

GREEN'S FUNCTIONS

A SHORT INTRODUCTION

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OUTLINE

- 1 Basic idea
- 2 Finding the Green's function
- 3 Constructing the solution
- 4 Special properties
- 5 Conclusion

BASIC IDEA

WHAT IS A GREEN'S FUNCTION?

Linear equation to solve:

$$\mathcal{L}u(x) = f(x)$$

Green's function is the **impulse response**:

$$\mathcal{L}G(x, x') = \delta(x - x')$$

WHAT IS A GREEN'S FUNCTION?

Linear equation to solve:

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Green's function is the **impulse response**:

$$\mathcal{L}G(x, x') = \delta(x - x')$$

WHY IS IT USEFUL?

$$\delta(x - x') \xrightarrow{\mathcal{L}^{-1}} G(x, x')$$

$$f(x) = \int \delta(x - x') f(x') dx \xrightarrow{\mathcal{L}^{-1}} \int G(x, x') f(x') dx$$

WHY IS IT USEFUL?

$$\delta(x - x') \xrightarrow{\mathcal{L}^{-1}} G(x, x')$$

$$f(x) = \int \delta(x - x') f(x') \, dx \xrightarrow{\mathcal{L}^{-1}} \int G(x, x') f(x') \, dx$$

WHY IS IT USEFUL?

$$\mathcal{L}u(x) = f(x)$$

$$\mathcal{L}G(x, x') = \delta(x - x')$$

$$u(x) = \int G(x, x')f(x') \, dx$$

(Some conditions apply.)

FAMILIAR GREEN'S FUNCTIONS

Impulse response of a linear time-invariant system:

$$y(t) = \int_{-\infty}^{\infty} x(t')h(t-t') dt'$$

FAMILIAR GREEN'S FUNCTIONS

Poisson's equation:

$$\nabla^2 V(\mathbf{r}) = -\frac{\rho(\mathbf{r})}{\epsilon_0}$$

$$V(\mathbf{r}) = \iiint \frac{1}{4\pi\epsilon_0 |\mathbf{r} - \mathbf{r}'|} \rho(\mathbf{r}') d^3\mathbf{r}'$$

FAMILIAR GREEN'S FUNCTIONS

Helmholtz equation:

$$(\nabla^2 + k^2) A_z(\mathbf{r}) = -J_z(\mathbf{r})$$

$$A_z(\mathbf{r}) = \iiint \frac{e^{-jk|\mathbf{r}-\mathbf{r}'|}}{4\pi |\mathbf{r}-\mathbf{r}'|} J_z(\mathbf{r}') d^3\mathbf{r}'$$

FAMILIAR GREEN'S FUNCTIONS

Green's functions let us:

- Derive and understand these expressions.
- Generalize to other problems and boundary conditions.

FINDING THE GREEN'S FUNCTION

A SIMPLE EXAMPLE

Original problem:

$$\frac{d^2 u(x)}{dx^2} - k^2 u(x) = f(x)$$

Green's function problem:

$$\frac{d^2 G(x, x')}{dx^2} - k^2 G(x, x') = \delta(x - x')$$

A SIMPLE EXAMPLE

For $x \neq x'$

$$\frac{d^2 G(x, x')}{dx^2} - k^2 G(x, x') = 0$$

So we have

$$G(x, x') = \begin{cases} Ae^{+k(x-x')} & \text{for } x < x' \\ Be^{-k(x-x')} & \text{for } x > x' \end{cases}$$

A SIMPLE EXAMPLE

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A SIMPLE EXAMPLE

$$\frac{d^2 G(x, x')}{dx^2} - k^2 G(x, x') = \delta(x - x')$$

Continuity of the Green's function:

$$\lim_{\epsilon \rightarrow 0} [G(x' + \epsilon, x') - G(x' - \epsilon, x')] = 0$$

A SIMPLE EXAMPLE

$$\frac{d^2 G(x, x')}{dx^2} - k^2 G(x, x') = \delta(x - x')$$

Continuity of the Green's function:

