GREEN'S FUNCTIONS

A short introduction

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- This is intended as a quick overview of Green's functions for electrical engineers.
- Green's functions are a huge subject: it's easy to get overwhelmed by calculation techniques.
- Focus here will be on intuition/understanding and awareness of some key techniques.
- · Lots of further reading provided at the end.
- Tip: read a lot of different references. Different authors take totally different approaches and it's interesting to see them all.

OUTLINE

- 1 Introduction
- 2 Generalized functions
- 3 Direct solution
- 4 Boundary conditions
- 5 Spectral methods
- 6 3D problems
- 7 Properties of the Green's function
- 8 Advanced topics

INTRODUCTION

WHAT IS A GREEN'S FUNCTION?

Linear equation to solve:

$$\mathcal{L}u(x) = f(x)$$

Green's function is the **impulse response**:

$$\mathcal{L}G(x,x')=\delta(x-x')$$

- Fortunately, the basic idea of Green's functions is really simple.
- You've actually used them before!
- What's far more interesting is how to calculate/interpret them.

- Most EM problems are described by linear (differential) equations with some source/driving function f(x).
- The Green's function is the solution when the source f(x) is an impulse located at x'.
- Can think of it as a generalization of the impulse response from signal processing.

$$\delta(X-X') \xrightarrow{\mathcal{L}^{-1}} G(X,X')$$

$$f(x) = \int \delta(x - x') f(x') dx \xrightarrow{\mathcal{L}^{-1}} \int G(x, x') f(x') dx$$

 $\mathcal{L}u(x) = f(x)$

$$\mathcal{L}G(x,x')=\delta(x-x')$$

$$u(x) = \int G(x, x') f(x') \, \mathrm{d} x$$

4

- Once we know the Green's function for a problem, we can find the solution for any source f(x).
- Impulses $\delta(x x')$ produce a response G(x, x').
- We can split the source f(x) up into a sum (integral) of impulses $\delta(x-x')$.
- Then the response to f(x) is just a weighted sum (integral) of impulse responses.

• Once we know the Green's function, we have an explicit formula for the solution u(x) for any source function f(x).

FAMILIAR GREEN'S FUNCTIONS

Impulse response of a LTI system:

$$y(t) = \int_{-\infty}^{\infty} x(t')h(t-t') dt'$$

E.g., for an RL-circuit:

$$G(t, t') = h(t - t') = u(t - t')e^{-\alpha(t - t')}$$

- In electrical engineering, we've seen Green's functions before.
- Impulse response h(t-t') from linear system theory is an example of a Green's function.

$$G(t,t')=h(t-t')$$

• Usually find h(t-t') using Fourier transform of the transfer function.

FAMILIAR GREEN'S FUNCTIONS

Poisson's equation:

$$\nabla^2 V(\mathbf{r}) = -\frac{\rho(\mathbf{r})}{\epsilon_0}$$

$$V(\mathbf{r}) = \iiint \frac{1}{4\pi\epsilon_0 |\mathbf{r} - \mathbf{r}'|^2} \rho(\mathbf{r}') d^3 \mathbf{r}'$$

· Green's function for Poisson's equation is

$$G(\mathbf{r},\mathbf{r}') = \frac{1}{4\pi\epsilon_0 \left|\mathbf{r} - \mathbf{r}'\right|^2}$$

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FAMILIAR GREEN'S FUNCTIONS

FAMILIAR GREEN'S FUNCTIONS

Helmholtz equation:

$$(\nabla^2 + k^2) A_z(\mathbf{r}) = -J_z(\mathbf{r})$$

$$A_z(\mathbf{r}) = \iiint \frac{e^{-jk|\mathbf{r}-\mathbf{r}'|}}{4\pi |\mathbf{r}-\mathbf{r}'|} J_z(\mathbf{r}') d^3 \mathbf{r}'$$

Our goal:

- · Derive these expressions.
- · Generalize to other problems and boundary conditions.

