

# The Price of Curiosity:

## *Information Recovery in de Sitter Space*



Based on upcoming work with Watse Sybesma

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Seminar Series on String Phenomenology  
March 30, 2021



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# Cosmological horizons

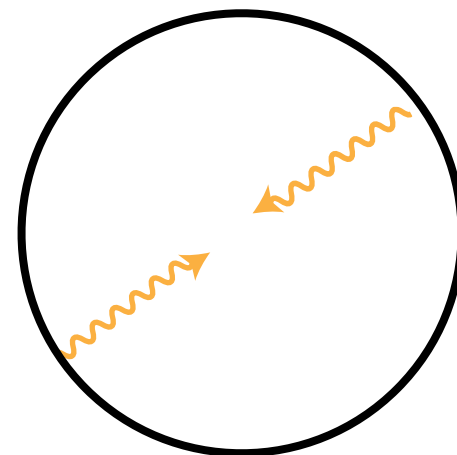
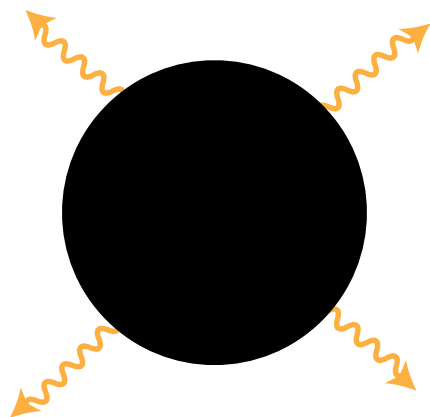
A static observer in de Sitter is surrounded by a **cosmological horizon**.

The horizon..

- Has a temperature:  $T_{\text{dS}} = \frac{1}{2\pi\ell}$
- And an entropy:  $S_{\text{dS}} = \frac{\text{Area}}{4G_N}$

[Gibbons, Hawking '77]

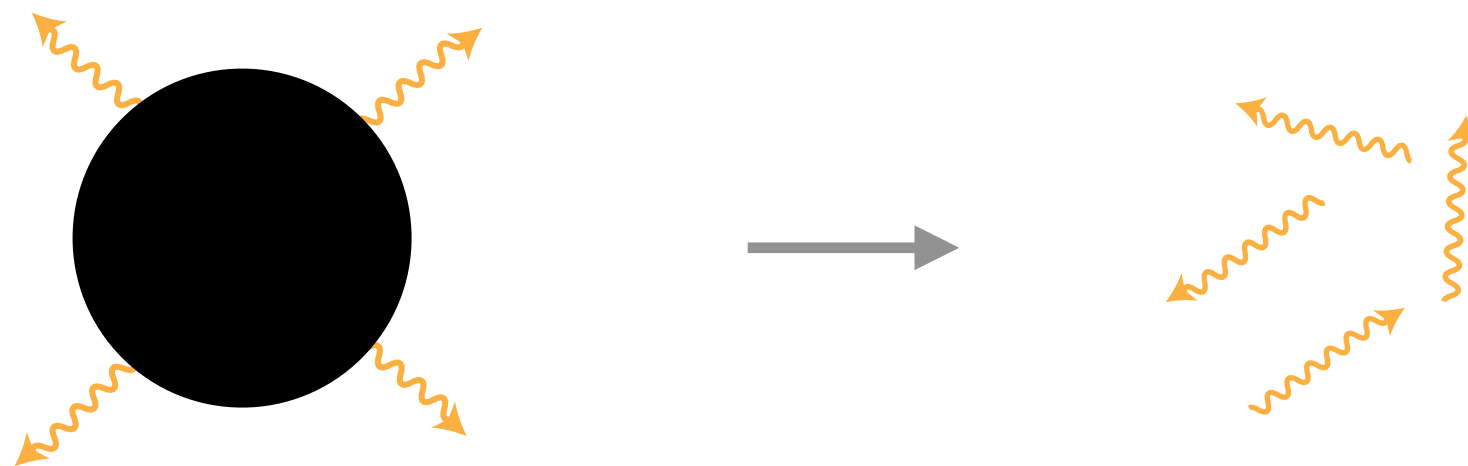
And shares many similarities with a **black hole horizon**.



# Entropy constraints

The **finite black hole entropy** can be associated to **microstates**. Is there a similar interpretation for the **de Sitter entropy**?

Together with **unitarity**, finiteness of the entropy places **constraints** on black hole evaporation.



To resolve the **black hole information paradox**, new effects should correct the EFT describing Hawking radiation.

# Corrections to BHs and dS

In last ~2 years, it has been realised what **new effects** should be included to resolve the information paradox.

