

WKB approximate solutions for standard baroclinic modes

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In this notes I derive approximate WKB solutions to the standard baroclinic modes of physical oceanography. The elementary textbook example with constant buoyancy frequency is recovered as a special case.

1 Pressure modes

The standard baroclinic modes for pressure, here denoted $\mathbf{p}_n(z)$, is defined by the regular Sturm-Liouville eigenproblem

$$\mathbf{L}\mathbf{p}_n = -\kappa_n^2 \mathbf{p}_n, \quad (1)$$

with homogeneous Neumann boundary conditions

$$@z = -h, 0 : \quad \mathbf{p}'_n = 0, \quad (2)$$

and the self-adjoint Linear operator

$$\mathbf{L} \stackrel{\text{def}}{=} \frac{d}{dz} \frac{f_0^2}{N^2} \frac{d}{dz}, \quad (3)$$

is a self-adjoint linear operator. Hence the eigenmodes, \mathbf{p}_n , are orthogonal. The real eigenvalues, κ_n , are the deformation wavenumber of the n 'th mode. It is convenient to normalize the eigenmodes to have the unit L^2 -norm

$$\frac{1}{H} \int_{-h}^0 \mathbf{p}_n \mathbf{p}_m dz = \delta_{mn}, \quad (4)$$

where δ_{mn} is the Dirac delta. Equation (1) can be rewritten as

$$\left(\frac{f_0}{N}\right)^2 \mathbf{p}_n'' + \left[\left(\frac{f_0}{N}\right)^2\right]' \mathbf{p}_n' + \kappa_n^2 \mathbf{p}_n = 0. \quad (5)$$

Introducing the following definitions

$$\epsilon \stackrel{\text{def}}{=} \frac{1}{\kappa_n} \quad \text{and} \quad S^2(z) \stackrel{\text{def}}{=} \left(\frac{N}{f_0}\right)^2. \quad (6)$$

we have the renotated equation

$$\epsilon^2 \mathbf{p}_n'' - \epsilon^2 [\log S^2(z)]' \mathbf{p}_n' + S^2(z) \mathbf{p}_n = 0. \quad (7)$$

In the WKB spirit we assume that $S^2(z)$ is slowly varying i.e., the buoyancy frequency $N^2(z)$ does not vary very fast. (This assumption may be problematic near the base of the mixed-layer.)

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We also assume that ϵ is small; the accuracy of the WKB solution improves with mode number. We now make the exponential approximation (e.g., Bender and Orszag)

$$\mathbf{p}_n^e \stackrel{\text{def}}{=} e^{Q(z)/\epsilon} . \quad (8)$$

Hence

$$\mathbf{p}_n^{e'} = \frac{Q'(z)}{\epsilon} \mathbf{p}_n^e , \quad (9)$$

and

$$\mathbf{p}_n^{e''} = \left[\left(\frac{Q'(z)}{\epsilon} \right)^2 + \frac{Q''(z)}{\epsilon} \right] \mathbf{p}_n^e , \quad (10)$$

Next we expand $Q(z)$ in powers of ϵ

$$Q(z) = Q_0(z) + \epsilon Q_1(z) + \epsilon^2 Q_2(z) + \mathcal{O}(\epsilon^3) . \quad (11)$$

Substituting (11) in (7) we obtain, to lowest order, $\mathcal{O}(\epsilon^0)$,

$$Q_0'^2 + S^2(z) = 0 . \quad (12)$$

Thus

$$Q_0 = \pm i \int^z S(\xi) d\xi = \pm i \frac{1}{f_0} \int^z N(\xi) d\xi . \quad (13)$$

At next order, $\mathcal{O}(\epsilon)$, we have

$$2 Q_0' Q_1' + Q_0'' - Q_0' [\log S^2(z)]' = 0 . \quad (14)$$

Hence

$$Q_1 = \frac{1}{2} \log S^2(z) - \frac{1}{2} \log \pm i S(z) + \text{const} . \quad (15)$$

Notice that the imaginary part in the log in (15) just contributes an irrelevant constant. Thus

$$Q_1 = \log \sqrt{S(z)} + \text{const} . \quad (16)$$

In the most common WKB approximation (a.k.a “physical optics”) we truncate (11) at $\mathcal{O}(\epsilon)$. The solution to (7), consistent with the bottom boundary condition (2), is

$$\mathbf{p}_n^{po} = A_n \sqrt{N(z)} \cos \left(\frac{\kappa_n}{f_0} \int_{-h}^z N(\xi) d\xi \right) , \quad (17)$$

where A_n is a constant. By imposing the boundary condition at $z = 0$ (2), we obtain the eigenvalues κ_n :

$$\kappa_n = \frac{n\pi f_0}{\overline{N} h} , \quad n = 0, 1, 2, \dots , \quad (18)$$

where the mean buoyancy frequency is

$$\overline{N} \stackrel{\text{def}}{=} \frac{1}{h} \int_{-h}^0 N(\xi) d\xi . \quad (19)$$

