

# Behavior of weakly adsorbing impurities in flow-through ion-exchange chromatography

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## Abstract

Flow-through ion-exchange chromatography is frequently used in polishing biotherapeutics, but the factors that contribute to impurity persistence are incompletely understood. A large number of dilute impurities may be encountered that exhibit physicochemical diversity, making the flow-through separation performance highly sensitive to process conditions. The analysis presented in this work develops two novel correlations that offer transferable insights into the chromatographic behavior of weakly adsorbing impurities. The first, based on column simulations and validated experimentally, delineates the relative contributions of thermodynamic, transport, and geometric properties in dictating the initial breakthrough volume of dilute species. The Graetz number for mass transfer was found to generalize the transport contributions, enabling estimation of a threshold in the equilibrium constant below which impurity persistence is expected. Impurity adsorption equilibria are needed to use this correlation, but such data are not typically available. The

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second relationship presented in this work may be used to reduce the experimental burden of estimating adsorption equilibria as a function of ionic strength. A correlation between stoichiometric displacement model parameters was found by consolidating isocratic retention data for over 200 protein-pH-resin combinations from the extant literature. Coupled with Yamamoto’s analysis of linear gradient elution data, this correlation may be used to estimate retentivity approximately from a single experimental measurement, which could prove useful in predicting host-cell protein chromatographic behavior.

*Keywords:* Flow-through, Ion-exchange, Impurity, Clearance, Breakthrough, Stoichiometric displacement model

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## 1. Introduction

The biopharmaceutical market has grown substantially over the past few decades, with over 300 biologics having received regulatory approval and hundreds more in development pipelines [1, 2]. Monoclonal antibodies (mAbs) comprise the majority of these biologics, partly because platform purification processes have enabled their rapid development [3, 4]. One stage in a typical platform process is polishing, where trace impurities such as host-cell proteins (HCPs) are removed prior to product formulation using one or more chromatographic operations [5]. Despite their low concentrations, some of these impurities may pose a risk to therapeutic safety and stability, and polishing operations are designed to remove them completely [6, 7].

Polishing may be performed using ion-exchange (IEX), hydrophobic interaction or multimodal resins. Since the majority of secreted HCPs are more

acidic than the typical mAb [8], anion-exchange (AEX) resins are more naturally suited to flow-through operations [9], in which impurities are intended to adsorb to the resin, but they may be used in bind-and-elute mode as well [10]. The inverse is true of cation-exchange (CEX) stationary phases. The flow-through operational mode offers the advantage of relatively high throughput; as a heuristic, 1 g of mAb usually requires 1 ml of resin in flow-through IEX [11]. For mAb concentrations on the order of 20 mg/ml, this corresponds to an apparent processing capacity on the order of 50 column volumes (CV). Such large capacities can enable high productivity operations with relatively small columns. If convective media are used instead of resin particles, there is the potential to increase productivity beyond the limitations inherently associated with packed column flow rate constraints ~~[12]~~[12, 13]. These process intensification benefits have led to an increasing consideration of flow-through operations for applications outside of polishing, particularly in the development of continuous downstream processes [11, 14, 15].

Despite the utility of flow-through IEX steps, relatively few studies have systematically investigated their development and limitations. This is partly because impurity behavior is essential to understanding flow-through separations, but it is difficult to interrogate experimentally. Total HCP concentrations encountered in mAb polishing are usually on the order of 1000 ppm [16], rendering in-line detection and quantification infeasible [17, 18]. Hundreds of HCPs may be present, and their biophysical diversity leads to heterogeneous chromatographic behavior. Although process conditions are tuned to maximize HCP adsorption, they are generally less conducive to strong retention than in bind-and-elute mode. This may make impurity clearance highly

39 sensitive to the mobile phase composition and the resin.

40 Coupled with the plethora of available resins, the high sensitivity of sep-  
41 aration performance typically necessitates screening studies during process  
42 development. Commercial IEX resins are available that differ in particle  
43 size, pore size and morphology, base matrix, ligand chemistry, ionic capac-  
44 ity, and surface functionalization [19, 20]. Screening studies to determine  
45 the best resin among several alternatives can provide valuable application-  
46 specific data, and perhaps furnish general resin retentivity heuristics [21], but  
47 they offer limited transferable insights into individual HCP chromatographic  
48 retention and capacity. Column modeling has therefore been suggested as a  
49 complementary technique to elucidate the factors that contribute to impurity  
50 persistence [17, 18]. A better understanding of this phenomenon could po-  
51 tentially improve flow-through process design, expedite development, reduce  
52 costs, and increase consistency with quality by design principles. It could also  
53 help direct future *in silico* optimization efforts, which to date have focused  
54 primarily on product molecules rather than impurities.

55 Two HCP persistence mechanisms have been hypothesized for flow-through  
56 processes: product association and weak adsorption [5]. The relative im-  
57 portance of these mechanisms remains unclear. Product association is ex-  
58 pected to vary mechanistically with the therapeutic molecule and the impu-  
59 rity species, which may hinder a general understanding of the phenomenon.  
60 However, it may be tractable to analyze generally the breakthrough of weakly  
61 adsorbing impurities. Two attributes of the flow-through process simplify the  
62 analysis, namely that the process is isocratic, and that the adsorption equi-  
63 lbria of weakly adsorbing species are expected to be approximately linear

64 and independent. Assuming HCP concentrations on the order of 1000 ppm,  
65 column capacities on the order of 100 mg/ml, and the load heuristic of 1 g  
66 mAb/ml of column, only  $\sim 1\%$  of the column is expected to saturate with  
67 HCPs. The majority of the column is therefore available for weakly retained  
68 HCPs to adsorb with negligible competition for the IEX surface.

