1. Consider the differential equation:

$$u_{xx} + u_{yy} + \lambda u = 0$$

in the strip $\{(x,y): 0 < y < \pi, -\infty < x < +\infty\}$ with boundary conditions

$$u(x,0) = 0, \quad u(x,\pi) = 0.$$

Find all bounded solutions of the boundary value problem when (a) $\lambda = 0$, (b) $\lambda > 0$, and (c) $\lambda < 0$.

By separation of variables, the solution has the form

$$u(x,y) = \sum_{n=1}^{\infty} X_n(x)\sin(ny).$$

For all n,

$$0 = X_n''(x)\sin(ny) - n^2X_n(x)\sin(ny) + \lambda X_n(x)\sin(ny)$$
$$0 = X_n''(x) + (\lambda - n^2)X_n(x)$$
$$X_n(x) = a_n e^{\sqrt{n^2 - \lambda}x} + b_n e^{-\sqrt{n^2 - \lambda}x}.$$

(a) If $\lambda = 0$, then

$$X_n(x) = a_n e^{nx} + b_n e^{-nx}.$$

The only bounded solution is $a_n = b_n = 0$, so $u \equiv 0$.

(b) If $\lambda > 0$, then for $n^2 > \lambda$,

$$X_n(x) = a_n e^{i\sqrt{\lambda - n^2}x} + b_n e^{-i\sqrt{\lambda - n^2}x}$$
$$= a'_n \sin\sqrt{\lambda - n^2}x + b'_n \cos\sqrt{\lambda - n^2}x.$$

Thus bounded solutions have the form

$$u(x,y) = \sum_{n^2 < \lambda} \left(a'_n \sin \sqrt{\lambda - n^2} \, x + b'_n \cos \sqrt{\lambda - n^2} \, x \right) \sin ny.$$

(c) If $\lambda < 0$, then similar to part (a),

$$X_n(x) = a_n e^{\sqrt{n^2 - \lambda} x} + b_n e^{-\sqrt{n^2 - \lambda} x}.$$

Thus the only bounded solution is $a_n = b_n = 0$, $u \equiv 0$.

2. Let $C^2(\overline{\Omega})$ be the space of all twice continuously differentiable functions in the bounded smooth closed domain $\overline{\Omega} \subset \mathbb{R}^2$. Let $u_0(x,y)$ be the function that minimizes the functional

$$D(u) = \int_{\Omega} (u_x^2 + u_y^2 + fu) + \int_{\partial \Omega} au^2$$

where f and a are given continuous functions. Find the differential equation and boundary condition that u_0 satisfies.

In other words, we must find the Euler-Lagrange equations associated with the minimization problem. Let $u=u_0, v\in C^2(\overline{\Omega})$, and define $g(\epsilon)=D(u+\epsilon v)$. Then

$$0 = g'(0) = \int_{\Omega} (2u_x v_x + 2u_y v_y + fv) + \int_{\partial\Omega} a 2u v$$

$$= 2 \int_{\Omega} \nabla u \cdot \nabla v + \int_{\Omega} fv + 2 \int_{\partial\Omega} a u v$$

$$= 2 \left[-\int_{\Omega} \Delta u v + \int_{\partial\Omega} (n \cdot \nabla u)v \right] + \int_{\Omega} fv + 2 \int_{\partial\Omega} a u v$$

$$= \int_{\Omega} (-2\Delta u + f)v + \int_{\partial\Omega} 2(n \cdot \nabla u + au)v.$$

Since this holds for any $v \in C^2(\overline{\Omega})$, it implies

$$\begin{cases} \Delta u = \frac{1}{2}f & \text{in } \Omega \\ n \cdot \nabla u + au = 0 & \text{on } \partial \Omega \end{cases}$$