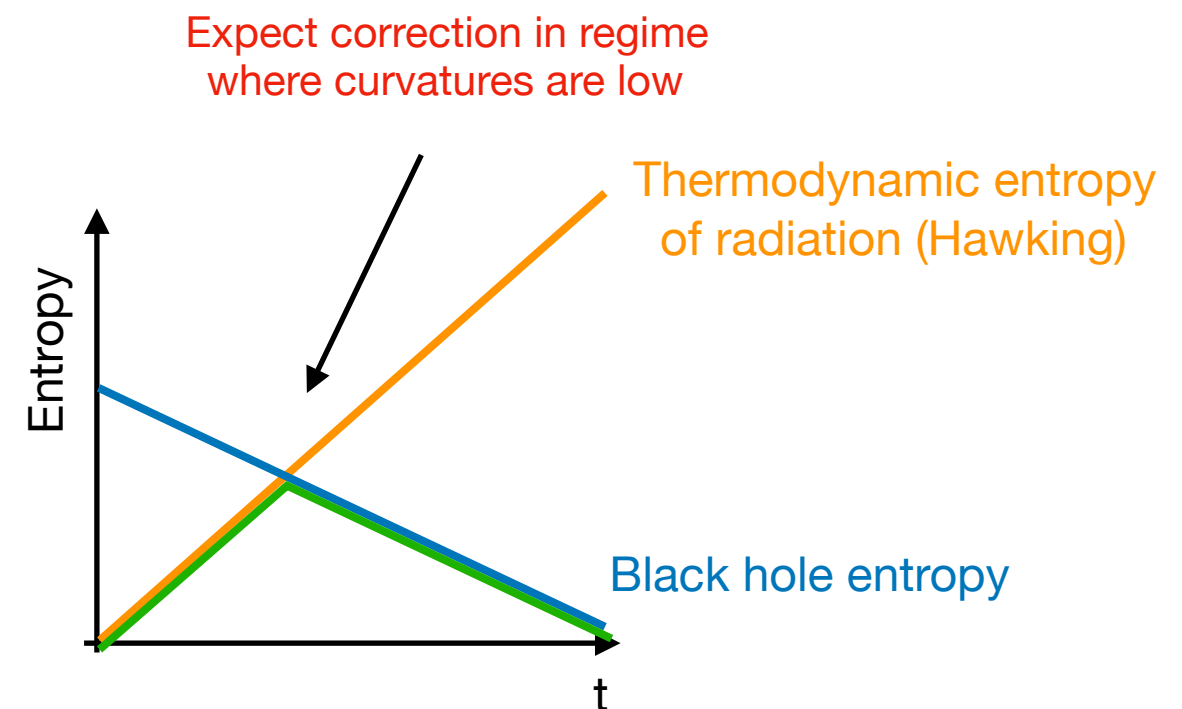
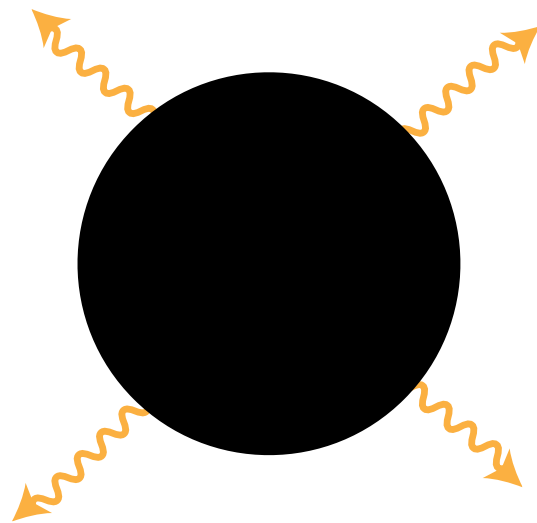
These **new contributions** do not rely on the specific background. Can similar effects occur in **de Sitter space**?

## Today we'll see:

- A review of **recent developments** resolving the info paradox.
- How these results can be applied to **JT gravity in  $dS_2$** .
- This implies **information can be recovered** behind the cosmological horizon.
- This experiment does not end well for a **static observer**.

# BH information paradox

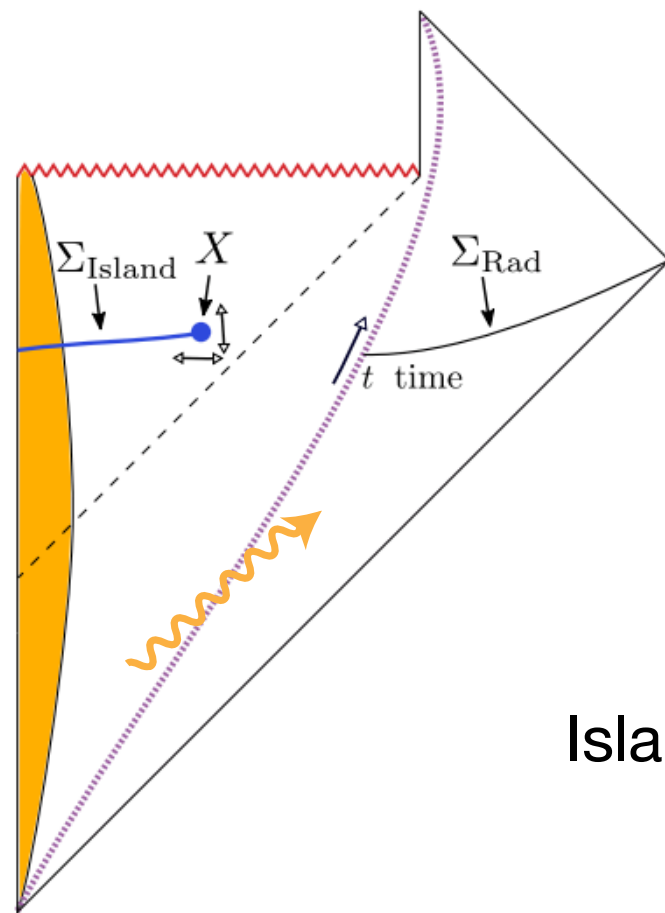
The fact that Hawking radiation is thermal, leads to a paradox.



- Unitary evolution requires the radiation entropy to follow the green (Page) curve. [Page '93]
- Recently, a new "island" contribution has been found that reproduces the Page curve. [Almheiri, Hartman, Engelhardt, Maldacena, Marolf, Maxfield, Penington, Shaghoulian, Tajdini + many more works]

# The island formula

The key insight is that a **disconnected** (island) region can contribute to the entropy of radiation.



$$S_{\text{Rad}} = \min \text{ext}_X \left[ \frac{\text{Area}(X)}{4G_N} + S_{\text{vN}}(\Sigma_{\text{Rad}} \cup \Sigma_{\text{Island}}) \right]$$

Extremize X to find a quantum extremal surface and pick minimum entropy

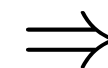
Island contribution appears as a new saddle in the Euclidean path integral

1) At early times, trivial island dominates:

$$S_{\text{Rad}} \simeq S_{\text{vN}}(\Sigma_{\text{Rad}})$$

2) After Page time, island starts dominating:

