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Statistical approach to muography as a non-destructive testing technique for industry problem solving

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Abstract

In this work, cosmic muons have been used to perform a muography experiment, a non-destructive method allowing us to study the internal properties of several physical objects with various geometries. Such study is extremely useful in the industry, in order for example to study the degradation of pipes, and finds many application in our daily lives.

In particular, a new C++ framework has been developed in order to study the results of such muography experiments and study their results from a statistical point of view, to find the optimal geometrical parameters of different objects placed between two muon detectors.

Key words : cosmic muons, muography, geometry, likelihood

Resumen

Palabras claves :

Acknowledgments

I would like to mainly thank Pablo for his outstanding work as the director of this work over the past few months. He has been incredibly helpful and present all along the way, teaching me the basic of muography and helping and guiding me when needed. I know it has not always been easy but I am very grateful for the time he spent with me, and I am more than happy with the final result and glad I decided to choose this particular subject.

Thank you to Muon Systems and its CEO Carlos Díez in particular, the co-director of this work, for allowing me be a part of their experiment during a few months and allowing me to get a privileged access to actual data collected. Having the opportunity to work in collaboration with a real industry has been an extremely interesting experience that will for sure help me a lot in the future. I obviously wish to the company all the best since the technology developed definitely show a lot of potential.

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Chapter 1

Introduction

Nowadays, muon tomography is an active field of research since it is a non destructive method allowing to map the inside of large objects difficult and/or dangerous to access, without any contact or damage, and without even having physical access to them. To do so, this technique uses muons, elementary particles similar to electrons but with a much greater mass, which allows them to penetrate much deeper and probe matter in a more efficient way, since they suffer less from the bremsstrahlung radiation affecting all leptons. Such technology is relatively well-know and presents several advantages over other techniques such as X-ray imaging since it is globally safe and clean and uses natural radiation, cosmic muons, while providing an excellent penetration in matter in order to study it.

In this particular work, muon tomography will be applied to industry and used in order to probe several kinds of physical objects, such as pipes, to study their degradation without damaging them. This is the main objective of the company called Muon systems, founded in 2017 and based in Bilbao.

This company already built powerful muon detectors allowing to perform such experiment and developed over the years a powerful framework allowing to study the degradation of pipes using convolutional neural networks, but this method is limited in the sense that training such networks is computing intensive and time consuming, and therefore does not allow to study several properties of such pipes at once. In this work, an alternative method based on a statistical analysis of the problem is presented, allowing a greater generalization of the previous method to other object geometries and to a higher dimension parameters phase space.

This project is highly relevant in today's society because it allows to combine data science algorithms and high-performing statistical methods that have been developed in this context to solver this particular muography problem, putting them in practice in an actual industry.

Chapter 2

Muons and muography

The main particles used throughout this work, muons, will be briefly introduced in Section 2.1 of this Chapter. Muons are particles produced naturally by cosmic rays, introduced in Section 2.2 and, once produced, they obviously interact when traversing the atmosphere and matter in general in ways described in Section 2.3 that actually need to be understood extremely well for the muon tomography process introduced in Section 2.4 to be useful and applicable to industry. Finally, the actual experimental setup used for this work will be explained in Section 2.5.

2.1 Particle physics and muons

Particle physics is the field which studies the matter surrounding us, along with the fundamental interactions between the particles. In this context, the Standard Model of particle physics [1] is nowadays the most accepted mathematical model used to describe the elementary particles and three of the 4 fundamental forces of nature (electromagnetic, weak and strong interactions, while the gravitational interaction is out of reach of this model). Even though quite simple in concept, it has been able to describe most of the phenomena observed in nature so far with an incredible level of precision, and has made a lot of predictions that have now been proven to be true, such as the discovery of the top quark [2] in 1995, the tau neutrino [3] in 2001 and the Higgs boson itself [4, 5], the last missing piece of the Standard Model, announced at CERN in July 2012.

According to this model, 12 different fermions (along with their 12 corresponding anti-particles) exist in nature, as shown in Figure 2.1, most of them being unstable. These fermions can be divided into two fundamentally different categories, the quarks and the leptons, containing each 6 particles and sensitive to different forces. Even though quite interesting, the quarks do not play a fundamental role in the muon tomography detailed in this work, so only leptons will be considered from now on. In particular, leptons can be divided even more into three different generations of particles, and the muon, one particular lepton belonging to the second generation, will be the main focus of this work.

Muons μ^- [9] are therefore fundamental particles having a negative charge and quite similar in nature to electrons, even though they have a quite high mass (200 times larger than the electron), which implies that they are not stable particles: they have a lifetime of approximately $2.2\mu\text{s}$, and typically decay into an electron and a pair of neutrinos. However, this lifetime is actually quite long with respect to other fundamental particles and muons are on average able to travel more than 700 meters, allowing us to consider them to be stable particles in many processes, such as the one presented in this work. Muons also have a relatively small interaction cross-section with ordinary matter, even though they do interact with baryonic matter by several processes described in Section 2.3.

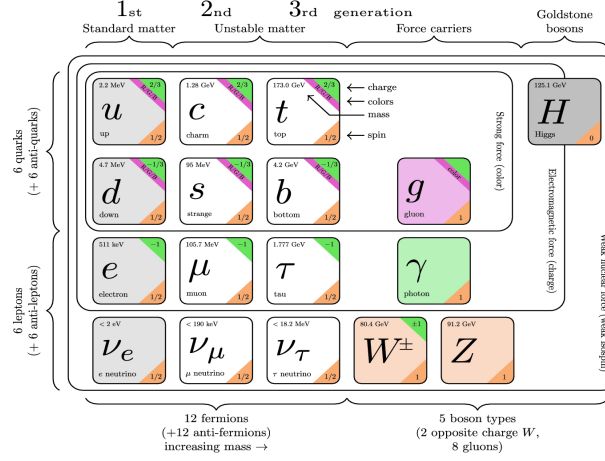


Figure 2.1: Representation of the 12 fermions of the Standard Model [8] along with the main force carriers and the Higgs boson, discovered in 2012 and completing this model.

2.2 Cosmic rays

Being unstable by nature, once produced, muons decay almost instantly by a weak process into an electron and a pair of neutrinos. However, it is possible to observe them in nature, since they are continually produced, mainly thanks to cosmic rays [10], constant flux of high energy particles (mostly protons and atomic nuclei) coming mostly from supernovae explosions and AGN emissions and reaching the Earth every day. Indeed, as they impact our atmosphere, these particles start a chain reaction, as shown in Figure 2.2: first of all, several neutral and charged pions are produced, decaying themselves into a pair of photons (and, later on, electron and positron pairs) and muons, respectively.

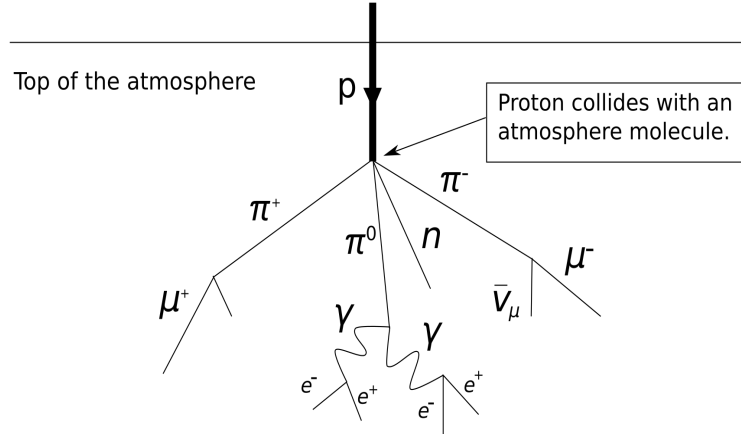


Figure 2.2: Typical chain of decays induced by highly energetic cosmic rays when reaching the Earth's atmosphere.

Muons are the most abundant charged particles produced by these processes actually reaching the sea level, as shown in Figure 2.3. Even though they are unstable and have a limited lifetime, around 0.06% of muons produced by such processes indeed do manage to reach the sea level thanks to the temporal distorsion induced by their high energy and relativistic speed. As a rule of thumb, one can expect to observe 10.000 muons per square meter and per minute at the sea level. Muons were actually discovered thanks to cosmic rays in 1936 [11].

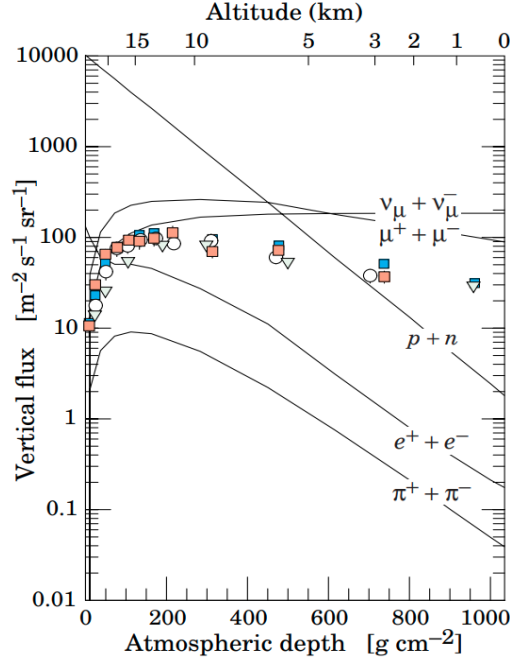


Figure 2.3: Abundance of particles observed at the sea level, due to cosmic rays [10].

With this work, millions of cosmic muons have been generated in order to simulate an actual experiment.

2.3 Muons interaction with matter

Muons typically interact with ordinary matter through two main processes: ionization and multiple scattering, both resulting in different effects on the incident particle.

2.3.1 Ionization process

First of all, cosmic muons can interact with matter through **ionization**, when the incident muon gives some of its energy to the electrons of the absorber. This process is well described by the famous Bethe-Bloch formula shown in Equation 2.1 [9], relating the average loss of energy over a distance $\frac{dE}{dx}$ (typically referred to as the **mass stopping power**) of material with several parameters, such as the charge number of incident particle z (for cosmic muons, $z = 1$), the atomic mass and charge of absorber A and Z , the relativistic factors β and γ , the maximum possible energy transfer to an electron in a single collision W_{\max} and the mean excitation energy I .

$$-\left\langle \frac{dE}{dx} \right\rangle = K z^2 \frac{Z}{A} \frac{1}{\beta^2} \left[\frac{1}{2} \ln \left(\frac{2m_e c^2 \beta^2 \gamma^2 W_{\max}}{I^2} - \beta^2 - \frac{\delta(\beta\gamma)}{2} \right) \right] \quad (2.1)$$

This previous equation gives an accuracy of a few percent in the range $0.1 < \beta < 1000$ and we can easily see that the quantity of energy lost by in a muon when traversing any given medium actually depends on the energy of the incident muon, as shown in Figure 2.4. In practice, this means that cosmic muons, having a mean energy of a few GeV, have mean energy loss rates actually close to the

minimum: they are usually called for this reason *minimum ionizing particles* or MIPs. Energy losses of muons due to this ionization process are for this reason quite small and difficult to detect, so this interaction process will not be considered in this work.

