

# CICY orientifolds and GV invariants

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# Based on...

- 1 A landscape of orientifold vacua
  - F.C, J. Moritz, A. Westphal. 2020
- 2 Gopakumar-Vafa hierarchies in winding inflation and uplifts
  - F.C, A. Mininno, N. Righi, A. Westphal. 2021

# Motivation

- One of the main goals of String Phenomenology is to make concrete string models which:
  - 1 Reproduce the established low-energy physics (SM,  $\Lambda$ CDM...)
  - 2 Extend the established low-energy physics (DM, SUSY, Inflation...)
- Type II string theory on  $\mathbb{R}^{1,3} \times X_6$ , with  $X_6$  compact CY 3-fold.
- 4d EFT is  $\mathcal{N} = 2$  SUGRA. More realistic  $\mathcal{N} = 1$  EFTs arise from modding out by an orientifold action  $\Omega\mathcal{R}(-1)^{F_L}$
- Focus on IIB, with orientifold action allowing for  $O3/O7$  planes. Typical setup for the flux landscape (GKP '01. KKLT '03)
- Properties of the CY orientifold fix properties of the low energy 4d EFT.

# Even and odd cohomology

The orientifold action induces a splitting in cohomology

$$H^{p,q}(X) = H_+^{p,q}(X) \oplus H_-^{p,q}(X)$$

$h_{\pm}^{p,q}$  counts the number of fields in the 4d EFT (Grimm, Louis '04).

- $h_-^{2,1}$  complex structure moduli. Stabilized by fluxes.
- $h_-^{1,1}$   $C_2$  and  $B_2$  axions. Usually light (DM? Inflation?)
- $h_+^{2,1}$  U(1) vector multiplets. (Dark photons?)
- $h_+^{1,1}$  kahler moduli. Stabilized by non-perturbative effects.

Few orientifolds with  $h_-^{1,1} \neq 0$  appear in the literature (Gao, Shukla '13). This motivated us to systematically look for such cases, on a large scale search.

# Complete intersection CYs

- Defined as the zero locus of a set of homogeneous polynomials  $p_i[z]$ , ( $i = 1, \dots, K$ ) in an ambient space  $\mathcal{A} = \mathbb{P}^{n_1} \times \dots \mathbb{P}^{n_s}$
- Multidegrees of  $p_i[z]$  encoded in the configuration matrix.

$x^i$	$\mathbb{P}^2$	0	2	0	1
$y^i$	$\mathbb{P}^2$	2	1	0	0
$w^i$	$\mathbb{P}^3$	1	1	1	1

$$\sum_{i=1}^s n_i - K = 3$$

$$\begin{aligned} p_1[z] &= a_{ijk} y^i y^j w^k, \\ p_2[z] &= b_{ijkl} x^i x^j y^k w^l, \\ p_3[z] &= c_i w^i, \\ p_4[z] &= d_{ij} x^i w^j. \end{aligned} \quad (1)$$

- List of at most 7890 (possibly) distinct CY constructed in this way.  
(Candelas, Dale, Lutken, Schimmrigk '87) (Green, Hubsch, Lutken '89)  
(Anderson, Gao, Gray, Lee '17)

# Looking for geometric involutions $\mathcal{R}$

- Consider all ambient space involutions
  - 1 Swap of identical  $\mathbb{P}^n$  factors
  - 2 Invert some coordinates ( $z \rightarrow -z$ ) of non-swapped  $\mathbb{P}^n$  factors.
- Consider all the ways in which the  $p_i[z]$  can map into each other.
  - 1  $p_i[z] \rightarrow \pm p_i[z]$
  - 2  $p_i[z] \rightarrow p_j[z]$

Most choices give a singular manifold. Many others are equivalent to each other. Ex. Take  $\mathbb{P}^2$  and consider the two  $\mathbb{Z}_2$  actions

$$[z_0, z_1, z_2] \rightarrow [-z_0, z_1, z_2] \quad (2)$$

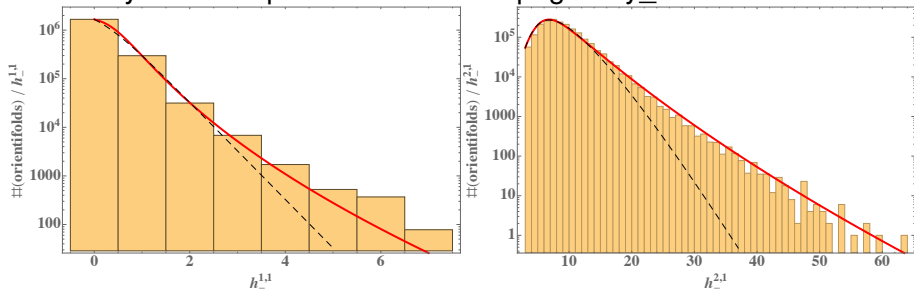
$$[z_0, z_1, z_2] \rightarrow [z_0, -z_1, -z_2] \quad (3)$$

They are equivalent because of the usual  $\mathbb{P}^2$  scaling.

