Development of a statistical analysis framework in the context of Muography applied to the industry

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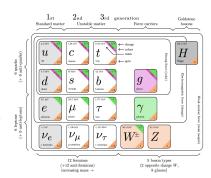
Section I General introduction

Develop a new framework allowing us to study the results from a muography experiment to characterize the inner properties of physical objects using data science and advanced statistical models.

Particle physics and muons

The Standard Model describes the fundamental particles existing and their interactions:

- Introduced in the 1970s and still considered to be valid, but probably incomplete
- Simple in concept but extremely precise
- Lots of successful predictions made over the years, such as the existence of the top quark and the Higgs boson



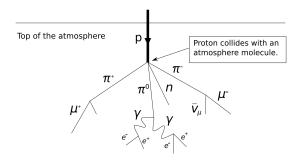
Muons

- ullet Muons μ^- are one of the 12 fundamental particles existing
- They have a relatively small interaction cross-section with ordinary matter, allowing them to cross material without being stopped, making them interesting.

Cosmic rays

Comic rays are a **constant flux of high energy particles** reaching the Earth:

- Mostly made out of protons and atomic nuclei
- Trigger a decay chain by interacting with the atmosphere, producing muons
- Muons are not stable ($\tau \simeq 2.2 \mu s$) but relativity makes them live long enough to reach the ground $\to 10.000$ cosmic muons are observed per m^2 and per minute at sea level.



Muons interactions with matter

Ionization

When the incident muon collides and gives some of its energy to the electrons of the absorber, but quite small for MIPs such as cosmic muons.

→ Basis for the absorption muography, not covered here.

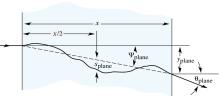
Multiple Coulomb scattering

Inducing a stochastic deviation whose central angular deviation can be described by a Gaussian of width θ_0 under our experimental conditions.

$$\theta_0 = \frac{13.6 \text{ MeV}}{\beta \textit{cp}} \sqrt{\frac{\textit{x}}{\textit{X}_0}} \left[1 + 0.038 \ln \left(\frac{\textit{x}}{\textit{X}_0 \beta^2} \right) \right]$$

Deviation depends on the number of radiation lengths X_0 and on the medium crossed.

ightarrow Basis for the scattering muography.



Muon tomography

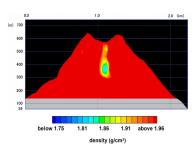
Instead of *calculating* the deviation expected for a cosmic muon, we can *measure* the positional and angular deviation suffered to estimate the properties of the medium crossed.

Several advantages over other imaging techniques:

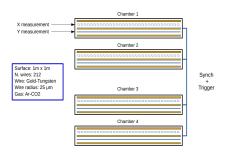
- Non-destructive technique
- High penetrating capabilities allowing to probe large and dense objects
- Completely safe, by using natural cosmic rays for the measurement.

Muography can be used in many different fields:

- Find hidden rooms in Pyramids
- In volcanology, to know whether a pocket is empty or full of lava [1]
- Nuclear waste/facilities inspection



Experimental setup

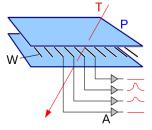


Quite simple experimental setup:

- Two $1m^2$ detectors placed below and above the object under investigation
- Two chambers in each detector, to measure the position and direction of muons along the x and y axes
- Each chamber is filled with a mixture of Argon and CO2
- More than 200 wires separated by 4mm make up each chamber

A muon crosses the detector leaves small electric charges behind, collected by the wires while leaving a signal. The combination of the signals on the different wires give us information regarding the muon.

The data is collected from a USB stick and goes through a complete reconstruction process before being available in a rootfile, as detailed in the backup.



Section II Statistical basis

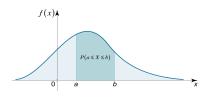
The algorithm developed heavily relies on several important statistical concepts that we can now define.