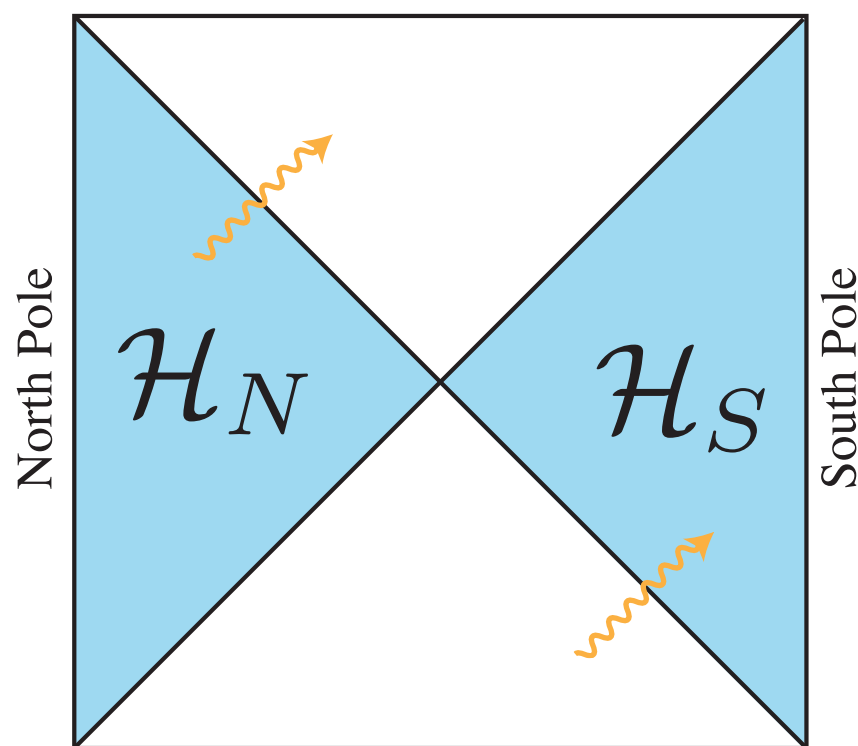
$$S_{\text{Rad}} \simeq \frac{\text{Area}(X)}{4G_N}$$



**Results in Page curve**

# Entanglement in de Sitter space

The natural vacuum state in de Sitter space is the **Bunch-Davies state**.



Side remark:  $\langle \mathcal{O} \rangle = \text{tr}_S(\mathcal{O}\rho_S)$

Tracing does not break the de Sitter isometries.

$$\rho = \frac{1}{Z} \sum_i e^{-\beta E_i} |E_i\rangle_N |E_i\rangle_S \langle E_i|_N \langle E_i|_S$$

$\beta = 2\pi\ell$  (pointing to  $\beta$ )  
 Pure state. (pointing to  $\rho$ )

Tracing over NP:

$$\rho_S = \text{tr}_N(\rho) = \frac{1}{Z} \sum_i e^{-\beta E_i} |E_i\rangle_S \langle E_i|_S$$

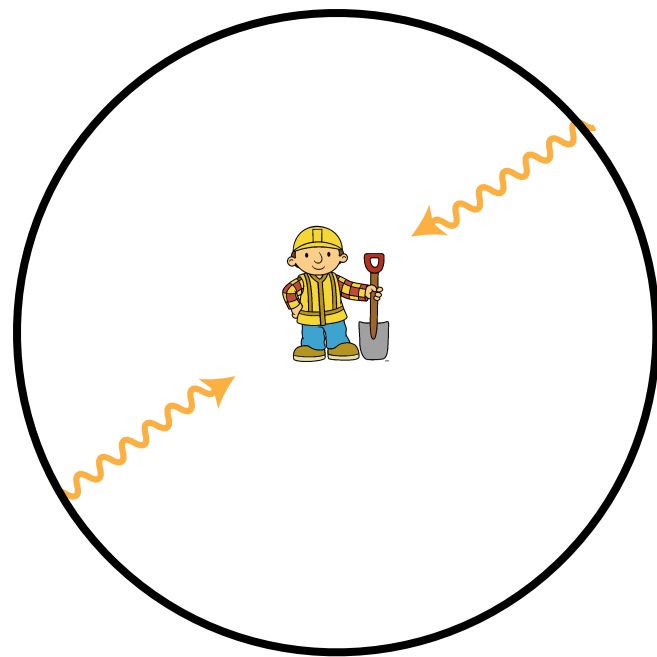
Maximally entangled state. (pointing to  $\rho$ )

Large entanglement entropy suggests possibility of information recovery.

# Islands in de Sitter space?

To decode, **islands** should contribute to the entropy in de Sitter space.

[Hartman, Jiang, Shaghoulian '20] [Chen, Gorbenko, Maldacena '20] [Chen, Gorbenko, Maldacena '20]  
[Balasubramanian, Kar, Ugajin '20][Sybesma '20][Geng, Nomura, Sun '21]



Decoding information requires  
radiation to be stored in the  
static patch.

Positive energy shrinks the horizon,  
breaking the thermal equilibrium!

We will go beyond the **thermal equilibrium** studied so far and  
compute **backreaction** in JT gravity.