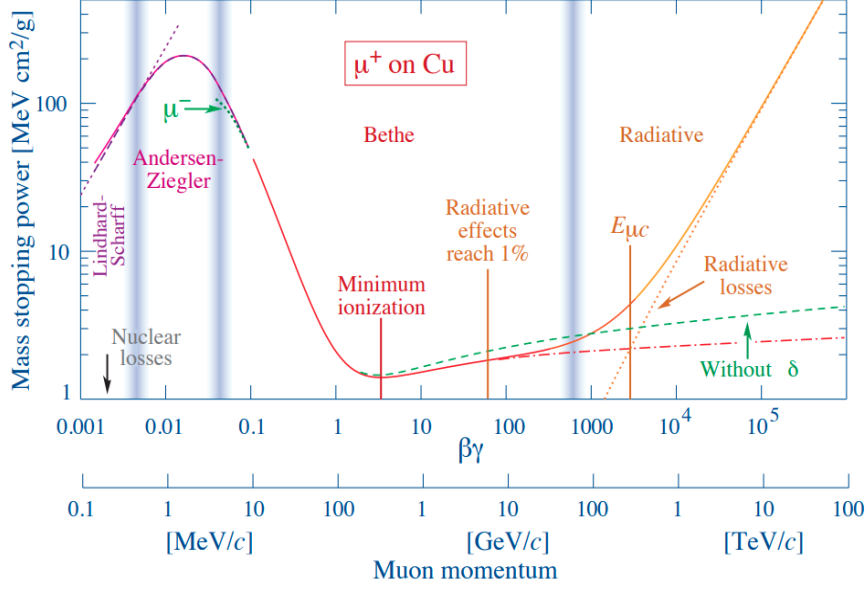


Figure 2.4: Mass stopping power of copper energy lost by a muon in copper due to the ionization process with respect to its momentum [9].

2.3.2 Multiple scattering process

Muons also interact with matter through another process, called **multiple scattering**. Since muons have a negative electric charge, by getting close to the nuclei of the absorber, they are suffering from Coulomb scattering. Given the high number of nuclei in matter, this process is repeated many times, deflecting each time the muon by a small angle in a stochastic way, meaning that there is no way of calculating this deviation exactly, but only using probabilities and the so-called theory of Molière [13].

According to this theory, the central 98% of the projected angular distribution due to Coulomb scattering can be described with a Gaussian function, whose width is given by the θ_0 parameter shown in Equation 2.2 [12], where p is the momentum of the incident particle and X_0 is the radiation length, defined as the characteristic amount of matter traversed by the incident particle for a particular interaction. The second term of the relation is actually typically neglected, being much smaller and therefore having a small impact on the final results obtained.

$$\theta_0 = \frac{13.6 \text{ MeV}}{\beta c p} \sqrt{\frac{x}{X_0}} \left[1 + 0.038 \ln \left(\frac{x}{X_0 \beta^2} \right) \right] \quad (2.2)$$

This formula is expected to be valid for distances up to $\sim 100X_0$, giving an error smaller than 11%. Corrections do exist though in order to get slightly better results, especially in the tails of the distribution, but this theory is precise enough for our needs, given the small deviation angles we expect to observe given the experimental conditions considered.

Even though this deviation angle is the most important parameter when considering the multiple scattering, other parameters can be important as well, as shown in Figure 2.5, such as the value of

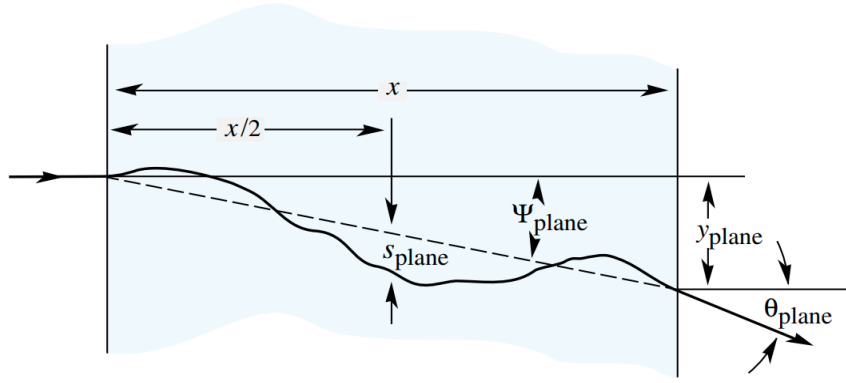


Figure 2.5: Schematic representation of the deviation induced by an absorber to an incident muon because of the multiple scattering effect [12].

y_{plane} , corresponding to the deviation observed with respect to the initial expected position. All these additional parameters are typically highly correlated to the deviation angle previously defined but cannot be easily described with this simple theory and not relevant to this work anyway, since the deviation angle will be the main parameters studied with the muon tomography method used.

2.4 Muon tomography

Given Molière's theory, it is obvious to see that, instead of *calculating* the expected deviation of a muon crossing a given material, we can instead try and *measure* it, by inverting the two previous relations. Since this deviation depends on several parameters related to the absorber itself, such as its width and radiation length X_0 , we can then infer such parameters experimentally and determine the properties of the medium crossed by cosmic muons. This is the so-called muon tomography, or **muography**, taking advantage of the interaction of muons with matter to use it in practice.

Muography is therefore a non-destructive imaging technique producing a density map of the inside of an object by measuring a flux of muons. Such method presents many different advantages over other imaging techniques such as X-rays, since it usually uses natural cosmic rays to make the measurements, being therefore completely safe. Additionally, muons interact lightly with matter, meaning that they typically have high penetrating capabilities and can therefore probe even large and/or dense objects. Muography can in this sense be applied to many different fields: it has for example even been used in 1970 in order to try and find hidden cavities of pyramids in Egypt [14] and can also be used in vulcanology, to determine whether a pocket inside of a volcano is empty or full of lava, among many other practical applications.

Such imaging techniques can be divided into two categories:

- **Absorption muography.** In this case, the observed muon flux in a given direction is compared to what is expected from cosmic rays, trying to determine the inner structure of the absorber as discrepancies between these two values. Only one detector is needed in this case, making this technique useful mostly to study large objects, even though the time to make a single measurement can last up to a few months.
- **Scattering muography.** On the other hand, the multiple scattering of muons can be used, by placing one detector on each side of the object being studied to determine the deviation of

the flux of incoming muons. The denser the material put in between, the larger the observed deviation will be, as shown in Equation 2.2. This technique is mostly used to study smaller objects and is able to make quick measurements.

In this particular work, scattering muography is being applied to industry in order to try and determine the degradation of the interior of industrial equipment, as we will now see.

2.5 Experimental setup

2.5.1 Muon detectors

If we want to work with cosmic muons, we need to be able to build some detectors able to spot them and give us back useful data, such as their energy and/or direction of propagation. Many different technologies exist nowadays in order to detect muons but in this particular case, the typical **multi-wire proportional chambers** have been used.

Invented at CERN in 1968, these detectors use an array of high-voltage wires (playing the role of the anode), running through a chamber filled with gas and whose walls are typically grounded (the cathode), as shown in Figure 2.6. Such an experimental setup therefore creates an electric field inside of the chamber, that needs to be made as large and uniform as possible.

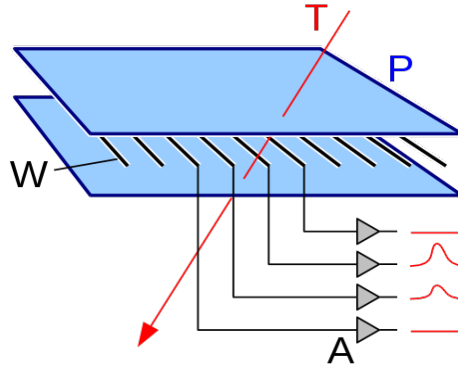


Figure 2.6: Schematic representation of a wire chamber muon detector.

When a charged particle such as a muon crosses this chamber, it will ionize the gas of the chamber and leave small electric charges along its path which will start to drift until reaching one of the wire. This drift induces an electric signal proportional to the ionisation effect in the different wires surrounding the particle path, and the combination of all the signals collected is able to give information regarding the actual path followed by the incoming particle.

Several properties are extremely important when designing a detector:

- First of all, the **spatial resolution**, the precision by which we can tell the position of the muon, should be ideally as small as possible, depending on the actual problem faced.
- The **efficiency** is also an important parameter, since we want to be able to detect as many muons as possible, to make the measurement faster and more precise.
- Finally, we typically want the detector to be large enough to avoid any **acceptance** issues.

2.5.2 Working setup

For this particular work, four muon chambers of 1 meter by 1 meter have been build using these principles, as shown in Figure 2.7. As we can see, more than 200 wires connected to the high voltage and made out of gold and tungsten have been placed every 4 mm in two different planes rotated by 90 degrees, to measure the x and y position of cosmic muons.

These chambers, filled while a mixture of Argon and CO_2 , are then setup in pairs above and below the object being studied, in order to determine the position and the direction of the incoming and outgoing cosmic muon. This setup allows to measure with a good spatial resolution the deviation of muons, to determine as precisely as possible the properties of the object put in between both detectors.