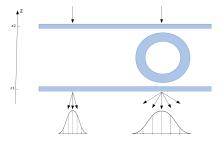
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Probability density functions

PDFs are mathematical expressions defining probability distributions which represent the likelihood of any given outcome.

The area below the PDF in an interval can be interpreted as the value of the probability of a random variable occurring.



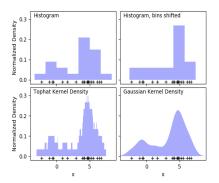


The multiple scattering is a stochastic process, making PDFs extremely important.

ightarrow A thicker and denser object results in a higher expected deviation and a larger standard deviation σ of the Gaussian PDF.

Kernel density estimation

This method allows us to estimate the shape of an unknown PDF f of a random variable X from a set of N observations.



The usual way to proceed is to simply put the observations in an histogram, but this results in a non continuous function with possible gaps.

We then define $\hat{f}_h(x)$, an estimator of f defined as the sum of continuous functions instead.

$$\hat{f}_h(x) = \frac{1}{Nh} \sum_{i=1}^{N} K\left(\frac{x - x_i}{h}\right)$$

Two important parameters

- The kernel K, chosen as Gaussian functions in this case
- And the bandwidth h, a smoothing parameter.

Monte-Carlo simulations

Monte-Carlo simulations are obtained from algorithms developed to **compute approximate numerical values to stochastic problems** using random processes and probabilistic techniques.

Instead of relying on actual data collected from the experiment, we can:

- Use CRY, a cosmic ray generator to simulate thousands of incident muons
- Make these muons go through a complete simulation of the detector done with Geant4
- Go through the same post-processing process as actual data, simulating thousands of experiments without the need to perform them
- Build the PDF for a given experiment from these simulations

Simulating an experiment is cheaper and faster than running it and allows to compare results obtained from both channels. In this work, the dependence on Geant4 will be removed as well, to make this process even faster.

Maximum likelihood estimation

We still need a method allowing us to estimate the geometrical parameters of an object from a given measurement, such as the thickness of a pipe.

The likelihood measures the goodness of a fit with respect to a sample of data for one or several unknown parameters. If θ are the parameters of the model and x is the measurement of a random variable X defined from a PDF f, then:

$$\mathcal{L}(\boldsymbol{\theta}|\mathbf{x}) = f_{\boldsymbol{\theta}}(\mathbf{x}) = P(\mathbf{X} = \mathbf{x}|\boldsymbol{\theta})$$

The likelihood can be described as an hypersurface whose peak gives the optimal set of parameters maximizing the probability of drawing the actual sample measured.

ightarrow The objective is then to find the set of parameters minimizing the log-likelihood $I(\theta|x) = -2\log(\mathcal{L}(\theta|x))$. In this work, the parameter to be optimized will be the thickness of a steel pipe placed between the detectors.

Section III Algorithm implementation

A C++ framework has been developed in order to solve this problem, relying on three main parts: a PipeReconstructor, a Generator and a Plotter, now described.

General idea

PipeReconstructor

General set of classes allowing us to:

- Define the geometry of the problem (mainly, the detector and object position and size)
- Compute the intersection points between cosmic muons and this geometry
- Propagate these muons through the geometry using the multiple scattering
- Calculate the likelihood of a given measurement for a given geometry
- Finally, find the optimal object parameters from this likelihood.

To reach all these goals, several classes have been defined as described in the next slides.

Generator

Another class allowing us to generate muons by performing Monte-Carlo simulations using some of the functions of the PipeReconstructor and without relying on Geant4.

Plotter

Simple set of scripts used in order to plot all the results shown in the next slides.

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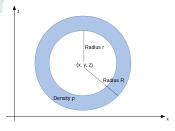
MuonStates, surfaces and volumes

Surfaces and Volumes

General virtual classes used to define the geometry of the problem; in particular, the subclasses Cylinder and Pipe are mostly used in this work.

