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# **Statistical approach to muography as a non-destructive testing technique for industry problem solving**

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# Abstract

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# Chapter 1

## Introduction

Nowadays, muon tomography is an active field of research since it is a non destructive method allowing to map the inside of large objects difficult and/or dangerous to access, without any contact or damage, and without even having physical access to them. To do so, this technique uses muons, elementary particles similar to electrons but with a much greater mass, which allows them to penetrate much deeper and probe matter in a more efficient way, since they suffer less from the bremsstrahlung radiation affecting all leptons. Such technology is relatively well-know and present several advantages over other techniques such as X-ray imaging since it is globally safe and clean and uses natural radiation, cosmic muons, while providing an excellent penetration in matter in order to study it.

In this particular work, muon tomography will be applied to industry and used in order to probe several kinds of physical objects, such as pipes, to study their degradation in a non-destructive way. This is the main objective of the company Muon systems, founded in Spain in 2017 and based in Bilbao.

This project is highly relevant in today's society because it allows to combine data science algorithms and high-performing statistical methods that have been developped in this context, putting them in practice in an actual industry. **INSIST ABOUT DATA SCIENCE RELEVANCE**

## Chapter 2

# Muons and muography

The main particles used throughout this work, muons, will be briefly introduced in Section 2.1 of this Chapter. Muons are particles produced naturally by cosmic rays, introduced in Section 2.2 and, once produced, they obviously interact when traversing the atmosphere and matter in general in ways described in Section 2.3 that actually need to be understood extremely well for the muon tomography process introduced in Section 2.4 to be useful and applicable to industry. Finally, the actual experimental setup used for this work will be explained in Section 2.5.

### 2.1 Particle physics and muons

Particle physics is the field which studies the matter surrounding us, along with the fundamental interactions between the particles. In this context, the Standard Model of particle physics [1] is nowadays the most accepted mathematical model used to describe the elementary particles and three of the 4 fundamental forces of nature (electromagnetic, weak and strong interactions, while the gravitational interaction is out of reach of this model). Even though quite simple in concept, it has been able to describe most of the phenomena observed in nature so far with an incredible level of precision, and has made a lot of predictions that have now been proven to be true, such as the discovery of the top quark [2] in 1995, the tau neutrino [3] in 2001 and the Higgs boson itself [4, 5], the last missing piece of the Standard Model, announced at CERN in July 2012.

According to this model, 12 different fermions (along with their 12 corresponding anti-particles) exist in nature, as shown in Figure 2.1, most of them being unstable. These fermions can be divided into two fundamentally different categories, the quarks and the leptons, containing each 6 particles and sensitive to different forces. Even though quite interesting, the quarks do not play a fundamental role in the muon tomography detailed in this work, so only leptons will be considered from now on. In particular, leptons can be divided even more into three different generations of particles, and the muon, one particular lepton belonging to the second generation, will be the main focus of this work.

Muons  $\mu^-$  [9] are therefore fundamental particles having a negative charge and quite similar in nature to electrons, even though they have a quite high mass (200 times larger than the electron), which implies that they are not stable particles: they have a lifetime of approximately  $2.2\mu\text{s}$ , and typically decay into an electron and a pair of neutrinos. However, this lifetime is actually quite long with respect to other fundamental particles and muons are on average able to travel more than 700 meters, allowing us to consider them to be stable particles in many processes, such as the one presented in this work. Muons also have a relatively small interaction cross-section with ordinary matter, even though they do interact with baryonic matter by several processes described in Section 2.3.



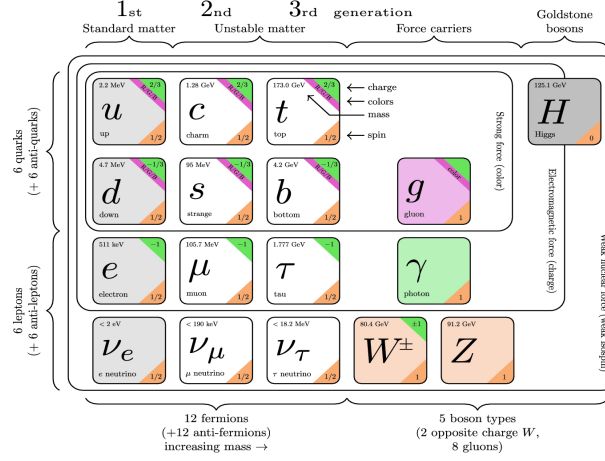


Figure 2.1: Representation of the 12 fermions of the Standard Model [8] along with the main force carriers and the Higgs boson, discovered in 2012 and completing this model.

## 2.2 Cosmic rays

Being unstable by nature, once produced, muons decay almost instantly by a weak process into an electron and a pair of neutrinos. However, it is possible to observe them in nature, since they are continually produced, mainly thanks to cosmic rays [10], constant flux of high energy particles (mostly protons and atomic nuclei) coming mostly from supernovae explosions and AGN emissions and reaching the Earth every day. Indeed, as they impact our atmosphere, these particles start a chain reaction, as shown in Figure 2.2: first of all, several neutral and charged pions are produced, decaying themselves into a pair of photons (and, later on, electron and positron pairs) and muons, respectively.

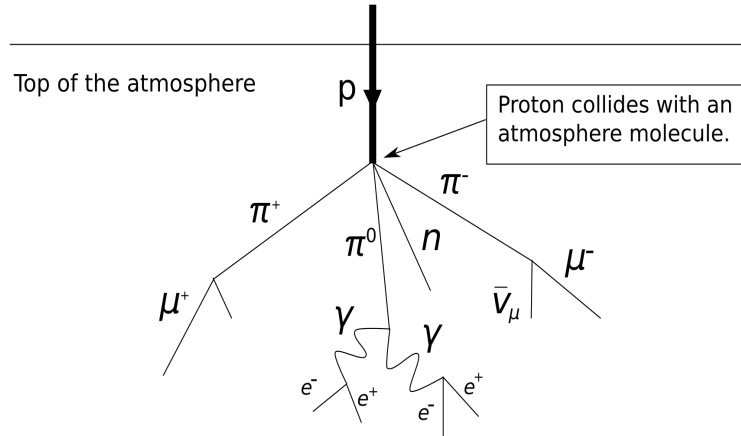


Figure 2.2: Typical chain of decays induced by highly energetic cosmic rays when reaching the Earth's atmosphere.