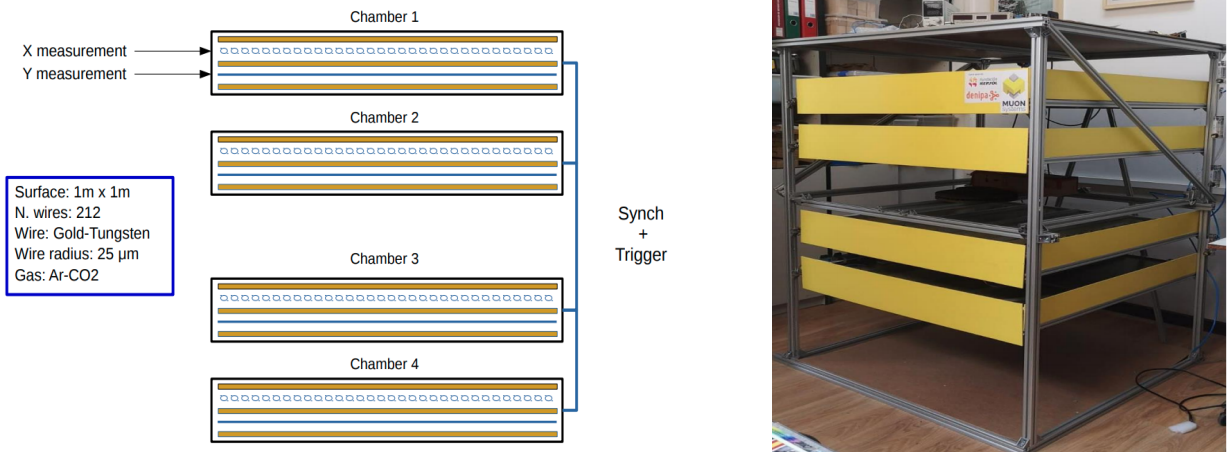


Figure 2.7: Actual working setup used for this work, with two multi-wires muon chambers of 1m^2 placed on each side of the object being studied.

2.5.3 Data flow

Once collected, the data needs to be pass through several different layers in order to convert electric signals into files that can be processed, as shown on Figure 2.8. As we can see, the data follows a different path depending on its nature, if it has been collected by the detector or if we are considering Monte-Carlo simulations that will be described in details in the next Chapter. The stream of data is collected by a simple USB, and then sent to a DAQ/DQM system and to an unpacker, which takes into account the detector geometry and translates into a physical position the information received from the DAQ, typically only telling that a certain wire has been activated.

Both streams then join at the 1D histogram format level, gathering a collection of all the hits measured for every event. All these hits are then processed and reconstructed into two trajectories, one above and one below the object, using advanced techniques that will also be described in next Chapter. Finally, 4D segments are made available from this reconstruction process, having the position and direction information needed, from which the actual analysis can take place.

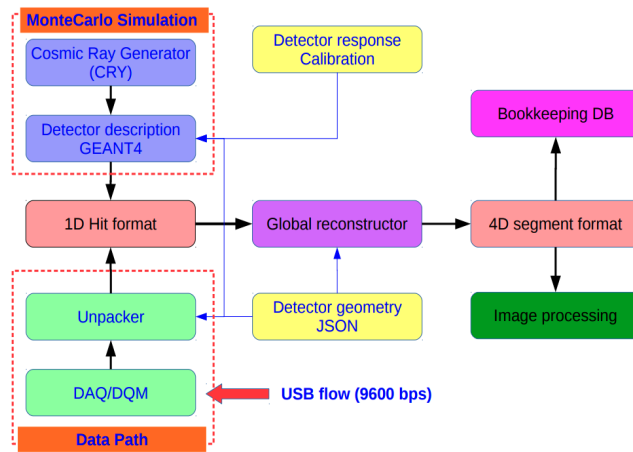


Figure 2.8: Schematic representation of the detector data flow.

Chapter 3

Statistical parameters

The general objective of this work is to find a way to determine the geometry of an eventual object placed between the two muon detectors by measuring the deviation affecting incident cosmic muons, a typical non-deterministic process which therefore needs to be addressed using advanced statistical parameters and methods, such as probability and kernel density functions (Sections 3.1 and 3.4), Monte-Carlo simulations (Section 3.2), p-values (Section 3.3) and likelihood function (Section 3.5).

3.1 Probability density functions

The probability density functions, or PDFs, are statistical expressions defining probability distributions representing the likelihood of any given outcome. They are typically represented as curves in the two dimensional space (x, y) , as shown in Figure 3.1, in which the total area below the curve in an interval is equal to the probability of a discrete random variable occurring.

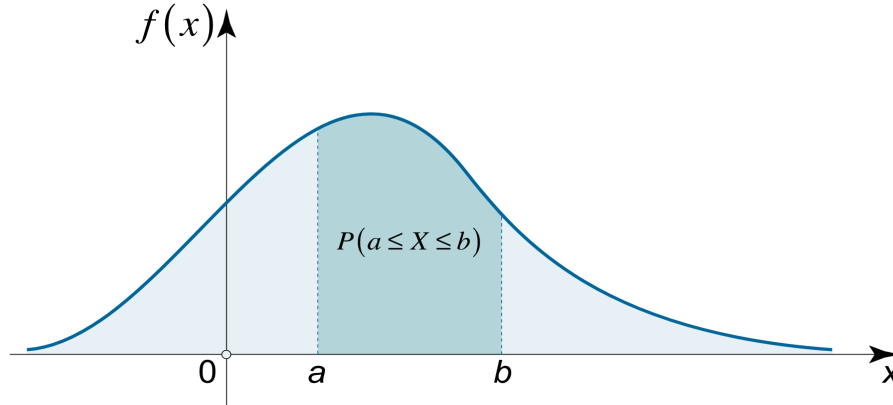


Figure 3.1: Schematic representation of a random PDF [15].

Gaussian distributions are in this sense probability density functions having a particular expression defined in Equation 3.1, where σ is the so-called **standard deviation** of the distribution, representing its width: a low σ indicates that the values of the distribution tend to be close to its mean value μ , while a high σ indicates that the values are spread out over a wider range.

$$f(x) = \frac{1}{\sqrt{2\pi}\sigma} e^{-\frac{1}{2}\left(\frac{x-\mu}{\sigma}\right)^2} \quad (3.1)$$

These definitions are important because we already know that the multiple scattering process which affects the incident cosmic muons is a stochastic process: this means that two muons having similar kinematics can leave the detector with very different output directions and positions, whose actual distribution can be approximated by a Gaussian function for small deviation angles (at larger angles, the distribution is behaving like Rutherford scattering, with slightly larger tails).

The standard deviation of the distribution of such output parameters is actually highly dependant on the geometrical parameters θ_i of the eventual object put between the two detectors, as we saw in Equation 2.2. Indeed, the mean deviation observed is expected to depend on the number of radiation lengths crossed, typically much higher when a muon crosses denser objects, such as a steel pipe, as shown in Figure 3.2. This means that in first approximation, the thicker the object investigated is, the higher the expected deviation will be and this simple observation will actually be used as the driving process of the statistical study performed in this work.

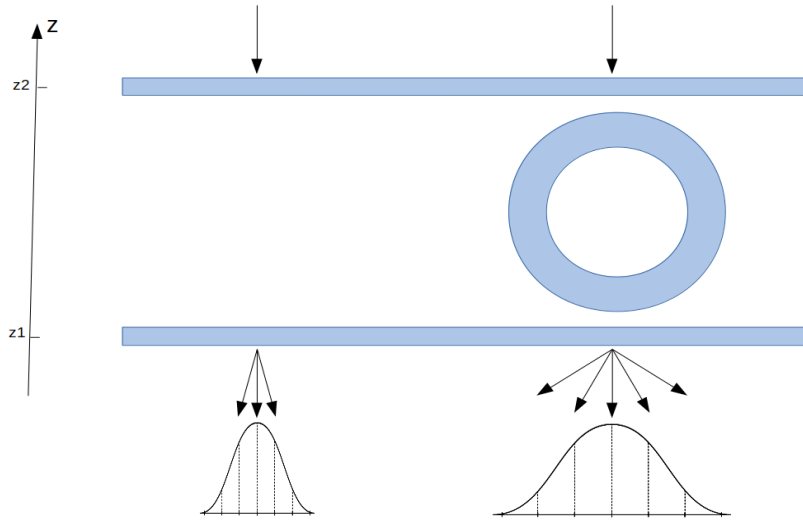


Figure 3.2: Schematic representation of the deviation expected for an incoming muon and PDFs associated without (on the left) and with (on the right) an object placed between the detectors.

3.2 Monte-Carlo simulations

Monte-Carlo simulations are algorithms developed in order to compute approximate numerical values to stochastic problems using random processes and advanced probabilistic techniques. In this particular case, such computation methods apply extremely well since, by being a stochastic process, the actual probability density function associated to a given geometry cannot be computed analytically (and if we were able to do it, a simple gradient descent method would be enough to solve this problem and find the optimal parameters θ_i of the object).

The only way available to actually estimate these parameters is by using thousands of Monte-Carlo simulated toys, simulating thousands of different incident muons for different geometries. The principle is simple: for each incident muon measured, a large number N of Monte-Carlo simulations will be performed, propagating the muon using an algorithm described in Chapter 4 until reaching the lower detector. Repeating this experiment over and over again always gives different results, giving us the opportunity to build the expected PDF for a given object geometry and for the incident muon that has been measured and that will be used later on.

At the end of the day, the main objective of this process is to be able to estimate the probability

of observing a certain deviation given an object geometry and the input trajectory of the muon, $P(\text{deviation}|\text{input})$ in simulation. Then, this process will be reversed using actual data in order to try and obtain the geometry of the object from the actual measurement of the input and output muon trajectories and positions, using a likelihood minimization method.

3.3 p-values

So far, we developed a technique allowing us to reconstruct the expected probability density function for a given incident muon using Monte-Carlo techniques. However, we still need a way to estimate the goodness of the actual measurement with this simulated PDF: this is where the so-called **p-values** enter, a key ingredient for this work since they allow to relate the two main parts of the problem: the simulations performed and the actual data collected.

Probability values, or p-values, are based on the concept of **null hypothesis** H_0 , a general statement or default position telling that there is no actual relationship between two measured phenomena and assumed to be true until an evidence indicates otherwise. This hypothesis is then typically opposed to the **alternative hypothesis** H , usually more interesting by being the interesting hypothesis of the experiment being performed. The main goal is to compare the data with both hypotheses: if the data is consistent with H_0 , then the null hypothesis can simply not be rejected. However, we can reject H_0 (and therefore accept H , its exact opposite) if the data collected is significantly unlikely to have occurred if the null hypothesis were to be true, according to a **confidence level** previously defined.

From these definitions, a p-value is then defined as the probability for a variable to be observed equal or more extreme (lower or higher, depending on the cases) than the actual value x observed under the null hypothesis. This definition, can be represented in a schematic way in Figure 3.3.

