

Math 269B, 2012 Winter, Homework 1 (Solutions)

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1 Theory

1. (Strikwerda 1.1.3.) Solve the initial value problem for

$$u_t + \frac{1}{1 + \frac{1}{2} \cos x} u_x = 0$$

Show that the solution is given by $u(t, x) = u_0(\xi)$, where ξ is the unique solution of

$$\xi + \frac{1}{2} \sin \xi = x + \frac{1}{2} \sin x - t.$$

Solution

We use the method of characteristics, in which we use the change of variables $\tau = \tau(t, x)$, $\xi = \xi(t, x)$, and $\tilde{u}(\tau, \xi) = u(t(\tau, \xi), x(\tau, \xi))$. Then

$$\tilde{u}_\tau = u_t t_\tau + u_x x_\tau = u_t + a(t, x) u_x \equiv 0$$

if

$$t_\tau \equiv 1, \quad x_\tau = a(t, x),$$

where

$$a(t, x) := \frac{1}{1 + \frac{1}{2} \cos x}.$$

We may solve the equation $t_\tau \equiv 1$ as $t = \tau$, while the equation for x_τ is an ordinary differential equation in $x = x(\tau)$:

$$\frac{dx}{d\tau} = \frac{1}{1 + \frac{1}{2} \cos x}, \quad x(\tau = 0) = \xi.$$

Via separation of variables, we obtain an implicit relation among x , $\tau = t$, and ξ :

$$x + \frac{1}{2} \sin x - \xi - \frac{1}{2} \sin \xi = \tau = t.$$

For each fixed ξ , we have $\tilde{u}_\tau = 0$, hence $\tilde{u}(\tau, \xi) = \tilde{u}(0, \xi) = u(0, \xi) = u_0(\xi)$. It follows that $u(t, x) = u_0(\xi)$, where $\xi = \xi(t, x)$ satisfies

$$\xi + \frac{1}{2} \sin \xi = x + \frac{1}{2} \sin x - t.$$

Note that the mapping $\xi \mapsto \xi + \frac{1}{2} \sin \xi$ is an automorphism on \mathbb{R} (its derivative is uniformly bounded away from 0), hence such a ξ always exists and is unique.

2. Solve the initial value problem

$$u_t + (\sin t) u_x = \frac{1}{1+t^2}, \quad u(0, x) = u_0(x), \quad x \in \mathbb{R}, \quad t > 0.$$

We use the method of characteristics, in which we use the change of variables $\tau = \tau(t, x)$, $\xi = \xi(t, x)$, and $\tilde{u}(\tau, \xi) = u(t(\tau, \xi), x(\tau, \xi))$. Then

$$\tilde{u}_\tau = u_t t_\tau + u_x x_\tau = u_t + a(t, x) u_x = f(t, x)$$

if

$$t_\tau \equiv 1, \quad x_\tau = a(t, x),$$

where

$$a(t, x) := \sin t, \quad f(t, x) := \frac{1}{1+t^2}.$$

We may solve the equation $t_\tau \equiv 1$ as $t = \tau$, while the equation for x_τ is an ordinary differential equation in $x = x(\tau)$:

$$\frac{dx}{d\tau} = \sin t = \sin \tau, \quad x(\tau = 0) = \xi.$$

This easily solves to

$$x(\tau, \xi) = \xi + \int_0^\tau \sin \tau' d\tau' = 1 - \cos \tau + \xi.$$

For each fixed ξ , we have

$$\tilde{u}_\tau = \frac{1}{1+t^2} = \frac{1}{1+\tau^2}, \quad \tilde{u}(0, \xi) = u(0, \xi) = u_0(\xi),$$

which easily solves to

$$\tilde{u}(\tau, \xi) = \tilde{u}(0, \xi) + \int_0^\tau \frac{1}{1+(\tau')^2} d\tau' = \arctan \tau + u_0(\xi).$$

It follows that

$$u(t, x) = \arctan t + u_0(\cos t + x + 1).$$

3. Consider the first order system of PDEs of the form

$$\vec{u}_t + A \vec{u}_x = 0, \quad \vec{u}(0, x) = \vec{u}_0(x), \quad x \in [0, 1], \quad t > 0.$$

(a) Give the solution to the initial value problem when

$$A = \begin{pmatrix} 2 & 1 \\ 1 & 2 \end{pmatrix}.$$

(b) Describe appropriate boundary conditions at $x = 0$ and/or $x = 1$, if possible, which make the initial boundary value problem in (a) well-posed. Try to be as general as possible. How should such boundary conditions be presented to put the solution in a simple form?