# JT Gravity in $dS_2$

[Maldacena, Turiaci, Yang '19]

JT gravity can be obtained as the  $dS_2 \times S^2$  near-horizon limit of a **near-Nariai black hole**.

$$I = I_0(\Phi_0) + \int d^2x \Phi \left( R - \frac{2}{\ell^2} \right) + I_{\text{CFT}}$$

Topological term

Deviation from Nariai limit

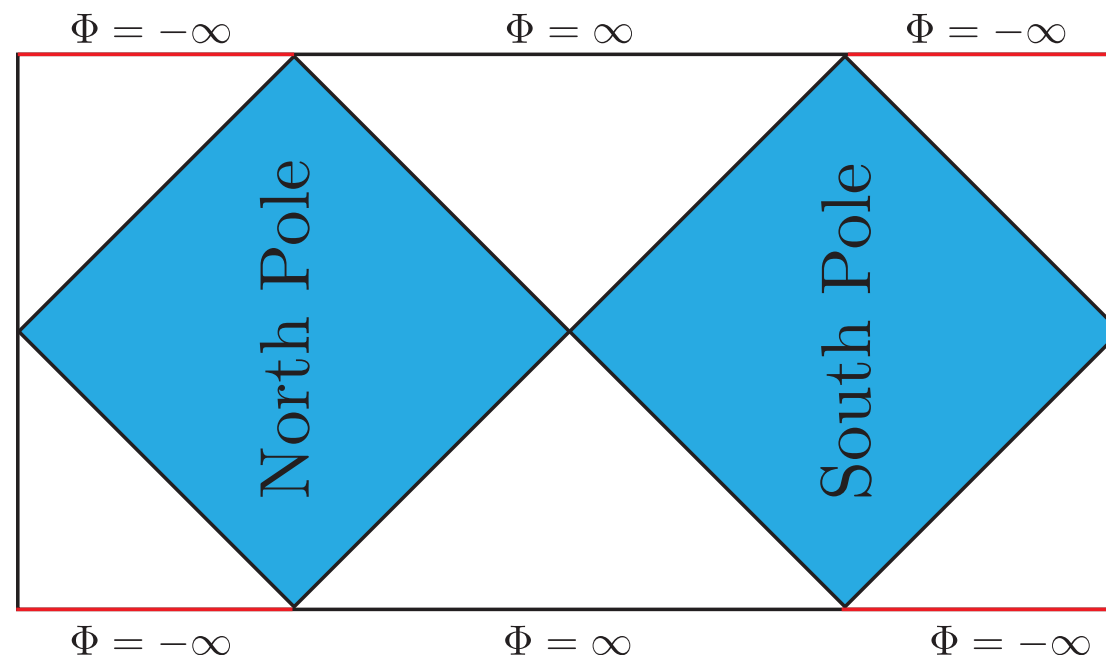
Large c matter sector

Picking coordinates that cover the static patch:  $ds^2 = -\sin^2 \theta dt^2 + \ell^2 d\theta^2$

$$\Phi = \frac{\Phi_s}{24} \cos \theta$$

$$S_{\text{dS}} = 2\Phi_0 + 2\Phi(\theta = 0)$$

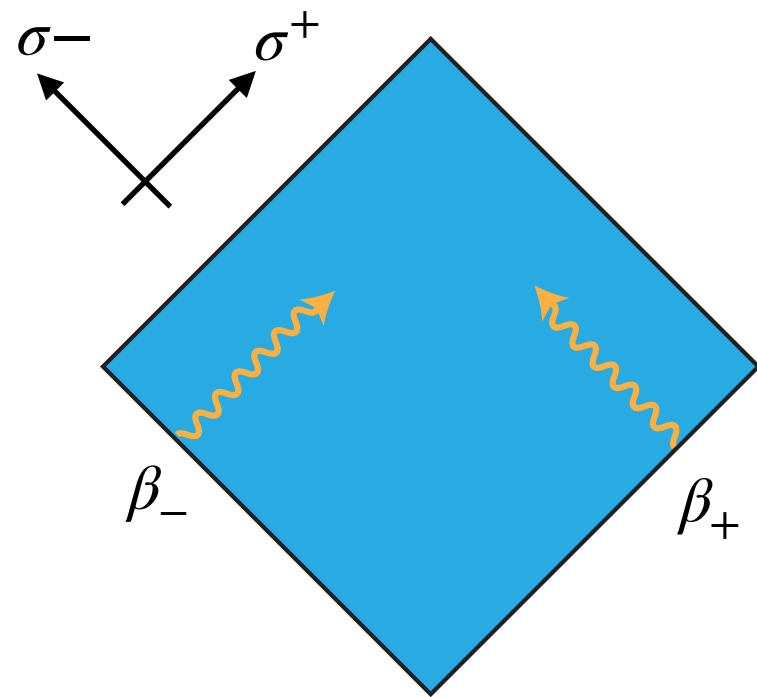
Large dilaton corresponds to weak gravitational coupling



# Including matter

Including conformal matter, we can compute the **quantum-corrected geometry**.

Using null coordinates  $\sigma^\pm = t \pm r_*$ , we find in the **Bunch-Davies state**:



$$\langle T_{\pm\pm}(\sigma^\pm) \rangle = \text{tr}_S (T_{\pm\pm}(\sigma^\pm) \rho_S) = \frac{\pi c}{12\beta_\pm^2}$$
$$\Phi = \frac{\Phi_s}{24} \cos \theta + \frac{c}{24}$$

$\beta_\pm = 2\pi\ell$

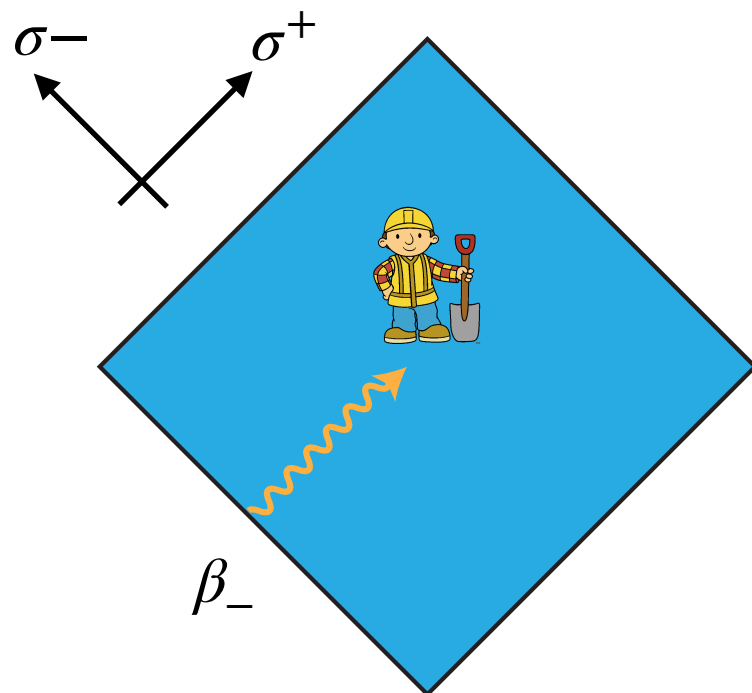
This describes a thermal equilibrium with **no backreaction**.

