## PH4606 - Lecture 2

Claus-Dieter Ohl

February 23, 2016

## 1 Harmonic Plane Waves

In the following we will consider a constant speed of sound and start with the description of a plane wave propagation. Planar waves have a constant amplitude and phase perpendicular to the direction of propagation. Thus if the wave moves in the x-plane we have:

$$\frac{\partial^2 p}{\partial x^2} = \frac{1}{c^2} \frac{\partial^2 p}{\partial t^2} \quad . \tag{2.1}$$

The complex solution to the harmonic wave equation is

$$p = A e^{i(\omega t - kx)} + B e^{i(\omega + kx)} \quad . \tag{2.2}$$

Using the linear Euler Eq. (1.11) we can obtain the particle velocity in a plane wave as

$$\mathbf{u} = u \,\vec{e}_1 = \left[ \frac{A}{\rho_0 c} e^{i(\omega t - kx)} - \frac{B}{\rho_0 c} e^{i(\omega t + kx)} \right] \vec{e}_1 \tag{2.3}$$

Please note that the constants A and B, pressures and velocities are now all complex. We can separate the wave travelling in positive x-direction from the the wave travelling in negative x-direction and designate them with the subscript + and -, respectively.

The wave travelling in positive x-direction is

$$p_{+} = Ae^{i(\omega t - kx)} \tag{2.4}$$

while in negative x-direction is

$$p_{-} = Be^{i(\omega t + kx)} \tag{2.5}$$

These two waves can be formally written for the other acoustic variables:

$$u_{\pm} = \pm \frac{p_{\pm}}{\rho_0 c} \tag{1}$$

$$u_{\pm} = \pm \frac{p_{\pm}}{\rho_0 c}$$

$$s_{\pm} = \frac{p_{\pm}}{\rho_0 c^2}$$
(1)

$$\Phi_{\pm} = -\frac{p_{\pm}}{i\omega\rho_0} \tag{2.8}$$

Plane waves travelling in an arbitrary direction can be written as

$$p = Ae^{i(\omega t - k_x x - k_y y - k_z z)} \tag{2.9}$$

If we insert Eq. (2.9) into the wave equation, Eq. (1.17) we obtain the condition

$$\left(\frac{\omega}{c}\right)^2 = k_x^2 + k_y^2 + k_z^2 \quad . \tag{2.10}$$

Now we can introduce the propagation vector,  $\mathbf{k}$  of the wave as

$$\mathbf{k} \equiv k_x \mathbf{e}_1 + k_y \mathbf{e}_2 + k_z \mathbf{e}_3 \tag{2.11}$$

with

$$|\mathbf{k}| = \frac{\omega}{c} \quad . \tag{2.12}$$

Thus using the position vector  $\mathbf{r} = x\mathbf{e}_1 + y\mathbf{e}_2 + z\mathbf{e}_3$  we can rewrite the plane wave solution, Eq. (2.9), as

$$p = Ae^{i(\omega t - \mathbf{kr})} (2.13)$$

The propagation vector points in the direction of wave propagation which can be simply shown. Surfaces of constant phase are defined by the scalar product  $\mathbf{k} \cdot \mathbf{r} = \text{const.}$  The vector  $\nabla(\mathbf{k} \cdot \mathbf{r})$  is a vector in the normal direction to the constant phase surface. But  $\mathbf{k} = \nabla(\mathbf{k} \cdot \mathbf{r})$ , thus  $\mathbf{k}$  is pointing normal to the plane of constant phase and thus in the direction of propagation.

## 1.1 Example

Let a plane wave have a constant phase parallel to the z-axis, thus the wave travels in the xy-plane as sketched in Fig. 2.1. Formally the surface of constant phase is  $\mathbf{k} \cdot \mathbf{r} = 0$ ; thus if the phase is independent of z,  $\mathbf{k}$  can only have  $k_x$  and  $k_y$  components. We can thus write Eq. (2.9) as

$$p = Ae^{i(\omega t - k_x x - k_y y)} \tag{2.14}$$

and the equation for the surface of constant phase is

$$k_x x + k_y y = \text{const.} (2.15)$$

or

$$y = -\frac{k_x}{k_y}x + \text{const.} . (2.16)$$

Equation (2.16) is line with a slope  $-k_x/k_y$ . The surfaces are thus oriented as sketched in Fig. 2.1.

Figure 2.1

Now let's study this wave in the x-direction while y=0

$$p = Ae^{i(\omega t - k_x x)} \tag{2.17}$$

where  $k_x = 2\pi/\lambda_x$  with  $\lambda_x$  being the apparent wavelength in x-direction. Note that this wavelength is larger than the wavelength in  $\vec{k}$ -direction. We can relate  $\lambda_x$  with the help of Fig. 1 to  $\cos \phi = \lambda/\lambda_x$  and thus  $k_x = k \cos \phi$ . Applying the same arguments in the y-direction gives

$$\mathbf{k} = k\cos\phi\,\mathbf{e}_1 + k\sin\phi\,\mathbf{e}_2 \quad . \tag{2.18}$$