69 This work focuses on developing transferable insights into the break-  
70 through of weakly adsorbing impurities in flow-through IEX. We placed em-  
71 phasis on understanding contributions to the initial breakthrough volume, as  
72 the goal of flow-through polishing is the complete removal of trace impuri-  
73 ties. To maintain generalizability across diverse sets of HCPs, we simulated  
74 the chromatographic behavior of dilute species with a variety of transport  
75 and thermodynamic properties. Simulated initial breakthrough volumes were  
76 analytically related to the Graetz number for mass transfer, and this rela-  
77 tionship was validated experimentally. Transport and thermodynamic pa-  
78 rameters need to be estimated to use this relationship, but HCP adsorption  
79 equilibrium constants are usually unknown. To gain a better understanding  
80 of IEX adsorption equilibria, we consolidated dilute isocratic retention data  
81 from the extant literature on model species. From these data, a correlation  
82 was observed between stoichiometric displacement model (SDM) parameters  
83 that are commonly used to relate the adsorption equilibrium constant to  
84 ionic strength [22]. This thermodynamic correlation was corroborated by  
85 linear gradient elution data, suggesting a way to estimate SDM parameters  
86 approximately from one experimental measurement. These relationships pro-  
87 vide novel insights into properties that contribute to flow-through impurity  
88 clearance.

## 2. Theory and simulation

### 2.1. Column chromatography model

The 1D general rate model of column chromatography was used in this work, which describes the transport of solutes within the column interstitial volume as [20, 23, 24]:

$$\frac{\partial c}{\partial t} + v \frac{\partial c}{\partial z} = D_{ax} \frac{\partial^2 c}{\partial z^2} - \frac{3}{r_p} \left( \frac{1 - \varepsilon_c}{\varepsilon_c} \right) k_f \left( c - c_p|_{r=r_p} \right) \quad (1a)$$

with Danckwerts' boundary conditions:

$$v (c_{in} - c|_{z=0}) = -D_{ax} \frac{\partial c}{\partial z} \Big|_{z=0} \quad (1b)$$

$$\frac{\partial c}{\partial z} \Big|_{z=L_{col}} = 0 \quad (1c)$$

where  $c$  is the solute concentration in the interstitial volume,  $c_p$  is the solute concentration in the resin bead pore space,  $t$  is time,  $z$  is the column axial coordinate,  $r$  is the resin bead radial coordinate,  $c_{in}$  is the column inlet concentration,  $L_{col}$  is the column length,  $r_p$  is the resin bead radius,  $v$  is the interstitial velocity,  $D_{ax}$  is the axial dispersion coefficient,  $k_f$  is the film mass transfer coefficient, and  $\varepsilon_c$  is the column porosity ([also known as the bed void fraction](#)). Solute transport within the resin beads is described by:

$$\varepsilon_p \frac{\partial c_p}{\partial t} + (1 - \varepsilon_p) \frac{\partial q}{\partial t} = \varepsilon_p D_p \left( \frac{\partial^2 c_p}{\partial r^2} + \frac{1}{r} \frac{\partial c_p}{\partial r} \right) + (1 - \varepsilon_p) D_s \left( \frac{\partial^2 q}{\partial r^2} + \frac{1}{r} \frac{\partial q}{\partial r} \right) \quad (2a)$$

with boundary conditions of spherical symmetry at the bead center and a mass balance at the bead edge, respectively:

$$\frac{\partial c_p}{\partial r} \Big|_{r=0} = 0, \quad \frac{\partial q}{\partial r} \Big|_{r=0} = 0 \quad (2b)$$

105

$$\varepsilon_p D_p \frac{\partial c_p}{\partial r} \Big|_{r=r_p} + (1 - \varepsilon_p) D_s \frac{\partial q}{\partial r} \Big|_{r=r_p} = k_f \left( c - c_p \Big|_{r=r_p} \right) \quad (2c)$$

106 where  $q$  is the adsorbed solute concentration,  $\varepsilon_p$  is the resin bead porosity  
 107 accessible to the solute,  $D_p$  is pore diffusivity, and  $D_s$  is surface diffusivity.  
 108 Typically, the relevant initial conditions are:

$$c|_{t=0} = 0, \quad c_p|_{t=0} = 0, \quad q|_{t=0} = 0 \quad (3)$$

109 An isotherm relating  $q$  to  $c_p$  is needed to complete the column model,  
 110 along with a description of the adsorption kinetics. Since adsorption and des-  
 111 orption steps are usually much faster than the other transport processes, in-  
 112 stantaneous equilibration was assumed in this work. The Langmuir isotherm  
 113 was used:

$$q = q_{max} \left( \frac{K_L c_p}{K_L c_p + 1} \right) \quad (4)$$

114 where  $q_{max}$  is the resin saturation capacity and  $K_L$  is a thermodynamic pa-  
 115 rameter (equal to the ratio of the adsorption and desorption rate constants,  
 116  $k_a/k_d$ ). The adsorption equilibrium constant,  $K_{eq}$ , is defined by the initial  
 117 isotherm slope, which for the Langmuir model is:

$$K_{eq} \equiv \lim_{c_p \rightarrow 0} \frac{q}{c_p} = q_{max} K_L \quad (5)$$

118 Dimensions of volume concentration were used for  $q$  in this work (as opposed  
 119 to excess surface concentration), making  $K_{eq}$  a dimensionless parameter.

120 Column model equations were solved with the Chromatography Analysis  
 121 and Design Toolkit (CADET, versions 4.1.0 and 4.2.0) [25].