# State of a curious observer

[LA, Sybesma]

[LA, Parikh, van der Schaar '19]

Now consider an observer at the South Pole that collects right-moving radiation with the aim of **decoding information**.



$$\langle T_{--}(\sigma^-) \rangle = \frac{\pi c}{12\beta_-^2}$$

$$\langle T_{++}(\sigma^+) \rangle = 0$$

Effectively sets left-moving temperature to zero.

Working in  $x^\pm$  (Kruskal) coordinates, the **backreacted** solution is:

$$x^\pm = \pm \ell e^{\pm \sigma^\pm / \ell}$$

$$\Phi = \frac{c}{48} \left[ 1 + \frac{2\Phi_s}{c\ell} r - \frac{r}{\ell} \log(x^+ / \ell) \right]$$

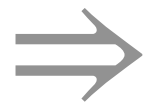
Explicit time dependence.

# dS information paradox

The entropy of the **collected radiation** increases, while the **horizon entropy** decreases.

1st law:  $dS_{\text{dS}} = -\frac{c}{24\ell}dt$

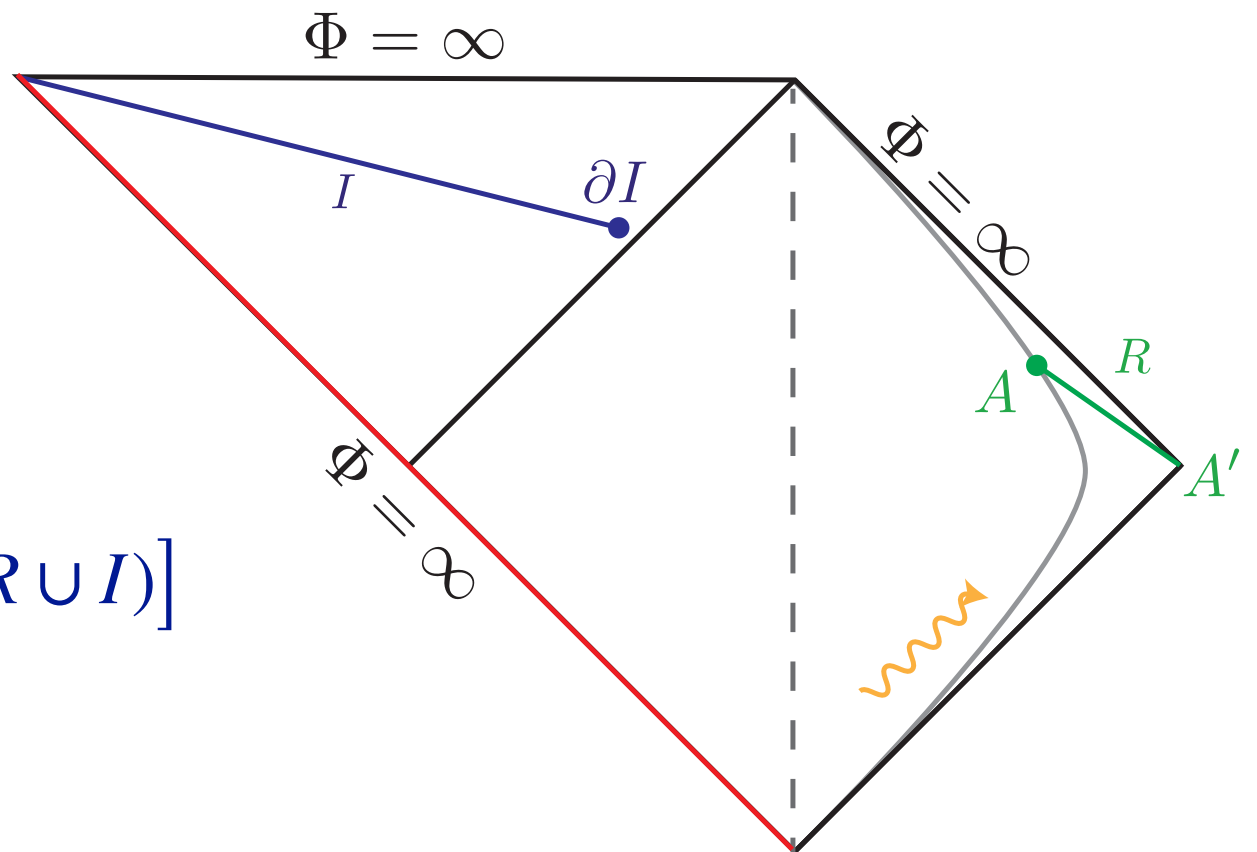
CFT:  $dS_{\text{Rad}} = \frac{c}{12\ell}dt$



Without island contribution,  
violates unitarity.