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A SIMPLE EXAMPLE

$$\frac{d^2 G(x, x')}{dx^2} - k^2 G(x, x') = \delta(x - x')$$

Discontinuity condition:

$$\lim_{\epsilon \rightarrow 0} \left[\left. \frac{dG}{dx} \right|_{x=x'+\epsilon} - \left. \frac{dG}{dx} \right|_{x=x'-\epsilon} \right] = 1$$

A SIMPLE EXAMPLE

$$\int_{x'-\epsilon}^{x'+\epsilon} \left[\frac{d^2 G(x, x')}{dx^2} - k^2 G(x, x') \right] dx = \int_{x'-\epsilon}^{x'+\epsilon} \delta(x - x') dx$$

Discontinuity condition:

$$\lim_{\epsilon \rightarrow 0} \left[\left. \frac{dG}{dx} \right|_{x=x'+\epsilon} - \left. \frac{dG}{dx} \right|_{x=x'-\epsilon} \right] = 1$$

A SIMPLE EXAMPLE

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A SIMPLE EXAMPLE

$$G(x, x') = \begin{cases} Ae^{+k(x-x')} & \text{for } x < x' \\ Be^{-k(x-x')} & \text{for } x > x' \end{cases}$$

Continuity of $G(x, x')$:

$$A = B$$

Discontinuity of $\frac{dG(x, x')}{dx}$:

$$kA + kB = 1$$

A SIMPLE EXAMPLE

Solving, our Green's function is

$$G(x, x') = \frac{1}{2k} \begin{cases} e^{+k(x-x')} & \text{for } x < x' \\ e^{-k(x-x')} & \text{for } x > x' \end{cases}$$

Or, more compactly:

$$G(x, x') = \frac{e^{k|x-x'|}}{2k}$$

A SIMPLE EXAMPLE

Original problem:

$$\frac{d^2 u(x)}{dx^2} - k^2 u(x) = f(x)$$

Solution:

$$u(x) = \int_{-\infty}^{\infty} f(x') \frac{e^{k|x-x'|}}{2k} dx'$$

GENERAL APPROACH

Direct solution:

- $G(x, x')$ obeys source-free equation for $x \neq x'$.
- $G(x, x')$ and its derivatives are continuous or discontinuous at $x = x'$.

CONSTRUCTING THE SOLUTION

$$u(x) = \int G(x, x') f(x') dx'$$

Can we prove/generalize this?

ADJOINT OPERATORS

Inner product:

$$\langle u, v \rangle = \int_a^b u(x) \bar{v}(x) \, dx$$

Adjoint operator \mathcal{L}^* :

$$\langle \mathcal{L}u, v \rangle = \langle u, \mathcal{L}^*v \rangle$$

ADJOINT BOUNDARY CONDITIONS

$$\langle \mathcal{L}u, v \rangle = \langle u, \mathcal{L}^*v \rangle \longrightarrow \text{Only for certain } u, v !$$

Adjoint boundary conditions:

$$\mathcal{B}_i[u] = 0; \quad \mathcal{B}_i^*[v] = 0$$

Otherwise

$$\langle \mathcal{L}u, v \rangle = \langle u, \mathcal{L}^*v \rangle + \underbrace{J(u, \bar{v})}_{\text{Conjunct}} \Big|_a^b$$

ADJOINT OPERATORS: EXAMPLE

$$\mathcal{L}u(x) = \left[\frac{d^2}{dx^2} + k^2 \right] u(x)$$

Want \mathcal{L}^* so that

$$\langle \mathcal{L}u, v \rangle = \langle u, \mathcal{L}^*v \rangle + J(u, \bar{v}) \Big|_a^b$$