A Pipe is therefore defined from 7 parameters:

- Its central position (x, y, z)
- Its inner r and outer radii R
- ullet Its constant density ho
- And its length L along the axis of the cylinder.



MuonStates

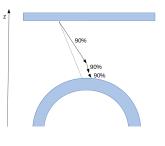
A MuonState is a class defining a muon from two vectors representing its position (x, y, z) and direction (v_x, v_y, v_z) , and its momentum value p.

Two MuonStates are defined for each experiment: one measured at the top detector $(x_1, y_1, z_1, v_{x1}, v_{y1}, v_{z1}, p_1)$, and one at the bottom, with the index 2.

Propagator

The propagator is the object allowing to propagate a MuonState through a Volume, defined as a vector of a fixed number of Surfaces:

- First, the distances between a MuonState and all the Surfaces of the Volume are computed and the first intersection point is kept
- The muon is propagated 90% of the distance to this cut point using the multiple scattering, and this process is repeated several times until being closer to 0.1mm to the first surface
- The Surface is manually crossed by slightly moving the MuonState along its direction
- This process is repeated for all the Surfaces of the Volume and then one last time, until reaching the bottom detector, where the MuonState is returned.

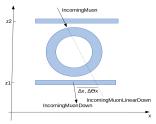


Eventual muons which do not actually cross the pipe or out of the acceptance of any of the detectors are rejected at this stage.

Likelihood

The Likelihood class takes as input a Volume and a rootfile containing either data or Monte-Carlo simulations, to tell us **how likely it is that we obtained exactly these measurements** for the geometry given:

- First, four MuonStates are computed:
 - ▶ The IncomingMuon and OutgoingMuon, defined as the actual measurements made by the top/bottom detectors
 - The IncomingMuonLinearDown, defined as the linear propagation of the IncomingMuon, used as reference
 - And the IncomingMuonDown, by propagating the initial state with our Propagator through the Volume



- This process is repeated n_{iter} times for each input event, obtaining each time values for the Δx , $\Delta \theta_x$, Δy and $\Delta \theta_y$ parameters, filling 2 bi-dimensional PDFs histograms
- The histograms obtained are smoothen using the kernel density estimation method
- The probability to observe the actual measurement (OutgoingMuon) is computed, and summed over all the events in the file, returning the value:

$$\mathcal{L} = \left(\frac{1}{N}\right) \sum_{i=1}^{N} -2.0 \left(\log(\mathsf{value}_{x,i}) + \log(\mathsf{value}_{y,i}) \right)$$



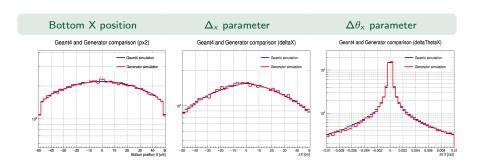
Section IV Results obtained

First of all, the results obtained by our Generator whill be shown and compared to Geant4, before moving on to the actual results obtained with our algorithm.

Generator validation

The Generator allows us to generate Monte-Carlo experiments in a faster way than using Geant4, which is more complex, relying on a complete description of the detector.

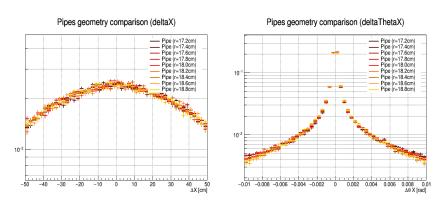
The first step was then to validate our Monte-Carlo simulations by comparing measurements simulated by both Geant4 and our own Generator for a given geometry.



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Pipes geometries

Once validated, 12 different Monte-Carlo files have been generated with our Generator with 10 to 50.000 events each, for different pipe geometries, having an outer radius of 20cm and inner radii ranging from 16.8 to 19.0cm.