We now allow for the  
contribution of an **island**.

$$S(R) = \mathbf{min\,ext}_{\partial I} [2\Phi_0 + 2\Phi(\partial I) + S_{\text{vN}}(R \cup I)]$$



# Island contributions

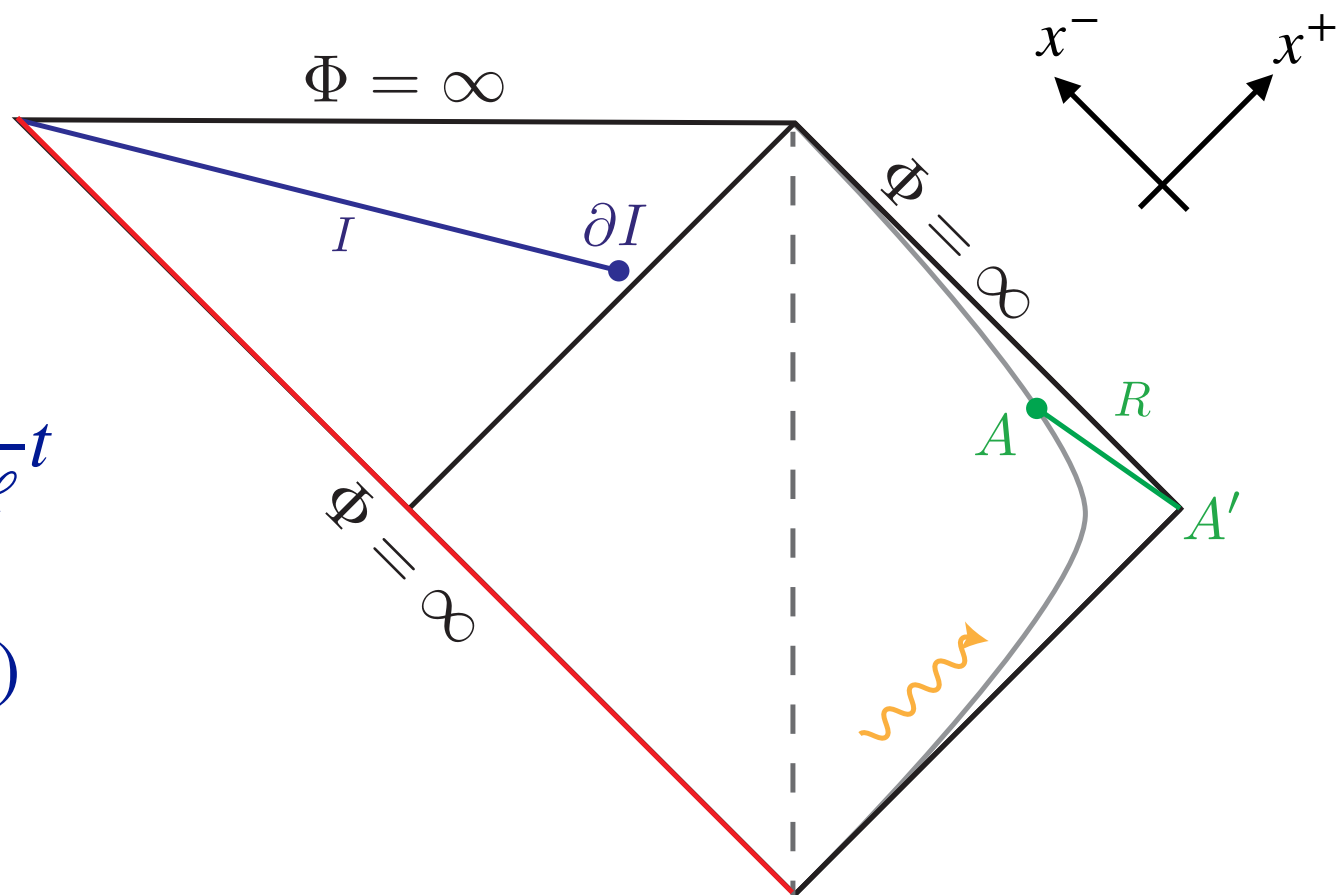
[LA, Sybesma]

**Extremizing** the fine-grained entropy, we find in addition to the trivial island an additional contribution:

Island location:  $x_{\partial I}^- \gtrsim 0$        $x_{\partial I}^+ \simeq \ell e^{1+2\Phi_s/c} + 2\ell e^{t/\ell}$

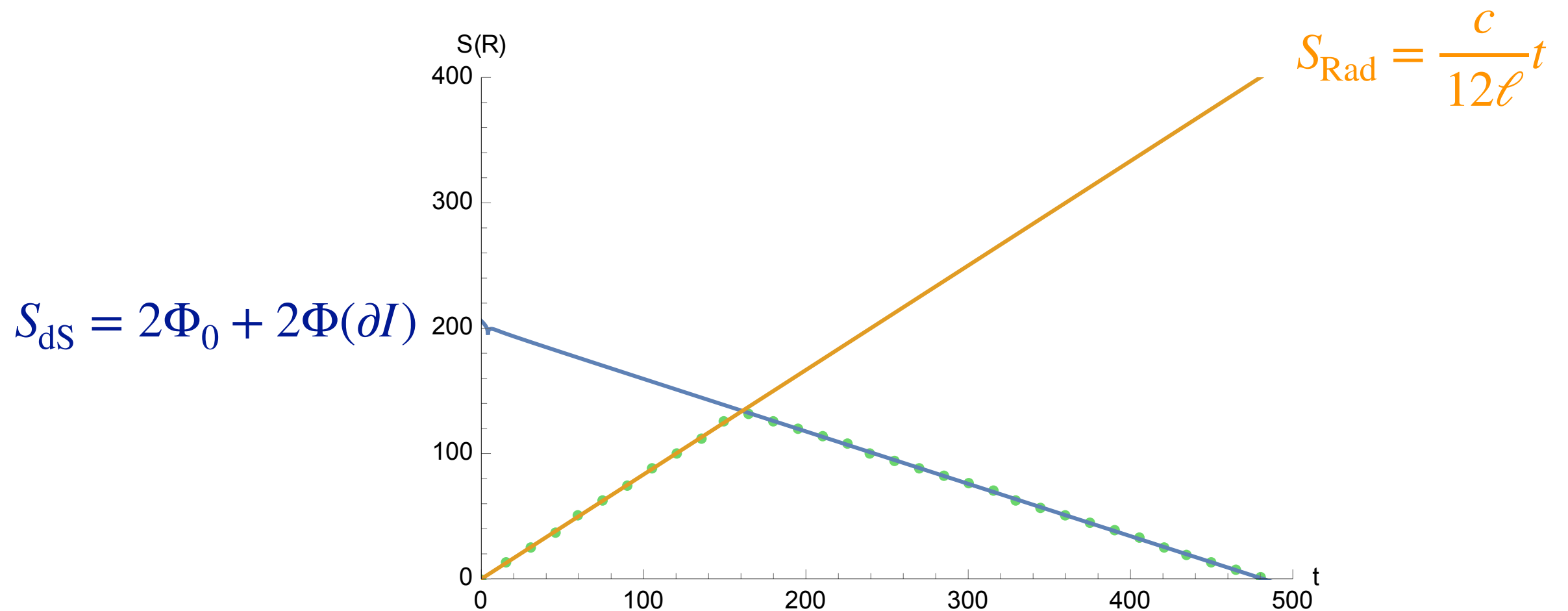
Trivial island:  $S(R) \simeq S_{\text{vN}}(R) = \frac{c}{12\ell} t$