ADJOINT OPERATORS: EXAMPLE

$$\langle \mathcal{L}u, v \rangle = \int_a^b [u'' + k^2 u] \bar{v} \, dx$$

$$\langle \mathcal{L}u, v \rangle = \int_a^b [-u' \bar{v}' + k^2 u \bar{v}] \, dx + [u' \bar{v}]_a^b$$

$$\langle \mathcal{L}u, v \rangle = \int_a^b u [\bar{v}'' + k^2 \bar{v}] \, dx + [u' \bar{v} - u \bar{v}']_a^b$$

ADJOINT OPERATORS: EXAMPLE

$$\langle \mathcal{L}u, v \rangle = \int_a^b u [\bar{v}'' + k^2 \bar{v}] dx + [u' \bar{v} - u \bar{v}']_a^b$$

Looks like

$$\langle \mathcal{L}u, v \rangle = \langle u, \mathcal{L}^* v \rangle + J(u, \bar{v}) \Big|_a^b$$

with

$$\mathcal{L}^* = \frac{d^2}{dx^2} + \overline{k^2}$$

$$J(u, \bar{v}) = u' \bar{v} - u \bar{v}'$$

ADJOINT OPERATORS: EXAMPLE

For

$$\langle \mathcal{L}u, v \rangle = \langle u, \mathcal{L}^*v \rangle$$

must have

$$\begin{aligned} J(u, \bar{v}) \Big|_a^b &= 0 \\ [u'\bar{v} - u\bar{v}']_a^b &= 0 \end{aligned}$$

$$u'(b)\bar{v}(b) - u(b)\bar{v}'(b) - u'(a)\bar{v}(a) + u(a)\bar{v}'(a) = 0$$

ADJOINT OPERATORS: EXAMPLE

$$u'(b)\bar{v}(b) - u(b)\bar{v}'(b) - u'(a)\bar{v}(a) + u(a)\bar{v}'(a) = 0$$

Boundary conditions:

$$\begin{array}{l} \mathcal{B}_1[u] = u(a) = 0 \\ \mathcal{B}_2[u] = u(b) = 0 \end{array} \implies \begin{array}{l} \mathcal{B}_1^*[v] = v(a) = 0 \\ \mathcal{B}_2^*[v] = v(b) = 0 \end{array}$$

$$\mathcal{B}_i = \mathcal{B}_i^*$$

ADJOINT OPERATORS: EXAMPLE

$$u'(b)\bar{v}(b) - u(b)\bar{v}'(b) - u'(a)\bar{v}(a) + u(a)\bar{v}'(a) = 0$$

Initial conditions:

$$\begin{array}{ll} \mathcal{B}_1[u] = u(a) = 0 & \mathcal{B}_1^*[v] = v(b) = 0 \\ \mathcal{B}_2[u] = u'(a) = 0 & \mathcal{B}_2^*[v] = v'(b) = 0 \end{array} \implies$$

$$\mathcal{B}_i \neq \mathcal{B}_i^*$$

ADJOINT OPERATORS: SUMMARY

Adjoint operator:

$$\langle \mathcal{L}u, v \rangle = \langle u, \mathcal{L}^*v \rangle$$

Adjoint boundary conditions:

$$\begin{aligned} \mathcal{B}_i[u] = 0 \\ \mathcal{B}_i^*[v] = 0 \end{aligned} \implies J(u, \bar{v}) \Big|_a^b = 0$$

Otherwise:

$$\langle \mathcal{L}u, v \rangle = \langle u, \mathcal{L}^*v \rangle + J(u, \bar{v}) \Big|_a^b$$

THE ADJOINT GREEN'S FUNCTION

Original problem:

$$\mathcal{L}u(x) = f(x); \quad \mathcal{B}_i[u(x)] = \alpha_i$$

Green's problem:

$$\mathcal{L}G(x, x') = \delta(x - x'); \quad \mathcal{B}_i[G(x, x')] = 0$$

Adjoint Green's problem:

$$\mathcal{L}^*H(x, x') = \delta(x - x'); \quad \mathcal{B}_i^*[H(x, x')] = 0$$

CONSTRUCTING SOLUTIONS: DERIVATION

$$\langle \mathcal{L}u(x), H(x, x') \rangle = \langle u(x), \mathcal{L}^*H(x, x') \rangle + J(u(x), \bar{H}(x, x')) \Big|_a^b$$