(c) Give the solution to the initial value problem when

$$A = \begin{pmatrix} 2 & 3 \\ 3 & 2 \end{pmatrix}.$$

(d) Describe appropriate boundary conditions at $x = 0$ and/or $x = 1$, if possible, which make the initial boundary value problem in (c) well-posed. Try to be as general as possible. How should such boundary conditions be presented to put the solution in a simple form?

Solution

- (a) The eigenvalues and corresponding eigenvectors of A are

$$\lambda_{\pm} = 2 \pm 1, \quad \vec{u}_{\pm} = \begin{pmatrix} 1 \\ \pm 1 \end{pmatrix},$$

leading to the general solution

$$\vec{u}(t, x) = \begin{pmatrix} \vec{u}_+ & \vec{u}_- \end{pmatrix} \begin{pmatrix} c_+(x - \lambda_+ t) \\ c_-(x - \lambda_- t) \end{pmatrix},$$

where the scalar functions c_{\pm} are determined by the initial conditions and boundary conditions.

- (b) Since both $\lambda_{\pm} > 0$, boundary conditions must only be specified at the $x = 0$ boundary. A simple way to present these boundary conditions is to simply specify c_{\pm} when $x = 0$, i.e., $c_{\pm}(-\lambda_{\pm}t) = g_{\pm}^0(t)$ for $t > 0$.
- (c) A here has the same eigenvectors as in (a) but now with eigenvalues $\lambda_{\pm} = 2 \pm 3$. Otherwise, the solution takes the same general form.
- (d) This time, $\lambda_+ > 0$ but $\lambda_- < 0$, so boundary conditions must be specified at both the $x = 0$ boundary and the $x = 1$ boundary. A simple way to present these boundary conditions is to specify c_+ when $x = 0$ in terms of t and c_- , and vice versa, to specify c_- when $x = 1$ in terms of t and c_+ :

$$\begin{aligned} c_+(-\lambda_+ t) &= g_+^0(t, c_-(-\lambda_- t)), \\ c_-(1 - \lambda_- t) &= g_-^1(t, c_+(1 - \lambda_+ t)), \end{aligned}$$

for $t > 0$.

4. Derive the leading term of the local truncation error for the following finite difference schemes used to approximate solutions to the equation $u_t + au_x = 0$.

- (a)

$$\frac{1}{k} (v_m^{n+1} - v_m^n) + a \frac{1}{2h} (v_{m+1}^n - v_{m-1}^n) = 0.$$

- (b)

$$\frac{1}{k} \left(v_m^{n+1} - \frac{1}{2} (v_{m+1}^n + v_{m-1}^n) \right) + a \frac{1}{2h} (v_{m+1}^n - v_{m-1}^n) = 0.$$

Solution

Unless otherwise noted, all functions are evaluated at a common point (t, x) .

- (a) We use the following Taylor expansions about (t, x) :

$$\begin{aligned} \phi(t + k, x) &= \phi + k\phi_t + \frac{1}{2}k^2\phi_{tt} + O(k^3); \\ \phi(t, x \pm h) &= \phi \pm h\phi_x + \frac{1}{2}h^2\phi_{xx} \pm \frac{1}{6}h^3\phi_{xxx} + \frac{1}{24}h^4\phi_{xxxx} + O(h^5). \end{aligned}$$

Thus,

$$\begin{aligned} P_{k,h}\phi &= \frac{1}{k} (\phi(t + k, x) - \phi) + a \frac{1}{2h} (\phi(t, x + h) - \phi(t, x - h)) \\ &= \phi_t + a\phi_x + \frac{1}{2}k\phi_{tt} + \frac{1}{6}h^2\phi_{xxx} + O(k^2 + h^4), \end{aligned}$$

which agrees with $P\phi = (\partial_t + a\partial_x)\phi$ up to a truncation error of

$$P_{k,h}\phi - P\phi = \frac{1}{2}k\phi_{tt} + \frac{1}{6}h^2\phi_{xxx} + O(k^2 + h^4).$$

(b) Using the same Taylor expansions as above,

$$\begin{aligned} P_{k,h}\phi &= \frac{1}{k} \left(\phi(t+k, x) - \frac{1}{2} (\phi(t, x+h) + \phi(t, x-h)) \right) + a \frac{1}{2h} (\phi(t, x+h) - \phi(t, x-h)) \\ &= \phi_t + a\phi_x + \frac{1}{2}k\phi_{tt} - \frac{1}{2}\frac{h^2}{k}\phi_{xx} + O(k^2 + h^2), \end{aligned}$$

which agrees with $P\phi$ up to a truncation error of

$$P_{k,h}\phi - P\phi = \frac{1}{2}k\phi_{tt} - \frac{1}{2}\frac{h^2}{k}\phi_{xx} + O(k^2 + h^2).$$