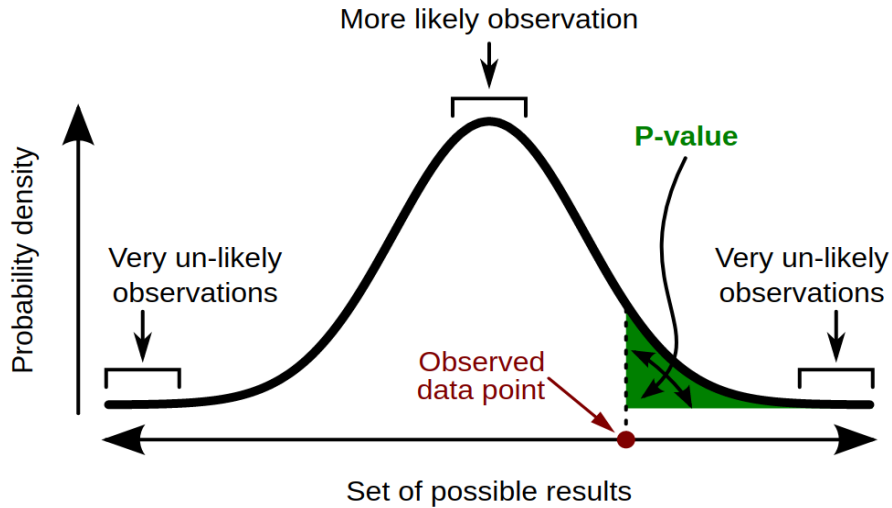


Figure 3.3: Schematic representation of the statistical concept of p-value.

In general, a smaller the p-value tells us that the hypothesis under consideration may not adequately explain the observation, leading to a higher statistical significance. At the end of the day, the null hypothesis is rejected if the p-value obtained is smaller than a previously defined, arbitrary and fixed threshold α , the **level of significance** of the experiment. This parameter can take a large range of values, typically ranging from 0.05 to 0.001.

In this work, the p-value will then allow us to compare a single data measurement with the probability

density functions obtained from Monte-Carlo and quantify the probability that such muon was actually observed for a given object geometry. This value will then be fed to a likelihood function, as explained in Section 3.5.

3.4 Kernel density estimations

In statistics, the kernel density estimation is a method allowing to estimate the shape of the density probability f , the PDF defined in Section 3.1, of a random variable \mathbf{x} , from N observations drawn from this unknown function PDF. The usual way to proceed is to place each observation made in an histogram, where the density in each point x can be therefore estimated as the proportion of observations close to x . However, this method depends strongly on the binning used and the function obtained by this process is non-continuous by definition. The **kernel method** has been developed in this context, to try and solve the non-continuity of the PDF obtained, by simply using various continuous functions instead of bins.

Mathematically, we can define a function called the **density estimator** of f , $\hat{f}_h(x)$, defined as a sum of such general functions K over the number of observations drawn, as shown in Equation 3.2.

$$\hat{f}_h(x) = \frac{1}{Nh} \sum_{i=1}^N K\left(\frac{x - x_i}{h}\right) \quad (3.2)$$

In this last equation, two parameters that need to be defined by hand before using this method are actually extremely important:

- First, the **kernel** K , is a non-negative window function defined depending on the problem, as the "building block" of the final function that needs to be estimated: indeed, as we can see, the function $\hat{f}_h(x)$ will be given as a simple sum of these kernels. In most of the problems, including this work, a standard gaussian kernel is used, taking the particular shape shown in Equation 3.3.

$$K(x) = \frac{1}{\sqrt{2\pi}} e^{-\frac{1}{2}x^2} \quad (3.3)$$

- The **bandwidth** h is another interesting parameter, affecting the final smoothness of the resulting curve. The optimal selection of such parameter is not obvious at all [16] but, as a rule of thumb, it can be shown that for Gaussian kernels of standard deviation σ , it should be approximately equal to $1.06 \sigma N^{-1/5}$ in order to minimize the mean integrated square error obtained.

A comparison of the different methods so far and different possible is shown in Figure 3.4.

In summary, this method is used in order to model the distribution of an arbitrary input dataset as a superposition of Gaussian kernels, one for each data point, each contributing $1/N$ to the total integral of the PDF. This method allows us to construct continuous functions and recover for the eventual low statistics and possibles holes in the bi-dimensional histograms for the position and direction of the muons obtained in x and y , following the method that will be described in Chapter 4.

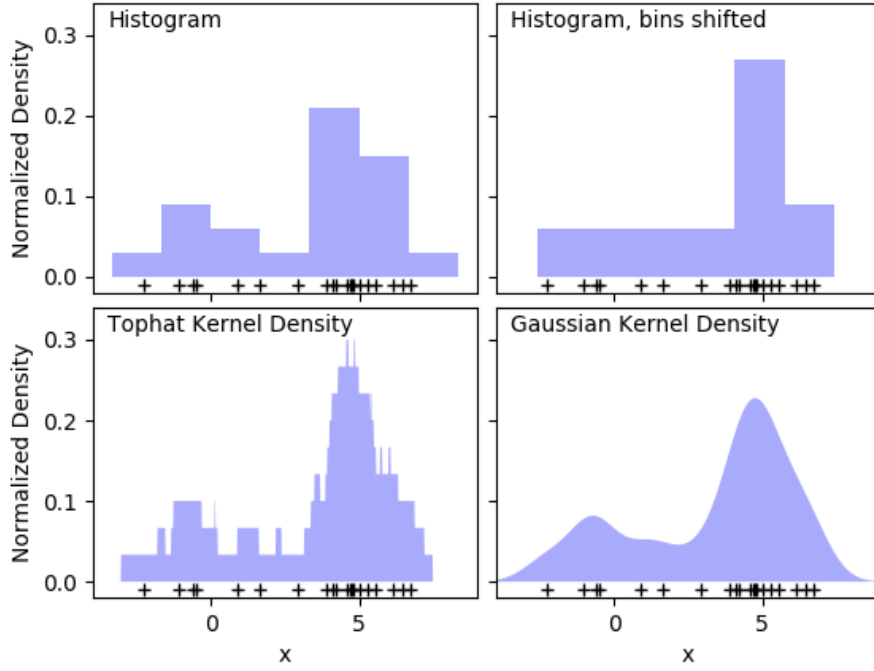


Figure 3.4: Comparison of the histogram and the kernel density estimations methods for different kernels in order to estimate a PDF [17].

3.5 Likelihood

So far, we have been able to generate thousands of different Monte-Carlo experiments for a given muon and for a given object geometry, but we still need to find a way to reverse this process to estimate the geometry of the object from a given measurement.

This is where the likelihood becomes useful, defined as a function that measures the goodness of a fit with respect to a sample of data, for several unknown parameters of a mathematical model. In general, this function can be defined in Equation 3.4, where $\boldsymbol{\theta} = \{\theta_1, \dots, \theta_i\}$ are the parameters of the model and x is the actual measurement of a random variable X following a density function f .

$$\mathcal{L}(\boldsymbol{\theta}|x) = f_{\boldsymbol{\theta}}(x) = P(X = x|\boldsymbol{\theta}) \quad (3.4)$$

It is important to note at this point that the likelihood is a function of the parameters $\boldsymbol{\theta}$ but not a probability density function itself. Additionally, it should in general not be confused with the probability $p(\boldsymbol{\theta}|x)$ since it is equal to the probability of observing a given outcome x when the true values of the parameters are $\boldsymbol{\theta}$: this means that the likelihood is equal to a probability density over the outcome x , not over the parameters $\boldsymbol{\theta}$.

To understand this concept better, the simple example of an unfair coin toss can be considered, where the fairness of the coin is the parameter of the model, represented by the probability of obtaining head p_H and taking values between 0 and 1 (for a regular coin, $p_H = 0.5$). The likelihood is then defined

from a given observation, such as obtaining two heads in a row HHT in the following way:

$$\mathcal{L}(p_H|HHT) = P(HHT|p_H) = P(H|p_H) \cdot P(H|p_H) \cdot P(T|(1 - p_H)) \quad (3.5)$$

For each value of p_H , the value of the likelihood can then be computed and ultimately plotted, as shown in Figure 3.5, showing a minimum value for a particular value of p_H .

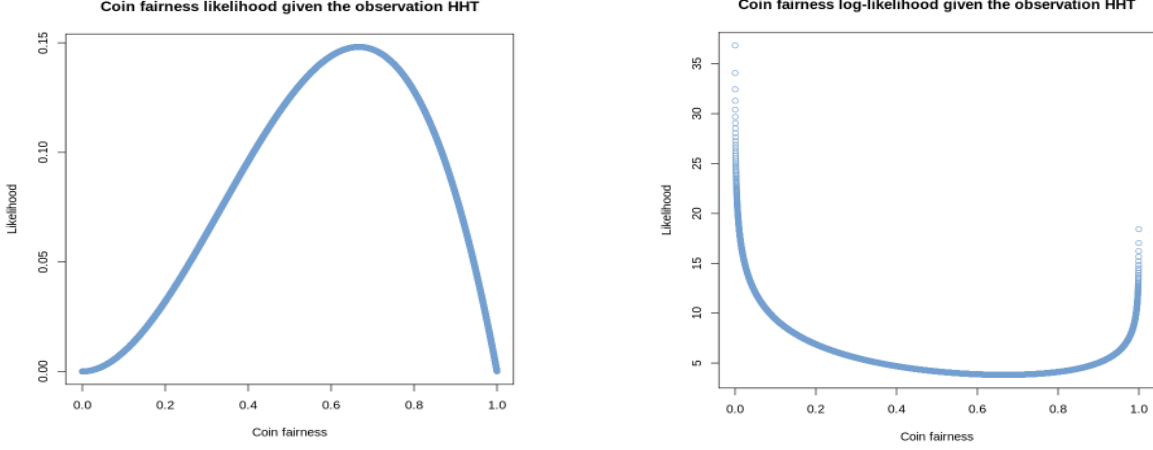


Figure 3.5: Likelihood (on the left) and log-likelihood (on the right) obtained for our particular example of the fairness of a coin given the observation HHT .

The likelihood is in this sense interesting because it can be described as a hypersurface whose peak gives the optimal set of parameters maximizing the probability of drawing the actual sample measured. Often, the **log-likelihood** $l(\theta|x) = -2\log(\mathcal{L}(\theta|x))$ is actually used instead, being a bit more convenient to deal with the special part concavity plays in the maximization process. Given the properties of the logarithm, maximizing the likelihood is equivalent to minimizing this particular log-likelihood.

Another important property of the likelihood that will be used extensively and already assumed in the previous example is the fact that, for two independent measurements, the likelihood of both measurements is equal to the product of both likelihoods independently computed for each event.

In this work, a likelihood function will be obtained as the product of the thousands of muons collected, but for a fixed geometry of the object. A minimization method will then be performed in order to try and minimize the value obtained by trying out different parameters θ , until finding the optimal set of parameters able to describe best the object put in between the detectors.