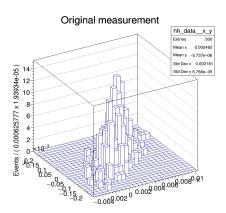


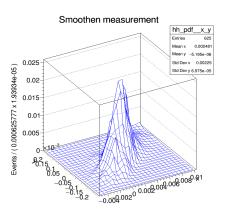
It took 8 seconds to produce 50.000 events with our Generator, more than two orders of magnitude faster than a complete Geant4 simulation.

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Kernel density functions

Since the likelihood needs to be computed for every single event of the input simulation file by performing another computing extensive loop, the $n_{\rm iter}$ parameter is kept as small as possible using the kernel density estimation method.

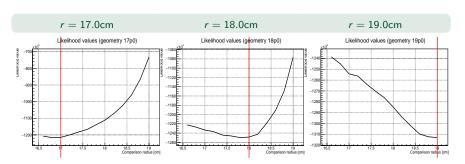




Likelihood curves

Finally, we estimated the value of the total likelihood obtained for different pipe geometries, characterized by different inner radii, ranging from 16.6 to 19.0cm, by steps of 0.2cm.

The idea was to plot the likelihood obtained comparing one by one the generated file with the different geometries available and try to figure out which pipe geometry, in the x-axis, is more likely to give rise to the file considered.



We do observe the minimum exactly where it is supposed to most of the geometries, as seen in the backup.

Conclusions

In conclusion, we developed a new framework allowing us to:

- Quickly generate thousands of Monte-Carlo experiments for different pipe geometries without relying on Geant4
- Compare the output measurement distributions expected at the bottom detector after propagating a muon through our Volume, using our own Propagator
- Compute the likelihood of a given measurement with respect to different pipe geometries, in order to find a way to estimate the thickness of a steel pipe by studying the deviation of incident cosmic muons
- And plot all the results obtained.

With this framework, we were able to determine the thickness of such pipes with a precision of the order of the millimiter by considering 10.000 events, which is equivalent to 1 minute of data taking only, therefore solving the initial problem solved.

Future improvements

This exercise is only a first approach to the problem, and different improvements can be considered to improve and/or generalize the results obtained:

- Consider more general geometries that a pipe, or perform the likelihood minimization with respect to additional parameters, not only the thickness of the pipe
- The analysis can be repeated using actual data collected by the detector
- We could also consider the ionization process to improve these results
- The interaction between the detector and the cosmic muons could be considered in our Generator to make it more reliable and precise
- A further analysis estimating the Hessian of the likelihoods and the impact of systematic uncertainties would definitely improve the algorithm
- Finally, we have been limited computionally in this case, taking a few hours to produce
 a single plot with few statistics. Accessing to computers with higher capacities will be
 extremely interesting to improve the results obtained.

Thank you for your attention!

Backup slides

Ionization

lonization happens when the incident muon gives some of its energy to the electrons of the absorber, as described by the Bethe-bloch formula.

$$-\left\langle \frac{\text{d}E}{\text{d}x}\right\rangle = \text{K}z^2 \frac{Z}{A} \frac{1}{\beta^2} \left[\frac{1}{2} \ln \left(\frac{2m_e c^2 \beta^2 \gamma^2 W_{\text{max}}}{I^2} - \beta^2 - \frac{\delta(\beta\gamma)}{2} \right) \right]$$

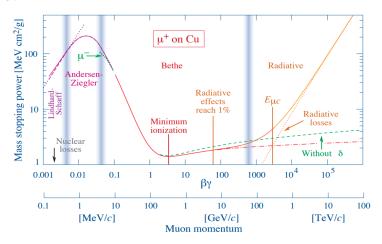
The mass stopping power of material depends on:

- The charge number of incident particle z
 - The atomic mass and charge of absorber A and Z
 - ullet The relativistic factors eta and γ
 - ullet The maximum possible energy transfer to an electron in a single collision $W_{ extsf{max}}$
 - And the mean excitation energy I.