Muons are the most abundant charged particles produced by these processes actually reaching the sea level, as shown in Figure 2.3. Even though they are unstable and have a limited lifetime, around 0.06% of muons produced by such processes indeed do manage to reach the sea level thanks to the temporal distortion induced by their high energy and relativistic speed. As a rule of thumb, one can expect to observe 10.000 muons per square meter and per minute at the sea level.

Muons were actually discovered thanks to cosmic rays in 1936 [11].

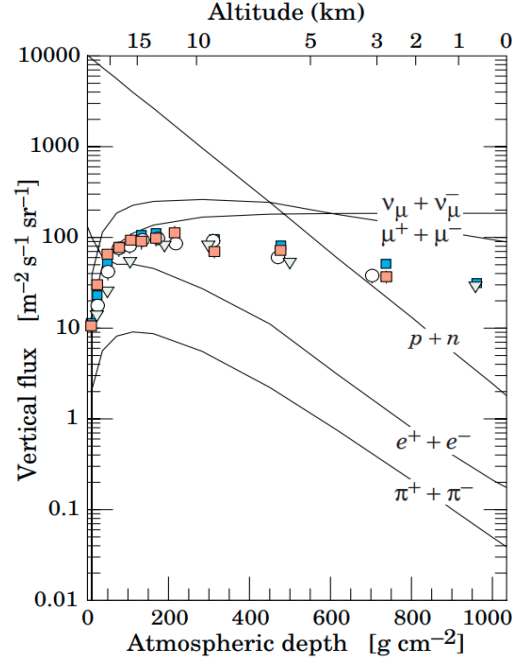


Figure 2.3: Abundance of particles observed at the sea level, due to cosmic rays [10].

## 2.3 Muons interaction with matter

Muons typically interact with ordinary matter through two main processes: ionization and multiple scattering, both resulting in different effects on the incident particle.

### 2.3.1 Ionization process

The most frequent interaction process of cosmic muons is through **ionization**, when the incident muon is giving some of its energy to the electrons of the absorber. This process is described by the Bethe-Bloch formula, shown in Equation 2.1, and describing the average loss of energy over a distance  $\frac{dE}{dx}$  of material, depending on several parameters, such as the charge number of incident particle  $z$ , the atomic mass and charge of absorber  $A$  and  $Z$ , the relativistic factors  $\beta$  and  $\gamma$ , the maximum possible energy transfer to an electron in a single collision  $W_{\max}$  and the mean excitation energy  $I$ .

$$-\left\langle \frac{dE}{dx} \right\rangle = K z^2 \frac{Z}{A} \frac{1}{\beta^2} \left[ \frac{1}{2} \ln \left( \frac{2m_e c^2 \beta^2 \gamma^2 W_{\max}}{I^2} - \beta^2 - \frac{\delta(\beta\gamma)}{2} \right) \right] \quad (2.1)$$

This previous equation gives an accuracy of a few percent in the range  $0.1 < \beta < 1000$  and we can easily see that the quantity of energy lost by in a muon when traversing any given medium actually depends on the energy of the incident muon, as shown in Figure 2.4.

In practical cases, most relativistic particles, such as muons coming from cosmic rays, have mean energy loss rates actually close to the minimum. They are usually called for this reason *minimum ionizing particles* or MIPs.

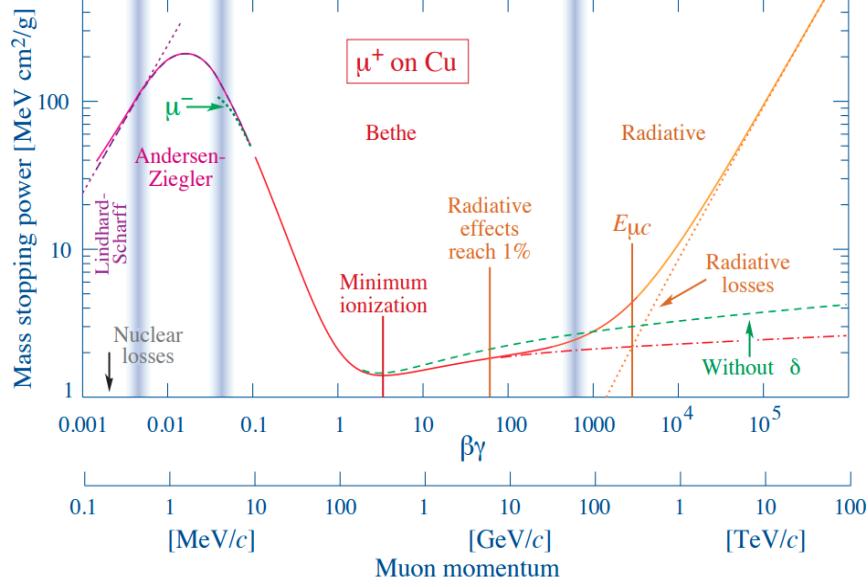


Figure 2.4: Average energy lost a muon due to ionization depending on its momentum [9].

### 2.3.2 Multiple scattering process

Muons also interact with matter through another process, called **multiple scattering**. Since muons have a negative electric charge, by getting close to the nuclei of the absorber, they are suffering from Coulomb scattering. Given the high number of nuclei in matter, this process is repeated many times, deflecting each time the muon by a small angle in a stochastic way, meaning that there is no way of calculating this deviation exactly, but only using probabilities and the so-called theory of Molière [12], shown in Equation 2.2, where  $\theta_0$  is the expected angle of deviation shown in Figure 2.5,  $p$  is the momentum of the incident particle and  $X_0$  is the radiation length, defined as the characteristic amount of matter traversed by the incident particle for a particular interaction.

$$\theta_0 = \frac{13.6 \text{ MeV}}{\beta c p} z \sqrt{\frac{x}{X_0}} \left[ 1 + 0.038 \ln \left( \frac{x z^2}{X_0 \beta^2} \right) \right] \quad (2.2)$$

This formula is expected to be valid for distances up to  $\sim 100X_0$ , giving an error smaller than 11% [9]. Corrections do exist though in order to get slightly better results, but this theory is precise enough for our needs, given the experimental conditions considered in this work.