Write a code to list all inequivalent involutions.

# Results

Explicit database of **2.004.513 orientifolds**. List available at [www.desy.de/~westphal/orientifold\\_webpage/cicy\\_orientifolds.html](http://www.desy.de/~westphal/orientifold_webpage/cicy_orientifolds.html)



Histograms for the number of orientifolds with non-zero  $h_{-}^{1,1}$  and  $h_{-}^{2,1}$ .  $h_{-}^{1,1}$  appears to be quite generic! Good for model building with  $B_2$  and  $C_2$  axions!

# Singularities at codimension 3: frozen conifolds

- Singularities at codimension 1, and codimension 3.
- We identify and drop from the database the cases of singularities of codimension 1. Non-trivial implementation for this.
- Codimension 3 singularities are interesting: conifold points that lie on the fixed surfaces, i.e. conifold points on top of the O-planes.
- Let  $f_-^i$  be a set of antisymmetric polynomials across the fixed locus, whenever  $df_-^1[z] \wedge df_-^2[z] \wedge \dots = 0$  we have such singularities.
- The orientifold projects out the deformation branch.

$$\det \begin{pmatrix} x & v \\ u & y \end{pmatrix} = \epsilon \quad (4)$$

$(v, y) \rightarrow (-v, -y)$  is a symmetry only for  $\epsilon = 0$ .



# Resolving frozen conifolds

- Let  $[\alpha, \beta]$  be the homogeneous coordinates of the resolution  $\mathbb{P}^1$ .
- **A-type resolution** (locally)

$$\begin{pmatrix} x & v \\ u & y \end{pmatrix} \begin{pmatrix} \alpha \\ \beta \end{pmatrix} = 0 \quad (5)$$

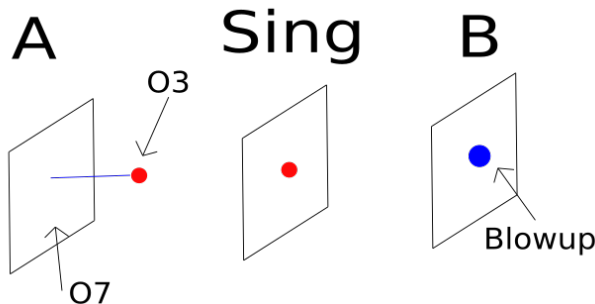
- **B-type resolution** (locally)

$$\begin{pmatrix} x & u \\ v & y \end{pmatrix} \begin{pmatrix} \alpha \\ \beta \end{pmatrix} = 0 \quad (6)$$

- A-type resolution. Action of the orientifold  $[\alpha, \beta] \rightarrow [-\alpha, \beta]$ . Two fixed points:  $[1, 0]$ ,  $[0, 1]$ . Now,  $[1, 0]$  lies on the fixed divisor, while  $[0, 1]$  is isolated fixed point. This is an  $O7$ -plane on  $\mathbb{C}^2$ , plus a distant  $O3$ .
- B-type resolution. Action of the orientifold  $[\alpha, \beta] \rightarrow [\alpha, \beta]$ . Whole resolution  $\mathbb{P}^1$  is fixed. This is an  $O7$ -plane on  $\mathbb{C}^2$  blown up at a point.