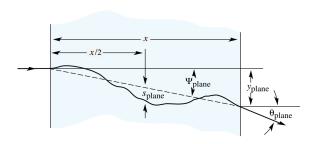


Ionization

Cosmic muons have an energy of the order of the GeV and are therefore referred to as minimum ionizing particles, so ionization is not considered in this work.



Multiple Coulomb scattering



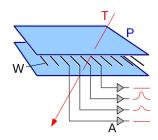
Highly correlated deviation parameters ($\rho_{\theta_{\text{plane}}, y_{\text{plane}}} = \sqrt{3}/2$). The pairs ($\theta_{\text{plane}}, y_{\text{plane}}$) are approximately distributed by a bi-dimensional Gaussian distribution with a given covariance:

$$Cov(\theta_{plane}, y_{plane}) = \begin{bmatrix} \theta_0^2 & \frac{\theta_0^2 x}{2} \\ \frac{\theta_0^2 x}{2} & \frac{\theta_0^2 x^2}{3} \end{bmatrix}$$
 (1)

Muon detectors

Multiwire proportional chambers use an array of highvoltage wires, placed within a chamber filled with a gas, in which an electric field is created.

A muon crosses the detector leaves small electric charges behind, collected by the wires while leaving a signal. The combination of the signals on the different wires give us information regarding the muon.

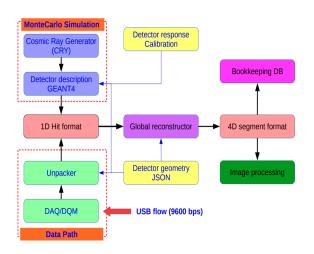


Most important parameters of a muon detector:

- The spatial resolution, ideally as small as possible
- The acceptance, related to the size of the detector
- The efficiency, which should be high to make the measurement reliable and fast.

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Data flow



Positional and angular deviation parameters

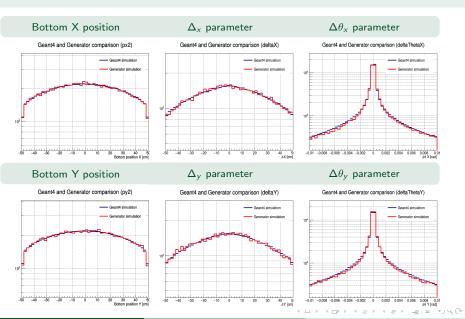
Main parameters used throughout this work:

$$\begin{cases} \Delta x = x_2 + d(v_{x2} - x_1) \\ \Delta y = y_2 + d(v_{y2} - y_1) \end{cases}$$

$$\begin{cases} \Delta \theta_x = \arctan\left(\frac{v_{x2}}{v_{z2}}\right) - \arctan\left(\frac{v_{x1}}{v_{z1}}\right) \\ \Delta \theta_y = \arctan\left(\frac{v_{y2}}{v_{z2}}\right) - \arctan\left(\frac{v_{y1}}{v_{z1}}\right) \end{cases}$$

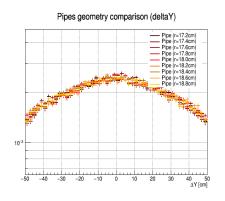
Bi-dimensional histograms (Δx vs $\Delta \theta_x$ and Δy vs $\Delta \theta_y$) are filled with these values, smoothened and used for the computation of the likelihood.

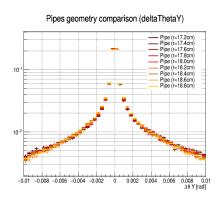
Generator validation



Pipes geometries

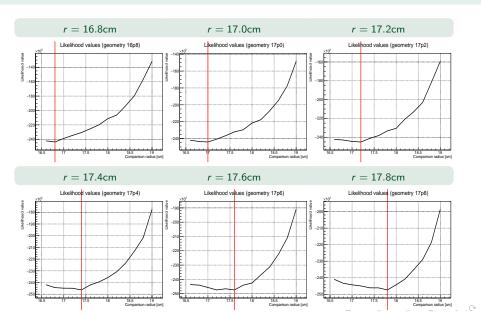
The same comparison has been performed along the Y-axis.