## 122 2.2. Simulations for analyzing breakthrough volumes

123 To ascertain qualitative differences between the breakthrough of concen-  
 124 trated and dilute species, simulations were performed for single-component  
 125 systems with load concentrations between 10 mg/ml and 1  $\mu\text{g}/\text{ml}$  with  $q_{max}$   
 126 fixed at 100 mg/ml of packed column,  $K_{eq}$  varied from 1 to 10000, and trans-  
 127 port parameters selected from representative calibrations to model species.  
 128 Exploratory simulations were then performed for a load concentration of 1  
 129  $\mu\text{g}/\text{ml}$  at variable  $v$ ,  $D_p$ ,  $L_{col}$ , and  $r_p$ .  $D_{ax}$  was estimated as a function of  $v$   
 130 from correlations for beds packed with solid particles [26], using a dilute solu-  
 131 tion diffusivity of  $7.5 \times 10^{-11} \text{ m}^2/\text{s}$  to represent general protein behavior [27].  
 132 Film mass transfer was assumed to ~~be relatively fast~~never be rate-limiting,  
 133 and  $k_f$  was set to  $1 \times 10^{-3} \text{ m/s}$  ~~for~~to essentially remove this degree of freedom  
 134 from all simulations.  $D_s$  was estimated as a function of  $K_{eq}$  using results re-  
 135 ported for a mAb of the form  $D_s = aK_{eq}^b$ , where the power law coefficient  
 136  $a = 1.66 \times 10^{-12} \text{ m}^2/\text{s}$  and  $b = -0.24$  [28]. All simulation parameters are  
 137 summarized in Supplementary Table S1.

138 Results from these simulations were correlated to relate breakthrough  
 139 volume with system parameters, making use of an effective intraparticle dif-  
 140 fusivity ( $D_{eff}$ ) to describe the combined effects of pore and surface diffusion  
 141 [29]:

$$D_{eff} \equiv D_p + \left( \frac{1 - \varepsilon_p}{\varepsilon_p} \right) K_{eq} D_s \quad (6)$$

142 This expression may be derived from Equation 2a by assuming instantaneous  
 143 equilibration and linear chromatographic conditions, as given by:

$$q = K_{eq} c_p \quad (7)$$



144 A linear isotherm was used for comparing simulations with experimental  
 145 data, and the extra-column volume, which consisted of valves and 0.75 mm  
 146 ID tubing, was described as a continuously stirred tank reactor (CSTR) in  
 147 series with a dispersed plug flow reactor (PFR) [30]. The fraction of the  
 148 extra-column volume allocated to the CSTR was fit as a function of flow  
 149 rate, and the PFR axial dispersion coefficient was set to  $1 \times 10^{-12}$  m<sup>2</sup>/s.

### 150 2.3. Models for ion-exchange adsorption equilibria

151 A number of models have been proposed to relate protein IEX adsorp-  
 152 tion equilibria to solution conditions, with perhaps the simplest and most  
 153 commonly cited being the stoichiometric displacement model (SDM), which  
 154 describes adsorption as a strict ion-exchange process [22]. From the law of  
 155 mass action, the adsorption equilibrium constant is related to ionic strength,  
 156  $I$ , as:

$$K_{eq} = \alpha I^{-\nu} \quad (8)$$

157 where  $\alpha$  is a measure of adsorption strength and  $\nu$  is the protein charac-  
 158 teristic charge, which represents the stoichiometry of ion exchange. These  
 159 parameters are commonly estimated by fitting isocratic retention data at  
 160 different ionic strengths or regressing linear gradient elution data according  
 161 to Yamamoto’s method [31]. The fundamental equation of ideal linear chro-  
 162 matography predicts that retention factors should also exhibit a power law  
 163 dependence on ionic strength [23]:

$$k' = \phi K_{eq} = \phi \alpha I^{-\nu} \quad (9)$$

164 where  $k' = (V_R - V_0)/V_0$  is the retention factor for the solute and solution  
 165 conditions of interest,  $V_R$  is the retention volume, and  $V_0$  is the ~~flow-through~~

166 retention volume under non-adsorbing conditions. For dimensionless  $K_{eq}$ ,  
 167  $\phi = (1 - \varepsilon_t)/\varepsilon_t$  is a dimensionless phase ratio where  ~~$\varepsilon_t = \varepsilon_c + (1 - \varepsilon_c)\varepsilon_p$~~   
 168  ~~$\varepsilon_t = \varepsilon_c + (1 - \varepsilon_c)\varepsilon_p = V_0/V_{column}$~~  represents the total column porosity.

169 More sophisticated colloidal models that are based on statistical thermo-  
 170 dynamics have also been proposed for estimating  $K_{eq}$  from the free energy  
 171 change of molecular adsorption,  $\Delta F$ , using some form of the general expres-  
 172 sion [32]:

$$K_{eq}(I, \text{pH}) = \int_{\Omega} \int_{z_0}^{\infty} (e^{-\Delta F(z, \Omega, I, \text{pH})/k_B T} - 1) dz d\Omega \quad (10)$$

173 where  $\Omega$  represents the protein orientation with respect to the resin and  $z$  is  
 174 the separation distance. Previous models have used continuum electrostatics  
 175 and van der Waals equations to estimate  $\Delta F$ , often making use of mean-field  
 176 approximations [33, 34]. Notably, the  $K_{eq}$  estimates from these models are  
 177 quite sensitive to  $\Delta F$  due to the Boltzmann weighting.

### 178 3. Materials and methods

#### 179 3.1. Chemicals and buffers

180 Sodium hydroxide, hydrochloric acid, sodium chloride, monobasic sodium  
 181 phosphate, dibasic sodium phosphate, acetic acid, sodium acetate, boric  
 182 acid, and sodium bicarbonate were purchased from Fisher (Fair Lawn, NJ).  
 183 Sodium carbonate, ethanolamine, and 2000 kDa blue dextran were pur-  
 184 chased from Sigma (St. Louis, MO). Fluorescein isothiocyanate (FITC) and  
 185 dimethylformamide were purchased from Thermo Scientific (Rockford, IL).  
 186 All chemicals were used without any further purification.