Thus an harmonic plane wave in 2-dimensions propagating under an angle  $\phi$  to the x-axis can be written as

$$p = Ae^{i(\omega t - kx\cos\phi - ky\sin\phi)} \tag{2.19}$$

```
## Simulation of a 2-dimensional wave
    %matplotlib notebook
   import math as m
   import numpy #array operations
   import matplotlib.pyplot as plt #plotting
   from ipywidgets import widgets #for the widgets
   #from ipykernel.client import display
    from IPython import display #for continous display
    #from PIL import Image #to export images
   nimg=0
10
    def savemyimage(visual):
12
      global nimg
13
      visual = (visual +2.)/4.
14
      #result = Image.fromarray((visual * 255).astype(numpy.uint8))
15
      #result.save('out{:03d}.bmp'.format(nimg))
16
      nimg=nimg+1
    def plotwave(u,time,px,py,pp,pt):
19
      plt.figure(1)
20
      plt.clf()
21
      plt.subplot2grid((4,4),(0,0), colspan = 4, rowspan = 3)
22
      plt.imshow(u, origin='upper', extent=[0., 2., 0., 2.], vmax=2, vmin=-2) #plot the wave field
23
      plt.text(0.1,1.8,"time {0:.5f}".format(time)) #annotate the time
      plt.plot(w_probex.value,1.,'o') #position of probe
      plt.subplot2grid((4,4),(3,0), colspan = 4)
26
      plt.plot(pt,pp) #plot the pressure at the probe
27
      plt.gca().set_ylim([-2,2])
28
      display.clear_output(wait=True)
29
      display(plt.gcf())
30
31
    def solvewave(b):
32
      tabs.visible=False
33
      #computational domain
34
      nx = ny = 381
35
      size=2. #size of the domain
36
      #parameters of the wave
      c = 5. #speed of sound
      l=w_wavelength.value #wavelength
39
      nu=c/l #frequency
40
      omega=nu*2.*m.pi #angular frequency
41
      duration=w_sourceduration.value/nu #duration of source
42
      #position
43
      emissionlength=(w_sizeemit.value/100.)*nx
44
      startx=int(nx/2-emissionlength/2)
45
```

endx=int(nx/2+emissionlength/2)

46

```
47
       #further variables
 48
       dx = size/(nx-1)
 49
       CFL=0.1 #CFL number <1
 50
       dt = CFL*dx/c
       nt=int(w_simduration.value/dt) #number of time steps
 52
       if w_position.value=='Left':
 53
           sourcepos=0
54
       else:
 55
           sourcepos=int(nx/2)
 56
       #arrays for measuring the pressure at a position
       pt=numpy.arange(nt+1)*dt
       pp=numpy.zeros(nt+1)
 60
       px=int(w_probex.value*(nx-1)/2.)
 61
       py=int(ny/2)
 62
 63
       #every xx times over the total nt timesteps an output should be generated
 64
       output=map(int,list(numpy.linspace(1,nt,int(nt/50))))
 65
 66
       u = numpy.zeros((nx,ny)) #pressure at t
67
       un = numpy.zeros((nx,ny)) #pressure at t-dt
 68
       unn= numpy.zeros((nx,ny)) #pressure at t-2*dt
 69
       row, col = u.shape
       #Assign initial conditions, for a sine wave
       un[startx:endx,0]=0.
                                \#amplitude is sin(omega*t) with t=0
 73
       unn [startx:endx,0]=1.
                                #velocity is cos(omega*t) with t=0
 74
 75
       C=c*c*dt*dt/dx/dx
 76
 77
       plt.figure(1, figsize=(8, 8), dpi=300)
 78
 79
       for n in range(nt+1): ##loop across number of time steps
 80
           #this line computes the finite differences of the wave equation
 81
           u[1:-1,1:-1]=2.*un[1:-1,1:-1]-unn[1:-1,1:-1]+C*(un[1:-1,:-2]+
 82
                            un[:-2,1:-1]+un[2:,1:-1]+un[1:-1,2:]-4.*un[1:-1,1:-1])
           #hard reflective boundary conditions
 85
           u[0,:] = u[1,:]
86
           u[-1,:] = u[-2,:]
 87
           u[:,0] = u[:,1]
 88
           u[:,-1] = u[:,-2]
 89
 90
           #pressure source
91
92
           #if float(n*dt<duration):</pre>
           u[startx:endx,sourcepos]=m.sin(omega*n*dt)*float(n*dt<duration)
93
94
           pp[n]=u[py,px]
95
           #save values for the time derivative
97
           unn=un.copy() \#n-1 time steop
           un=u.copy()
                         #n time step
100
           if (n in output):
101
               plotwave(u,n*dt,px,py,pp,pt)
102
               if (w_saveplots.value):
103
104
                   savemyimage(u)
105
```

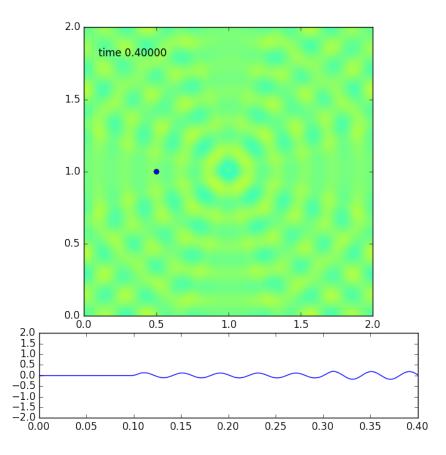


Figure 1:

#and plot the last figure
plotwave(u,n\*dt,px,py,pp,pt)
tabs.visible=True

1