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**1 Local wind regime induced by the presence of giant  
2 linear dunes.**

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**7**  
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**9 Abstract**

**10** In large-scale studies of arid areas, sediment fluxes and resulting geomorpho-  
**11** logical quantities (sediment pathways, erosion rates, dune morphodynamics)  
**12** are often derived from wind data. However, the feedback of nearby topogra-  
**13** phy on the flow, such as kilometric giant dunes, is rarely taken into account  
**14** because barely studied. Here, we compare locally measured wind data to the  
**15** predictions of the Era5Land climate reanalysis in four different places across  
**16** and around the Namib sand sea. In places with flat environments, we show  
**17** that the two wind datasets agree with each other. In contrast, within the  
**18** sand sea, we show that they significantly differ due to the presence of giant  
**19** linear dunes. We quantify this discrepancy, and link it to dynamics of the  
**20** atmospheric boundary layer. Aside setting limits to the applicability of the  
**21** Era5land dataset, this study highlights the giant dunes feedback on the wind,  
**22** and discuss its implications for smaller scale bedforms. All codes used, from  
**23** the raw data to the figures of the paper, are publicly available with full doc-  
**24** documentation at <https://github.com/Cgadal/GiantDunes> (will be made public  
**25** upon acceptance of this manuscript for publication).

**26** **Keywords** Boundary layer · Turbulent flow · Sand dunes · Fluide-structures  
**27** interactions

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**28 1 Introduction**

29 As a flow encounters an obstacle, different interactions can arise depending on  
30 the different time and length scales involved. In the case of atmospheric flows,  
31 this mainly depends on the part of the atmosphere, schematically composed  
32 of a turbulent boundary layer topped by a stably stratified one, with which  
33 the obstacle interacts (Stull 1988). At the largest scale, the feedback of moun-  
34 tains on the stratified free atmosphere results in wave generation as well as  
35 significant wind disturbances, such as downslope winds in the lee side (Dur-  
36 ran 1990). Inside the boundary layer, the interaction between a turbulent flow  
37 and hilly surfaces is for example key to the understanding ocean surface wind-  
38 driven waves (Sullivan and McWilliams 2010), or eolian bedforms in desert  
39 (Courrech du Pont 2015).

40 Looking at the flow close to the surface, two different components of the  
41 topography feedback can be isolated. First, the flow accelerates on the upwind  
42 slope, and slows down on the downwind one, with a maximum velocity slightly  
43 upwind of the crest (Jackson and Hunt 1975; Sykes 1980; Hunt et al. 1988).  
44 Several measurements are available, performed in subaqueous flumes (Zilker  
45 et al. 1977; Zilker and Hanratty 1979; Frederick and Hanratty 1988) and in  
46 eolian field conditions (Claudin et al. 2013; Lü et al. 2021). Second, deflection  
47 occurs when the incident flow direction is not perpendicular to the ridge crest.  
48 While predicted to be small (less than 10°) in the linear regime (Gadal et al.  
49 2019), significant flow steering has been reported in the field on the downwind  
50 side of larger aspect ratios obstacles, such as mountain ranges (Kim et al. 2000;  
51 Lewis et al. 2008; Fernando et al. 2019) or sand dunes (Walker et al. 2009;  
52 Walker and Shugar 2013; Hesp et al. 2015; Smith et al. 2017).

53 In the case of aeolian bedforms, the measurements mentioned above were  
54 made on dunes a few meters high only. The resulting wind disturbances have  
55 then been shown to significantly impact the sediment pathways of the whole  
56 system (Hesp et al. 2015), but also to affect the collective behaviour of dune  
57 populations through long-range interactions (Bacik et al. 2020). However, eo-  
58 lian dunes coexist in sand seas at different scales, each representative of their  
59 different evolution stages, from emergence at twenty meters wavelengths to gi-  
60 ant size, corresponding to kilometric wavelengths and heights of several dozen  
61 meters (McKee 1979). Quantifying these two effects for giant dunes is then  
62 all the more important for several reasons. First, larger wind disturbances  
63 are expected from larger obstacles. Second, the flow–structure coupling is  
64 not straightforward due to vertical structure of the atmosphere with which  
65 the giant dunes interact due to their sizes (Andreotti et al. 2009). Finally,  
66 the morphodynamics of aeolian bedforms are strongly dependent on the wind  
67 regime (Livingstone and Warren 1996). Thus, small-scale bedforms are sen-  
68 sitive to wind disturbances induced by larger-scale bedforms. This has been  
69 demonstrated for impact ripples on dunes (Howard 1977; Hood et al. 2021),  
70 but has never been studied for medium-scale dunes near giant dunes.