$$\langle f(x), H(x, x') \rangle = \langle u(x), \delta(x - x') \rangle + J(u(x), \bar{H}(x, x')) \Big|_a^b$$

$$\int_a^b f(x) H^*(x, x') dx = \int_a^b u(x) \delta(x - x') dx + J(u(x), \bar{H}(x, x')) \Big|_a^b$$

$$u(x') = \int_a^b f(x) \bar{H}(x, x') dx - J(u(x), \bar{H}(x, x')) \Big|_a^b$$

CONSTRUCTING SOLUTIONS: DERIVATION

How are $G(x, x')$ and $H(x, x')$ related?

$$\langle \mathcal{L}G(x, x'), H(x, x'') \rangle = \langle G(x, x'), \mathcal{L}^*H(x, x'') \rangle$$

$$\langle \delta(x - x'), H(x, x'') \rangle = \langle G(x, x'), \delta(x - x'') \rangle$$

$$\int_a^b \delta(x - x') \bar{H}(x, x'') dx = \int_a^b G(x, x') \delta(x - x'') dx$$

$$\bar{H}(x', x'') = G(x'', x')$$

$$G(x, x') = \bar{H}(x', x)$$

CONSTRUCTING SOLUTIONS: DERIVATION

$$u(x') = \int_a^b f(x) \bar{H}(x, x') \, dx - J(u(x), \bar{H}(x, x')) \Big|_a^b$$

and

$$G(x, x') = \bar{H}(x', x)$$

so

$$u(x) = \int_a^b f(x') G(x, x') \, dx' - J(u(x'), G(x, x')) \Big|_{x'=a}^b$$

EXAMPLE: 1D POISSON EQUATION

Original problem:

$$\frac{d^2V(x)}{dx^2} = -\frac{\rho(x)}{\epsilon_0}; \quad \begin{aligned} V(a) &= V_a \\ V(b) &= V_b \end{aligned}$$

Green's problem:

$$\frac{d^2G(x, x')}{dx^2} = \delta(x - x'); \quad \begin{aligned} G(a, x') &= 0 \\ G(b, x') &= 0 \end{aligned}$$

EXAMPLE: 1D POISSON EQUATION

Solution:

$$V(x) = \int_a^b -\frac{\rho(x')}{\epsilon_0} G(x, x') dx' - J(V(x'), G(x, x')) \Big|_{x'=a}^b$$

Take $\rho(x) = 0$ for now.

$$V(x) = -J(V(x'), G(x, x')) \Big|_{x'=a}^b$$

EXAMPLE: 1D POISSON EQUATION

Can show

$$\langle \mathcal{L}u, v \rangle = \int_a^b u'' \bar{v} \, dx = \int_a^b u \bar{v}'' \, dx + [u' \bar{v} - u \bar{v}']_a^b$$

Comparing with

$$\langle \mathcal{L}u, v \rangle = \langle u, \mathcal{L}^* v \rangle + J(u, \bar{v}) \Big|_a^b$$

we see

$$\mathcal{L}^* = \frac{d^2}{dx^2} = \mathcal{L}$$

$$J(u, \bar{v}) = u' \bar{v} - u \bar{v}'$$

EXAMPLE: 1D POISSON EQUATION

$$V(x) = -J(V(x'), G(x, x')) \Big|_{x'=a}^b$$

$$V(x) = \left[V(x') \frac{dG(x, x')}{dx'} - \frac{dV(x')}{dx'} G(x, x') \right]_{x'=a}^b$$

$$\begin{aligned} V(x) = & V_b \frac{dG(x, b)}{dx'} - \frac{dV(b)}{dx'} G(x, b) - \\ & - V_a \frac{dG(x, a)}{dx'} + \frac{dV(a)}{dx'} G(x, a) \end{aligned}$$