Non-trivial island:  $S(R) \simeq 2\Phi_0 + 2\Phi(\partial I)$



# Page curve for dS

Taking the **minimum** of both contributions, we find a **Page curve** for the radiation in de Sitter space.



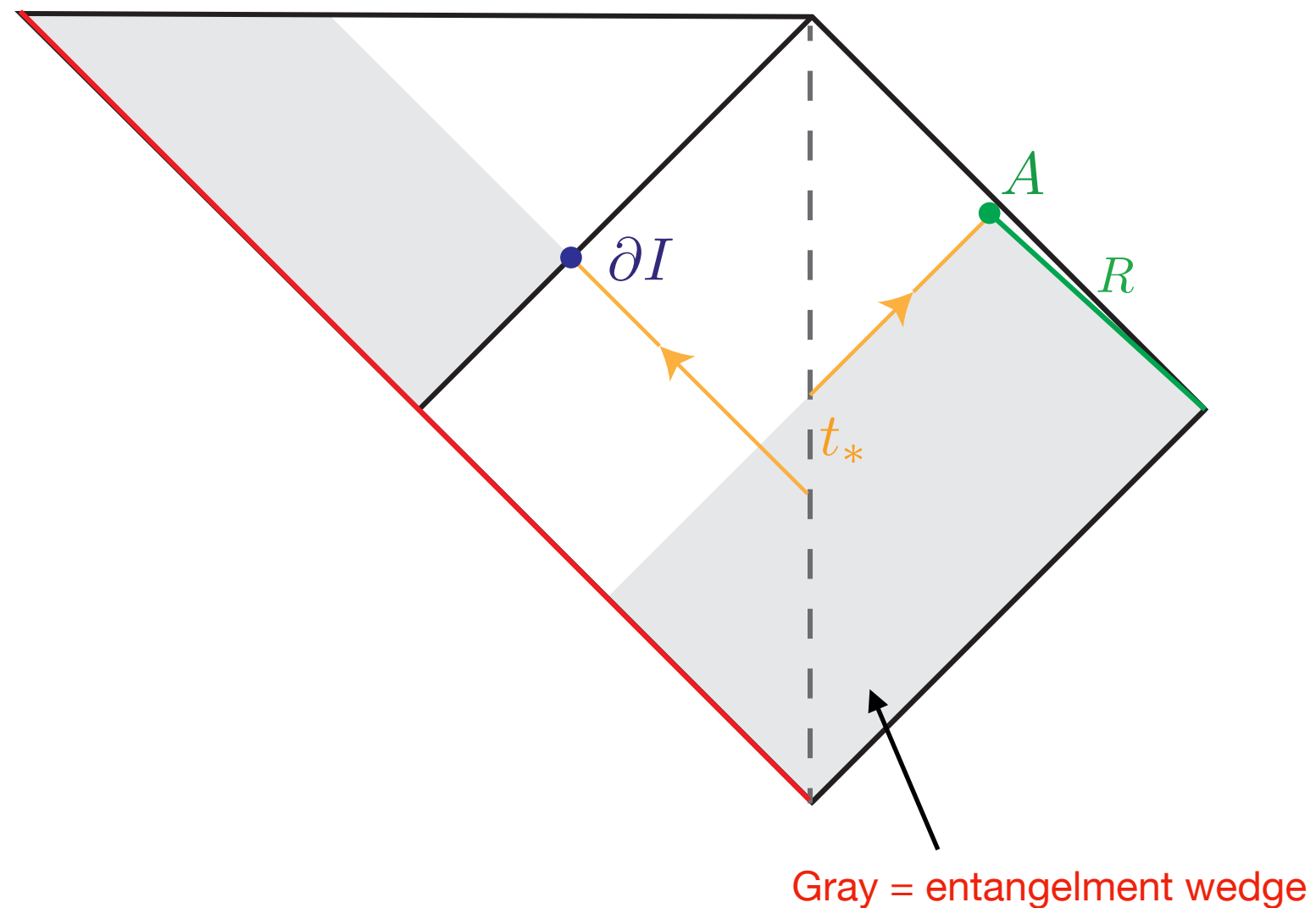
The Page time is given by:  $t_{\text{Page}} = \frac{16\ell\Phi_0}{c}$ .

# Information recovery in dS

After the Page time, information recovery becomes possible in a **scrambling time**.

Compute time difference between a lightray reaching the **island** and arrival in **R**:

$$t_* = \ell \log(S_{\text{dS}})$$



Agrees with computation using **OTOCs** [LA, Shiu '20] [Blommaert '20] .

