# Math 269B, 2012 Winter, Homework 3 (Solutions)

Professor Joseph Teran Jeffrey Lee Hellrung, Jr.

February 17, 2012

### 1 Theory

1. (Strikwerda 5.1.2.) Show that the modified leapfrog scheme (5.1.6) is stable for  $\epsilon$  satisfying

$$0 < \epsilon \le 1 \quad \text{if} \quad 0 < a^2 \lambda^2 \le \frac{1}{2}$$

and

$$0 < \epsilon \le 4a^2\lambda^2 \left(1 - a^2\lambda^2\right)$$
 if  $\frac{1}{2} \le a^2\lambda^2 < 1$ .

Note that these limits are not sharp. It is possible to choose  $\epsilon$  larger than these limits and still have the scheme be stable.

#### Solution

Continuing from the text, we find the amplification factors to be

$$g_{\pm}(\theta) = -ia\lambda \sin \theta \pm \sqrt{1 - a^2\lambda^2 \sin^2 \theta - \epsilon \sin^4 \frac{1}{2}\theta}.$$

If the expression under the  $\sqrt{\cdot}$  is nonnegative, then

$$|g_{\pm}(\theta)|^2 = 1 - \epsilon \sin^4 \frac{1}{2}\theta \le 1,$$

hence the scheme is stable. We thus wish to satisfy

$$0 \le 1 - a^2 \lambda^2 \sin^2 \theta - \epsilon \sin^4 \frac{1}{2} \theta =: \alpha(\theta).$$

We compute that

$$\alpha'(\theta) = -\frac{1}{2}\sin\theta \left( \left( 4a^2\lambda^2 - \epsilon \right)\cos\theta + \epsilon \right),\,$$

and hence the extrema of  $\alpha$  occur when  $\sin \theta = 0$  or  $\cos \theta = \epsilon / \left(\epsilon - 4a^2\lambda^2\right)$ . Values of  $\theta$  satisfying  $\sin \theta = 0$  give  $\alpha = 1$  or  $\alpha = 1 - \epsilon$ , requiring that  $\epsilon \leq 1$ . Values of  $\theta$  satisfying  $\cos \theta = \epsilon / \left(\epsilon - 4a^2\lambda^2\right)$  exist if and only if  $\left|\epsilon / \left(\epsilon - 4a^2\lambda^2\right)\right| \leq 1$ , which is equivalent to  $\epsilon \leq 2a^2\lambda^2$ . For such  $\theta$ , we get  $\alpha = 1 - 4a^4\lambda^4 / \left(4a^2\lambda^2 - \epsilon\right)$ , and for this to be nonnegative, we must have  $\epsilon \leq 4a^2\lambda^2 \left(1 - a^2\lambda^2\right)$ . In particular, we must have  $|a\lambda| < 1$ .

So far, we have deduced that, at a minimum,  $0 < \epsilon \le 1$ . Furthermore, if  $\epsilon \le 2a^2\lambda^2$ , then we must additionally satisfy  $\epsilon \le 4a^2\lambda^2\left(1-a^2\lambda^2\right)$ . Now, in the instance that  $2a^2\lambda^2 \le 4a^2\lambda^2\left(1-a^2\lambda^2\right)$ , we would automatically satisfy the second condition, and this latter inequality is equivalent to  $a^2\lambda^2 \le \frac{1}{2}$ . It follows that

– If  $0 < a^2 \lambda^2 \le \frac{1}{2}$ , it is sufficient to take  $0 < \epsilon \le 1$ .

- If 
$$\frac{1}{2} \le a^2 \lambda^2 < 1$$
, it is sufficient to take  $0 < \epsilon \le 4a^2 \lambda^2 (1 - a^2 \lambda^2)$ .

2. Derive the stability condition for the backward-time forward-space scheme

$$\frac{1}{k} \left( v_m^{n+1} - v_m^n \right) + \frac{a}{h} \left( v_{m+1}^{n+1} - v_m^{n+1} \right) = 0$$

used to approximate solutions to  $u_t + au_x = 0$  with, say,  $x \in [0, 1]$  and periodic boundary conditions. Give an example of an initial condition  $v_m^0$  and an explicit expression for  $v_m^n$  that demonstrate unstable behavior for a particular  $\lambda$  (your choice) which fails to satisfy the stability condition. Does the growth in your example agree with your theoretical amplification factor?

#### Solution

Our difference operator is

$$P_{k,h}v_m^n = \frac{1}{k} \left( v_m^{n+1} - v_m^n \right) + \frac{a}{h} \left( v_{m+1}^{n+1} - v_m^{n+1} \right)$$

which has symbol

$$\begin{split} p_{k,h}(s,\xi) &= P_{k,h} \left( e^{skn + imh\xi} \right) / e^{skn + imh\xi} \\ &= \frac{1}{k} \left( e^{sk} - 1 \right) + \frac{a}{h} e^{sk} \left( e^{ih\xi} - 1 \right). \end{split}$$

We determine stability by finding the roots of the symbol as a function of  $g := e^{sk}$ , yielding

$$g = \frac{1}{1 + a\lambda \left(e^{i\theta} - 1\right)}$$

where  $\lambda := k/h$  and  $\theta := h\xi$ . We find that

$$|g|^{-2} = 1 + 2a\lambda (a\lambda - 1) (1 - \cos \theta),$$

hence the scheme is stable ( $|g| \le 1$ ) if and only if  $a \le 0$  or  $a\lambda \ge 1$ .