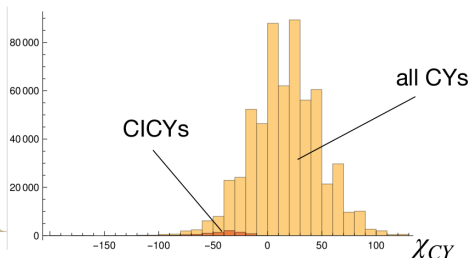
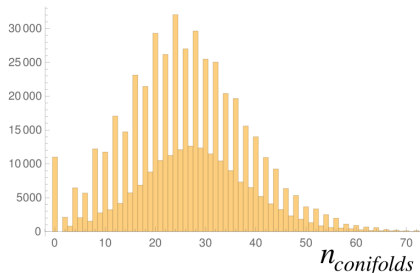
# Frozen conifold transitions

- A distant  $O3$  approaches the  $O7$ . We are in the resolved  $A$ -side.
- The  $O3$  on top of the  $O7$ . Singular frozen conifold point on the  $O7$ .
- The  $O7$  eats the  $O3$ , converting it into geometry. Now the topology of the fixed divisor wrapped by the  $O7$ s, there is one blowup. We are in the resolved B-side.
- Notice it is consistent with modification of  $D3$  tadpole.



# Frozen conifold transitions

- One can access the resolution (Higgs) branch, in two ways.
- How generic? Almost all orientifolds have non-deformable conifolds, whose resolution branches are often outside the CICY set.



- Notice the  $\chi$  plot is now symmetric. Are we connecting CICYs to mirror CICYs in such a way? Or maybe a different and new class of CYs?

# CICY redundancies

- Some of the configuration matrices in the original CICY list define the same CY
- **Wall's Theorem.** If two CY  $X$  and  $Y$  are such that
  - 1  $h^{1,1}(X) = h^{1,1}(Y)$
  - 2  $h^{2,1}(X) = h^{2,1}(Y)$
  - 3 There exist a choice of bases for  $H^2(X, \mathbb{Z})$  and  $H^2(Y, \mathbb{Z})$  such that  $\int_{D_i} c_2(X) = \int_{D_i} c_2(Y)$ ,  $\forall i = 1, \dots, h^{1,1}$
  - 4 In the same base as above,  $K_{ijk}(X) = K_{ijk}(Y)$then the two CY are diffeomorphic as real manifolds.
- Idea: pick two CICYs with same  $h^{1,1}$  and  $h^{2,1}$  and naively different  $\int_{D_i} c_2$  and  $K_{ijk}$ . Find a change of basis to make these quantities equal, or show it can't exist.
- Write a code for this. Find and list 1169 redundancies. Extending the results of (Candelas, He, '90) and (Anderson, He, Lukas, '08)

# Genus zero GV invariants

- Topological invariants of the CY 3-fold  $X$ . (Gopakumar, Vafa '98)
- Count the number of holomorphic maps from a Riemann surface  $\Sigma_g$  to a 2-cycle class  $\beta \in H_2(X)$
- Good to know the value, for pheno model building. I.e. Kahler potential for cs moduli  $z^i$  at LCS is

$$\begin{aligned} K = & -\ln \left( -\frac{4}{3} k_{ijk} \operatorname{Im}(z^i) \operatorname{Im}(z^j) \operatorname{Im}(z^k) + ic \right. \\ & - 2 \sum_{\beta}^{\infty} n_{\beta}^0 \left( \operatorname{Li}_3(e^{i\beta_i z^i}) + \operatorname{Li}_3(e^{-i\beta_i \bar{z}^i}) \right) + \\ & \left. - 2 \sum_{\beta}^{\infty} n_{\beta}^0 \beta_i \operatorname{Im}(z^i) \left( \operatorname{Li}_2(e^{i\beta_i z^i}) + \operatorname{Li}_2(e^{-i\beta_i \bar{z}^i}) \right) \right) \end{aligned} \quad (7)$$

# Computing the genus 0 GV by Mirror symmetry

(Hosono, Klemm, Theisen, Yau, '94)

- $w_0(z) = \sum_{n_1 \geq 0} \cdots \sum_{n_{h^2,1} \geq 0} c(n) \prod_{i=1}^{h^2,1} z_i^{n_i}.$

Generic solution of the PF equation for the first entry of the period vector, in terms of data in the configuration matrix.

- $\Pi(z) = \left( w_0(z), \frac{\partial}{\partial \rho_i} w_0(z, \rho) \big|_{\rho=0}, \dots \right)^t$

- $w_i(z) = \sum_{n_1 \geq 1} \cdots \sum_{n_{h^2,1} \geq 0} \frac{1}{2\pi i} \frac{\partial}{\partial \rho^i} c(n + \rho) \big|_{\rho=0} \prod_{i=1}^{h^2,1} z_i^{n_i} + w_0(z) \frac{\ln z_i}{2\pi i}$

- $t^i = \frac{w_i(z)}{w_0(z)}$  mirror map. Relation between Kahler moduli at large radius in the A-model side, with complex structure moduli at LCS in the B-model side.