## 2.4 Muon tomography

Given Molière's theory, it is obvious to see that, instead of *calculating* the deviation of a muon crossing a given material, we can instead try and *measure* it, by inverting the two previous relations. Since this deviation depends on several parameters related to the absorber itself, such as its width and radiation length  $X_0$ , we can in this case actually infer such parameters experimentally and determine the properties of the medium crossed by cosmic muons. This is the so-called muon tomography, or **muography**, taking advantage of the interaction of muons with matter to use it in practice.

Muography is therefore a non-destructive imaging technique producing a density map of the inside

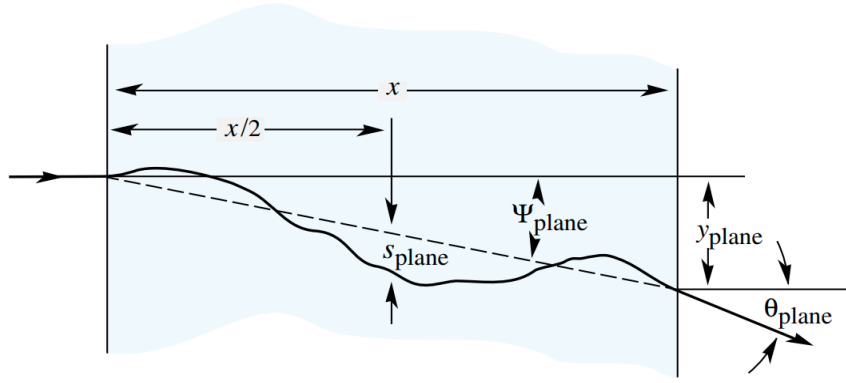


Figure 2.5: Schematic representation of the deviation induced by an absorber to an incident muon because of the multiple scattering effect [9].

of an object by measuring a flux of muons. Such method presents many different advantages over other imaging techniques such as X-rays, since it actually uses natural cosmic rays to make the measurements, being therefore completely safe. Additionally, muons interact lightly with matter, meaning that they typically have high penetrating capabilities and can therefore probe even large and/or dense objects. Muography can in this sense be applied to many different fields: it has for example even been used in 1970 in order to try and find hidden cavities of pyramids in Egypt [13] and can also be used in vulcanology, to determine whether a pocket inside of a volcano is empty or full of lava, among many other practical applications.

Such imaging techniques can be divided into two categories:

- **Absorption muography.** In this case, the observed muon flux in a given direction is compared to what is expected from cosmic rays, trying to determine the inner structure of the absorber as discrepancies between these two values. Only one detector is needed in this case, making this technique useful mostly to study large objects, even though the time to make a single measurement can last up to a few months.
- **Scattering muography.** On the other hand, the multiple scattering of muons can be used, by placing one detector on each side of the object being studied to determine the deviation of the flux of incoming muons. The denser the material put in between, the larger the observed deviation will be, as shown in Equation 2.2. This technique is mostly used to study smaller objects and is able to make quick measurements.

In this particular work, scattering muography is being applied to industry in order to try and determine the degradation of the interior of industrial equipment, as we will now see.

## 2.5 Experimental setup

### 2.5.1 Muon detectors

If we want to work with cosmic muons, we need to be able to build some detectors able to spot them and give us back useful data, such as their energy and/or incident direction. Many different

technologies exist nowadays in order to detect muons, even though in this particular case, the typical **multi-wire proportional chambers** have been used.

Invented at CERN in 1968, these detectors use an array of high-voltage wires (playing the role of the anode), running through a chamber filled with gas and whose walls are typically grounded (the cathode), as shown in Figure 2.6. Such an experimental setup therefore creates an electric field inside of the chamber, that needs to be made as large and uniform as possible.

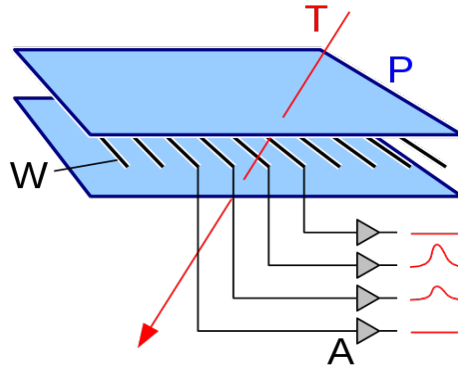


Figure 2.6: Schematic representation of a wire chamber muon detector.

When a charged particle such as a muon crosses this chamber, it will ionize the gas of the chamber and leave small electric charges along its path which will start to drift until reaching one of the wire. This drift induces an electric signal proportional to the ionisation effect in the different wires surrounding the particle path, and the combination of all the signals collected is able to give information regarding the actual path followed by the incoming particle.

Several properties are extremely important when designing a detector:

- First of all, the **spatial resolution**, the precision by which we can tell the position of the muon, should be ideally as small as possible, depending on the actual problem faced.
- The **efficiency** is also an important parameter, since we want to be able to detect as many muons as possible, to make the measurement faster and more precise.
- Finally, we typically want the detector to be large enough to avoid any **acceptance** issues.

### 2.5.2 Working setup

For this particular work, four muon chambers of 1 meter by 1 meter have been build using these principles, as shown in Figure 2.7. As we can see, more than 200 wires connected to the high voltage and made out of gold and tungsten have been placed every 4mm in two different planes rotated by 90 degrees, to measure the  $x$  and  $y$  postion of cosmic muons.

These chambers, filled while a mixture of Argon and  $\text{CO}_2$ , are then setup in pairs above and below the object being studied, in order to determine the position and the direction of the incoming and outgoing cosmic muon. This setup allows to measure with a good spatial resolution the deviation of muons, to determine as precisely as possible the properties of the object put in between both detectors.