187 Buffer solutions were prepared to the desired constituent concentrations  
188 at room temperature with deionized water from an EMD Millipore Milli-Q  
189 system ( $> 18.2 \text{ M}\Omega \text{ cm}$ ). Low and high ionic strength buffers were prepared  
190 with 0 and 1 M added NaCl, respectively. Buffer pH and conductivity were  
191 measured with a Cole-Parmer PC200 meter, and pH adjustments were made  
192 with concentrated sodium hydroxide and hydrochloric acid solutions. Prior  
193 to use, buffer solutions were filtered with Fisher  $0.2 \mu\text{m}$  aPES membranes  
194 (Pittsburgh, PA).

### 195 3.2. *Proteins*

196 Lyophilized hen egg white lysozyme was purchased from Sigma (St. Louis,  
197 MO). A mAb sample was obtained from the protein A eluate pool of a  
198 manufacturing process at Bristol Myers Squibb (Devens, MA). The mAb  
199 was supplied at 35 mg/ml and stored at  $-80^\circ\text{C}$ . It was thawed prior to use  
200 and exchanged into the desired buffer using a GE Sephadex G25 desalting  
201 column, and protein solution concentrations were measured with a Thermo  
202 Scientific NanoDrop Lite spectrophotometer.

### 203 3.3. *Resins and chromatography equipment*

204 All chromatographic measurements were performed using an Amersham  
205 Biosciences Äkta Explorer equipped with a P-960 sample pump, fraction  
206 collector, 10 mm UV flow cell, and in-line pH and conductivity meters. SP  
207 Sepharose FF, a strong cation-exchanger with a nominal particle diameter  
208 of  $90 \mu\text{m}$ , was purchased from GE Healthcare (Uppsala, Sweden). It was  
209 exchanged three times into a pH 7.0 (25 mM sodium phosphate) buffer by  
210 centrifugation and decantation prior to slurry packing (50%) into a  $0.5 \times$

211 10 cm Waters AP chromatography column to a final bed volume of 2.0 ml.  
212 Column packing was validated at pH 7.0 by injecting a 100  $\mu$ l pulse of a 1 M  
213 NaCl buffer and verifying that the conductivity trace asymmetry fell within  
214 commonly accepted limits [19].

### 215 3.4. *Linear gradient elution*

216 The retention of lysozyme on SP Sepharose FF was measured with lin-  
217 ear gradient elution at pH 5.0, 7.0, and 9.0 (in 75 mM sodium acetate, 25  
218 mM sodium phosphate, and 25 mM ethanolamine buffers, respectively) at a  
219 superficial velocity of 300 cm/h. Feed solutions were prepared by dissolving  
220 lysozyme in the appropriate low ionic strength buffer at  $\sim$ 20 mg/ml, and  
221 lyophilized excipients were removed with three buffer exchange cycles using  
222 EMD Millipore Amicon centrifugal filter units (Cork, Ireland). Lysozyme  
223 solutions were filtered using Thermo Scientific 0.2  $\mu$ m PVDF membranes  
224 (Rockwood, TN) and diluted to  $\leq$  10 mg/ml prior to use. The SP Sepharose  
225 FF column was equilibrated in the relevant low ionic strength buffer for 5  
226 ~~CV~~column volumes (CV), and gradient elution from 0 to 1 M NaCl began  
227 shortly after a 100  $\mu$ l sample injection. The gradient elution volume was  
228 varied from 10 to 50 CV in increments of 10 CV, and the conductivity at  
229 peak elution was determined. This was used with a correlation for the in-  
230 strument's conductivity response to estimate the peak elution ionic strength.  
231 Columns were regenerated at least once every five runs via 15 minutes of  
232 contact with 0.5 M NaOH.

233 *3.5. Breakthrough measurements during isocratic elution*

234 To validate corresponding simulation results, frontal loading chromatog-  
235 raphy was performed with dilute lysozyme solutions on SP Sepharose FF  
236 at pH 7.0 (20 mM sodium phosphate) under conditions of weak to moder-  
237 ate retention. The requisite ionic strength for such retention was estimated  
238 from linear gradient elution measurements and tuned empirically to target  
239 a breakthrough volume of  $\sim 10$  CV. A feedstock of 10  $\mu\text{g}/\text{ml}$  lysozyme was  
240 used to mimic the low impurity concentrations encountered in flow-through  
241 polishing, and outlet concentrations were estimated from the absorbance at  
242 215 nm with a computed extinction coefficient [35]. At such low protein con-  
243 centrations, appreciable variability in component separation was observed  
244 between replicate feedstock preparations. For this reason, feedstock batches  
245 were prepared at the desired ionic strength in sufficient volume to service  
246 the entire set of measurements, and a batch exhibiting minimal component  
247 separation was selected for use. For each measurement, a 10 CV equilibra-  
248 tion period preceded a 20 CV load through the Äkta's sample pump, after  
249 which the column was washed as described previously. This was performed  
250 at superficial velocities ranging from 30 to 305 cm/h in a randomized order,  
251 and comparable breakthrough measurements were made under non-adsorbing  
252 conditions (at high ionic strength).

253 To qualitatively validate simulated trends under conditions of greater  
254 relevance to flow-through polishing, breakthrough behavior was also demon-  
255 strated with fluorescently labelled lysozyme in the presence of a mAb. Lysozyme  
256 was conjugated with FITC according to the manufacturer's protocol, produc-  
257 ing a complex product mixture with different label ratios and conjugation

258 sites. A fraction of the conjugation products exhibiting relatively homo-  
259 geneous chromatographic behavior was required to effectively represent the  
260 breakthrough of an individual impurity. The conjugation product mixture  
261 was therefore subjected to repeated fractionation on SP Sepharose FF at  
262 pH 9.5 (20 mM sodium carbonate) using shallow ionic strength steps during  
263 elution. Average label ratios and FITC-lysozyme concentrations in the col-  
264 lected fractions were estimated from the absorbance profiles at 280 and 495  
265 nm. Once a sufficiently homogeneous fraction was obtained, it was spiked  
266 to  $\sim 10 \mu\text{g/ml}$  in a mAb solution at pH 9.5, where the mAb was observed  
267 to have negligible retention on SP Sepharose FF. As with lysozyme break-  
268 through measurements, an ionic strength was identified to achieve weak to  
269 moderate FITC-lysozyme retention, and breakthrough profiles were acquired  
270 over a range of superficial velocities in randomized order. The mAb concen-  
271 tration was dilute due to material constraints, and the eluate absorbance at  
272 495 nm was used to distinguish FITC-lysozyme breakthrough from the mAb  
273 flow-through.

