

## Quantum Mechanics: Worksheet 6

*Klein-Gordon equation and the Klein paradox*

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### 1 Klein-Gordon equation

The Klein-Gordon equation is given by

$$\left( \partial^2 + \frac{m^2 c^2}{\hbar^2} \right) \phi(x) = 0,$$

where  $\phi(x)$  is a scalar function of the space-time coordinates  $x^\mu = (ct, \mathbf{r})$ . Here  $\partial^2$  is the d'Alembertian,

$$\partial^2 \equiv \partial_\mu \partial^\mu = \frac{1}{c^2} \frac{\partial^2}{\partial t^2} - \nabla^2,$$

with  $\partial^\mu = \eta^{\mu\nu} \partial_\nu = (c^{-1} \partial_t, -\nabla)$ . For stationary solutions  $\phi(x) = e^{-\frac{i}{\hbar} Et} \varphi(\mathbf{r})$ , the time-independent Klein-Gordon equation becomes

$$\left( -\frac{\hbar^2}{2m} \nabla^2 - \frac{E^2 - m^2 c^4}{2mc^2} \right) \varphi(\mathbf{r}) = 0.$$

External electromagnetic fields are included in the same way as for the Schrödinger equation, with minimal coupling:  $p_\mu \rightarrow p_\mu - \frac{q}{c} A_\mu$  (in Gaussian units).

The Klein-Gordon equation has a conserved U(1) charge with a conserved four-current

$$j^\mu = \frac{i\hbar}{2m} (\phi^* \partial^\mu \phi - \phi \partial^\mu \phi^*) = (c\rho, \mathbf{j}),$$

where the prefactor is chosen such that it matches the Schrödinger current for positive energy states. Show that  $\partial_\mu j^\mu = 0$  and  $j^\mu = (\hbar/m) k^\mu$  for a plane wave. Also calculate  $j^\mu$  for a general stationary solution. Also show that

$$\frac{dQ}{dt} = - \oint_{\partial\mathcal{V}} \mathbf{j} \cdot d\mathbf{a}, \quad Q = \int_{\mathcal{V}} \rho d\tau,$$

with  $Q$  the total “charge” in a volume  $\mathcal{V}$  with boundary  $\partial\mathcal{V}$ .

Why does the Klein-Gordon equation not admit a consistent interpretation in terms of a relativistic single-particle wave equation?

### 2 Klein paradox

The Klein paradox illustrates that, unlike the Dirac equation, the Klein-Gordon equation cannot be consistently interpreted as a relativistic single-particle equation. To this end, consider scattering at a potential step:

$$V(x) = V_0 \theta(x),$$

with  $V_0 \geq 0$ . This potential can be introduced in the time-independent Klein-Gordon equation by substituting  $E \rightarrow E - V(x)$ . This problem has translational symmetry in the  $yz$  plane. For simplicity, we consider the case where  $k_y = k_z = 0$  (normal incidence).

- (a) Write down the general solution  $\varphi^<(x)$  for  $x < 0$  and  $\varphi^>(x)$  for  $x > 0$ .

- (b) Now consider a solution for an incoming plane wave for  $x < 0$  with positive energy  $E > mc^2$  and wave vector  $k > 0$ . Call the amplitude of the reflected mode  $r$  and of the transmitted mode  $t$ . To ensure causality, the transmitted mode always needs to have positive group velocity. Hint: make a drawing of the energy dispersion in both regions and consider three cases:  $V_0 < E - mc^2$ ,  $E - mc^2 < V_0 < E + mc^2$ , and  $V_0 > E + mc^2$ .
- (c) Calculate the Klein-Gordon current for  $x < 0$  and  $x > 0$ . Show that  $\partial_x j = 0$ . Now separate the current in incoming ( $j_i$ ), reflected ( $j_r$ ), and transmitted ( $j_t$ ) parts.
- (d) From the previous question, we find that

$$j_i + j_r = j_t.$$

Use the continuity of  $\phi$  and  $\partial_x \phi$  to find an expression for the coefficients  $r$  and  $t$ . Check that your solution satisfies current conservation.

- (e) Show that  $j_t = 0$  for  $E - mc^2 < V_0 < E + mc^2$ ? Hint: write down the wave vector  $q$  for  $x > 0$  and show that  $q$  is imaginary in this case.
- (f) Consider the case  $V_0 > E + mc^2$ . Show that  $j_t < 0$  and thus  $j_r < -j_i$ . This is the Klein paradox. Argue why there is no paradox if we interpret  $j$  as a charge current. Namely, explain when  $j_t < 0$  is consistent with a right-moving wave with density  $\rho$ .