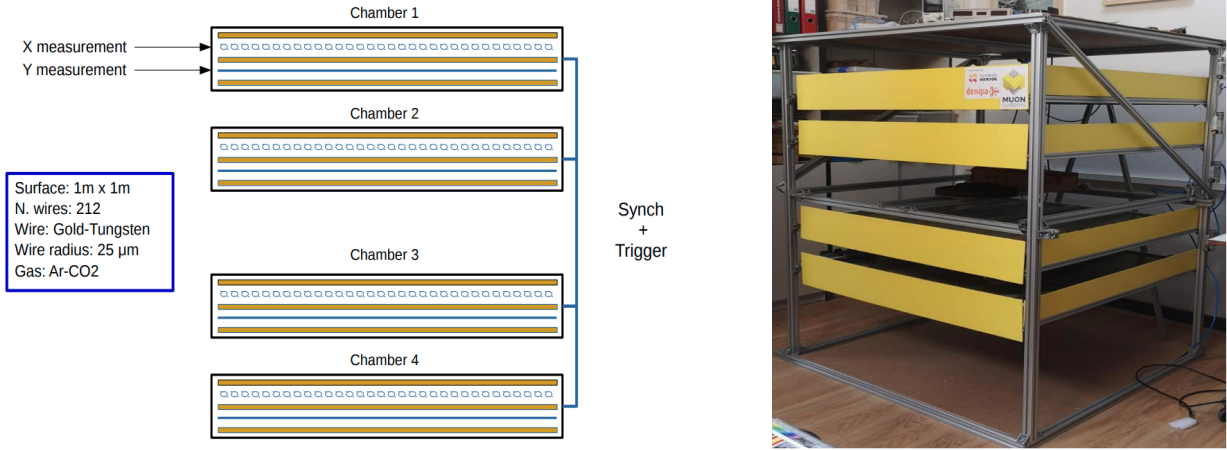


Figure 2.7: Actual working setup used for this work, with two multi-wires muon chambers of  $1\text{m}^2$  placed on each side of the object being studied.

### 2.5.3 Data flow

Once collected, the data needs to be pass through several differents in order to convert electric signals into files that can be processed, as shown on Figure 2.8. As we can see, the data follows a different path depending on its nature, if it has been collected by the detector or if we are considering Monte-Carlo simulations that will be described in details in the next Chapter. The stream of data is collected by a simple USB, and then sent to a DAQ/DQM system and to an unpacker, which takes into account the detector gemetry and translates into a physical position the information received from the DAQ, typically only telling that a certain wire has been activated.

Both streams then join at the 1D histogram format level, gathering a collection of all the hits measured for every event. All these hits are then processed and reconstructed into two trajectories, one above and one below the object, using advanced techniques that will also be described in next Chapter. Finally, 4D segments are made available from this reconstruction process, having the position and direction information needed, from which the actual analysis can take place.

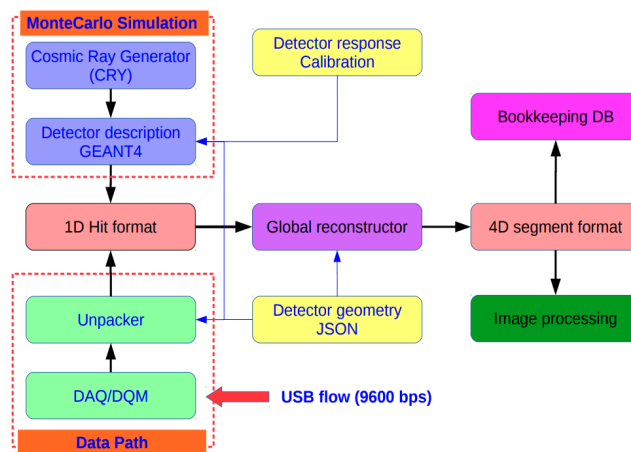


Figure 2.8: Schematic representation of the detector data flow.

## Chapter 3

# Statistical parameters

The general objective of this work is to find a way to determine the geometry of an eventual object placed between the two muon detectors by measuring the deviation affecting incident cosmic muons, a typical non-deterministic process which therefore needs to be addressed using advanced statistical parameters and methods, such as probability and kernel density functions (Sections 3.1 and 3.4), Monte-Carlo simulations (Section 3.2), p-values (Section 3.3) and likelihood function (Section 3.5).

### 3.1 Probability density functions

The probability density functions, or PDFs, are statistical expressions defining probability distributions representing the likelihood of any given outcome. They are typically represented as curves in the two dimensional space  $(x, y)$ , as shown in Figure 3.1, in which the total area below the curve in an interval is equal to the probability of a discrete random variable occurring.

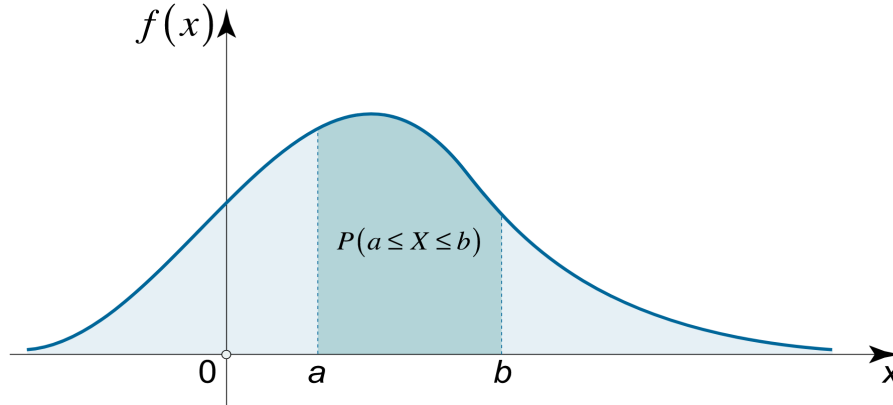


Figure 3.1: Schematic representation of a random PDF [14].

Gaussian distributions are in this sense probability density functions having a particular expression defined in Equation 3.1, where  $\sigma$  is the so-called **standard deviation** of the distribution, representing its width: a low  $\sigma$  indicates that the values of the distribution tend to be close to its mean value  $\mu$ , while a high  $\sigma$  indicates that the values are spread out over a wider range.

$$f(x) = \frac{1}{\sqrt{2\pi}\sigma} e^{-\frac{1}{2}\left(\frac{x-\mu}{\sigma}\right)^2} \quad (3.1)$$

These definitions are important because we already know that the multiple scattering process which affects the incident cosmic muons is a stochastic process: this means that two muons having similar kinematics can leave the detector with very different output directions and positions, whose actual distribution can be approximated by a gaussian function for small deviation angles (at larger angles, the distribution is behaving like Rutherford scattering, with slightly larger tails).

The standard deviation of the distribution of such output parameters is actually highly dependant on the geometrical parameters  $\theta_i$  of the eventual object put between the two detectors, as we saw in Equation 2.2. Indeed, the mean deviation observed is expected to depend on the number of radiation lengths crossed, typically much higher when a muon crosses denser objects, such as a steel pipe, as shown in Figure 3.2. This means that in first approximation, the thicker the object investigated is, the higher the expected deviation will be and this simple observation will actually be used as the driving process of the statistical study performed in this work.