3. Let $f(x_1, x_2)$ be a continuous function with compact support. Define

$$u(x_1, x_2) = \frac{1}{2\pi} \iint_{\mathbb{R}^2} \frac{f(y_1, y_2)}{z - w} \, dy_1 \, dy_2$$

where $z = x_1 + ix_2$ and $w = y_1 + iy_2$. Prove that

$$\frac{\partial u}{\partial x_1} + i \frac{\partial u}{\partial x_2} = f(x_1, x_2) \qquad \text{in } \mathbb{R}^2.$$

Consider $v_{x_1} + iv_{x_2} = f$. Then

$$f_{x_1} = v_{x_1x_1} + iv_{x_1x_2}$$

$$f_{x_2} = v_{x_1x_2} + iv_{x_2x_2}.$$

Therefore, $\Delta v=f_{x_1}-if_{x_2}$. Let $g=\frac{1}{2\pi}\log\sqrt{x_1^2+x_2^2}$, then since g is a fundamental solution of $\Delta v=f$,

$$v = f_{x_1} * g - i f_{x_2} * g$$

$$= f * (g_{x_1} - i g_{x_2})$$

$$= \frac{1}{2\pi} f * \left(\frac{x_1}{x_1^2 + x_2^2} - \frac{i x_2}{x_1^2 + x_2^2}\right)$$

$$= \frac{1}{2\pi} f * \frac{1}{x_1 + i x_2}$$

$$v(x_1, x_2) = \frac{1}{2\pi} \iint_{\mathbb{R}^2} \frac{f(y_1, y_2)}{z - w} \, \mathrm{d}y_1 \, \mathrm{d}y_2.$$

So u=v and hence $u_{x_1}+iu_{x_2}=f$.

4. Consider the boundary value problem on $[0, \pi]$:

$$\begin{cases} y''(x) + p(x)y(x) = f(x), & 0 < x < \pi \\ y(0) = 0, & y'(\pi) = 0. \end{cases}$$

Find the smallest λ_0 such that the boundary value problem has a unique solution whenever $p(x) > \lambda_0$ for all x. Justify your answer.

Let y_1 and y_2 be two solutions and let $w = y_1 - y_2$, then

$$\begin{cases} w''(x) + p(x)w(x) = 0, & 0 < x < \pi \\ w(0) = 0, & w'(\pi) = 0. \end{cases}$$

So

$$0 = \int_0^{\pi} (-w'' + pw)w = -w'w\Big|_0^{\pi} + \int_0^{\pi} (w')^2 + pw^2 = \int_0^{\pi} (w')^2 + pw^2.$$

Therefore,

$$\int_0^{\pi} (w')^2 = -\int_0^{\pi} pw^2.$$

The left side is nonnegative. If p>0, then the right side is nonpositive and $w\equiv 0$. Therefore, $\lambda_0=0$.

5. Consider the Laplace equation

$$u_{xx} + u_{yy} = 0, \quad y > 0, \quad -\infty < x < +\infty$$

with the boundary condition

$$u_{y}(x,0) - u(x,0) = f(x),$$

where $f(x) \in C_0^{\infty}(\mathbb{R})$. Find a bounded solution u(x,y) and show that $u(x,y) \to 0$ when $|x| + y \to \infty$.

Apply the Fourier transform in the x variable to obtain

$$\begin{cases} -\xi^2 \hat{u}(\xi, y) - \hat{u}(\xi, y) = 0 \\ \hat{u}_y(\xi, 0) - \hat{u}(\xi, 0) = \hat{f}(\xi) \end{cases}$$

where $(\hat{\cdot})$ denotes the Fourier transform. For each fixed frequency, $-\xi^2\hat{u} + \hat{u}'' = 0$ is an ODE in y with solutions $e^{\xi y}$ and $e^{-\xi y}$. For bounded \hat{u} , we keep only the solution such that the exponent is negative,

$$\hat{u}(\xi, y) = c(\xi) e^{-|\xi|y}.$$

The initial conditions determine $c(\xi)$:

$$\hat{u}_{y}(\xi,0) - \hat{u}(\xi,0) = \hat{f}(\xi) -|\xi|c(\xi) - c(\xi) = \hat{f}(\xi) c(\xi) = -\frac{\hat{f}(\xi)}{1 + |\xi|},$$

and we have

$$\hat{u}(\xi, y) = \underbrace{\frac{-e^{-|\xi|y}}{1 + |\xi|}}_{\hat{H}(\xi, y)} \hat{f}(\xi).$$

