# Towards MDP.jl: The Julia Library of MD Potentials

April 5, 2021

#### Overview

MDP.jl: the Julia library of MD potentials

Descriptor and force calculation

Next steps...

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#### MDP.jl: the Julia library of MD potentials

- Molecular Dynamics (MD) simulations require a potential energy function to describe the force field governing the interaction among atoms.
- ► Force calculation often takes between 50% and 90% of the overall complexity, thus, **determining an adequate potential is a crucial task**.
- ▶ MDP.jl will provide fast and accurate potentials for classical MD simulations on exascale supercomputers.
  - Coupling empirical and machine learning (ML) potentials
  - Quantifying uncertainties in trained ML potentials.
  - Identifying near-optimal configurations to include in the training data.

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- ▶ **Present limitations** when solving complex study cases as ultrahigh temperature ceramics in hypersonic flows.
  - Existing EP, ReaxFF and COMB, do not produce satisfactory results since they require retraining.
  - MLP demand a significant amount of DFT data for the training process.

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$$\mathbf{f}(\mathbf{r}^{N_j}, \mathbf{c}, j, i) = \mathbf{f}_E(\mathbf{r}^{N_j}, \mathbf{c}, j, i) + \mathbf{f}_{ML}(\mathbf{r}^{N_j}, \mathbf{c}, j, i)$$

## Empirical component of the force

Empirical component of the force:

$$\mathbf{f}_{E}(\mathbf{r}^{N_{j}}, \mathbf{c}, j, i) = \mathbf{f}_{E,PS}(\mathbf{r}^{N_{j}}, \mathbf{c}, j, i) + \mathbf{f}_{E,BS}(\mathbf{r}^{N_{j}}, \mathbf{c}, j, i)$$

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The **power spectrum** component (long range interaction):

$$\mathbf{f}_{E,PS}(\mathbf{r}^{N_j}, \mathbf{c}, j, i) = \sum_{t=1}^{N_z} \sum_{k=1}^K \sum_{k'=k}^K \sum_{l=0}^L c_{tkk'l} \frac{\partial d_{tkk'l}^{PS}(\mathbf{r}^{N_j}, j, i)}{\partial \mathbf{r}_i^{N_j}}$$

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The **bispectrum** component (short range interactions):

$$\mathbf{f}_{E,BS}(\mathbf{r}^{N_j}, \mathbf{c}, j, i) = \sum_{t=1}^{N_z} \sum_{k=1}^{K} \sum_{k'=k}^{K} \sum_{l=0}^{L} \sum_{l_1=0}^{L} \sum_{l_2=0}^{L} c_{tkk'll_1 l_2} \frac{\partial d_{tkk'll_1 l_2}^{BS}(\mathbf{r}^{N_j}, j, i)}{\partial \mathbf{r}_i^{N_j}}$$

Derivative of the power spectrum basis functions

$$\begin{split} \frac{\partial d_{tkk'l}^{BS}(\boldsymbol{r}^{N_{j}},j,i)}{\partial \boldsymbol{r}_{i}^{N_{j}}} &= \sum_{s \in \Omega'_{jit}} p_{iskk'l}^{\partial}(\boldsymbol{r}^{N_{j}},j) - \sum_{s \in \Omega'_{jit}} p_{sikk'l}^{\partial}(\boldsymbol{r}^{N_{j}},j), \\ p_{i_{0}i_{1}kk'l}^{\partial}(\boldsymbol{r}^{N_{j}},j) &= \sum_{m=-l}^{l} \left( \frac{\partial u_{klm}(\boldsymbol{r}_{i_{0}}^{N_{j}} - \boldsymbol{r}_{i_{1}}^{N_{j}})}{\partial (\boldsymbol{r}_{i_{0}}^{N_{j}} - \boldsymbol{r}_{i_{1}}^{N_{j}})} \sum_{s \in \Omega_{j,i_{1}}} \left( u_{k'lm}(\boldsymbol{r}_{s}^{N_{j}} - \boldsymbol{r}_{i_{1}}^{N_{j}}) \right) \right) + \\ \sum_{m=-l}^{l} \left( \frac{\partial u_{k'lm}(\boldsymbol{r}_{i_{0}}^{N_{j}} - \boldsymbol{r}_{i_{1}}^{N_{j}})}{\partial (\boldsymbol{r}_{i_{0}}^{N_{j}} - \boldsymbol{r}_{i_{1}}^{N_{j}})} \sum_{s \in \Omega_{j,i_{1}}} \left( u_{klm}(\boldsymbol{r}_{s}^{N_{j}} - \boldsymbol{r}_{i_{1}}^{N_{j}}) \right) \right) \end{split}$$

► The descriptors are based on products of radial basis functions  $g_{lk}(r)$  and spherical harmonics  $Y_{lm}(\theta, \phi)$ .