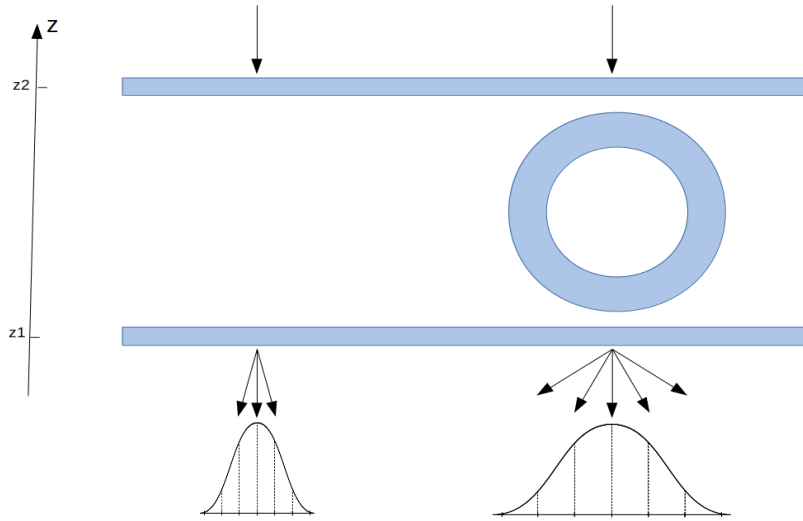


Figure 3.2: Schematic representation of the deviation expected for an incoming muon and PDFs associated without (on the left) and with (on the right) an object placed between the detectors.

### 3.2 Monte-Carlo simulations

Monte-Carlo simulations are algorithms developed in order to compute approximate numerical values to stochastic problems using random processes and advanced probabilistic techniques. In this particular case, such computation methods apply extremely well since, by being a stochastic process, the actual probability density function associated to a given geometry cannot be computed analytically (and if we were able to do it, a simple gradient descent method would be enough to solve this problem and find the optimal parameters  $\theta_i$  of the object).

The only way available to actually estimate these parameters is by using thousands of Monte-Carlo simulated toys, simulating thousands of different incident muons for different geometries. The principle is simple: for each incident muon measured, a large number  $N$  of Monte-Carlo simulations will be performed, propagating the muon using an algorithm described in Chapter 4 until reaching the lower detector. Repeating this experiment over and over again always gives different results, giving us the opportunity to build the expected PDF for a given object geometry and for the incident muon that has been measured and that will be used later on.

At the end of the day, the main objective of this process is to be able to estimate the probability



of observing a certain deviation given an object geometry and the input trajectory of the muon,  $P(\text{deviation}|\text{input})$  in simulation. Then, this process will be reversed using actual data in order to try and obtain the geometry of the object from the actual measurement of the input and output muon trajectories and positions, using a likelihood minimization method.

### 3.3 p-values

So far, we developed a technique allowing us to reconstruct the expected probability density function for a given incident muon using Monte-Carlo techniques. However, we still need a way to estimate the goodness of the actual measurement with this simulated PDF: this is where the so-called **p-values** enter, a key ingredient for this work since they allow to relate the two main parts of the problem: the simulations performed and the actual data collected.

Probability values, or p-values, are based on the concept of **null hypothesis**  $H_0$ , a general statement or default position telling that there is no actual relationship between two measured phenomena and assumed to be true until an evidence indicates otherwise. This hypothesis is then typically opposed to the **alternative hypothesis**  $H$ , usually more interesting by being the interesting hypothesis of the experiment being performed. The main goal is to compare the data with both hypotheses: if the data is consistent with  $H_0$ , then the null hypothesis can simply not be rejected. However, we can reject  $H_0$  (and therefore accept  $H$ , its exact opposite) if the data collected is significantly unlikely to have occurred if the null hypothesis were to be true, according to a **confidence level** previously defined.

From these definitions, a p-value is then defined as the probability for a variable to be observed equal or more extreme (lower or higher, depending on the cases) than the actual value  $x$  observed under the null hypothesis. This definition, can be represented in a schematic way in Figure 3.3.

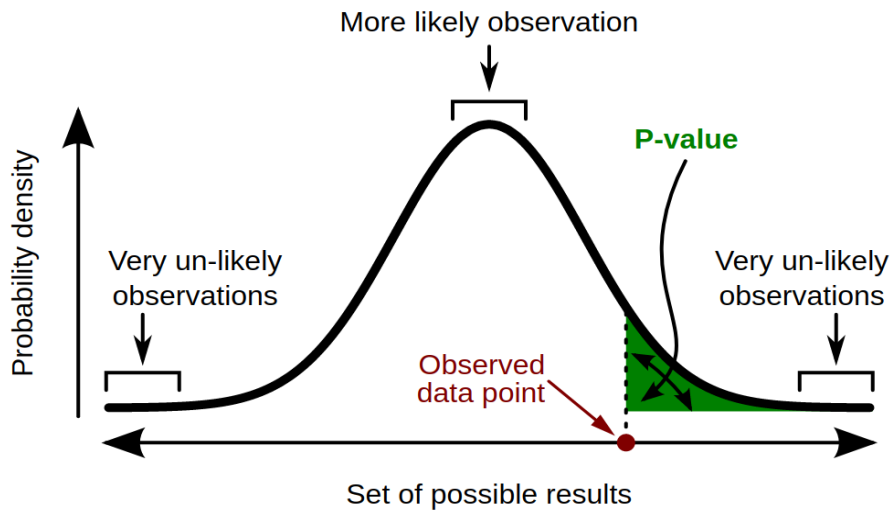


Figure 3.3: Schematic representation of the statistical concept of p-value.

In general, a smaller the p-value tells us that the hypothesis under consideration may not adequately explain the observation, leading to a higher statistical significance. At the end of the day, the null hypothesis is rejected if the p-value obtained is smaller than a previously defined, arbitrary and fixed threshold  $\alpha$ , the **level of significance** of the experiment. This parameter can take a large range of values, typically ranging from 0.05 to 0.001.

In this work, the p-value will then allow us to compare a single data measurement with the probability

density functions obtained from Monte-Carlo and quantify the probability that such muon was actually observed for a given object geometry. This value will then be feeded to a likelihood function, as explained in Section 3.5.

### 3.4 Kernel density functions

In statistics, the kernel density estimation is a method allowing to estimate the density probability of a random variable. The objective of this method is to make inference regarding the data in order to try and smooth the probability density function obtained experimentally, as shown in Figure 3.4.