8

• Green's function for the Helmholtz equation is

$$G(\mathbf{r}, \mathbf{r}') = \frac{e^{-jk|\mathbf{r} - \mathbf{r}'|}}{4\pi |\mathbf{r} - \mathbf{r}'|}$$

GENERALIZED FUNCTIONS

TYPICAL DELTA FUNCTION DEFINITION

Typical "definition" of $\delta(x - x_0)$:

$$\delta(x - x_0) = 0 \quad \text{for} \quad x \neq x_0$$

$$\int_{-\infty}^{\infty} \delta(x - x_0) = 1$$

- Delta functions play a key role in Green's functions (and electrical engineering in general), but tend to lead to hand-waving.
- Worth seeing how they can be rigorously defined before moving on.
- Machinery for this is Schwartz's theory of distributions (generalized functions).
- See Folland (1992), Fourier analysis and its applications, Chapter 9 for more.

- · Often see definitions like this one.
- Often said to imply that $\delta(x x_0) = \infty$ at $x = x_0$.
- Might be okay intuitively, but very imprecise mathematically.
- There is no true function which satisfies both of these requirements!

f(x) defines a linear operator $\phi(x)$ via

$$f[\phi] = \int_{-\infty}^{\infty} f(x)\phi(x) \, \mathrm{d} x$$

If we have $f[\phi]$, but no f(x), then f is a generalized function.

Symbolically, we write

$$f[\phi] \stackrel{s}{=} \int_{-\infty}^{\infty} f(x)\phi(x) dx$$

13

12

- Let's see if we can generalize the idea of a "function" so that it includes delta functions.
- Given a function f(x), we can use it to define a linear operator (a functional, to be exact) on other functions $\phi(x)$.
- $f[\cdot]$ is a linear operator. It takes a function $\phi(x)$ and returns the number

$$f[\phi] = \int_{-\infty}^{\infty} f(x)\phi(x) \, \mathrm{d} x$$

• If we ensure that $\phi(x)$ is very well-behaved, then every function f(x) defines an operator in this way.

- It's possible to have an operator $f[\phi]$, but we can't find an f(x) to implement it via an integral.
- Then f(x) is a generalized function. It is not a function in its own right, but it is defined purely by its action on other functions $f[\phi]$.
- · We still symbolically write

$$f[\phi] \stackrel{s}{=} \int_{-\infty}^{\infty} f(x)\phi(x) \,\mathrm{d}x$$

but this just suggestive notation. It is not actually an integral unless f(x) is a "proper" function!

 $\delta(x-x_0)$ is a generalized function defined by the sifting property

$$\delta_{x_0}[\phi] = \phi(x_0) \stackrel{s}{=} \int_{-\infty}^{\infty} \delta(x - x_0) \phi(x) \, \mathrm{d}x$$

We can define derivatives too:

$$\delta_{x_0}^{(n)}[\phi] = (-1)^n \phi^{(n)}(x_0) \stackrel{s}{=} \int_{-\infty}^{\infty} \delta^{(n)}(x - x_0) \phi(x) \, dx$$

14

- We can define a simple linear operator via the sifting property $\delta_{x_0}[\phi] = \phi(x_0)$.
- There is no actual function $\delta(x x_0)$ which gives

$$\int_{-\infty}^{\infty} \delta(x - x_0) \phi(x) \, \mathrm{d} \, x = \phi(x_0)$$

so $\delta(x-x_0)$ is a generalized function and the above integral is purely symbolic.

- Generalized function theory lets us make sense of the derivatives of the delta function too.
- $\delta_{x_0}^{(n)}$ is just an operator that picks out the value of the *n*th derivative of $\phi(x)$ at the point x_0 .