Let $\mathcal S$ denote the Schwartz space in one dimension. Since the Fourier transform maps $\mathcal S$ into itself, $f\in C_0^\infty(\mathbb R)\subset \mathcal S$ implies $\hat f\in \mathcal S$. Considering y as a fixed parameter, $\hat H$ is in $\mathcal S$ and hence $\hat u$ and u are in $\mathcal S$. Therefore, u is bounded and decays as $|x|+y\to\infty$.

6. Consider the first order system $u_t - u_x = v_t + v_x = 0$ in the diamond shaped region -1 < x + t < 1, -1 < x - t < 1. For each of the following boundary value problems state whether this problem is well-posed. If it is well-posed, find the solution.

(a)
$$\begin{cases} u(x+t) &= u_0(x+t) & \text{on } x-t=-1 \\ v(x-t) &= v_0(x+t) & \text{on } x+t=-1 \end{cases}$$
(b)
$$\begin{cases} v(x+t) &= v_0(x+t) & \text{on } x-t=-1 \\ u(x-t) &= u_0(x+t) & \text{on } x+t=-1 \end{cases}$$

(b)
$$\begin{cases} v(x+t) = v_0(x+t) & \text{on } x-t = -1 \\ u(x-t) = u_0(x+t) & \text{on } x+t = -1 \end{cases}$$

By method of characteristics,

$$u: \left\{ \begin{array}{ll} t' & = & 1 \\ x' & = & -1 \\ u' & = & 0 \end{array} \right. \quad \text{characteristics } x+t=c \qquad \qquad \left\{ \begin{array}{ll} t' & = & 1 \\ x' & = & 1 \\ v' & = & 0 \end{array} \right. \quad \text{characteristics } x-t=c$$

(a) The boundary conditions are along noncharacteristic curves, x-t=-1 for u and x+t=-1 for v, so the problem is well-posed. The solution is

$$\begin{cases} u(x,t) &= u_0(x+t) \\ v(x,t) &= v_0(x-t) \end{cases}$$
 in D .

(b) In this case, the boundary conditions are along characteristic curves, so it is not well-posed.

- 7. For the two-point boundary value problem $Lf = f_{xx} f$ on $-\infty < x < \infty$ with $\lim_{x \to \infty} f(x) = 0$ and $\lim_{x \to -\infty} f(x) = 0$, the Green's function G(x, x') solves $LG = \delta(x x')$ in which L acts on the variable x.
 - (a) Show that G(x, x') = G(x x').
 - (b) For each x', show that

$$G(x, x') = \begin{cases} a_- e^x & \text{for } x < x', \\ a_+ e^{-x} & \text{for } x' < x, \end{cases}$$

in which a_{\pm} are functions that depend only on x'.

- (c) Using (a), find the x' dependence of a_{\pm} .
- (d) Finish finding G(x, x') by using the jump conditions to find the remaining unknowns in a_{\pm} .
- (a) Define $\mathscr{G}f(x) = \int_{\mathbb{D}} G(x, x') f(x') dx'$, then

$$f(x) = \int_{\mathbb{R}} LG(x, x') f(x') dx' = L\mathscr{G}f(x)$$

$$= \int_{\mathbb{R}} G_{xx}(x, x') f(x') - G(x, x') f(x') dx' = \int_{\mathbb{R}} G(x, x') f''(x') - G(x, x') f(x') dx' = \mathscr{G}Lf(x).$$

Suppose f_0, f_τ are such that $f_\tau(x) = f_0(x - \tau)$ and let $g_0 = Lf_0$, $g_\tau = Lf_\tau$, then $g_\tau(x) = g_0(x - \tau)$ and $\mathscr{G}g_\tau(x) = f_\tau(x) = f_0(x - \tau) = \mathscr{G}g_0(x - \tau)$. So \mathscr{G} is shift-invariant, which implies G(x, x') = G(x - x').