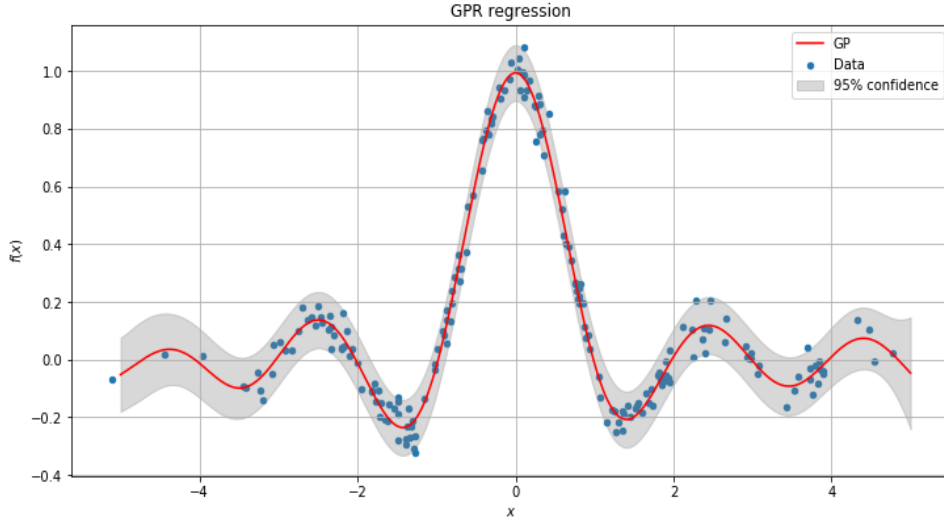


Figure 3.4: Academic example of a gaussian regression applied to a noisy and finite set of data points.

Mathematically, the objective is then to build a continuous function  $\hat{f}_h(x)$  defined in Equation 3.2 from a finite set of data samples  $\mathbf{x}$ , distributed along an unknown density function  $f$ , and where  $K$  is the so-called **kernel** and  $h$  is the **bandwidth** or **smoothing parameter** of the estimation.

$$\hat{f}_h(x) = \frac{1}{nh} \sum_{i=1}^n K\left(\frac{x - x_i}{h}\right) \quad (3.2)$$

In the previous equation the kernel  $K$  is a non-negative function defined depending on the problem, as the "building block" of the final function that needs to be estimated: indeed, as we can see, the function  $\hat{f}_h(x)$  will be given as a simple sum of kernels. In most of the problems, including this work, a standard gaussian kernel is used, taking the shape shown in Equation 3.3.

$$K(x) = \frac{1}{\sqrt{2\pi}} e^{-\frac{1}{2}x^2} \quad (3.3)$$

In this particular case, this method allows us to construct continuous functions and recover for the eventual low statistics and possible holes in the bi-dimensional histograms for the position and direction of the muons obtained in  $x$  and  $y$ , following the method that will be described in Chapter 4.

### 3.5 Likelihood

So far, we have been able to generate thousands of different Monte-Carlo experiments for a given muon and for a given object geometry, but we still need to find a way to reverse this process to estimate the geometry of the object from a given measurement.

This is where the likelihood becomes useful, defined as a function that measures the goodness of a fit with respect to a sample of data, for several unknown parameters of a mathematical model. In general, this function can be defined in Equation 3.4, where  $\theta = \{\theta_1, \dots, \theta_i\}$  are the parameters of the model and  $x$  is the actual measurement of a random variable  $X$  following a density function  $f$ .

$$\mathcal{L}(\theta|x) = f_{\theta}(x) = P(X = x|\theta) \quad (3.4)$$

It is important to note at this point that the likelihood is a function of the parameters  $\theta$  but not a probability density function itself. Additionally, it should in general not be confused with the probability  $p(\theta|x)$  since it is equal to the probability of observing a given outcome  $x$  when the true values of the parameters are  $\theta$ : this means that the likelihood is equal to a probability density over the outcome  $x$ , not over the parameters  $\theta$ .

To understand this concept better, the simple example of an unfair coin toss can be considered, where the fairness of the coin is the parameter of the model, represented by the probability of obtaining head  $p_H$  and taking values between 0 and 1 (for a regular coin,  $p_H = 0.5$ ). The likelihood is then defined from a given observation, such as obtaining two heads in a row  $HHT$  in the following way:

$$\mathcal{L}(p_H|HHT) = P(HHT|p_H) = P(H|p_H) \cdot P(H|p_H) \cdot P(T|(1 - p_H)) \quad (3.5)$$

For each value of  $p_H$ , the value of the likelihood can then be computed and ultimately plotted, as shown in Figure 3.5, showing a minimum value for a particular value of  $p_H$ .

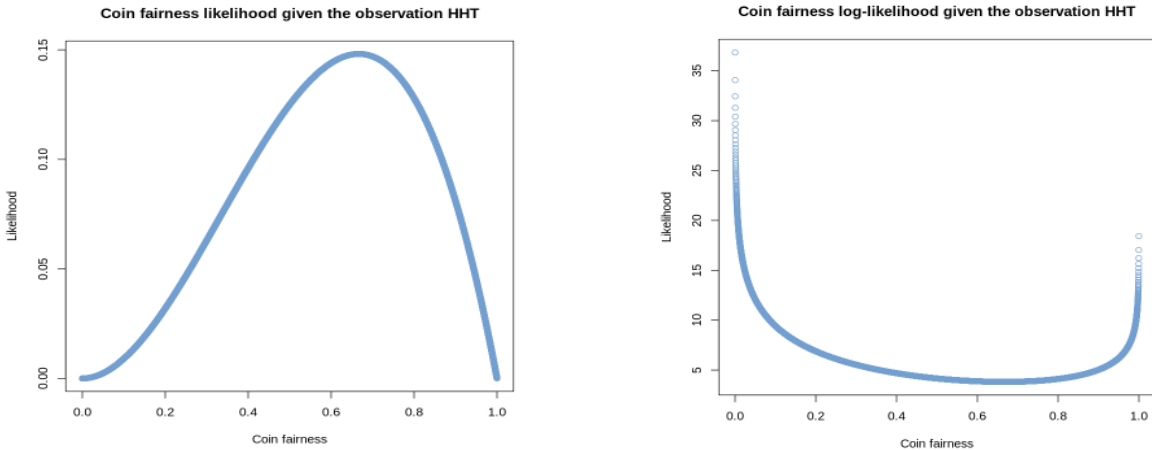


Figure 3.5: Likelihood (on the left) and log-likelihood (on the right) obtained for our particular example of the fairness of a coin given the observation  $HHT$ .