### 274 3.6. *Parameter estimation for test simulations*

275 The breakthrough of lysozyme fed in dilute solution on SP Sepharose FF  
276 was simulated in CADET as described in Section 2.2. Sequential parameter  
277 estimation was used to reduce the dimensionality of the optimization space  
278 when fitting the experimental data. Duplicate column bypass profiles were  
279 acquired with lysozyme at each flow rate used in breakthrough measurements.  
280 These were fit with a trust region optimization algorithm interfaced by the  
281 `Scipy least_squares` function to determine the extra-column volume (which  
282 is invariant with flow rate) and the fraction of that volume allocated to the

283 CSTR in the system model (which varies with flow rate). Triplicate blue  
 284 dextran pulses, monitored with the absorbance at 280 nm, were then passed  
 285 through the column to measure  $\varepsilon_c$ . Peak tailing was observed, which is  
 286 attributable to size heterogeneity in the blue dextran tracer. For this reason,  
 287 the average peak mode was used instead of the average first moment to  
 288 estimate  $\varepsilon_c$ .

289 Duplicate lysozyme pulses under non-adsorbing conditions (high ionic  
 290 strength) were then passed through the column at each flow rate. Retention  
 291 volumes under non-adsorbing conditions were computed from the first mo-  
 292 ment of the absorbance profiles and used to estimate  $\varepsilon_p$ . These profiles were  
 293 then fit to estimate  $D_p$  (which is invariant with flow rate) and  $D_{ax}$  (which  
 294 varies with flow rate). An evolutionary algorithm interfaced by the `Scipy`  
 295 `differential_evolution` function was used to estimate  $D_p$  and a coarse  
 296 dependence of  $D_{ax}$  on flow rate. The estimated value of  $D_p$  was retained,  
 297 and the coarse  $D_{ax}$  estimates were refined using the `least_squares` func-  
 298 tion. Without further refinement, the estimated parameters were validated  
 299 by modeling dilute lysozyme breakthrough profiles under non-adsorbing con-  
 300 ditions (at high ionic strength). Dilute breakthrough profiles under adsorbing  
 301 conditions (at lower ionic strength) were then simulated by fitting  $K_{eq}$  and  
 302  $D_s$  (which are invariant with flow rate) using the `differential_evolution`  
 303 function.

## 304 4. Results and discussion

### 305 4.1. Breakthrough volume correlation

306 In the context of polishing operations, the value of flow-through IEX is its  
307 ability to completely remove as many trace HCPs as possible from a mAb so-  
308 lution. Challenges inherent in this task may not be readily apparent because  
309 it is difficult to study the chromatographic behavior of dilute solutes em-  
310 pirically. A column simulation was therefore used to investigate differences  
311 between the breakthrough of overloaded products and dilute impurities. An  
312 initial set of single-component simulations were performed at different feed  
313 concentrations and  $K_{eq}$  values. Transport and geometric parameters were  
314 held fixed, and a Langmuir isotherm with constant  $q_{max} = 100$  mg/ml of  
315 column was used. Results for extreme feed concentrations are juxtaposed  
316 in Figure 1, and analogous plots for intermediate concentrations are shown  
317 in Supplementary Figure S1. As may be expected for the system with a 10  
318 mg/ml feed, retention increased with  $K_{eq}$  until the column saturated, and  
319 breakthrough profiles became sharper as column saturation was approached.  
320 This self-sharpening behavior did not occur for the 1  $\mu$ g/ml feed over the vol-  
321 ume scale simulated, revealing dilute solute breakthrough profiles to become  
322 more diffuse as retention increases due to the essentially linear isotherm. An  
323 order-of-magnitude difference can be observed between the load volumes cor-  
324 responding to the initial breakthrough and the inflection point in the more  
325 strongly retained profiles. This illustrates an appreciable challenge; even if  
326 trace impurities exhibit strong to moderate retention, their diffuse distribu-  
327 tion on the column may lead to relatively early breakthrough and preclude  
328 adequate clearance.



329 To better understand this challenge, the load volume corresponding to 1%  
 330 breakthrough was identified and plotted against  $K_{eq}$  (Figure 2). A couple  
 331 of trends may be observed. When column saturation is not approached, the  
 332 breakthrough volume appears to scale linearly with  $K_{eq}$ . An ideal limit on the  
 333  $K_{eq}$  dependence may be roughly estimated as  $CV_{breakthrough} = \varepsilon_t + (1 - \varepsilon_t)K_{eq}$   
 334 based on the fundamental equation of ideal linear chromatography. As  
 335 shown, the slope of the dilute solute series differs markedly from the ideal  
 336 limit, due largely to finite transport rates that limit the system's approach to  
 337 equilibrium. The discrepancy between the observed and ideal behavior there-  
 338 fore suggests that transport rate optimization may provide a useful means of  
 339 improving impurity clearance in flow-through processes.