71 In the literature, arid areas have been studied at the desert scale from  
72 climate reanalyses based on global atmospheric models (Blumberg and Greeley

73 1996; Livingstone et al. 2010; Ashkenazy et al. 2012; Jolivet et al. 2021; Hu  
74 et al. 2021), such as ERA-40, ERA-Interim or ERA-5 (Uppala et al. 2005;  
75 Dee et al. 2011; Hersbach et al. 2020). However, the spatial resolution (tens of  
76 kilometers) of these reanalyses implies average quantities that do not resolve  
77 the smaller scales, ranging from the study of individual dunes to the border of  
78 mountains (Livingstone et al. 2010). Lately, the release of ERA5-Land allow to  
79 push back this limitation by providing up to 70 years of hourly wind predictions  
80 at a 9 km spatial resolution (Muñoz-Sabater et al. 2021). However, due to its  
81 recent nature, the applicability limitations of this dataset remain to be studied.

82 Here, we compare local wind speeds and directions measured in four dif-  
83 ferent places across the Namib desert to the regional predictions of the ERA5-  
84 Land climate reanalysis. When the measurement stations are surrounded by a  
85 relatively flat environment, we show that local measurements and regional pre-  
86 dictions agree with each other. On the contrary, in the sand sea, we show that  
87 they differ due to the presence of giant dunes. Furthermore, we link the magni-  
88 tude of these discrepancies to the circadian cycle of the atmospheric boundary  
89 layer. Finally, we draw implications for smaller-scale eolian bedforms.

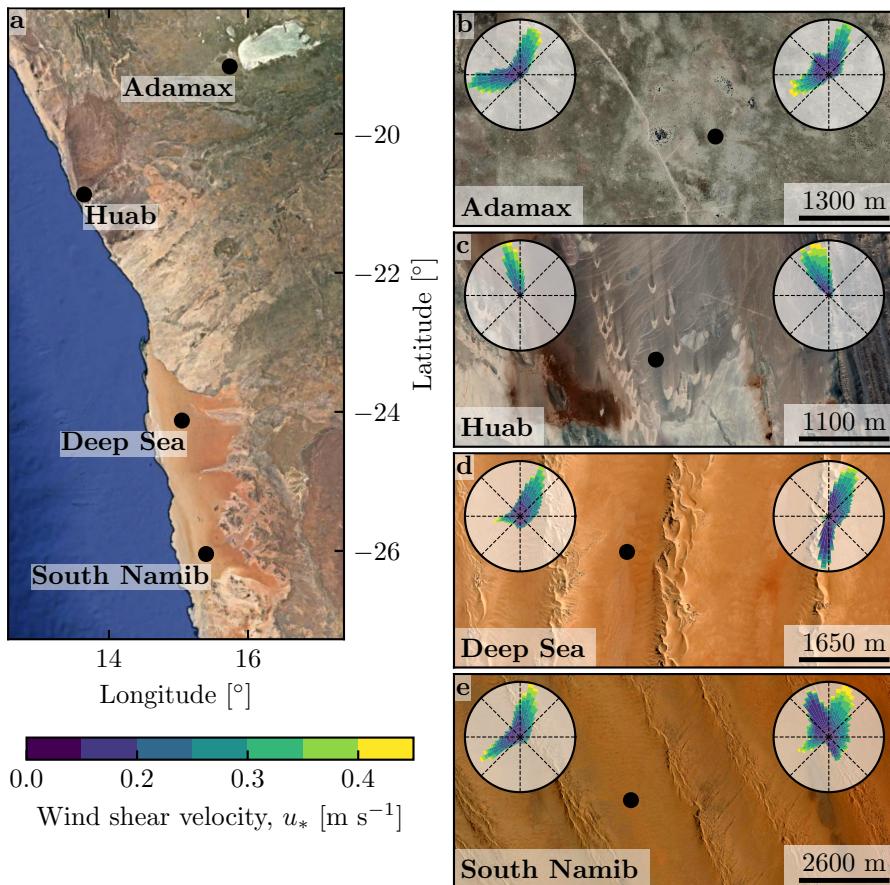
## 90 2 Wind regimes across the Namib Sand Sea