Chapter 4

The algorithm

The algorithm we developed to solve this particular problem has been written in C++ and can be divided in two categories. First of all, we have the so-called **PipeReconstructor** (Section 4.1), a set of classes able to perform several tasks:

- Define the geometry of the problem (the detector, and the volume under investigation are therefore defined using the smallest possible set of parameters at this stage);
- Compute the different intersection points of an incident cosmic muons with our geometry;
- Propagate the muons through the different medium of interaction encountered along their path;
- Calculate the likelihood of a given measurement for a given volume;
- Finally, probe different volumes and perform a descent method in order to minimize the likelihood encountered and find the optimal volume to solve our problem given the measurements collected.

This class is complemented with the **Generator** (Section 4.2), a small class allowing us to generate muons simulating our actual experiment, using some of the functions of the PipeReconstructor.

4.1 PipeReconstructor

The PipeReconstructor is a set of classes, allowing us to define different kind of objects:

- First of all, a muon is defined from the **MuonState** class and from 7 parameters: 3 parameters (x, y, z) representing the position of the muon, 3 additional parameters (v_x, v_y, v_z) representing its three-dimensional direction and the actual value of its momentum. The measurement performed at the top detector will be labelled with the index 1 $(x_1, y_1, z_1, v_{x1}, v_{y1}, v_{z1}, p_1)$ while the measurement performed by the lower detector will take the index 2.
- **Surfaces** and **Volumes**, virtual classes allowing us to define general surfaces (their spatial position and their geometrical center) and volumes, defined as vectors of several surfaces: a general volume is in this sense defined by its surfaces, central position and global density. Defining such virtual classes is extremely important since it would allow us to consider different geometries, not just a pipe put in between the two detectors. The constant size ($1m^2$) and z position of both the lower and upper detectors (-37 and 37 cm, respectively) is also defined at this stage.

- **Cylinders** and **Pipes**, particular subclasses of the virtual classes previously defined in order to define the geometry of this particular problem. A pipe is in this sense defined as a set of two cylinders and from 7 different parameters: its central position (x, y, z) , along with its density, the radii r and R of both cylinders and its length along the y-axis L , as shown in Figure 4.1. For simplicity, the pipes considered in this work are therefore considered to have a constant density.

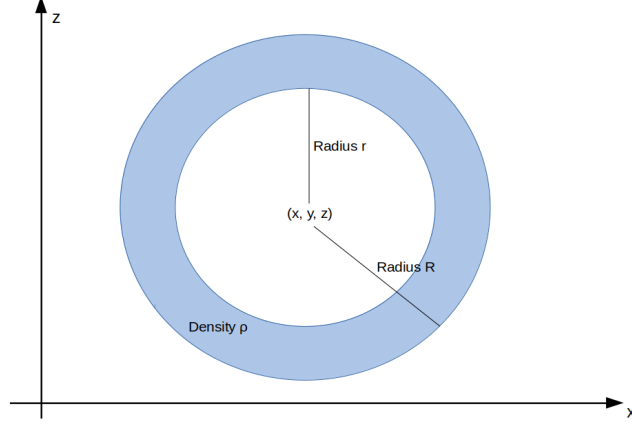


Figure 4.1: Schematic representation of the pipe considered in this work and its respective parameters.

One method computing the exact cut point between a `muonState` and a `Cylinder` has been defined at this point as well, from the general equation of a cylinder of radius R centered in $(X, Y = 0, Z)$ (given that the y-axis is defined along the cylinder, we can assume that the cylinder is actually centered in $Y = 0$ without any loss of generalization).

$$(x - X)^2 + (z - Z)^2 = R^2 \quad (4.1)$$

From this, computing the intersection between such cylinder and a general `MuonState` $(x_0, y_0, z_0, v_x, v_y, v_z, p)$ is trivial, since a `MuonState` can be represented by a straight line:

$$\begin{cases} x = x_0 + \lambda v_x \\ y = y_0 + \lambda v_y \\ z = z_0 + \lambda v_z \end{cases} \quad (4.2)$$

By putting this into the equation of the cylinder, we get the following expressions:

$$(x - X)^2 + \left(z_0 + \frac{v_z}{v_x}(x - x_0) - Z \right)^2 = R^2 \quad (4.3)$$

Once both the squared values applied, we get an equation having the shape $Ax^2 + Bx + C = 0$, where both cutting points will be given therefore by $x^{+/-} = \frac{-B \pm \sqrt{B^2 - 4AC}}{2A}$:

$$\begin{cases} A = \left(1 + \left(\frac{v_z}{v_x} \right)^2 \right) \\ B = -2X - 2x_0 \left(\frac{v_z}{v_x} \right)^2 + (2z_0 - 2) \left(\frac{v_z}{v_x} \right) \\ C = X^2 + Z^2 + z_0^2 - 2z_0Z + x_0^2 \left(\frac{v_z}{v_x} \right)^2 + x_0(2Z - 2z_0) \left(\frac{v_z}{v_x} \right) - R^2 \end{cases} \quad (4.4)$$

Finally, methods allowing us to determine whether a muon state is inside or outside of the pipe have also been written, taking into account the particular geometry of such objects in order to determine the medium in which the muon is propagating and therefore use the correct density value in all the calculations performed.

- A **Propagator** is then defined as the object allowing us to propagate a MuonState through a general Volume. The way it works is quite simple:
 - First of all, the distance between the initial MuonState and all the surfaces of the object considered is computed and the first intersection point is kept.
 - The muon is then propagated in a distance corresponding to 90% of the total distance to this first cut point by taking into account the multiple scattering effect and the Molière theory, and a new MuonState is therefore obtained. We cannot propagate the muon 100% of the distance immediately to avoid effects related to the actual shape of the Cylinder and due to the fact that the Equation 2.2 is only valid in case of normal incidence between the MuonState and the Cylinder, which is typically not verified, as shown in Figure 4.2.

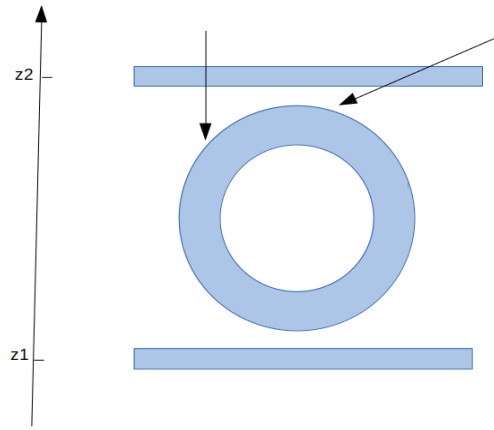


Figure 4.2: Typical example of two incident muons not respecting the normal incidence required by the Molière equation for the multiple scattering of muons.

- This propagation process is then several times, until reaching a position extremely close to the volume that is being studied (our tolerance parameter has been fixed to 0.1 mm). The muon is then manually propagated along its direction of propagation along a distance of 0.1 mm in order to force its crossing with respect to the first surface, so that the propagation can be repeated for the next surface.
- This whole process is repeated for all the different surfaces encountered by the muon (up to four times in total for our particular geometry), and one final time until reaching the bottom detector. At this point, the MuonState is kept as the second measurement ($x_2, y_2, z_2, v_{x2}, v_{y2}, v_{z2}, p_2$) that will be compared with the simulation using statistical methods.
- Eventual muons out of the acceptance of any of the detectors are of course rejected at this stage and simply not considered when computing the value of the likelihood.
- The **Likelihood** object itself can then be calculated. This class takes as input a pipe geometry and a file containing a set of measurements (either obtained from Monte-Carlo or actually measured by the detectors) and its main objective is to estimate how likely it is that we obtained exactly these measurements for the geometry given.

To reach this goal, several calculations are performed for each event found in the input file:

- First of all, three MuonStates are computed: the so-called *IncomingMuon* and *OutgoingMuon*, directly read as the measurement at the upper and lower detectors from the input file, and the *IncomingMuonLinearDown*, defined as the prolongation along a straight line of the *IncomingMuon*. Each *IncomingMuon* is then propagated along the geometry using our Propagator, defining a fourth *IncomingMuonDown* MuonState, as shown in Figure 4.3.

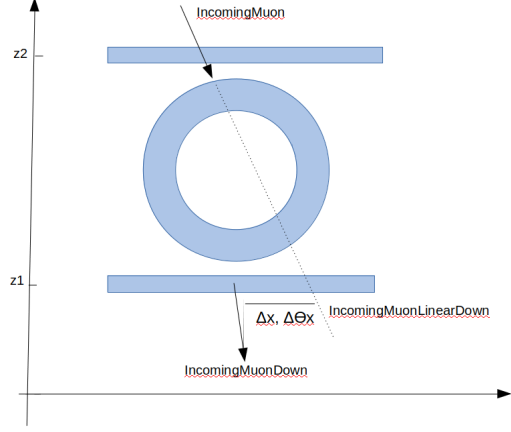


Figure 4.3: Schematic representation of the 4 four muonStates and main parameters used for the computation of the likelihood.

- From these muonStates, the Δx , $\Delta\theta_x$, Δy and $\Delta\theta_y$ parameters can easily be computed as the difference between the positions and angles measured between the *IncomingMuonDown* and *IncomingMuonLinearDown* muonStates. Such parameters are important and will be used a lot in this work since they represent the actual deviation suffered by the muon.
- This propagation process is repeated multiple times ($n_{\text{prop}} \sim 100 - 500$ times in total) since we know that the multiple scattering is a stochastic process. In each iteration of the loop, two bi-dimensional histograms (Δx vs $\Delta\theta_x$ and Δy vs $\Delta\theta_y$) are then filled with the actual values of the deviation parameters previously obtained.
- Keeping in mind that this process needs to be repeated for every single event found in the input file, and given the relatively high number of iterations needed to get statistically meaningful results for each event, the kernel density estimation method introduced in Section 3.4 becomes really helpful at this point. Indeed, this method allows us to keep the number of iterations (and therefore the computing time) to a minimum, while still obtaining smooth enough results to compute p-values in a reliable way.
- At this stage, all the ingredients needed to compute a p-value are available. A p-value for each *IncomingMuon* is then computed by using the bi-dimensional histograms previously defined and representing all the possible deviation values for a single event, and by comparing the actual measurement *OutgoingMuon* to such distributions.
- Two p-values, one along each of the x and y axes, are then obtained for each of the N events found in the input file. The total value returned by this function can then simply be computed from Equation 4.5.

$$\mathcal{L} = \sum_{i=1}^N -2.0 \cdot \log(\text{pvalue}_{x,i}) \cdot \log(\text{pvalue}_{y,i}) \quad (4.5)$$

4.2 Generator

As we have just seen, the computation of the p-values required for the likelihood takes as input a file of Monte-Carlo generated events, which has been previously produced thanks to this **Generator**. This class takes as input one rootfile containing the Monte-Carlo generation of incoming cosmic muons previously produced by the company using CRY [18], a dedicated generator for cosmic muons. To simplify the current problem, the goal is then to simply read the input parameters of such muons, such as the position of their impact with the upper detector, and to propagate them using the previously defined functions, to generate a distribution of the expected output.