The likelihood is in this sense interesting because it can be described as a hypersurface whose peak gives the optimal set of parameters maximizing the probability of drawing the actual sample measured. Often, the **log-likelihood**  $l(\theta|x) = -2\log(\mathcal{L}(\theta|x))$  is actually used instead, being a bit more convenient to deal with the special part concavity plays in the maximization process. Given the properties of the logarithm, maximizing the likelihood is equivalent to minimizing this particular log-likelihood.

Another important property of the likelihood that will be used extensively and already assumed in the previous example is the fact that, for two independant measurements, the likelihood of both measurements is equal to the product of both likelihoods independantly computed for each event.

In this work, a likelihood function will be obtained as the product of the thousands of muons collected, but for a fixed geometry of the object. A minimization method will then be performed in order to try and minimize the value obtained by trying out different parameters  $\theta$ , until finding the optimal set of parameters able to describe best the object put in between the detectors.

## Chapter 4

# The algorithm

The algorithm we developed to solve this particular problem has been written in C++ and can be divided in two categories. First of all, we have the so-called **PipeReconstructor** (Section 4.1), a set of classes able to perform several tasks:

- Define the geometry of the problem (the detector, and the volume under investigation are therefore defined using the smallest possible set of parameters at this stage);
- Compute the different intersection points of an incident cosmic muons with our geometry;
- Propagate the muons through the different medium of interaction encountered along their path;
- Calculate the likelihood of a given measurement for a given volume;
- Finally, probe different volumes and perform a descent method in order to minimize the likelihood encountered and find the optimal volume to solve our problem given the measurements collected.

This class is complemented with the **Generator** (Section 4.2), a small class allowing us to generate muons simulating our actual experiment, using some of the functions of the PipeReconstructor.

### 4.1 PipeReconstructor

The PipeReconstructor is a set of classes, allowing us to define different kind of objects:

- First of all, a muon is defined from the **MuonState** class and from 7 parameters: 3 parameters  $(x, y, z)$  representing the position of the muon, 3 additional parameters  $(v_x, v_y, v_z)$  representing its three-dimensional direction and the actual value of its momentum. The measurement performed at the top detector will be labelled with the index 1  $(x_1, y_1, z_1, v_{x1}, v_{y1}, v_{z1}, p_1)$  while the measurement performed by the lower detector will take the index 2.
- **Surfaces** and **Volumes**, virtual classes allowing us to define general surfaces (their spatial position and their geometrical center) and volumes, defined as vectors of several surfaces: a general volume is in this sense defined by its surfaces, central position and global density. Defining such virtual classes is extremely important since it would allow us to consider different geometries, not just a pipe put in between the two detectors. The constant  $z$  position of both the lower and upper detectors (-37 and 37 cm, respectively) is also defined at this stage.

- **Cylinders** and **Pipes**, particular subclasses of the virtual classes previously defined in order to define the geometry of this particular problem. A pipe is in this sense defined as a set of two cylinders and from 7 different parameters: its central position  $(x, y, z)$ , along with its density, the radii  $r$  and  $R$  of both cylinders and its length along the y-axis  $L$ , as shown in Figure 4.1. For simplicity, the pipes considered in this work are therefore considered to have a constant density.

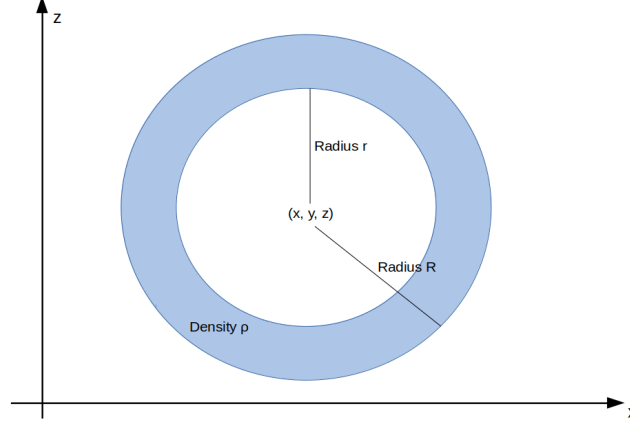


Figure 4.1: Schematic representation of the pipe considered in this work and its respective parameters.

One method computing the exact cut point between a `muonState` and a `Cylinder` has been defined at this point as well, from the general equation of a cylinder of radius  $R$  centered in  $(X, Y = 0, Z)$  (given that the y-axis is defined along the cylinder, we can assume that the cylinder is actually centered in  $Y = 0$  without any loss of generalization).

$$(x - X)^2 + (z - Z)^2 = R^2 \quad (4.1)$$

From this, computing the interesection between such cylinder and a general `MuonState`  $(x_0, y_0, z_0, v_x, v_y, v_z, p)$  is trivial, since a `MuonState` can be represented by a straight line:

$$\begin{cases} x = x_0 + \lambda v_x \\ y = y_0 + \lambda v_y \\ z = z_0 + \lambda v_z \end{cases} \quad (4.2)$$

By putting this into the equation of the cylinder, we get the following expression:

$$(x - X)^2 + \left( z_0 + \frac{v_z}{v_x}(x - x_0) - Z \right)^2 = R^2 \quad (4.3)$$

Once both the squared values applied, we get an equation having the shape  $Ax^2 + Bx + C = 0$ , where both cutting points will be given therefore by  $x^{+/-} = \frac{-B \pm \sqrt{B^2 - 4AC}}{2A}$ :

$$\begin{cases} A = \left( 1 + \left( \frac{v_z}{v_x} \right)^2 \right) \\ B = -2X - 2x_0 \left( \frac{v_z}{v_x} \right)^2 + (2z_0 - 2) \left( \frac{v_z}{v_x} \right) \\ C = X^2 + Z^2 + z_0^2 - 2z_0Z + x_0^2 \left( \frac{v_z}{v_x} \right)^2 + x_0(2Z - 2z_0) \left( \frac{v_z}{v_x} \right) - R^2 \end{cases} \quad (4.4)$$

Finally, methods allowing us to determine whether a muon state is inside or outside of the pipe have also been written, taking into account the particular geometry of such objects in order to determine the medium in which the muon is propagating and therefore use the correct density in all the calculations performed.