If, for example,  $a\lambda = \frac{1}{4}$ , then

$$|g|^{-2} = 1 - \frac{3}{8} (1 - \cos \theta) = \frac{5}{8} + \frac{3}{8} \cos \theta.$$

Choosing, for example,  $\theta = \pi$  ought to give an amplication factor of exactly g = 2 of the pure mode  $v_m = e^{i\theta m} = (-1)^m$ . Indeed, one can quickly verify that  $v_m^n = 2^n (-1)^m$  satisfies the difference equation:

$$kP_{k,h}v_m^n = v_m^{n+1} - v_m^n + a\lambda \left(v_{m+1}^{n+1} - v_m^{n+1}\right)$$

$$= 2^{n+1}(-1)^m - 2^n(-1)^m + \frac{1}{4}\left(2^{n+1}(-1)^{m+1} - 2^{n+1}(-1)^m\right)$$

$$= 2^n(-1)^m \left(2 - 1 + \frac{1}{4}(-2 - 2)\right)$$

$$= 0$$

One final remark: Notice that if  $a\lambda = \frac{1}{2}$ , |g| is unbounded near  $\theta = \pi$ . This corresponds to a null space in the resulting system of equations for  $v^{n+1}$  induced by the difference operator, and this null space is spanned precisely by the mode corresponding to  $\theta = \pi$ ,  $v_m = (-1)^m$ .

3. Prove that numerical solutions to the Lax-Friedrichs scheme

$$\frac{1}{k}\left(v_m^{n+1} - \frac{1}{2}\left(v_{m+1}^n + v_{m-1}^n\right)\right) + \frac{a}{2h}\left(v_{m+1}^n - v_{m-1}^n\right) = 0$$

converge to solutions to the corresponding modified equation

$$u_t + au_x = \frac{h^2}{2k} \left( 1 - \left( \frac{ak}{h} \right)^2 \right) u_{xx}$$

to second order accuracy in  $L^{\infty}$ . I.e., show that  $|v_m^n - u_{k,h}(t_n, x_m)| \to 0$  as  $h, k \to 0$  (according to the stability criterion), where the subscripts on  $u_{k,h}$  only indicate that the solution to the modified equation is parameterized by k, h.

#### Solution

The difference operator for the Lax-Friedrichs scheme is

$$P_{k,h}v_m^n = \frac{1}{k} \left( v_m^{n+1} - \frac{1}{2} \left( v_{m+1}^n + v_{m-1}^n \right) \right) + \frac{a}{2h} \left( v_{m+1}^n - v_{m-1}^n \right)$$

which has symbol

$$\begin{split} p_{k,h}(s,\xi) &= P_{k,h} \left( e^{skn + imh\xi} \right) / e^{skn + imh\xi} \\ &= \frac{1}{k} \left( e^{sk} - \cos h\xi \right) + i \frac{a}{h} \sin h\xi \\ &= \frac{1}{k} \left( 1 + sk + \frac{1}{2} s^2 k^2 - 1 + \frac{1}{2} h^2 \xi^2 \right) + i \frac{a}{h} \left( h\xi \right) + O\left( k^2 + h^2 + h^4 k^{-1} \right) \\ &= s + ia\xi + \frac{k}{2} s^2 + \frac{h^2}{2k} \xi^2 + O\left( k^2 + h^2 + h^4 k^{-1} \right). \end{split}$$

4. (Strikwerda 4.1.2.) Show that the (2,2) leapfrog scheme for  $u_t + au_{xxx} = f$  (see (2.2.15)) given by

$$\frac{v_m^{n+1} - v_m^{n-1}}{2k} + a\delta^2 \delta_0 v_m^n = f_m^n$$

with  $\nu = k/h^3$  constant, is stable if and only if

$$|a\nu| < \frac{2}{3^{3/2}}.$$

#### Solution

5. (Strikwerda 3.2.1.) Show that the (forward-backward) MacCormack scheme

$$\begin{split} &\tilde{v}_m^{n+1} = v_m^n - a\lambda \left( v_{m+1}^n - v_m^n \right) + kf_m^n, \\ &v_m^{n+1} = \frac{1}{2} \left( v_m^n + \tilde{v}_m^{n+1} - a\lambda \left( \tilde{v}_m^{n+1} - \tilde{v}_{m-1}^{n+1} \right) + kf_m^{n+1} \right) \end{split}$$

is a second-order accurate scheme for the one-way wave equation (1.1.1). Show that for f=0 it is identical to the Lax-Wendroff scheme (3.1.1).

#### Solution

## 2 Programming

1. For the one-way wave equation  $u_t + au_x = 0$ , investigate how close the numerical solution to a finite difference scheme is to the solution to the corresponding modified equation. To be concrete, suppose a convenient initial condition for which you can solve the modified equation explicitly with periodic boundary conditions. Take a = 1, k/h = 0.5, and final time T = 0.5. Compare the following finite difference schemes: upwinding, Lax-Friedrichs, and Lax-Wendroff. Also, include a derivation of the respective corresponding modified equations.