 $\lim_{\epsilon \to 0} f_{\epsilon}(x) = \delta(x)$

if and only if

$$\lim_{\epsilon \to 0} f_{\epsilon}[\phi] = \lim_{\epsilon \to 0} \int_{-\infty}^{\infty} f_{\epsilon}(x) \phi(x) \, \mathrm{d} \, x = \phi(0)$$

Limit of Gaussian functions:

$$\delta(x) = \lim_{\epsilon \to 0} \frac{1}{\sqrt{2\pi\epsilon}} e^{-x^2/2\epsilon^2}$$

Limit of Lorentzian functions:

$$\delta(x) = \lim_{\epsilon \to 0} \frac{1}{\pi} \frac{\epsilon}{t^2 + \epsilon^2}$$

16

- Often useful to show that some set of actual functions $f_{\epsilon}(x)$ "approach" the delta function in a limit.
- To do this, we need to show that the sifting property is obeyed in the limit.

- Two examples of delta function limits.
- Confirms our intuition of the delta function as a limit of sharply-peaked functions.
- In fact, basically any limit of sharply-peaked functions of area 1 will work: see Folland (Theorem 9.2).

A more interesting example:

$$\frac{1}{2\pi}\int\limits_{-\infty}^{\infty}e^{jxt}\,\mathrm{d}\,t=\delta(x)$$

because

$$\lim_{\epsilon \to 0} \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{-\epsilon^2 t^2} e^{jxt} dt = \delta(x)$$

 $\mathcal{L}G(x,x') = \delta(x-x')$

actually means

$$(\mathcal{L}G)[\phi] = \phi(X') \stackrel{s}{=} \int_{-\infty}^{\infty} (\mathcal{L}G(X,X')) \phi(X) dX$$

18

- Example of a common, but unintuitive expression for the Delta function.
- Can show that it's true by expressing it as a delta function limit. (If you want to go through it, use Theorem 9.2 from Folland.)

- Technically, the Green's function is a generalized function such that $\mathcal{L}G$ is the delta function (it has the sifting property).
- In practise, we'll keep using the less-precise way; just remember that there is a more correct way.

TAKEAWAY

If in doubt, think of $\delta(x - x_0)$ as an operator, not a function!

DIRECT SOLUTION

20

- In practise, thinking of $\delta(x-x_0)$ as a function is usually fine. (We'll even do that for the rest of this presentation.)
- But if anything starts to seem fishy, it's good to remember that $\delta(x-x_0)$ is actually an operator, and not a function.

· Back to Green's functions!

A SIMPLE EXAMPLE

Original problem:

$$\frac{\mathrm{d}^2 u(x)}{\mathrm{d} x^2} - k^2 u(x) = f(x)$$

Green's function problem:

$$\frac{\mathrm{d}^2 G(x,x')}{\mathrm{d} x^2} - k^2 G(x,x') = \delta(x-x')$$

22

A SIMPLE EXAMPLE

For $x \neq x'$

$$\frac{d^2 G(x, x')}{d x^2} - k^2 G(x, x') = 0$$

So we have

$$G(x,x') = \begin{cases} Ae^{+k(x-x')} & \text{for } x < x' \\ Be^{-k(x-x')} & \text{for } x > x' \end{cases}$$

· Let's look at a simple example now.

- This problem is similar to a simple harmonic oscillator, but the negative sign means we expect lossy behaviour rather than oscillation.
- We won't worry much about boundary conditions yet, we'll just look for solutions that don't blow up at $x = \pm \infty$.
- If we can find the Green's function, then we can find the solution to the original problem.
- But the Green's function problem looks pretty hard. The point of this example is to demonstrate that we can actually solve it.

• Key thing to notice is that the source is concentrated at x = x'.