The constant A_n is determined by the normalization condition (4). We have

$$A_n^2 \int_{-h}^0 N(z) \cos \left(\frac{\kappa_n}{f_0} \int_{-h}^z N(\xi) d\xi \right) dz = H , \quad n \geq 1 . \quad (20)$$

The integral in (20) can be evaluated exactly by making the change of variables

$$\eta \stackrel{\text{def}}{=} \frac{\kappa_n}{f_0} \int_{-h}^z N(\xi) d\xi \quad \Rightarrow \quad d\eta = \frac{\kappa_n}{f_0} N(z) dz , \quad (21)$$

and using the expression for the eigenvalues (18). We obtain

$$A_n = \left(2/\overline{N}\right)^{1/2}, \quad n \geq 1. \quad (22)$$

Thus the WKB approximate solution to the standard pressure modes is

$$\mathbf{p}_n^{po} = \left[\frac{2N(z)}{\overline{N}} \right]^{1/2} \cos \left(\frac{n\pi}{\overline{N}h} \int_{-h}^z N(\xi) d\xi \right), \quad n \geq 1. \quad (23)$$

The amplitude of the baroclinic modes at the boundaries is independent of the eigenvalue:

$$\mathbf{p}_n^{po}(z=0) = (-1)^n \left[\frac{2N(0)}{\overline{N}} \right]^{1/2}, \quad (24)$$

and

$$\mathbf{p}_n^{po}(z=-h) = \left[\frac{2N(-h)}{\overline{N}} \right]^{1/2}. \quad (25)$$

The barotropic mode is not recovered from the WKB solution because $\kappa_0 = 0$. From (1) we have that with $\kappa_0 = 0$, the barotropic mode is constant, independent of the stratification. With the normalization (4) we obtain $\mathbf{p}_0 = 1$.

Constant buoyancy frequency

With $N = \text{const.}$ the modes are simple sinusoids. That exact result is recovered as a special case of the WKB solution

$$\mathbf{p}_n^{po} = \sqrt{2} \cos [n\pi(1 + z/h)]. \quad (26)$$

2 Density modes

Similarly the baroclinic modes for density, here denoted by \mathbf{r}_n , are defined via the eigenproblem

$$\mathbf{r}_n'' = -\kappa_n^2 \left(\frac{N}{f_0} \right)^2 \mathbf{r}_n, \quad (27)$$

with homogeneous Dirichlet boundary conditions

$$@z = -h, 0 : \quad \mathbf{r}_n = 0, \quad (28)$$

and normalization

$$\frac{1}{h} \int_{-h}^0 \mathbf{r}_n \mathbf{r}_m dz = \delta_{mn}. \quad (29)$$

Alternatively, we can work on the approximation from the beginning. The WKB approximate solution to (27)-(28), consistent with the bottom boundary conditions (28), is

$$\mathbf{r}_n^{po} = \frac{B_n}{\sqrt{N(z)}} \sin \left(\frac{\kappa_n}{f_0} \int_{-h}^z N(\xi) d\xi \right), \quad (30)$$

The eigenvalues κ_n are the same as before (18). (This should be no surprise because it follow from the definition of \mathbf{p}_n and \mathbf{r}_n . Nonetheless, the verification is a good sanity check.) To find B_n we use the normalization (29)

$$B_n^2 \int_{-h}^0 \frac{1}{N(z)} \sin^2 \left(\frac{\kappa_n}{f_0} \int_{-h}^z N(\xi) d\xi \right) dz = h, \quad n \geq 1. \quad (31)$$

We use a similar trick as above i.e., we change variables with

$$\eta \stackrel{\text{def}}{=} \frac{\kappa_n}{N^2(z)f_0} \int_{-h}^z N(\xi) d\xi \quad \Rightarrow \quad d\eta = \frac{\kappa_n}{N(z)f_0} dz, \quad (32)$$

where, in the WKB spirit, we used the fact that $N(z)$ is slowly varying when differentiating the relation above. We obtain

$$B_n = \left(2/\overline{N}\right)^{1/2}. \quad (33)$$

Thus the WKB approximate solution to the density modes is

$$r_n^{po} = \left(\frac{2}{\overline{N}N(z)}\right) \sin\left(\frac{n\pi}{\overline{N}h} \int_{-h}^z N(\xi) d\xi\right), \quad n \geq 1. \quad (34)$$

Finally, note that the modes are simply related

$$\frac{dr_n^{po}}{dz} = \underbrace{\frac{n\pi}{\overline{N}h}}_{=\kappa_n/f_0} p_n^{po}. \quad (35)$$

Constant buoyancy frequency

Again we recover the $N = \text{const.}$ special case from (34):

$$r_n^{po} = \frac{\sqrt{2}}{\overline{N}} \sin[n\pi(1 + z/h)]. \quad (36)$$