The way this algorithm works is quite simple: a general volume matching the one that needs to be studied (in this case, a general Pipe located in the center $(0, 0, 0)$ of the detector, having an inner radius variable, an outer radius of 20 cm and a length of 50 cm) is first of all defined. Then, for each event we want to generate ($n_{\text{gen}} \sim 100.000$), a loop is performed: in each iteration of this loop, the incident muon is propagated throughout the geometry, histograms for the parameters Δx , $\Delta\theta_x$, Δy and $\Delta\theta_y$ are obtained and, more importantly, measurements in the lower detector are artificially generated to test our framework without having to consider detector effects for now.

One last important notion to introduce at this point is the difference between the actual measurement and the Monte-Carlo truth value, that can both be computed by the generator. The existence of both variables is related to the fact that the detector is made out of 200 discrete wires, each separated by a distance of 4 mm, meaning that the actual position measured is typically a discrete number. However, when considering simulation files, both the real and the measured position can be obtained and kept in different MuonStates labelled as $(px_i, py_i, pz_i, pv_{xi}, pv_{yi}, pv_{zi}, p_i)$ and $(mx_i, my_i, mz_i, mv_{xi}, mv_{yi}, mv_{zi}, m_i)$ respectively. Depending on whether we are working with Monte-Carlo simulations or data, one or the other MuonState will then be chosen.

Chapter 5

Results obtained

Now that the experimental setup and main statistical parameters used have been fully described and that the algorithm developed for this particular exercise has been introduced, we can present the results obtained. This chapter is divided into several sections: first of all, we are going to mention the different validation steps performed to make sure that this algorithm's results can be trusted in Section 5.1. Then, more general results obtained for different pipe geometries will be presented in Section 5.2, along with the final likelihood results of this work.

5.1 Generator validation

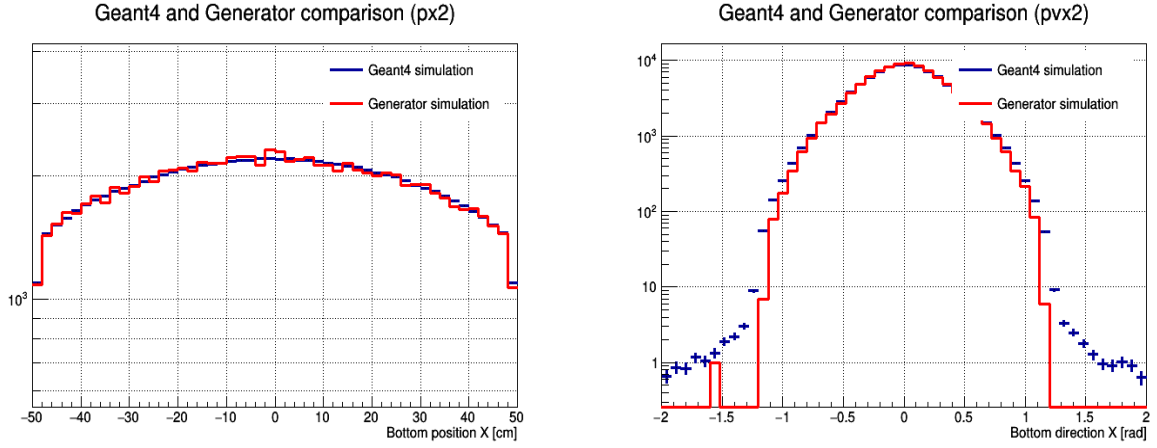
The validation step consists in comparing the results previously obtained with Geant4, a toolkit for the simulation of the passage of particles through matter [19], with a complete and realistic description of the interaction between the cosmic muons and the detectors, and the ones obtained with our Generator, assuming a perfect detector.

This step is actually extremely important, as it allowed us to find several bugs in the code. In this context, several files have been generated using both generators and different pipe geometries, allowing us to perform several different checks. In each case, the objective was to obtain the output MuonState $(x_2, y_2, z_2, v_{x2}, v_{y2}, v_{z2}, p_2)$ after propagating the initial MuonState throughout a given geometry.

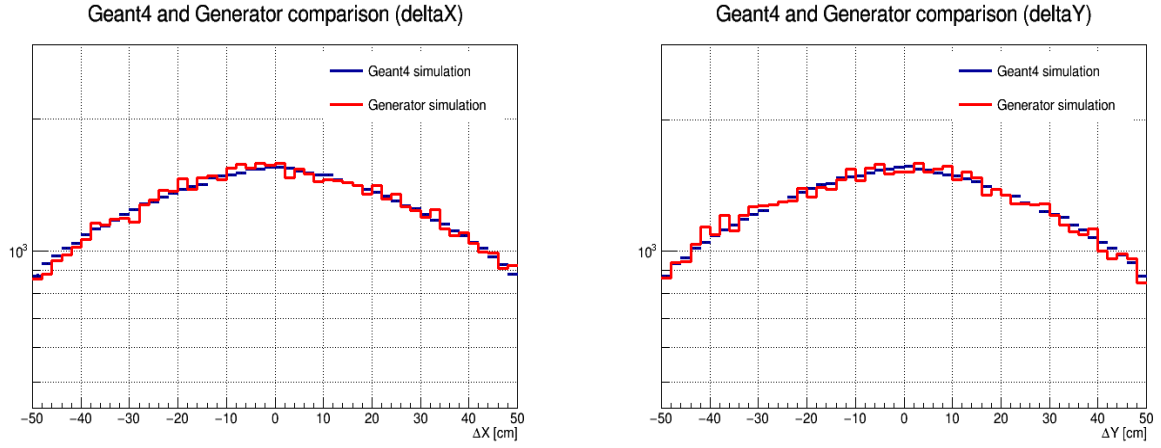
For a given pipe having a geometry $(r = 17.2 \text{ cm}, R = 20 \text{ cm}, L = 50 \text{ cm})$ and located in $(0, 0, 0)$, thousands of muons have then been generated using Geant4 and our Generator. The output distributions (the positions x and y and directions v_x and v_y after propagation, along with the deviations in position Δx and Δy , defined in Equation 5.1 where d is the vertical distance between the two detectors, and the deviation angles $\Delta\theta_x$ and $\Delta\theta_y$, defined in Equation 5.2) have then been obtained in this case in Figure 5.1. All these plots have been normalized to unity in order to account for the eventual different number of events generated in both cases.

$$\begin{cases} \Delta x = x_2 + d(v_{x2} - x_1) \\ \Delta y = y_2 + d(v_{y2} - y_1) \end{cases} \quad (5.1)$$

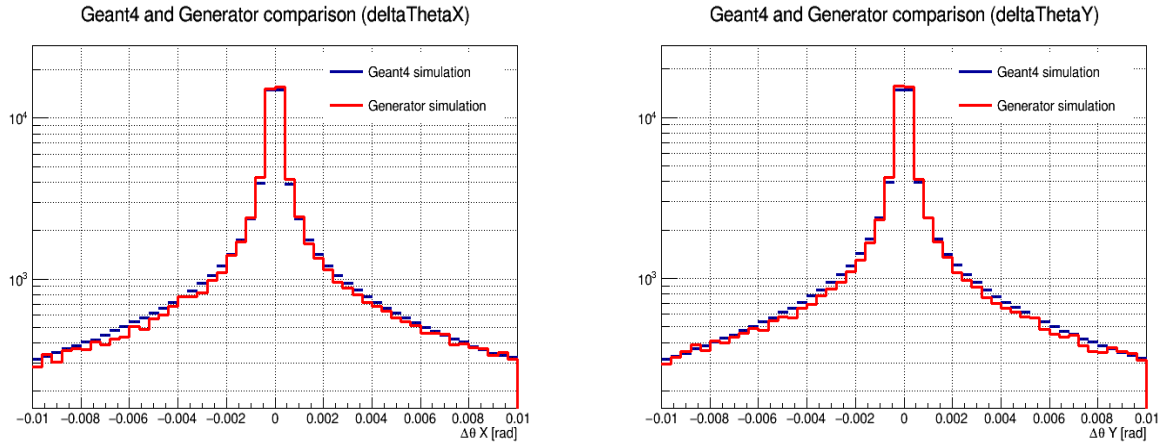
$$\begin{cases} \Delta\theta_x = \arctan\left(\frac{v_{x2}}{v_{z2}}\right) - \arctan\left(\frac{v_{x1}}{v_{z1}}\right) \\ \Delta\theta_y = \arctan\left(\frac{v_{y2}}{v_{z2}}\right) - \arctan\left(\frac{v_{y1}}{v_{z1}}\right) \end{cases} \quad (5.2)$$



(a) Muon position (on the left) and direction (on the right) along the x-axis at the bottom detector after propagation



(b) Muon Δx (on the left) and Δy (on the right) as measured after propagation



(c) Muon $\Delta \theta_x$ (on the left) and $\Delta \theta_y$ (on the right) as measured after propagation

Figure 5.1: Normalized variables measured on the bottom detector using the Geant4 (in blue) simulation process and our custom Generator (in red).

5.2 General results

5.2.1 Pipe geometries impact

Now that we know we can trust the results obtained from the Generator, the next step has then been to generate 9 Monte-Carlo files with 100.000 events each, using our Generator, corresponding to 9 possible pipe geometries characterized by a constant length $L = 50\text{cm}$, an outer radius of $R = 20\text{cm}$ and different inner radii r , ranging from 17.2 to 18.8cm, by steps of 0.2cm. Each file was generated with 100.000 events, taking a few minutes per geometry considered.

We can then start comparing the Δx , Δy , $\Delta\theta_x$ and $\Delta\theta_y$ distributions obtained for these different geometries, as shown in Figure 5.2.

