \quad \Rightarrow \quad Z = 1 - \frac{g_s}{12} (F_3)^2 y^2 \\ \text{singularities at } y^{-2} &= \frac{1}{12} g_s (F_3)^2 \quad \Rightarrow \quad \Delta^{-2} \sim \mathcal{T} \end{split}$$



singularities removed by O3's (possibly with D3's)





## Magnetized branes

- $\blacksquare$  type IIB on  $\mathsf{T}^2_{(1)}\,\times\mathsf{T}^2_{(2)}$  with O7's and D7's transverse to  $\mathsf{T}^2_{(1)}$
- Add **M** units of D7 worldvolume magnetic flux on  $T_{(2)}^2$  susy breaking and tadpole for dilaton and  $T_{(2)}^2$  Kähler modulus
- Solve by magnetic field  $F_2 = F(dz_2d\overline{z}_2 dz_3d\overline{z}_3)$

#### **Finite Distance**

lift to F-theory  $G_4 \sim F_2 \wedge \omega_2$  singular warp factor  $\Delta^{-2} \sim (F_2)^2 \sim \mathcal{T}$ 

## **Dynamical Cobordism**

compactification on extra  $\mathsf{T}_2/\mathsf{Z}_2$  with additional O7's and D7's



# 10d non susy USp(32) string

 $\blacksquare$  orientifold of IIB with  $O9^+$  and 32  $\overline{D9}\mbox{'s}$ 

[Sugimoto]

dilaton tadpole

$$S_E = \frac{1}{2\kappa^2} \int d^{10}x \sqrt{-G} \left[ R - \frac{1}{2} (\partial \phi)^2 \right] - T_9^E \int d^{10}x \sqrt{-G} \, 64 \, e^{\frac{3\phi}{2}}$$

■ running solution  $\phi = \frac{3}{4}\alpha_E y^2 + \frac{2}{3}\log|\sqrt{\alpha_E}y| + \phi_0$  [Dudas Mourad]  $ds_E^2 = |\sqrt{\alpha_E}y|^{\frac{1}{9}} e^{-\frac{\alpha_E y^2}{8}} \eta_{\mu\nu} dx^{\mu} dx^{\nu} + |\sqrt{\alpha_E}y|^{-1} e^{-\frac{3\phi_0}{2}} e^{-\frac{9\alpha_E y^2}{8}} dy^2$ 

singularities at  $y = 0, \infty$  separated by

Finite Distance 
$$\Delta \sim \int_0^\infty \sqrt{g_{yy}} \, dy \sim e^{-\frac{3\phi_0}{4}} \alpha_E^{-\frac{1}{2}} \Rightarrow \Delta^{-2} \sim \mathcal{T}$$

# 10d non susy USp(32) string

#### **Dynamical Cobordism**

strong coupling defect, that is able to gap chiral non anomalous content

■ Solve by magnetization:  $\mathsf{T}^6/(\mathsf{Z}^2{\times}\mathsf{Z}^2)$  with  $\mathsf{O9}^+$  and  $\mathsf{8}$   $\mathsf{O5}^-_i$ 

obj.	$N_{\alpha}$	$(n_{\alpha}^1, m_{\alpha}^1)$	$(n_\alpha^2,m_\alpha^2)$	$(n_{\alpha}^3, m_{\alpha}^3)$
O9 <sup>+</sup>	32	(1,0)	(1,0)	(1,0)
O5 <sub>1</sub>	-32	(1,0)	(0,1)	(0, -1)
$O5_{2}^{-}$	-32	(0,1)	(1,0)	(0, -1)
$O5_{3}^{-}$	-32	(0,1)	(0, -1)	(1,0)
D9	16	(-1,1)	(-1,1)	(-1,1)
D9'	16	(-1, -1)	(-1, -1)	(-1, -1)

 $n_{\alpha}^{i}$  wrapping on  $\mathsf{T}_{i}^{2}$   $m_{\alpha}^{i} \text{ magnetic flux}$  on  $\mathsf{T}_{i}^{2}$ 

# SM from spontaneous compactification

■ type I on  $(T^2 \times T^2)/Z^2$  with 32 magnetized D9's and O9, O5's

$N_{\alpha}$	$(n^1_{\alpha}, m^1_{\alpha})$	$(n_{\alpha}^2, m_{\alpha}^2)$
$N_{a+d} = 6 + 2$	(1,3)	(1, -3)
$N_{h_1} = 4$	(1, -3)	(1, -4)
$N_{h_2} = 4$	(1, -4)	(1, -3)
40	(0,1)	(0, -1)

 $n_{lpha}^{i}$  wrapping on  $\mathsf{T}_{i}^{2}$   $m_{lpha}^{i}$  magnetic flux on  $\mathsf{T}_{i}^{2}$ 

T dual to intersecting D7's

RR tadpoles 
$$\sum_{\alpha}N_{\alpha}n_{\alpha}^{2}n_{\alpha}^{3}=16$$
  $\sum_{\alpha}N_{\alpha}m_{\alpha}^{2}m_{\alpha}^{3}=-16$ 

non susy  $\Rightarrow$  dynamical tadpoles for inverse areas of  $T^2$ 's

# SM from spontaneous compactification

Solve by magnetization along two of the 6d spacetime dimensions

$$\Rightarrow~(\mathsf{T}^2{\times}\mathsf{T}^2{\times}\mathsf{T}^2)/(\mathsf{Z}^2{\times}\mathsf{Z}^2)$$
 with extra O5's, D5's

$N_{\alpha}$	$(n^1_{\alpha}, m^1_{\alpha})$	$(n_{\alpha}^2, m_{\alpha}^2)$	$(n_{\alpha}^3, m_{\alpha}^3)$
$N_{a+d} = 6 + 2$	(1,3)	(1, -3)	(1,0)
$N_b = 2$	(0,1)	(1,0)	(0,1)
$N_c = 2$	(-1,0)	(0, -1)	(0,1)
$N_{h_1} = 2$	(1, -3)	(1, -4)	(2,-1)
$N_{h_2} = 2$	(1, -4)	(1, -3)	(2,-1)
40	(0,1)	(0, -1)	(0,1)

## RR tadpoles

$$\sum_{\alpha} N_{\alpha} n_{\alpha}^{1} n_{\alpha}^{2} n_{\alpha}^{3} = 16$$

$$\sum_{\alpha} N_{\alpha} n_{\alpha}^{1} m_{\alpha}^{2} m_{\alpha}^{3} = 16$$

$$\sum_{\alpha} N_{\alpha} m_{\alpha}^{1} n_{\alpha}^{2} m_{\alpha}^{3} = 16$$

$$\sum_{\alpha} N_{\alpha} m_{\alpha}^{1} m_{\alpha}^{2} n_{\alpha}^{3} = -16$$

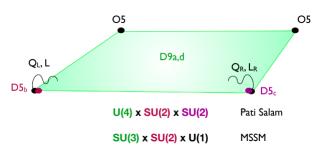
susy 
$$\chi_1=\chi_2$$
  $\chi_3=\frac{14\chi_1}{1-12\chi_1^2}$   $\Rightarrow$  no dynamical tadpoles



# SM from spontaneous compactification

3-family MSSM-like spectrum

[Marchesano Shiu]



Note that all MSSM but gluons/inos arise from cobordism branes

## Conclusions

- We have studied the properties of space-dependent solutions in theories with dynamical tadpoles
- Two lessons: Finite Distance and Dynamical Cobordism

Many open questions

Time-dependent backgrounds

More non susy examples

Links to other swampland conjectures

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