EXAMPLE: 1D POISSON EQUATION

Green's problem:

$$\frac{d^2 G(x, x')}{dx^2} = \delta(x - x'); \quad \begin{aligned} G(a, x') &= 0 \\ G(b, x') &= 0 \end{aligned}$$

Adjoint Green's problem:

$$\frac{d^2 H(x, x')}{dx^2} = \delta(x - x'); \quad \begin{aligned} H(a, x') &= 0 \\ H(b, x') &= 0 \end{aligned}$$

But $G(x, x') = \overline{H}(x', x)$ so

$$\frac{d^2 \overline{G}(x, x')}{dx'^2} = \delta(x - x'); \quad \begin{aligned} G(x, a) &= 0 \\ G(x, b) &= 0 \end{aligned}$$

EXAMPLE: 1D POISSON EQUATION

$$V(x) = V_b \frac{dG(x, b)}{dx'} - \frac{dV(b)}{dx'} G(x, b) - \\ - V_a \frac{dG(x, a)}{dx'} + \frac{dV(a)}{dx'} G(x, a)$$

With $G(x, a) = G(x, b) = 0$, we have

$$V(x) = V_b \frac{dG(x, b)}{dx'} - V_a \frac{dG(x, a)}{dx'}$$

EXAMPLE: 1D POISSON EQUATION

Full solution with $\rho(x)$:

$$V(x) = \int_a^b -\frac{\rho(x')}{\epsilon_0} G(x, x') dx' + V_b \frac{dG(x, b)}{dx'} - V_a \frac{dG(x, a)}{dx'}$$

SOLVING PROBLEMS WITH GREEN'S FUNCTIONS

$$\text{Given: } \mathcal{L}u(x) = f(x); \quad \mathcal{B}_i[u(x)] = \alpha_i$$

- ① Solve Green's problem

$$\mathcal{L}G(x, x') = \delta(x - x'); \quad \mathcal{B}_i[G(x, x')] = 0$$

- ② Find \mathcal{L}^* , \mathcal{B}_i^* , and $J(u, \bar{v})$ from $\langle \mathcal{L}u, v \rangle$.

- ③ Solution is

$$u(x) = \int_a^b f(x')G(x, x') dx' - J(u(x'), G(x, x')) \Big|_{x'=a}^b$$

- ④ Simplify using $\mathcal{B}^*[G(x, x')] = 0$ (with respect to x').

SOLVING 3D PROBLEMS WITH GREEN'S FUNCTIONS

$$\text{Given: } \mathcal{L}u(\mathbf{r}) = f(\mathbf{r}); \quad \mathcal{B}_i[u(\mathbf{r})] = \alpha_i$$

- 1 Solve Green's problem

$$\mathcal{L}G(\mathbf{r}, \mathbf{r}') = \delta(\mathbf{r} - \mathbf{r}'); \quad \mathcal{B}_i[G(\mathbf{r}, \mathbf{r}')] = 0$$

- 2 Find \mathcal{L}^* , \mathcal{B}^* , and $J(u, \bar{v})$ from $\langle \mathcal{L}u, v \rangle$.
- 3 Solution is

$$u(\mathbf{r}) = \int_V f(\mathbf{r}') G(\mathbf{r}, \mathbf{r}') d^3\mathbf{r}' - \oint_{\partial V} J(u(\mathbf{r}'), G(\mathbf{r}, \mathbf{r}')) \cdot d\mathbf{s}$$

- 4 Simplify using $\mathcal{B}^*[G(\mathbf{r}, \mathbf{r}')] = 0$ (with respect to \mathbf{r}').

Comments:

- Calculating Green's functions is not trivial.
- Adjoint approach is difficult, but offers clarity.