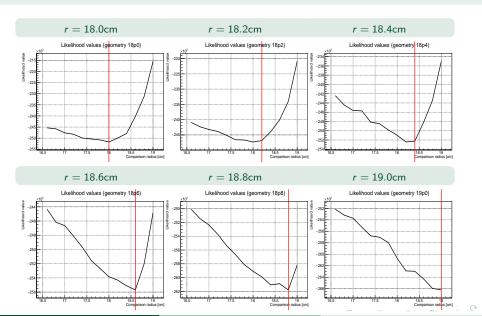


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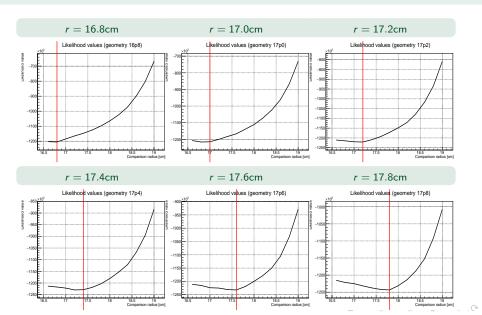
Likelihood curves (10.000 events, 100 iterations)



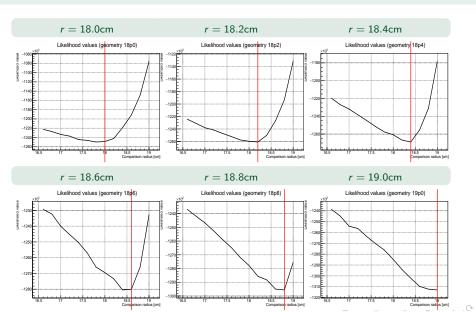
Likelihood curves (10.000 events, 100 iterations)



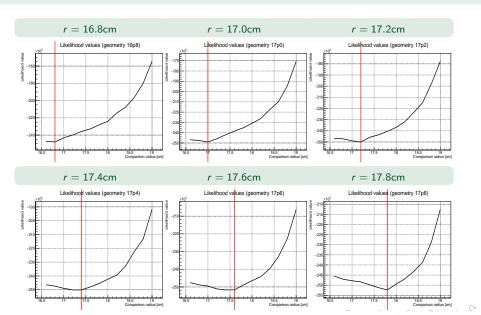
Likelihood curves (50.000 events, 100 iterations)



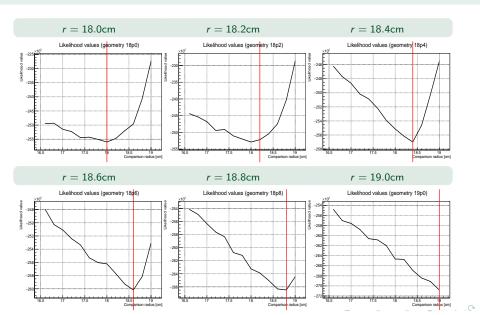
Likelihood curves (50.000 events, 100 iterations)



Likelihood curves (10.000 events, 250 iterations)

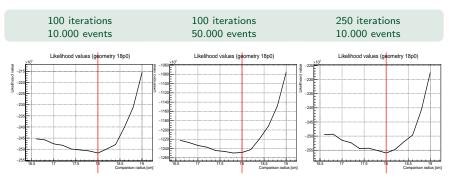


Likelihood curves (10.000 events, 250 iterations)



Likelihood curves

We also estimated the impact of the number of simulated events N_{MC} and the number of likelihood computation iterations N_{iter} on the likelihood curves, for the r = 18.0cm geometry.



In this case, it does seem that increasing the number of simulated events improve quite a lot the results (note that 10.000 events correspond to 1 minute of data taking only).

Additional references

[1] "A window into the Earth's interior", Earthquake Research Institute,2014

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