91 We focus on four places across and nearby the Namib desert, highlighting  
92 different environments (see Fig. 1). The Adamax station is located near the  
93 Adamax salt pan, in a highly vegetated area. The Huab station, on the coast  
94 at the outlet of the Huab river, is in an arid environment exhibiting 60-m scale  
95 barchan dunes. While the surroundings of these two stations are relatively flat,  
96 this is not the case for the Deep Sea and South Namib stations. Both are in  
97 the interdune between tens of meters high giant linear dunes with kilomet-  
98 ric wavelengths and superimposed patterns. In this section, we describe and  
99 compare winds from local measurements and climate reanalysis predictions.

### 100 2.1 Datasets

101 Local winds are provided by measurement stations located in the four different  
102 places (see black dots in Fig. 1). The wind strength and direction are sampled  
103 every 10 minutes by cup anemometers and wind vanes, at heights between  
104 2 m and 3 m depending on the station. The available period of measurements  
105 ranges from 1 to 5 discontinuous years distributed between 2012 and 2020 (see  
106 Fig. S1). We checked that at least one complete seasonal cycle is available at  
107 each station. Regional winds are extracted at the same locations and periods  
108 from the ERA5-Land dataset, which is a replay at a smaller spatial resolution  
109 of ERA5, the latest climate reanalysis from the ECMWF (Hersbach et al.  
110 2020; Muñoz-Sabater et al. 2021). It provides hourly predictions of the 10-m  
111 wind velocity and direction at a spatial resolution of  $\sim 9$  km ( $0.1^\circ \times 0.1^\circ$ ).

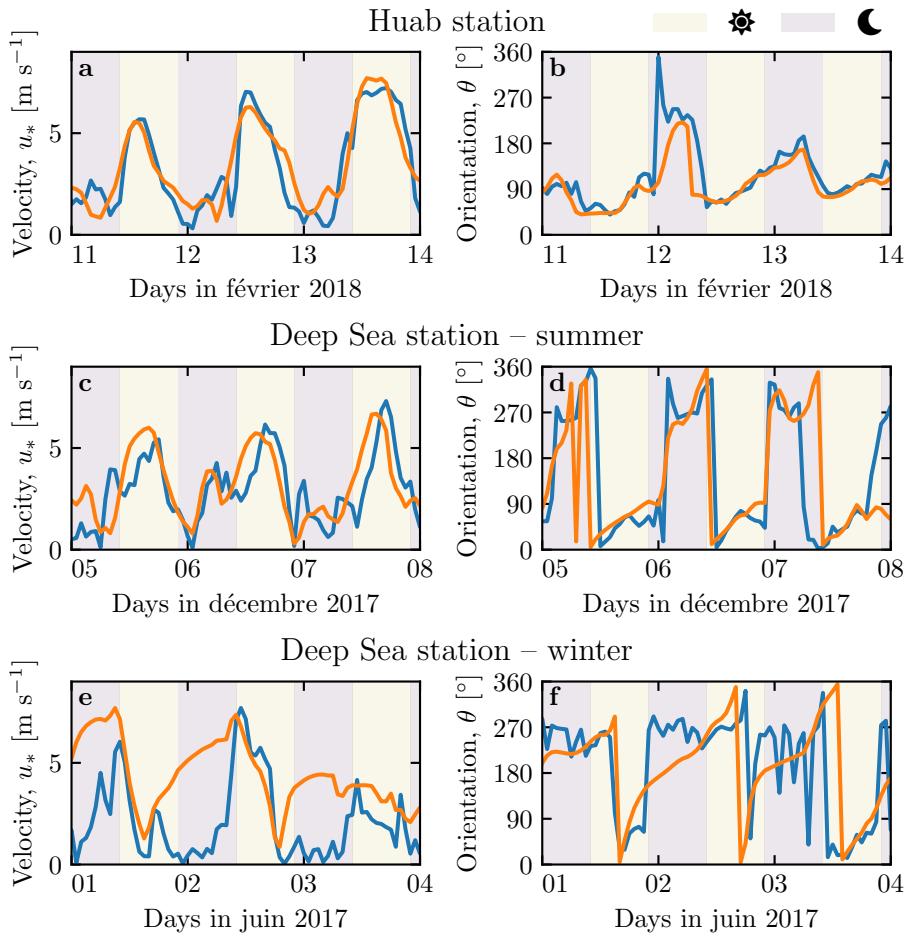
112 For comparison, the local measurements are averaged into 1-hr bins cen-  
113 tered on the temporal scale of the ERA5-Land estimates (see Fig. S2). As the