- A **Propagator** is then defined as the object allowing us to propagate a MuonState through a general Volume. The way it works is quite simple:
  - First of all, the distance between the initial MuonState and all the surfaces of the object considered is computed and the first intersection point is kept.
  - The muon is then propagated in a distance corresponding to 90% of the total distance to this first cut point by taking into account the multiple scattering effect, and a new MuonState is therefore obtained. We cannot propagate the muon 100% of the distance immediately to avoid effects related to the actual shape of the Cylinder and due to the fact that the Equation 2.2 is only valid in case of normal incidence between the MuonState and the Cylinder, which is typically not verified, as shown in Figure 4.2.

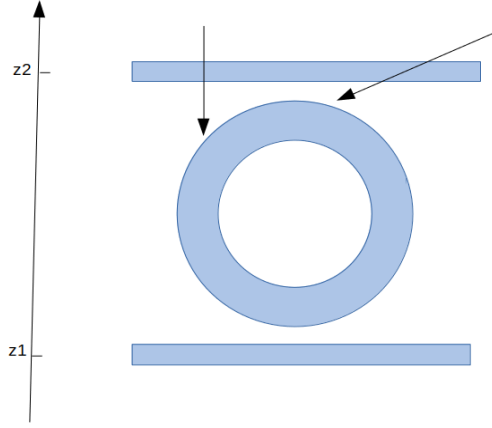


Figure 4.2: Typical example of two incident muons not respecting the normal incidence required by the Moliere equation for the multiple scattering of muons.

- This propagation process is then repeated 6 times in total, until reaching a position extremely close to the volume that is being studied. The muon is then manually propagated along a distance of around 0.1 mm in order to force its crossing with respect to the first surface, so that the propagation can be repeated for the next surface.
- This whole process is repeated for all the different surfaces encountered by the muon (up to four times in total for our particular geometry), and one final time until reaching the bottom detector. At this point, the MuonState is kept as the second measurement ( $x_2, y_2, z_2, v_{x2}, v_{y2}, v_{z2}, p_2$ ) that will be compared with the simulation using statistical methods.
- Eventual muons which not cross both the upper and lower detector are of course rejected and simply not considered when computing the value of the likelihood.
- The **Likelihood** object itself can then be calculated. This class reads a Monte-Carlo generated file (actual explanations regarding this generation process will be given in Section 4.2) and is able to perform several calculations for each generated event ( $n_{\text{gen}} \sim 100000$ ):
  - First of all, three MuonStates are computed: the so-called *IncomingMuon* and *OutgoingMuon*, directly read as the measurement at the upper and lower detectors, from the Monte-Carlo simulation, and the *IncomingMuonLinearDown*, defined as the prolongation along a straight line of the *IncomingMuon*.

- Each *IncomingMuon* is then propagated along the geometry using the method previously defined, defining a fourth so-called *IncomingMuonDown* MuonState. The  $\Delta x$ ,  $\Delta\theta_x$ ,  $\Delta y$  and  $\Delta\theta_y$  parameters shown in Figure 4.3 and representing the actual deviation suffered by the muon can then be obtained.

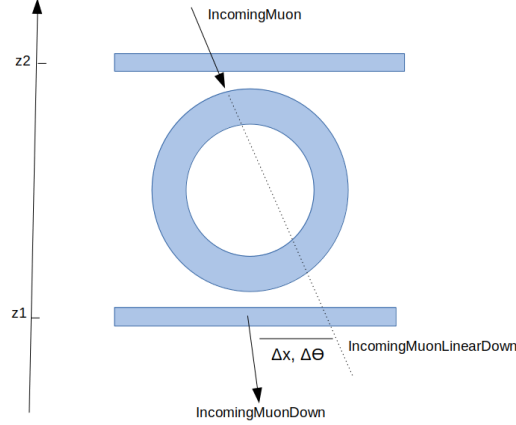


Figure 4.3: Typical example of two incident muons not respecting the normal incidence required by the Moliere equation for the multiple scattering of muons.

At this stage, all the ingredients needed to compute a p-value are available.

- This propagation process is repeated over and over again ( $n_{\text{prop}} \sim 100$  times in total) since we know that the multiple scattering is a stochastic process. Each time the propagation method is called for the same incident muon, different deviation values are therefore obtained, each kept in a vector, which can be represented in an histogram and smoothen using kernel density functions, as explained in Section 3.4.
- At this stage, all the ingredients needed to compute a p-value are available. A p-value for each *IncomingMuon* is then computed by using the vectors previously defined and representing an histogram of all the possible deviation values, and by comparing the actual measurement *OutgoingMuon* for such distributions. The p-values computed along the x and y-axis then enter the global likelihood value that is used for the minimization process.
- So far, we described a method allowing to compute the likelihood of observing a given measurement by using Monte-Carlo simulations for a given volume, but one last important step is still missing. **Talk about the minimization process once implemented.**

## 4.2 Generator

As we have just seen, the computation of the p-values required for the likelihood takes as input a file of Monte-Carlo generated events, which has been previously produced thanks to this **Generator**. This class takes as input one rootfile containing the Monte-Carlo generation of incoming cosmic muons previously produced by the company using CRY [15], a dedicated generator for cosmic muons. To simplify the current problem, the goal is then to simply read the input parameters of such muons, such as the position of their impact with the upper detector, and to propagate them using the previously defined functions, to generate a distribution of the expected output.

The way this algorithm works is quite simple: a general volume matching the one that needs to be studied (in this case, a general Pipe located in the center (0, 0, 0) of the detector, having an inner



radius of 18.2 cm, an outter radius of 20 cm is then defined and a lenght of 50 cm) is first of all defined. Then, for each event we want to generate ( $n_{\text{gen}} \sim 100000$ ), a loop is performed: in each iteration of this loop, the incident muon is propagated throughout the geometry and histograms for the same parameters  $\Delta x$ ,  $\Delta\theta_x$ ,  $\Delta y$  and  $\Delta\theta_y$  are obtained.

Talk about MC truth and measurement.

## Chapter 5

# Results obtained

### 5.1 Incident muons distribution

Prove the cos-square law.

### 5.2 Tests performed

## Chapter 6

# Conclusions

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