5. Determine the stability region Λ for each of the finite difference schemes in Problem 4.

Solution

(a) The amplification factor g satisfies the equation

$$\frac{1}{k}(g-1) + a\frac{1}{2h}(e^{i\theta} - e^{-i\theta}) = 0,$$

so

$$g = 1 - i\frac{ak}{h}\sin\theta,$$

so $|g|^2 = 1 + \frac{a^2k^2}{h^2}\sin^2\theta \leq 1 + O(k)$ if and only if $k = O(h^2)$.

(b) The amplification factor g satisfies the equation

$$\frac{1}{k}\left(g - \frac{1}{2}(e^{i\theta} + e^{-i\theta})\right) + a\frac{1}{2h}(e^{i\theta} - e^{-i\theta}) = 0,$$

so

$$g = \cos\theta - i\frac{ak}{h}\sin\theta,$$

so $|g| \leq 1$ if and only if $|ak/h| \leq 1$.

2 Programming

1. Implement the finite difference schemes in Problem 4. in the Theory section for $x \in [0, 1]$, $t \in [0, T]$ for some final time T , $u(x, 0) = u_0(x)$, and *periodic* boundary conditions.
2. Investigate the convergence of each scheme for $a = 1$ and $T = 1$. Set $k/h =: \lambda$ to be constant, and demonstrate which values of λ cause the scheme to converge and which to diverge. If no such λ gives convergence, find an alternate relation between k and h which does ensure convergence (if possible). Try using both a smooth initial condition (e.g., $u_0(x) = \sin(2\pi x)$); a non-smooth initial condition (e.g., $u_0(x) = 1 - 2|x - 1/2|$); and a discontinuous initial condition (e.g., $u_0(x) = 0$ if $|x - 1/2| > 1/4$ and $u_0(x) = 1$ if $|x - 1/2| < 1/4$). Use the discrete L^2 norm to measure the error between your numerical solution and the true solution:

$$\|w\|_h = \left(h \sum_m |w_m|^2 \right)^{1/2}.$$

[Note: Due to periodicity, be sure to avoid double-counting the contributions at $x = 0$ and $x = 1$]
Plot the numerical solutions from each scheme at $t = T$ when $h = 1/100$. Summarize your results. Which scheme do you think is better and why?

Solution

Using the included code, the results of the following statements investigate the convergence of the Lax-Friedrichs scheme with the various initial conditions and $\lambda = 0.99$.

```

test_convergence( ...
    1, 1, @(x) sin(2*pi*x), ...
    @lax_friedrichs, "Lax-Friedrichs", ...
    2.^(-(9:0.5:12)), @(h) 0.99*h);
test_convergence( ...
    1, 1, @(x) 1 - 2*abs(x - 1/2), ...
    @lax_friedrichs, "Lax-Friedrichs", ...
    2.^(-(9:0.5:12)), @(h) 0.99*h);
test_convergence( ...
    1, 1, @(x) (1 + sign(1/4 - abs(x - 1/2)))/2, ...
    @lax_friedrichs, "Lax-Friedrichs", ...
    2.^(-(9:0.5:12)), @(h) 0.99*h);

```

The numerical convergence rates, based on fitting a linear regression through a log-log plot of h vs. L^2 -error, are found to be, respectively, 1.06, 0.80, and 0.19. For the smooth initial condition $u_0(x) = \sin(2\pi x)$, the numerical convergence rate, 1.06, agrees well with the theoretical convergence rate of 1. The fact that the numerical convergence rates for the non-smooth initial conditions are significantly less than 1 is not surprising given that our theoretical convergence analysis depended on the existence of some number of derivatives.

Additionally, as predicted by theory, numerical evidence indicates that the Lax-Friedrichs scheme diverges if $\lambda > 1$ (and, still, $a = 1$), while the forward-time central-space scheme diverges for any λ , which is consistent with our theoretical analysis.

We do get the forward-time central-space scheme to converge if, e.g., we let $k = h^2$, as demonstrated by the results of the following statement.

```

test_convergence( ...
    1, 1, @(x) sin(2*pi*x), ...
    @ftcs, "forward-time central-space", ...
    2.^(-(4:0.5:7)), @(h) h*h);

```

The numerical convergence rate is found to be 2.01, which agrees with the theoretical rate of convergence of 2.