**Fig. 1** Wind data used in this study **a**: Location of the studied sites. **b–e**: Satellite images of the studied sites (Google-Earth, Maxar Technologies, CNES/Airbus). In each subplot, the left and right wind roses represent the data from the ERA5Land climate reanalysis and the local wind stations, respectively. Note that the bars show the direction towards which the wind blows. The black dots show the location of local wind stations.

114 wind velocities of both datasets are provided at different heights, we convert  
 115 them into shear velocities (see SI section 1), characteristic of the turbulent  
 116 wind profile, which are then used together with the wind direction for further  
 117 analysis. The resulting wind data are shown on the wind roses of Fig. 1(b–e).

118 Finally, the dune properties are computed using autocorrelation on the 30-  
 119 m Digital Elevation Models (DEMs) of the shuttle radar topography mission  
 120 (Farr et al. 2007). For the South Namib and Deep Sea stations, we obtain  
 121 respectively orientations of 85° and 125°, wavelengths of 2.6 km and 2.3 km  
 122 and amplitudes of 45 m and 20 m (see Fig. S4 for more details).



**Fig. 2** Temporal comparison between the wind data coming from the Era5Land climate reanalysis (orange lines) and from the local measurements (blue lines). Color swathes indicate day (between 1000 UTC and 2200 UTC) and night (before 1000 UTC or after 2200 UTC). **a–b:** Huab station. **c–d:** Deep Sea station in winter. **e–f:** Deep Sea station in summer.

## 123 2.2 Agreement between local and regional winds

124 The obtained wind regimes are shown in figure 1. In the Namib, the regional  
 125 wind patterns are essentially controlled by the see breeze, resulting in strong  
 126 northward components (sometimes slightly deviated by the large scale topog-  
 127 raphy) present in all regional wind roses (Lancaster 1985). These daytime  
 128 winds are dominant during the second-half of the year (Septembre–January).  
 129 In winter, an additional easterly component can be recorded during the night,  
 130 induced by the combination of katabatic winds forming on the mountains,  
 131 and infrequent ‘berg’ winds, which are responsible of the high wind velocities  
 132 observed (Lancaster 1984). The frequency of these easterly components de-

creases from the inland to the coast, resulting in bidirectional wind regimes within the Namib Sand Sea and at the Adamax salt pan (Fig. 1b, 1d and 1e) and a unidirectional wind regime on the coast at the outlet of the Huab River (Fig. 1c).

In the case of the Adamax and Huab stations, the wind roses from the regional predictions qualitatively match those corresponding to the local measurements. However, for the Deep Sea and South Namib stations, the measured wind roses exhibit additional components aligned with the giant dune orientation visible on the satellite images (Fig. 1c–d). The time series of wind speed and direction show that this agreement in the case of Adamax and Huab stations is always verified (Fig. 2a–b) and Fig. S5). In contrast, for the stations within the giant dune field, we observe that this agreement is limited to Septembre–January time periods (Fig. 2c–d).