- So for x > x' and x < x', we expect the solutions to look like those of the source-free equation.
- To keep the solutions finite, we expect exponential growth before x = x' and exponential decay afterward.
- Now, how do we find the constants A and B?

$$\frac{\mathrm{d}^2 G(x,x')}{\mathrm{d} x^2} - k^2 G(x,x') = \delta(x-x')$$

Continuity of the Green's function:

$$\lim_{\epsilon \to 0} \left[G(X' + \epsilon, X') - G(X' - \epsilon, X') \right] = 0$$

24

- · How continuous do we expect our Green's function to be?
- If G(x,x') is discontinuous (like a step function), then dG/dx will behave like a delta function and d^2G/dx^2 will behave like a delta function derivative. No good!
- So we expect G(x, x') to be continuous.
- That gives us one condition we can use to find A and B. (In fact, it tells us that A = B.)

 $\int_{x'-\epsilon}^{x'+\epsilon} \left[\frac{d^2 G(x,x')}{d x^2} - k^2 G(x,x') \right] dx = \int_{x'-\epsilon}^{x'+\epsilon} \delta(x-x') dx$

Discontinuity condition:

$$\lim_{\epsilon \to 0} \left[\frac{dG}{dx} \Big|_{x=y'+\epsilon} - \frac{dG}{dx} \Big|_{x=y'-\epsilon} \right] = 1$$

- But what if the derivative dG/dx is discontinuous?
- Then $d^2 G/dx^2$ is like a delta function. But that's fine, because we have a delta function on the right hand side too.
- We can find exactly how discontinuous the derivative is by integrating over a small interval around x'.
- In the limit of $\epsilon \to 0$, the second integral vanishes because G(x,x') is continuous.
- But, we expect dG/dx to be discontinous.
- Using fundamental theorem of calculus, we get at a discontinuity condition for the derivative. (Key idea for the direct solution method!)

A SIMPLE EXAMPLE

$$G(x, x') = \begin{cases} Ae^{+k(x-x')} & \text{for } x < x' \\ Be^{-k(x-x')} & \text{for } x > x' \end{cases}$$

Continuity of *G*:

$$A = B$$

Discontinuity of $\frac{dG}{dx}$:

$$kA + kB = 1$$

A SIMPLE EXAMPLE

At last, our Green's function is

$$G(x, x') = \begin{cases} \frac{e^{+k(x-x')}}{2k} & \text{for } x < x' \\ \frac{e^{-k(x-x')}}{2k} & \text{for } x > x' \end{cases}$$

or, more compactly

$$G(x,x')=\frac{e^{k|x-x'|}}{2k}$$

26

• Applying our two conditions, we can solve for A and B.

Solution:

$$u(x) = \int_{-\infty}^{\infty} f(x') \frac{e^{k|x-x'|}}{2k} dx'$$

Properties of G(x, x'):

- Looks like the sourceless solution except at x = x'.
- Function is continuous at x = x'.
- Derivative is discontinuous at x = x'.

28

• Now that we have the Green's function, we can construct the solution to our original problem for any forcing function f(t).

- Listed are the key things to note from that example.
- This approach works quite well for solving 1D Green's function problems.

BOUNDARY CONDITIONS SPECTRAL METHODS

3D PROBLEMS

PROPERTIES OF THE GREEN'S FUNCTION

ADVANCED TOPICS

INTRODUCTORY RESOURCES

Balanis (2012), Advanced engineering electromagnetics. Less rigorous, but good for getting the key ideas.

Folland (1992), Fourier analysis and its applications. Chapter on generalized functions is particularly nice.

Dudley (1994), Mathematical foundations for electromagnetic theory. Great introduction to 1D Green's functions: deals with subtleties that others ignore.

Byron and Fuller (1992), *Mathematics of classical and quantum physics*. Interesting alternative approach.

ADVANCED RESOURCES

Collin (1990), Field theory of guided waves. Huge chapter on Green's functions. Emphasis on dyadics.

Morse and Feshback, *Methods of theoretical physics*. Another big, detailed reference. Emphasis on theory and insights.

Warnick (1996), "Electromagnetic Green functions using differential forms." For the differential forms inclined.