340 Further simulations showed that the slope of the  $CV_{1\%}$  vs.  $K_{eq}$  curve for  
 341 dilute solutes varies with transport and geometric parameters. This obser-  
 342 vation indicates that the initial breakthrough volume may be approximately  
 343 described as:

$$CV_{\text{breakthrough}1\%} = \varepsilon_t + (1 - \varepsilon_t)K_{eq}f_{1\%} \quad (11)$$

344 where  $f_{1\%}$  is a function of transport parameters that modifies the break-  
 345 through volume dependence on  $K_{eq}$ , with  $0 \leq f \leq 1$  and  $0 \leq f_{1\%} \leq 1$ . The form  
 346 of this relationship has two implications: it suggests that thermodynamic  
 347 and transport contributions to impurity breakthrough may be independent,  
 348 and that transport contributions may be described by a single factor, which  
 349 can presumably be expressed generally in terms of a dimensionless group.

350 Based on this rationale, a second set of simulations were performed for  
 351 a 1  $\mu\text{g}/\text{ml}$  feed. A variety of thermodynamic, transport, and geometric pa-  
 352 rameters were used, and  $f_{1\%}$  was computed from the results by rearranging

Equation 11. However,  $\varepsilon_c$  was used in place of the  $\varepsilon_t$  intercept when computing  $f_{1\%}$  from Equation 11 because ~~this~~ it was observed to better describe the simulated breakthrough volumes in the limit of  $K_{eq} \rightarrow 0$ , which may be attributable to finite transport rates limiting the solute exploration of intraparticle void volumes. Based on the dynamic binding capacity correlation presented by Chen et al. [36],  ~~$f_{1\%}$~~  is plotted in Figure 3 against ~~the dimensionless group~~:

$$Gz_{eff} = \frac{vd_{part}^2}{D_{eff}L_{col}} = Péc_{p-eff} \frac{d_{part}}{L_{col}}$$

~~which represents a Graetz number for mass transfer, written as a product of two dimensionless subgroups. The first is:~~

$$\underbrace{Gz_{eff} = \frac{d_{part}^2/D_{eff}}{L_{col}/v}}_{\text{which represents the ratio of characteristic time scales for intraparticle diffusion and axial convection. This may be alternatively written as } Gz_{eff} = (d_{part}/L_{col}) Péc_{p-eff}, \text{ where the particle Péclet number, which represents the } Péc_{p-eff} = vd_{part}/D_{eff} \text{ represents a ratio of convective to diffusive transport rates on the length scale of the resin particle, when } D_{eff} \text{ is used to describe the diffusional transport. The second subgroup is a ratio of characteristic length scales, namely the resin particle diameter and the column length.}}$$

~~The second subgroup is a ratio of characteristic length scales, namely the resin particle diameter and the column length.~~

Figure 3 indicates that Equation 11 with a modified intercept may be written more specifically as:

$$CV_{\%} = \varepsilon_c + (1 - \varepsilon_t)K_{eq}f_{\%} \left( \frac{vd_{part}^2}{D_{eff}L_{col}} \underbrace{Gz_{eff}} \right) \quad (13)$$

371 where  $CV_{\%}$  is the load volume at a given percent breakthrough threshold, and  
 372  $f_{\%}$  is the corresponding function of ~~transport and geometric parameters~~[the](#)  
 373 [Graetz number](#). Shown are two series for illustrative breakthrough thresh-  
 374 olds, but any arbitrary threshold could be selected. Some noise is apparent,  
 375 but it is largely attributable to the approximation that  $CV_{\%}$  is invariant with  
 376 transport parameters when  $K_{eq} = 0$ . Within the noise, it is clear that  $CV_{\%}$   
 377 decays monotonically with the value of the Graetz number.

378 To maximize the load volume before breakthrough, which is approxi-  
 379 mated by  $CV_{1\%}$ , the Graetz number should be minimized. Unfortunately,  
 380 doing so comes at a cost to throughput if  $v$  is reduced or an increase in col-  
 381 umn pressure drop if  $d_{part}$  is decreased or  $L_{col}$  is increased. This correlation  
 382 does therefore not suggest a free way of improving flow-through processing.  
 383 What it does quantify, however, is the inherent balance between separa-  
 384 tion performance and throughput. Notably, the initial breakthrough volume  
 385 decays much more quickly with the Graetz number than the intermediate  
 386 breakthrough volume. The benefit of optimizing the system geometry and  
 387 transport rates is expected to be only marginal for typical process conditions,  
 388 which correspond to Graetz numbers on the flat part of the 1% breakthrough  
 389 curve in Figure 3. However, if the Graetz number could be decreased below  
 390  $\sim 5$ , the benefit would become much more pronounced.

391 An interesting subtlety that is implicit in this correlation is the effect  
 392 of thermodynamics on intraparticle transport rates, which is described by  
 393 the dependence of  $D_{eff}$  on  $K_{eq}$ .  $D_{eff}$  captures the combined effects of pore  
 394 and surface diffusion in the dilute solute limit by describing the effective  
 395 intraparticle diffusivity as a weighted sum, where  $K_{eq}$  weights the relative

importance of surface diffusion. IEX surface diffusivities have been shown to follow a power law of the form  $D_s = aK_{eq}^b$ , ~~where~~ [28], where the empirical power law coefficients  $a$  and  $b$  are expected to be on the order of  $D_p$  and  $-0.5$ , respectively, with  $b < 0$  [28][37, 38]. This leads to competing effects: increasing  $K_{eq}$  decreases the surface diffusivity but simultaneously increases the driving force for surface diffusion, with the effect that dominates being determined by whether ~~the power law exponent~~  $b > -1$ . If this is the case, increasing  $K_{eq}$  leads to an increase in  $D_{eff}$  and a reduction in the Graetz number, meaning that both thermodynamic and transport contributions to delaying impurity breakthrough are improved simultaneously.

