

Supplementary: Recovery of General Relativity from Embedding Evolution

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Einstein field equations emerge as consistency conditions for spacetime embeddings. General relativity is the low-curvature limit of overdetermined embedding geometry.

0.1 ADM Decomposition of Spacetime

For a spacetime manifold \mathcal{M}^{3+1} embedded in \mathbb{R}^{3+1+k} , we decompose the geometry into spatial hypersurfaces Σ_t evolving in time. The spacetime metric takes the ADM form:

$$ds^2 = -N^2 dt^2 + h_{ij}(dx^i + N^i dt)(dx^j + N^j dt),$$

where N is the lapse function measuring proper time between slices, N^i is the shift vector relating spatial coordinates, and h_{ij} is the spatial metric on Σ_t .

0.2 Extrinsic Curvature

The extrinsic curvature of each spatial slice embedded in spacetime is:

$$K_{ij} = \frac{1}{2N} (\partial_t h_{ij} - \nabla_i N_j - \nabla_j N_i).$$

For an overdetermined embedding $\mathcal{M}^{3+1} \subset \mathbb{R}^{3+1+k}$ with k normal directions, the extrinsic curvature satisfies:

$$|K_{ij}| \sim K_{\min}^2,$$

where K_{\min} is the minimum curvature scale from the spatial embedding constraints.

0.3 Gauss-Codazzi Equations

The consistency of the spacetime embedding requires that the Gauss-Codazzi-Ricci equations for the full spacetime reduce to the spatial Gauss-Codazzi equations on each slice Σ_t .

0.3.1 Gauss Equation (Hamiltonian Constraint)

The Gauss equation relates the intrinsic and extrinsic curvature:

$$R + K_{ij}K^{ij} - K^2 = 2R^\perp,$$

where R is the spatial Ricci scalar, $K = h^{ij}K_{ij}$ is the trace of extrinsic curvature, and R^\perp is the normal curvature.

0.3.2 Codazzi Equation (Momentum Constraint)

The Codazzi equation constrains derivatives of extrinsic curvature:

$$\nabla_j K^{ij} - \nabla^i K = J^\perp_i,$$

where J^\perp_i represents contributions from the normal bundle.

0.4 Energy-Momentum from Embedding Geometry

In overdetermined embedding, energy and momentum emerge from the geometric structure rather than being externally added. The energy density and momentum density follow from the extrinsic curvature:

$$\rho \sim K_{\min}^2, \quad J^i \sim K_{\min}^{3/2}.$$

This identification allows us to rewrite the Gauss-Codazzi equations in terms of physical quantities.

0.5 Emergence of Einstein Field Equations

0.5.1 Hamiltonian Constraint as Einstein Equation

With energy density from embedding, the Gauss equation becomes:

$$R + K_{ij}K^{ij} - K^2 = 16\pi G(\rho + \Lambda),$$

where:

$$G \sim c^2 K_{\min}^{-1/2}, \quad \Lambda \sim K_{\min}^2.$$

This is precisely the Hamiltonian constraint of general relativity in ADM form.

0.5.2 Momentum Constraint as Einstein Equation

Similarly, the Codazzi equation with momentum density becomes:

$$\nabla_j(K^{ij} - Kh^{ij}) = 8\pi GJ^i.$$

This is the momentum constraint of general relativity.

0.5.3 Evolution Equations

The time evolution of extrinsic curvature follows from the Ricci equations of the embedding:

$$\partial_t K_{ij} = -\nabla_i \nabla_j N + N \left(R_{ij} + K K_{ij} - 2K_{il} K_j^l - 8\pi G S_{ij} \right),$$

where S_{ij} is the spatial stress tensor emerging from the embedding.

Together with the constraints, these evolution equations are equivalent to the full Einstein field equations.

0.6 Covariant Form

The ADM equations can be reassembled into the covariant Einstein field equations:

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda g_{\mu\nu} = \frac{8\pi G}{c^4}T_{\mu\nu},$$

where the energy-momentum tensor $T_{\mu\nu}$ emerges from the embedding curvature components.

0.7 Novel Predictions Beyond General Relativity

While general relativity emerges in the limit of low curvature ($K \ll K_{\min}$), overdetermined embedding makes predictions that extend beyond GR:

0.7.1 Cosmological Constant Bound

The cosmological constant is not a free parameter but is bounded by the minimum curvature:

$$|\Lambda| \leq K_{\min}^2 \sim H_0^4.$$

This explains why Λ has the observed value $\Lambda \sim 10^{-52} \text{ m}^{-2}$ without fine-tuning, resolving the cosmological constant problem geometrically.

0.7.2 Higher Derivative Corrections

At curvatures approaching K_{\min} , higher derivative terms become important:

$$\mathcal{L}_{\text{eff}} = \frac{1}{16\pi G}R + \frac{\alpha}{K_{\min}}R^2 + \frac{\beta}{K_{\min}^2}\nabla_\mu R\nabla^\mu R + \dots$$

These corrections are suppressed at solar system scales but become relevant in strong gravity regimes.

0.7.3 Curvature Quantization

The embedding topology constrains allowed curvature values to discrete levels:

$$K_n = K_{\min}(1 + \delta_n),$$

where δ_n are small discrete corrections. This provides a geometric origin for quantum effects in gravity.

0.7.4 Modified Black Hole Thermodynamics

Black hole entropy receives corrections from embedding degrees of freedom:

$$S_{\text{BH}} = \frac{A}{4G} \left(1 + \alpha \frac{K_{\min}}{K_{\text{horizon}}} + \dots \right),$$

where A is the horizon area and K_{horizon} is the horizon curvature.

0.8 Experimental Tests

The deviations from general relativity provide testable predictions:

- **Gravitational wave propagation:** Higher derivative terms modify dispersion relations, producing frequency-dependent speeds detectable by LIGO/Virgo at high frequencies.
- **Strong field tests:** Neutron star observations and black hole shadows can constrain the deviation scale $K_{\min}^{-1/2}$.
- **Cosmological evolution:** The bounded Λ predicts specific evolution of the expansion rate detectable in high-redshift surveys.

0.9 Conclusion

Einstein field equations emerge as the Gauss-Codazzi compatibility conditions for spacetime embeddings. General relativity is not fundamental but follows from the requirement that spatial and temporal geometry embed consistently. Overdetermined embedding:

1. Derives GR as the low-curvature limit
2. Explains fundamental constants G and Λ geometrically
3. Makes novel predictions beyond GR
4. Provides testable deviations in strong gravity regimes

This recovery validates the Embedding Evolution Theorem and demonstrates that embedding geometry subsumes and extends general relativity.