(b) For
$$x < x'$$
, $LG(x, x') = G_{xx}(x, x') - G(x, x') = 0$, so

$$G(x, x') = a_{-}(x')e^{x} + b_{-}(x')e^{-x}$$
.

The boundary condition $\lim_{x\to -\infty} G(x,x')=0$ requires that $b_-(x')=0$. Similarly for x>x',

$$G(x, x') = b_{+}(x')e^{x} + a_{+}(x')e^{-x},$$

where $\lim_{x\to\infty} G(x,x') = 0$ implies $b_+(x') = 0$.

(c) If x < x', then the shift-invariance of G implies G(x + t, x' + t) = G(x, x') and

$$a_{-}(x'+t)e^{x+t} = a_{-}(x')e^{x}$$

 $a_{-}(x'+t) = a_{-}(x')e^{-t}$
 $a_{-}(t) = c_{-}e^{-t}, c_{-} = a_{-}(0).$

Similarly, $a_+(t) = c_+ e^t$.

(d) Notice that

$$f(x) = \mathcal{G}Lf(x) = c_{+} \int_{-\infty}^{x} e^{x'-x} f''(x') dx' + c_{-} \int_{x}^{\infty} e^{x-x'} f''(x') dx' - \mathcal{G}f(x).$$

Using integration by parts,

$$\int_{-\infty}^{x} e^{x'-x} f''(x') dx' = \left[e^{x'-x} f'(x') \right]_{-\infty}^{x} - \left[e^{x'-x} f(x') \right]_{-\infty}^{x} + \int_{-\infty}^{x} e^{x'-x} f(x') dx'$$

$$\int_{-\infty}^{\infty} e^{x-x'} f''(x') dx' = \left[e^{x-x'} f'(x') \right]_{x}^{\infty} - \left[-e^{x-x'} f(x') \right]_{x}^{\infty} + \int_{-\infty}^{\infty} e^{x-x'} f(x') dx'$$

So the jump condition is $f(x)=(c_--c_+)f'(x)+(-c_--c_+)f(x)$, which yields $c_-=c_+=-\frac{1}{2}$. Therefore, the Green's function is $G(x,x')=-\frac{1}{2}e^{-|x-x'|}$.

8. For the ODE

$$\begin{cases} u_t = u - v^2 \\ v_t = v - u^2 \end{cases}$$

do all of the following:

- (a) Find all stationary points.
- (b) Analyze their type.
- (c) Show that u = v is an invariant set for this ODE; i.e., if u(0) = v(0), then u(t) = v(t) for all t.
- (d) Draw the phase plane for this system.
- (a) The stationary points satisfy $u = v^2$ and $v = u^2$, they are (0,0) and (1,1).
- (b) Linearize the system about a stationary point:

$$\left(\begin{array}{c} u \\ v \end{array}\right)_t = \left(\begin{array}{cc} 1 & -2v \\ -2u & 1 \end{array}\right) \left(\begin{array}{c} u \\ v \end{array}\right).$$

At (0,0), $J=\left(\begin{smallmatrix} 1&0\\0&1\end{smallmatrix}\right)$ with eigenvalues 1,1. So (0,0) is an unstable proper node.

At (1,1), $J=\begin{pmatrix} 1 & -2 \\ -2 & 1 \end{pmatrix}$ with eigenvalues 3,-1 and eigenvectors $\begin{pmatrix} 1 \\ -1 \end{pmatrix}$, $\begin{pmatrix} 1 \\ 1 \end{pmatrix}$. So (1,1) is a saddle node.

(c) Notice that $|u_t - v_t| = |u - v^2 - v + u^2| = |u - v||1 + u + v|$. If u(0) = v(0), we have by Gronwall's inequality

$$|u(t) - v(t)| = \int_0^t |u(s) - v(s)| |1 + u(s) + v(s)| ds \le 0, \quad t \ge 0.$$

(d)