Equation 13 also enables estimation of the problematic  $K_{eq}$  threshold below which impurities will persist via weak adsorption. To gain a sense for the threshold's order of magnitude, an illustrative calculation was performed with a spline fit to the  $f_{1\%}$  curve shown in Figure 3. Results for a 20 cm column with a 6 minute residence time are shown in Figure 4 as a function of load volume using  $\varepsilon_c = 0.35$ ,  $\varepsilon_p = 0.5$ ,  $d_{part} = 50 \mu\text{m}$ ,  $D_p = 10^{-11} \text{ m}^2/\text{s}$ , and  $D_s$  estimated from the power law discussed previously. The results show that, for load volumes on the order of 100 CV, HCPs with  $K_{eq} < \sim 400$  are liable to break through before the end of loading. This would be equivalent to a retention factor of  $\sim 200$  for an input feed pulse under identical conditions.

#### 4.2. Validation of the breakthrough volume correlation

To validate the breakthrough volume correlation that was developed from simulation results, frontal loading chromatography was performed ~~using dilute feeds~~ as described in Section 3.5. Using a highly pure feedstock was found to be essential for demonstrating the behavior of individual species in the dilute

421 limit. Various model proteins were tested with AEX and CEX resins, but  
 422 only lysozyme was found to be commercially available in sufficient purity. It  
 423 was therefore used with SP Sepharose FF, despite the fact that most mAb  
 424 flow-through purification processes are performed with AEX resins. Figure 5  
 425 shows breakthrough profiles acquired by loading lysozyme onto SP Sepharose  
 426 FF at 10  $\mu\text{g}/\text{ml}$  under high and low ionic strength conditions. The flow rate  
 427 was varied to change the value of the Graetz number, and an increase in  
 428 breakthrough volume was observed at lower flow rates. Equation 13 indi-  
 429 cates that transport parameters should have a more pronounced effect on  
 430 the initial breakthrough volume when retention is stronger, due to the  $K_{eq}$   
 431 weighting of  $f\%$ . This is consistent with the observed behavior at low ionic  
 432 strength, where  $K_{eq}$  is higher. Column models with fit parameters showed  
 433 excellent agreement at high ionic strength. Although not perfect, the agree-  
 434 ment was also quantitatively close at low ionic strength. Discrepancies from  
 435 simulation under these conditions may be attributable to feedstock compo-  
 436 nent separation, as well as inaccuracies in describing extra-column effects  
 437 with the simplified model of a CSTR in series with a PFR, which are more  
 438 pronounced at low flow rates.

439 The breakthrough volume correlation was also qualitatively validated us-  
 440 ing FITC-lysozyme in the presence of a mAb. As with lysozyme measure-  
 441 ments, the flow rate was varied to change the value of the correlation variable  
 442 under conditions of weak FITC-lysozyme retention. Figure 6 shows break-  
 443 through profiles that were computed from absorbance measurements at 495  
 444 and 280 nm, using the average FITC-lysozyme label ratio estimated during  
 445 fractionation of the conjugation products. Appreciable measurement noise

446 is apparent in the FITC-lysozyme profiles due to the low load concentra-  
 447 tion of  $\sim 10 \mu\text{g/ml}$ . Nonetheless, the observed trend is consistent with the  
 448 previous results: using lower flow rates delayed the onset of FITC-lysozyme  
 449 breakthrough. This was not modeled explicitly due to the multicomponent  
 450 nature of the FITC-lysozyme conjugation products, but it does support the  
 451 breakthrough volume correlation qualitatively.

#### 452 4.3. Correlation of SDM parameters

453 A knowledge of phenomenological properties is required to estimate break-  
 454 through volumes using Equation 13, specifically values of  $D_p$ ,  $D_s$ , and  $K_{eq}$ .  
 455 However, such information is typically unknown for impurities such as HCPs  
 456 on IEX resins. Methods of estimating these properties could therefore be  
 457 useful, and some heuristics already exist for estimating the two intraparti-  
 458 cle diffusivities. For instance,  $D_s$  is typically an order of magnitude smaller  
 459 than  $D_p$  under the relevant conditions of weak adsorption, and  $D_p$  should  
 460 theoretically scale with the free solution diffusivity,  $D_0$ , according to [19]:

$$D_p = \frac{\varepsilon_p \psi_p}{\tau_p} D_0 \quad (14)$$

461 where  $\psi_p$  and  $\tau_p$  are the diffusional hindrance coefficient and the resin tor-  
 462 tuosity factor, respectively. If a putative weakly adsorbing HCP has been  
 463 identified, its molar mass may be used with the Stokes-Einstein equation or  
 464 correlations to estimate  $D_0$  [27]. Appreciable uncertainty surrounds the value  
 465 of  $\psi_p/\tau_p$ , but 0.2 appears to be representative of its order of magnitude [39].  
 466 Assuming a value like this enables rough estimation of  $D_p$  and  $D_s$  based on  
 467 protein mass.

468 This leaves the estimation of  $K_{eq}$ , which varies with the solution condi-  
 469 tions, as the main obstacle to applying the breakthrough volume correlation  
 470 in practice. It is expected that  $K_{eq}$  will follow a power law in ionic strength  
 471 of the SDM form, but the dependence on pH is less well defined. Several elec-  
 472 trostatics models have been proposed for describing this behavior, but doing  
 473 so predictively remains an open problem. In the absence of such predictive  
 474 tools, insights may be drawn from experimental data on the retention of  
 475 model proteins. To facilitate observation of system-independent trends, iso-  
 476 cratic  $k'$  measurements at different ionic strengths were consolidated from the  
 477 extant literature on 230 protein-pH-resin combinations, as shown in Supple-  
 478 mentary Figure S2 [and Table S2](#) [21, 30, 40–47]. These data were regressed  
 479 according to Equation 9 to extract the SDM parameter  $\nu$  and the quasi-  
 480 SDM parameter  $\phi\alpha$ , which are plotted against each other in Figure 7 and  
 481 observed to follow a significant intrinsically linear correlation (as determined  
 482 by regressor t-tests).