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- $ightharpoonup g_{lk}(r)$  can vary...
  - Spherical Bessel
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- Providing flexibility when calculating the derivative of  $u_{klm}(\mathbf{r})$  using the finite difference method (FDM)

$$\frac{\partial u_{klm}(\mathbf{r})}{\partial(\mathbf{r})} = \left(\frac{u_{klm}(\mathbf{r} + \Delta x) - u_{klm}(\mathbf{r} - \Delta x)}{2|\Delta x|}, ..., ...\right)$$



Derivative of the bispectrum basis functions

$$\frac{\partial d_{tkk'll_1l_2}^{BS}(\boldsymbol{r}^{N_j},j,i)}{\partial \boldsymbol{r}_i^{N_j}} = \left(\frac{d_{tkk'll_1l_2}(\boldsymbol{r}^{N_j} + \Delta X_i,j,i) - d_{tkk'll_1l_2}(\boldsymbol{r}^{N_j} - \Delta X_i,j,i)}{2|\Delta X_i|},...,...\right)$$

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The bispectrum basis functions are formulated as:

$$d_{tkk'll_1l_2}^{BS}(\boldsymbol{r}^{N_j},j,i) = \sum_{s \in \Omega_{it}^{\prime\prime\prime}} b_{skk'll_1l_2}(\boldsymbol{r}^{N_j},j,i)$$

$$b_{skk'll_1l_2}(\boldsymbol{r}^{N_j},j,i) = \sum_{m=-l}^{l} \sum_{m_1=-l_1}^{l_1} \sum_{m_2=-l_2}^{l_2} \bar{a}_{sklm}(\boldsymbol{r}^{N_j},j) C_{m_1m_2m}^{l_1l_2l} a_{sk'l_1m_1}(\boldsymbol{r}^{N_j},j) a_{sk'l_2m_2}(\boldsymbol{r}^{N_j},j)$$

$$a_{iklm}(\mathbf{r}^{N_j},j) = \sum_{s \in \Omega_{ii}} u_{klm}(\mathbf{r}_s^{N_j} - \mathbf{r}_i^{N_j})$$

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    - ► Fast calculation using "PartialWaveFunctions.jl" ✓
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  - Calculate only in the neighbors
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- ► Integration with NBodysimulator and LAMMPS



## $\mathsf{MDP.jl} + \mathsf{NBodySimulator.jl}$

#### Argon study case.

- 1. Run MDP.jl:
  - Input (from the initial condition):
    - Force field (calculated through LJ, thus, the initial force field). Atomic positions.
  - Output:
    - Force field (this fitted function will be used in the whole MD simulation)
- Run NBodySimulator.jl

For each time step:

- Input (from the initial condition):
  - Force field from MDP.jl (the fitted force field expression is evaluated here). Atomic Velocities. Domain(box). Temperature. Mass. External forces.
- Output:
  - Atom positions. Atomic velocities. Temperature. Energy (kinetic, potential, and total). Radial distribution function

# $\mathsf{DFTK}.\mathsf{jl} + \mathsf{NBodySimulator}.\mathsf{jl}$

#### Argon study case.

- 1. Run NBodySimulator.jl
  - For each time step:
    - Input: Quantum force from DFTK.jl
      - ▶ DFTK.jl input: Atomic positions from NBodySimulator.jl
      - DFTK.jl output: Quantum force (a single set of forces evaluated at one set of positions). Atomic positions. Atomic velocities. Domain (box). Temperature. Mass. External forces
    - Output: Atom positions. Atomic velocities. Temperature. Energy (kinetic, potential, and total). Radial distribution function.

#### References

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Thanks :-)