- Now, invert the mirror map in order to get  $z(t)$ . Hardest step.

# Computing the genus 0 GV by Mirror symmetry

(Hosono, Klemm, Theisen, Yau, 94)

- Having  $z(t)$ , compute the quantum corrected triple intersection number

$$k_{ijk} = \frac{\partial}{\partial t_i} \frac{\partial}{\partial t_j} \frac{\frac{1}{2} k_{kab}^0 \frac{\partial}{\partial \rho_a} \frac{\partial}{\partial \rho_a} w_0(z, \rho) |_{\rho=0}}{w_0(z)}(t)$$

- Introduce  $q_i = \exp(2\pi i t_i)$ , and the general expression for the quantum corrected triple intersection number

$$k_{ijk} = k_{ijk}^0 + \sum_{n_1 \geq 1} \dots \sum_{n_{h^1,1} \geq 0} n_{d_1, \dots, d_{h^1,1}} d_i d_j d_k \frac{\prod_{l=1}^{\bar{h}^{1,1}} q_l^{d_l}}{1 - \prod_{l=1}^{\bar{h}^{1,1}} q_l^{d_l}}$$

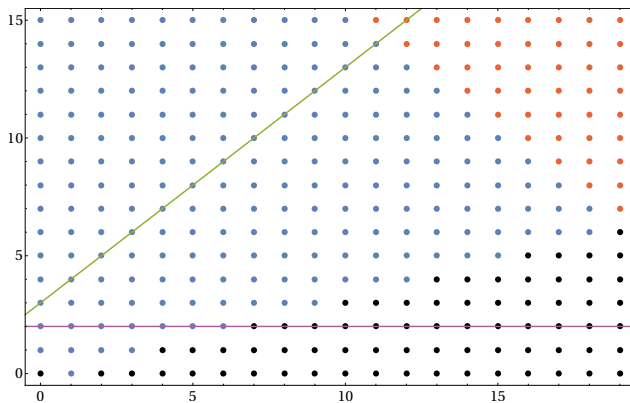
- Match and extract the GV.

# Instanton program, and the scan

- The above algorithm has been coded in a program called Instanton. (Klemm, Kreuzer)
- Made modification of their code, parallelized it, and let it run.
- Computing time  $\approx 6$  months, on two different clusters (DESY and Madrid IFT)
- List all the genus 0 GV, up to total degree 10, for all favourable CICY up to  $h^{1,1} = 9$ .
- We find directions in the Picard lattice in which the GV invariants grow at hierarchical rates, as well as "vanishing directions" (Demirtas, Kim, McAllister, Moritz '20) and "periodic directions".
- List at [www.desy.de/~westphal/GV\\_CICY\\_webpage/GVInvariants.html](http://www.desy.de/~westphal/GV_CICY_webpage/GVInvariants.html)



# Occupation sites for CICY 7858 ( $h^{1,1} = 2$ )



**Figure:** Blue=non-zero GV. Black = zero GV. Orange = not computed, believed to be non-zero GV. Green = non-vanishing direction. Purple = vanishing direction.

# Conclusions

- 1 We compute a database of all the  $O3/O7$  orientifolds of the CICYs that descend from the ambient space involutions. We compute quantities as the  $D3$  tadpole, the number of frozen conifolds, etc.
- 2 We observe the phenomenon of frozen conifold transitions, connecting the CICYs to a larger set of (potentially new) CYs.
- 3 We compute a database of all the genus 0 GV invariants for CICYs up to  $h^{1,1} = 9$  and total degree 10.
- 4 We observe the presence of directions in the Picard lattice where the GV growth is hierarchically dominant/suppressed with respect to other directions, as well as periodic directions.
- 5 We identify 1169 redundancies in the CICY list.

# Future directions

- Would be nice to do similar scans for Kreuzer-Skarke.
- Are frozen conifold transitions the generic  $\mathcal{N} = 1$  geometric transitions? Or are they special to CICYs? Can we see them in Kreuzer-Skarke?
- Compute the topology of the divisors for the CICY orientifolds (important to know which 4-cycles will support D7s for gaugino condensation or ED3)
- Suppose you realize a given CICY by a different, much larger configuration matrix. Do you find more orientifolds?
- Would be nice to do the orientifold scan for the CICYs, for IIA orientifolds. Much harder as the orientifold action is not now-holomorphic. i.e. fixed loci at real codimension 3.