SPECIAL PROPERTIES

If $\mathcal{L} = \mathcal{L}^*$ and $\mathcal{B}_i = \mathcal{B}_i^*$, then

$$G(x, x') = \overline{G}(x', x)$$

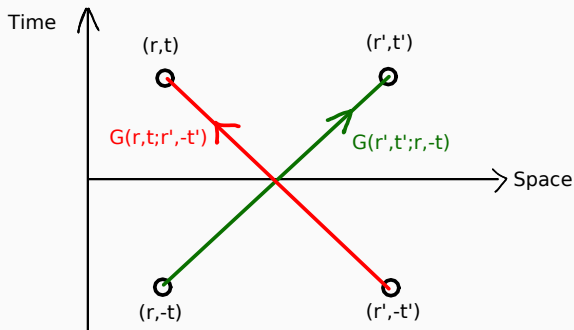
Reciprocity in the frequency-domain:

$$G(\mathbf{r}, \mathbf{r}') = G(\mathbf{r}', \mathbf{r})$$

RECIPROCITY

Reciprocity in the time-domain:

$$G(\mathbf{r}, t; \mathbf{r}', -t') = G(\mathbf{r}', t'; \mathbf{r}, -t)$$



Causality:

$$G(\mathbf{r}, t; \mathbf{r}', t') = 0 \quad \text{for } t < t'$$

Special relativity:

$$G(\mathbf{r}, t; \mathbf{r}', t') = 0 \quad \text{for } |\mathbf{r} - \mathbf{r}'| > c(t - t')$$

Time-invariance:

$$G(t, t') = G(t - t')$$

Spatial-invariance:

$$G(\mathbf{r}, \mathbf{r}') = G(\mathbf{r} - \mathbf{r}')$$

$$(\mathcal{L} - \lambda)u(x) = f(x); \quad \mathcal{B}[u] = 0$$

If $\mathcal{L} = \mathcal{L}^*$ and $\mathcal{B} = \mathcal{B}^*$ then

$$u(x) = \sum_n \frac{\langle f, \phi_n \rangle}{\lambda_n - \lambda} \phi_n(x)$$

$$u(x) = \int_a^b \left(\sum_n \frac{\phi_n(x) \phi_n^*(x')}{\lambda_n - \lambda} \right) f(x') dx'$$

$$G(x, x') = \sum_n \frac{\phi_n(x) \phi_n^*(x')}{\lambda_n - \lambda}$$

Green's function of $(\mathcal{L} - \lambda)$:

$$G(x, x'; \lambda) = \sum_n \frac{\phi_n(x) \phi_n^*(x')}{\lambda_n - \lambda}$$

$\lambda_n \longrightarrow$ Poles of $G(x, x'; \lambda)$

$\phi_n(x) \phi_n^*(x') \longrightarrow$ Residues of $G(x, x'; \lambda)$

CONCLUSION

TAKEAWAYS

- Green's function is the impulse response.
- Finding Green's function:
 - Source-free behaviour for $x \neq x'$.
 - Continuity/discontinuity requirements at $x = x'$.
- Constructing solutions:
 - Systematic method using adjoint equation.
 - Non-zero boundary conditions behave like sources.
- Lots of information in the Green's function.
- Green's functions \iff eigenvalues/eigenfunctions.

FURTHER READING

Dudley (1994), *Mathematical foundations for electromagnetic theory*. Great introduction to 1D Green's functions: deals with subtleties that others ignore.

Gerlach (2010), *Linear mathematics in infinite dimensions*. Nice set of online course notes for quick reference.

Balanis (2012), *Advanced engineering electromagnetics*. Not very rigorous, but decent for getting the key ideas.

Morse and Feshbach, *Methods of theoretical physics*. Big, detailed reference. Great resource for deeper insight and understanding.

FURTHER READING

Collin (1990), *Field theory of guided waves*. Huge chapter on Green's functions. Emphasis on dyadics.

Folland (1992), *Fourier analysis and its applications*. Rigorous math book. Chapter on generalized functions is particularly nice.

Byron and Fuller (1992), *Mathematics of classical and quantum physics*. Interesting alternative approach.

Warnick (1996), "Electromagnetic Green functions using differential forms." For the differential forms inclined.