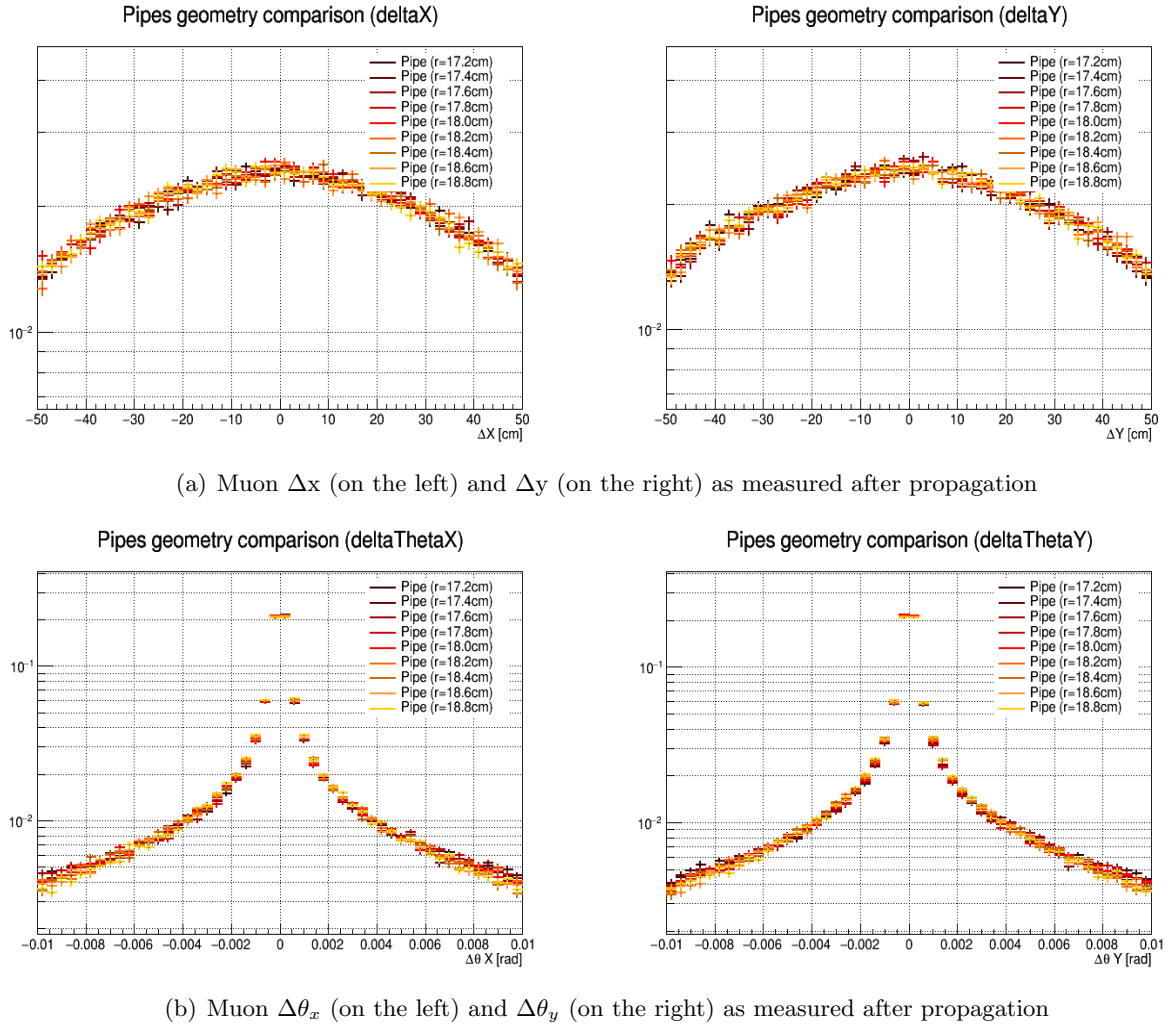


Figure 5.2: Normalized deviation variables generated using different pipe geometries.

As we can see, the geometry actually has only a little impact on the deviation observed in the position and direction, along both the x and y axes. We can however observe slightly larger tails for the pipes having a larger inner width, as expected. The study of these small differences between the different geometries will be the starting point of the determination of the properties of an unknown pipe placed between the two detectors using our Likelihood class.

5.2.2 Kernel density functions

As explained previously, we generated 100.000 events for each geometry with our Generator and, for every single event, a likelihood needs to be computed, which involves running another loop of at least several hundreds of iterations. In each iteration, the value of the deviation in position and angle between the MuonState obtained with our Propagator and the one defined as the linearly propagated MuonState to the bottom detector is computed, filling in this process two bi-dimensional histograms (Δx vs $\Delta\theta_x$ and Δy vs $\Delta\theta_y$) for each input event.

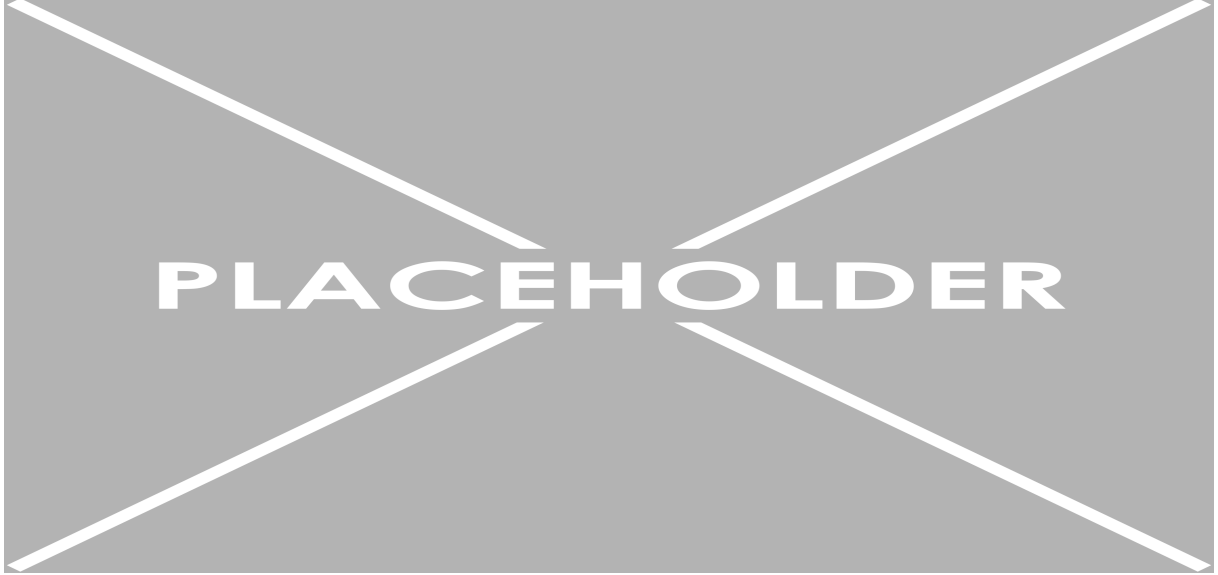
More precise results are expected to be obtained by increasing this number of iterations, but so does the computing time, so in order to keep the code running time manageable, we decided to use the kernel density estimation method described in Section 3.4. Both the original and the smoothen histograms obtained for a single event can be found in Figure 5.3.

5.2.3 Likelihoods

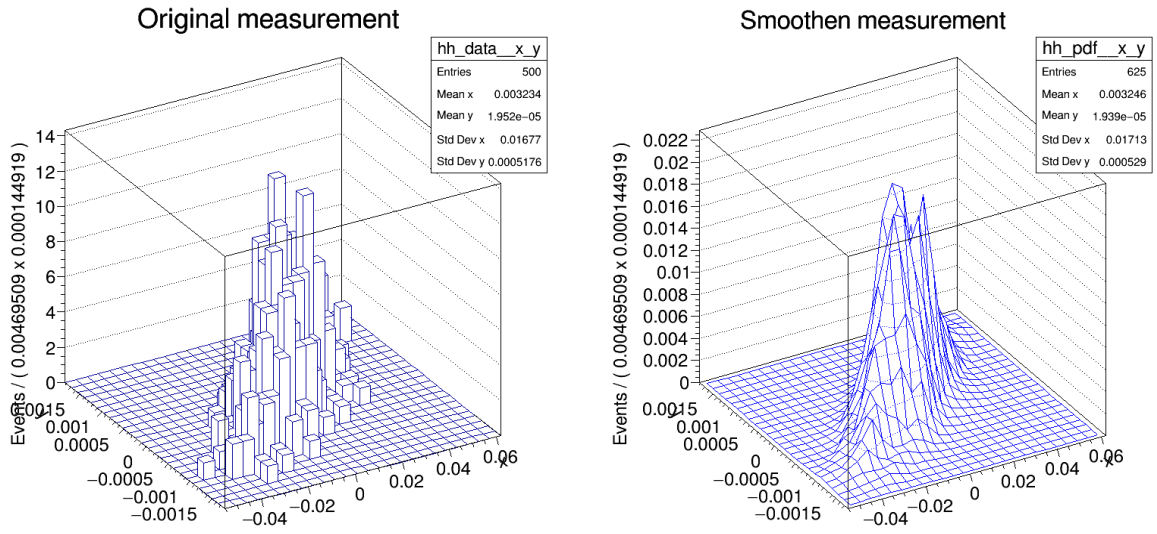
The goal is now to estimate the value of the total likelihood obtained for different pipe geometries, characterized by different inner radii, ranging from 17.2 to 18.8cm, by steps of 0.2cm.

The idea is then to compute the value of the likelihood obtained for each of these geometries by comparing an actual measurement obtained with our Propagator with these generated files one by one, in order to estimate the p-value in each case and minimize the likelihood obtained in each case to, at the end of the day, try to figure out which pipe geometry is more likely to give rise to such measurements. The results obtained are shown in Figures 5.4 and 5.5 for 100 and 500 iterations used for the computation of the likelihood value, and where we expect to see a minimum for the actual geometry used, represented by the vertical red line in each case.

As we can see in these results, the minimum of the likelihood curve is not always exactly at the place where it should be. However, and this is especially true for larger pipes, we do manage to reach in most of the cases a precision of the order of the step chosen, equal to 0.2cm. This means that using the algorithm developed here, we are capable of determining the actual inner width of a pipe using cosmic muon and without destructing this pipe, the initial objective of this work, with a precision of the order of a few millimeters.



(a) (Δx vs $\Delta \theta_x$) histogram along the x-axis, with (on the right) and without (on the left) smoothening applied.



(b) (Δy vs $\Delta \theta_y$) histogram along the y-axis, with (on the right) and without (on the left) smoothening applied.

Figure 5.3: Original and smoothen bi-dimensional histograms along the x and y axes, obtained for a single event generated with the Generator class.