483 This previously unobserved correlation is expected to be a consequence of  
 484 adsorption thermodynamics rather than resin morphology, so the true rela-  
 485 tionship represented in Figure 7 is suspected to be between  $\nu$  and  $\alpha$ . Inverse  
 486 size exclusion chromatography data are unavailable for the majority of the  
 487 resins studied, which precludes estimation of  $\phi$  values that are specific to  
 488 protein-resin pairs. However, all phase ratios are expected to be of compara-  
 489 ble magnitude, and  $\alpha$  may span several orders of magnitude, so the inclusion  
 490 of  $\phi$  is not expected to introduce much noise in the correlation. The fact  
 491 that  $\phi\alpha$  spans several orders of magnitude makes the prediction interval ap-  
 492 preciably broad in terms of absolute values, and explaining some of the noise

493 with additional factors would be advantageous. The ion-exchange type was  
494 tested for this purpose, and a significant difference was found between the  
495 correlations for AEX and CEX isocratic data (by including an interaction  
496 with the categorical variable). The mechanisms underlying this difference  
497 are not understood.

498 The utility of these relationships lies in their ability to remove one degree  
499 of freedom from the analysis of retention data. A series of measurements  
500 are usually performed to estimate SDM parameters, either from isocratic  
501 retention at different ionic strengths or elution peak ionic strength under  
502 linear gradient salt elutions of different volumes. The identified correlation  
503 could potentially enable the estimation of SDM parameters from only one  
504 experimental measurement. This concept is illustrated in Figure 8, which  
505 shows linear gradient elution data for lysozyme on SP Sepharose FF, collected  
506 as part of this study, plotted in the regression space for Yamamoto’s GH  
507 analysis [31]. Juxtaposed with these data are predictions from the correlation  
508 of isocratic CEX SDM parameters. Close agreement is observed between  
509 predictions and experiment at each of the pH values tested. One point in each  
510 series would be sufficient to estimate roughly the value of  $\nu$ , and therefore  
511 that of  $\phi\alpha$  when using the correlation. However, because the uncertainty in  
512  $\phi\alpha$  is with respect to its order of magnitude, the absolute value of its estimate  
513 needs to be interpreted with caution.

514 The linear gradient elution data were regressed according to Yamamoto’s  
515 method, and the SDM parameters are plotted in Figure 7. Analogous CEX  
516 gradient elution data for a mAb and its aggregates that had relatively large  
517 SDM parameters were also obtained from the literature and included in Fig-



518 ure 7 [48]. Both sets of gradient elution data are consistent with the correla-  
 519 tion of isocratic CEX data within the 95% prediction interval. This confirms  
 520 the correlation’s utility in removing a degree of freedom from the analysis of  
 521 retention data. Unfortunately, it does not eliminate the need for some ex-  
 522 perimental data in estimating  $K_{eq}$ , but it may reduce the burden of doing so.  
 523 As proteomic techniques advance, it may become possible to quantitatively  
 524 measure HCP retention. This would foreseeably be an expensive operation,  
 525 and it may be essential to minimize the number of requisite measurements.  
 526 The SDM parameter correlation may be useful in such an application.

527 From an intuitive perspective, the essential features of this correlation ap-  
 528 pear to be broadly consistent with theory. If the characteristic charge indeed  
 529 represents the number of attractive point charge interactions between the  
 530 protein and the IEX resin, it seems reasonable to hypothesize that increas-  
 531 ing this number at constant ionic strength would result in an approximately  
 532 proportional increase in the adsorption free energy. This would result in an  
 533 order-of-magnitude increase in  $K_{eq}$ , which at the unit concentration of ionic  
 534 strength is equivalent to  $\alpha$ ; such an argument may rationalize the positive  
 535 and log-linear correlation between  $\nu$  and  $\phi\alpha$ . The reality is necessarily more  
 536 complex, however, as the adsorption free energy would theoretically be a  
 537 function of the pH, resin, protein, and the protein’s adsorbed orientation.  
 538 These are incorporated in a statistical thermodynamic calculation of  $K_{eq}$  in  
 539 Equation 10, albeit implicitly with respect to the resin and protein species.  
 540 Perhaps the dependence on all four variables can be generalized approxi-  
 541 mately by the characteristic charge, but this would require evaluation and  
 542 proof using a rigorous biophysical model. Developing that proof, or at least

543 exhibiting behavior that is consistent with the empirical SDM parameter cor-  
544 relation (Figure 7), may provide a good validity test for molecular adsorption  
545 models.

## 546 **5. Conclusions**

547 Unlike concentrated solutes, the approximately linear chromatographic  
548 behavior of dilute impurities can lead to diffuse breakthrough. This can  
549 make it challenging to remove flow-through impurities completely, but it  
550 also enables breakthrough volumes to be analytically related to phenom-  
551 logical properties. Specifically, the Graetz number can describe transport  
552 contributions to dilute solute breakthrough volumes generally when an effec-  
553 tive intraparticle diffusivity is used. Knowledge of this relationship enables  
554 a problematic equilibrium constant threshold to be estimated, which is on  
555 the order of 400 for typical process conditions. However, for this to be used  
556 in practice, the IEX adsorption equilibria of impurities of interest must be  
557 known. The correlation of SDM model parameters may be useful in estimat-  
558 ing such data approximately from one linear gradient elution measurement,  
559 and it may serve as a validity test for molecular adsorption models as well.

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565     The authors declare that they have no known competing financial inter-  
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568 **CRedit authorship contribution statement**

569     **Chase E. Herman:** Conceptualization, Data curation, Formal analy-  
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571     nal draft. **Xuankuo Xu:** Conceptualization, Funding acquisition, Project  
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578 **Supplementary material**

579     Supplementary material is associated with this article.

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