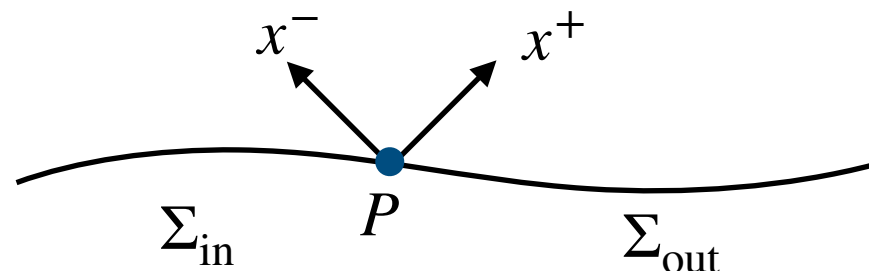
# Fate of the observer

Information recovery is possible, but what is the **fate of the observer**?

- Because energy is continuously collected in a finite volume, **backreaction** eventually becomes large.
- Does a **singularity** form?

To assess, use **quantum singularity theorem**: [Wall '10] (see also [Freivogel, Kontou, Krommydas '20])

**Quantum Singularity Theorem:**  
a singularity is unavoidable when a quantum trapped region forms.



$$S(P) = \frac{\text{Area}(\mathbf{P})}{4G_N} + S_{\text{vN}}(\Sigma_{\text{out}})$$

Quantum Trapped:  $\partial_{\pm} S(P) < 0$



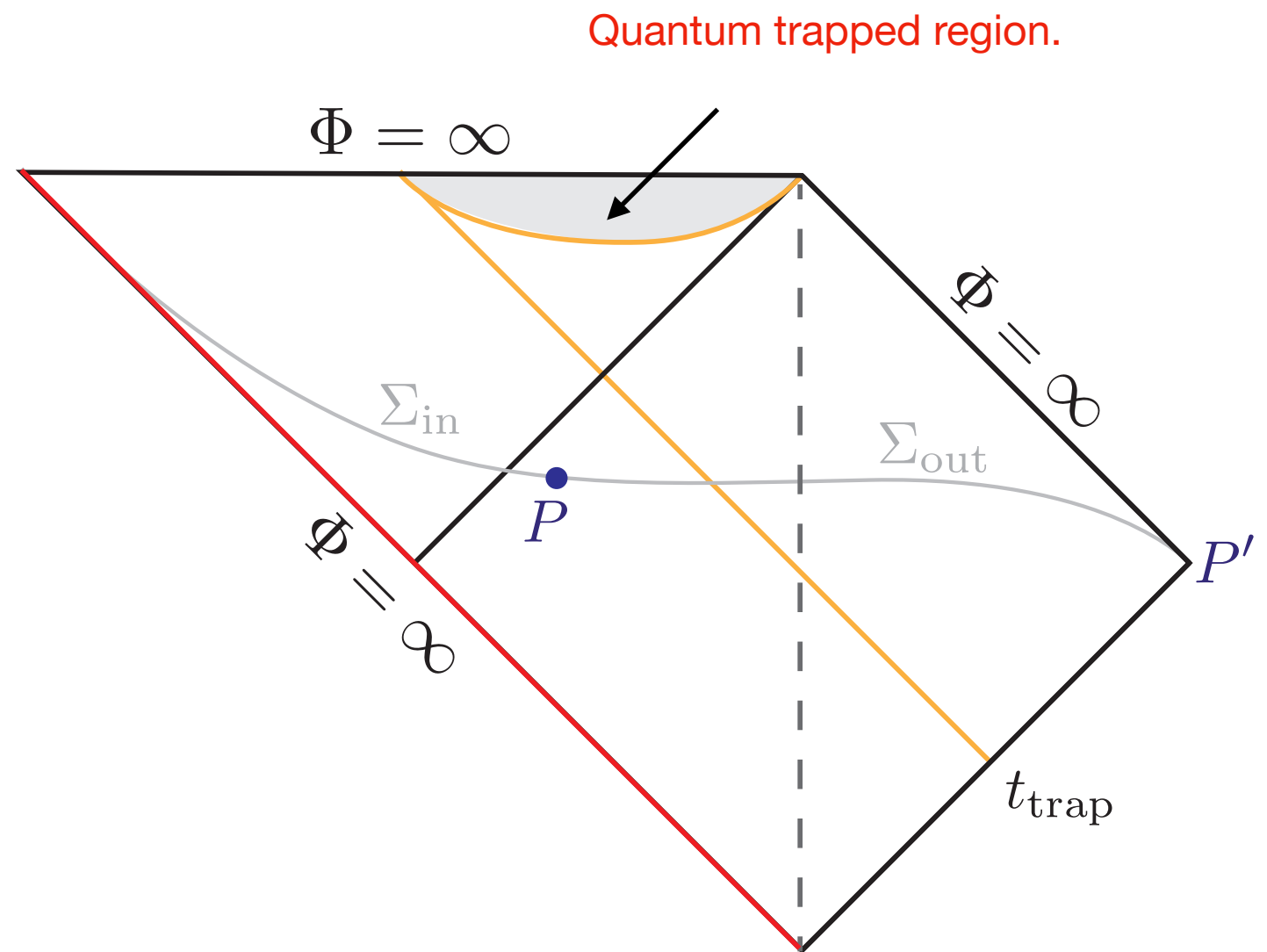
# Quantum singularity theorem

We now look for **quantum trapped regions** in the non-equilibrium state.

Solving for  $\partial_{\pm} S = 0$  we find a quantum trapped region after:

$$t > t_{\text{trap}} = \frac{2\Phi_s \ell}{c}$$

Hierarchy:  $t_* \ll t_{\text{trap}} \simeq t_{\text{Island}} \ll t_{\text{Page}}$



After  $t_{\text{trap}}$  the formulation of a singularity will be **unavoidable**.

# Conclusions

- Using recent developments in quantum information, we assessed if the static patch of de Sitter exhibits an **information paradox**.
- In JT gravity, we found that **quantum extremal islands** contribute to the fine-grained entropy.
- Gives an interpretation of the **de Sitter entropy** with respect to a single static observer.
- In de Sitter, curiosity comes at a price: when an island appears, formation of a **singularity is unavoidable**.
- Relation to **inflationary physics**?

# Thank you!

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