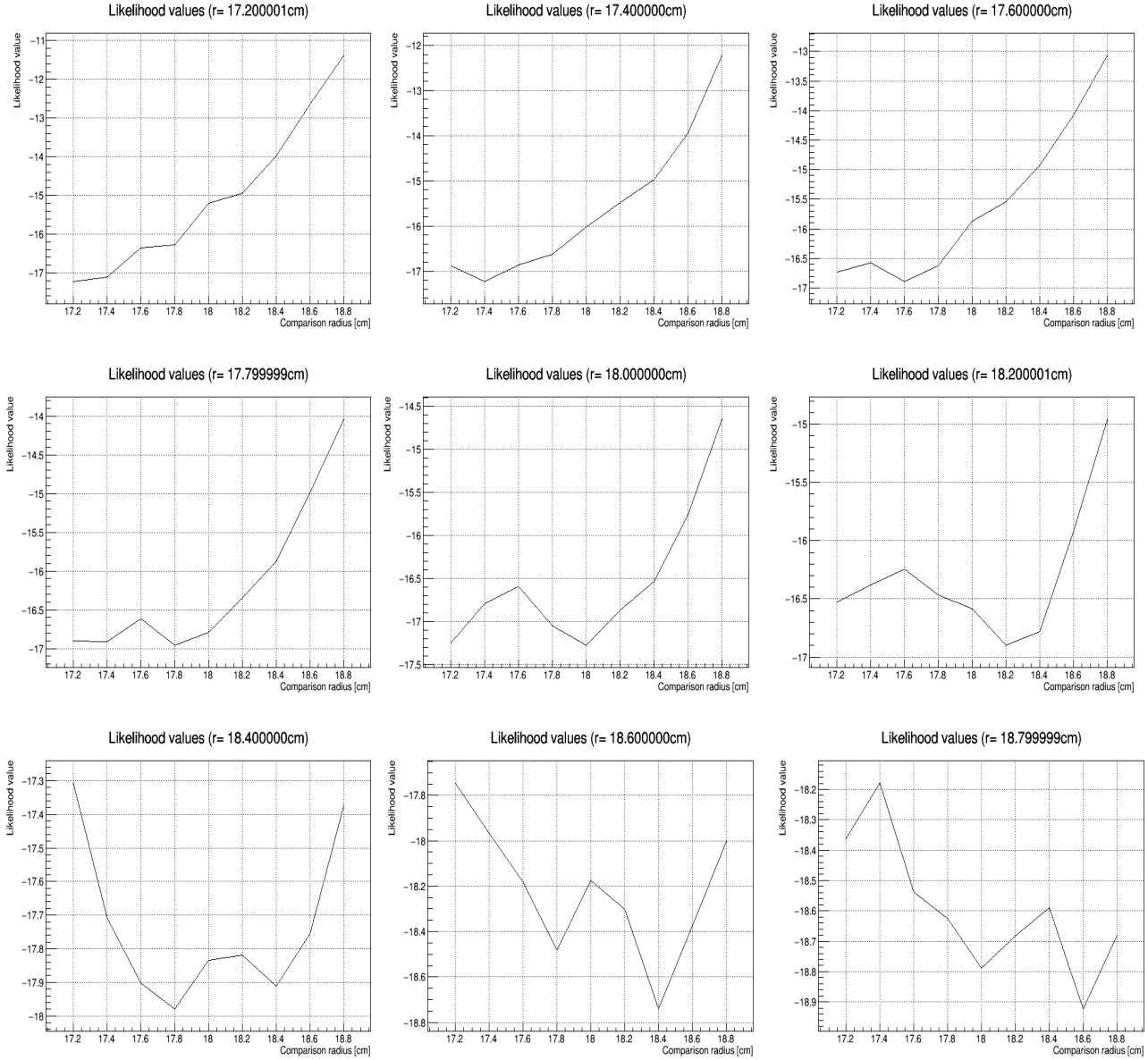


Figure 5.4: Likelihoods values obtained for different pipe geometries, ranging from 17.2 (on the top left) to 18.8cm of inner radius (on the bottom right), obtained for 100 iterations.

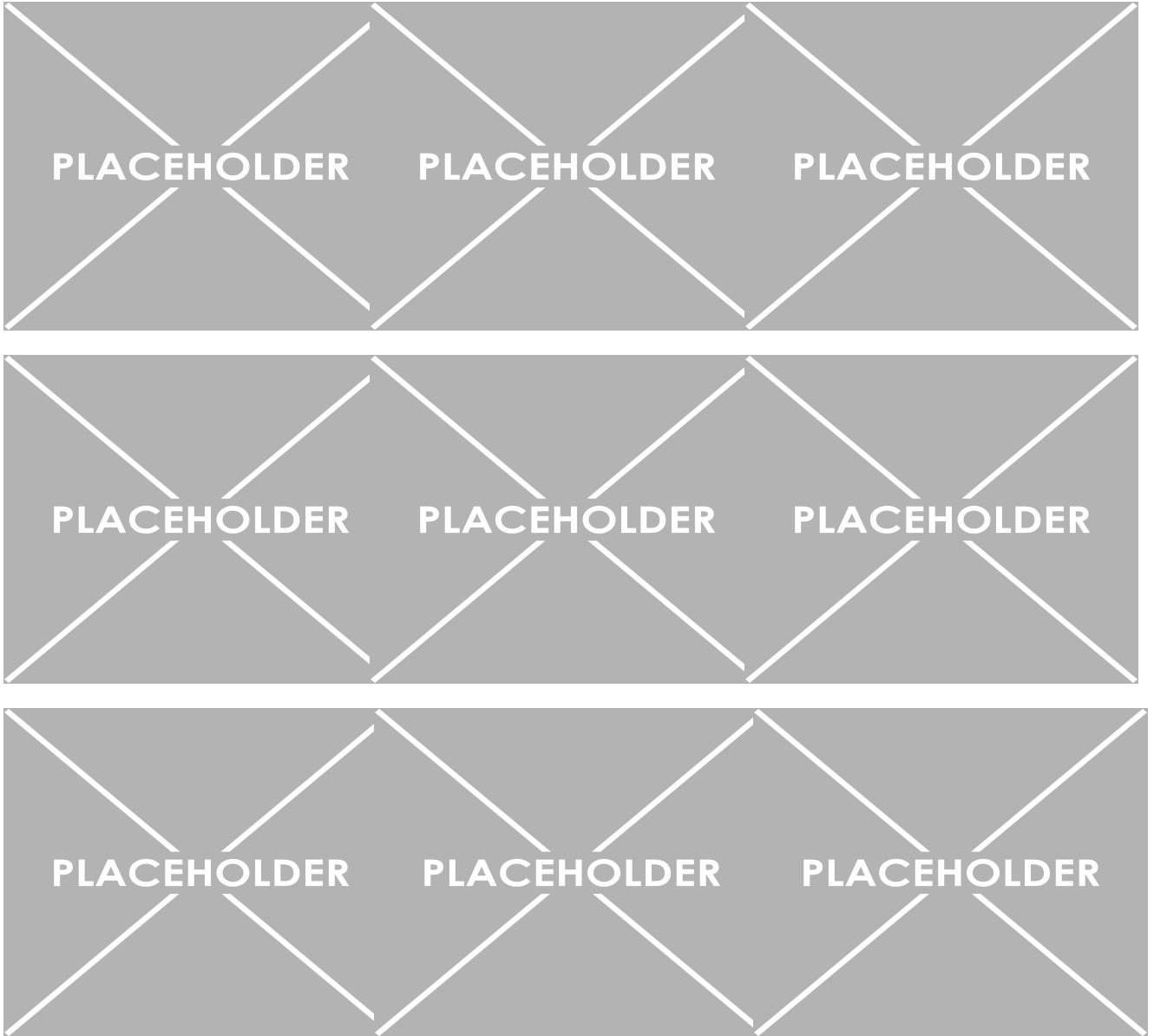


Figure 5.5: Likelihoods values obtained for different pipe geometries, ranging from 17.2 (on the top left) to 18.8cm of inner radius (on the bottom right), obtained for 500 iterations.

Chapter 6

Conclusions

In conclusion, this work was designed with one simple objective in mind: using the muon tomography technique in order to determine the inner properties of a general geometry (usually considered here to be a pipe made out of two cylinders, even though the code has been developed to be able to simply add more complex geometries in the future), using statistical methods to analyze the results obtained.

The principle of this statistical approach presented here is quite simple: in order to study the inner properties of a steel pipe, we can put a detector above the pipe, measuring the initial position and directions of cosmic muons, and another detector below the volume considered, to compute the actual deviation suffered by the cosmic muons when traversing the gap between the two detectors. Since this deviation in position and angle is expected to be larger when traversing the steel pipe than the air, given the difference in density of these two medium, we expected to be able to estimate the actual width of the pipe crossed from such measurements.

By developing our framework, able to generate Monte-Carlo simulations of such cosmic muons and to propagate them across our geometry using the Molière's theory of multiple scattering, and by implementing advanced statistical methods to solve this problem, we managed to estimate the actual width of such steel pipes with a precision of the order of 2 millimeters, therefore solving the problem which was originally presented. These results have been obtained after comparing our Generator with Geant4, in order to cross-check the results presentend throughout this work.

6.1 Suggested improvements

Of course, this work is only a first approach to this particular problem, and many different improvements could be added in the future in order to improve and/or generalize the results obtained:

- First of all, as previously stated, this work was focused on a typical approach related to steel pipes typically used in the industry in general. The framework is ready to include additional and more exotic geometries, to study them as well, even though this has not been done in this particular work. Even pipes could be quite easily studied in greater details using this framework, by for example considering additional parameters to characterize them, such as a possible dependence of the width of the pipe with respect to the position. Being able to find weak spots in particular locations in the pipe itself would be something of great value for the industry.
- Another interesting idea might be to consider ionization as well. In this work, we did not consider the loss of energy suffered by a muon when crossing the volume studied since it is expected to

be quite low, but such information could help us increase the precision of the measurements performed.

- In this work, we did not consider detector effects even though by interacting with the incident cosmic muons, they have an effect on them, even though this effect is supposed to be quite small. Simulating in a precise way this interaction could also improve the precision of the method.
- Finally, we have been quite limited in this work by the computing power available. Given the fact that the determination of the likelihood required us to do complex calculations, we had to limit as much as possible several parameters, such as the number of iterations performed to compute the likelihood. Increasing this parameter is expected to give better results, but also increase the computing time. Running the code on more powerful machines is therefore also expected to give better results.

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