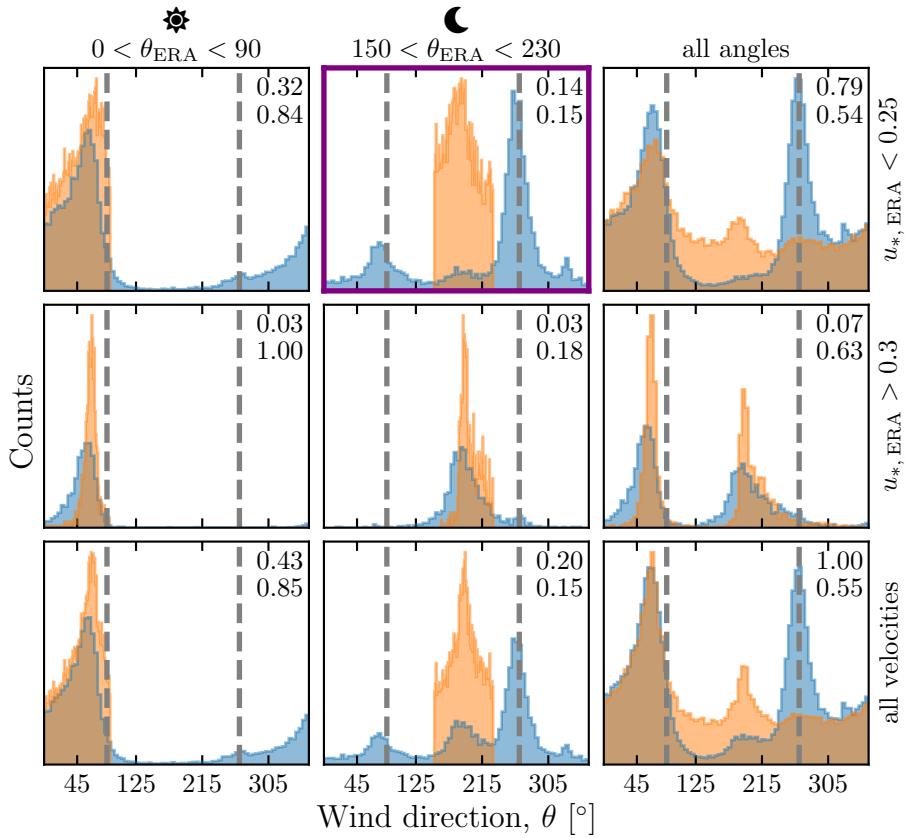
### 2.3 Influence of the giant dunes on local wind regimes

In the February–August period, when giant dunes are present, the local and regional winds match during daytime only, i.e when the southerly/southwesterly sea breeze dominates (see Fig. 2(e–f), Fig. 3 and Fig. S6). In the late afternoon and during the night, when the northwesterly ‘berg’ and katabatic winds blow, measurements and predictions differ. In this case, the angular wind distribution of the local measurements exhibits two additional modes separated of  $\simeq 180^\circ$ , each corresponding to the giant dune alignment (see the purple frame in Fig. 3 and Fig. S6, as well as Fig. S7). This deviation is also associated with a global attenuation of the wind strength (Fig. S8). Remarkably, all these figures show that this process occurs for low wind velocities, typically for  $u_* < 0.1 \text{ m s}^{-1}$ . For shear velocities larger than  $0.25 \text{ m s}^{-1}$ , this wind reorientation does not occur. Finally, for intermediate shear velocities, both reorientation along the dune crest and no reorientation are observed (Fig. S7).

## 3 Influence of the circadian cycle of the atmospheric boundary layer

For linear ridges, dune-induced flow disturbances have mainly been related to the angle between wind direction and crest orientation, with a maximum for angles between  $30^\circ$  and  $70^\circ$  (Walker et al. 2009; Hesp et al. 2015). In our case, the most deflected wind for both stations is the most perpendicular, such that the incident wind direction does not seem to be the dominant parameter controlling the wind deflection. In contrast, a different behavior is observed between low and high wind velocities, suggesting a change in hydrodynamical regime.

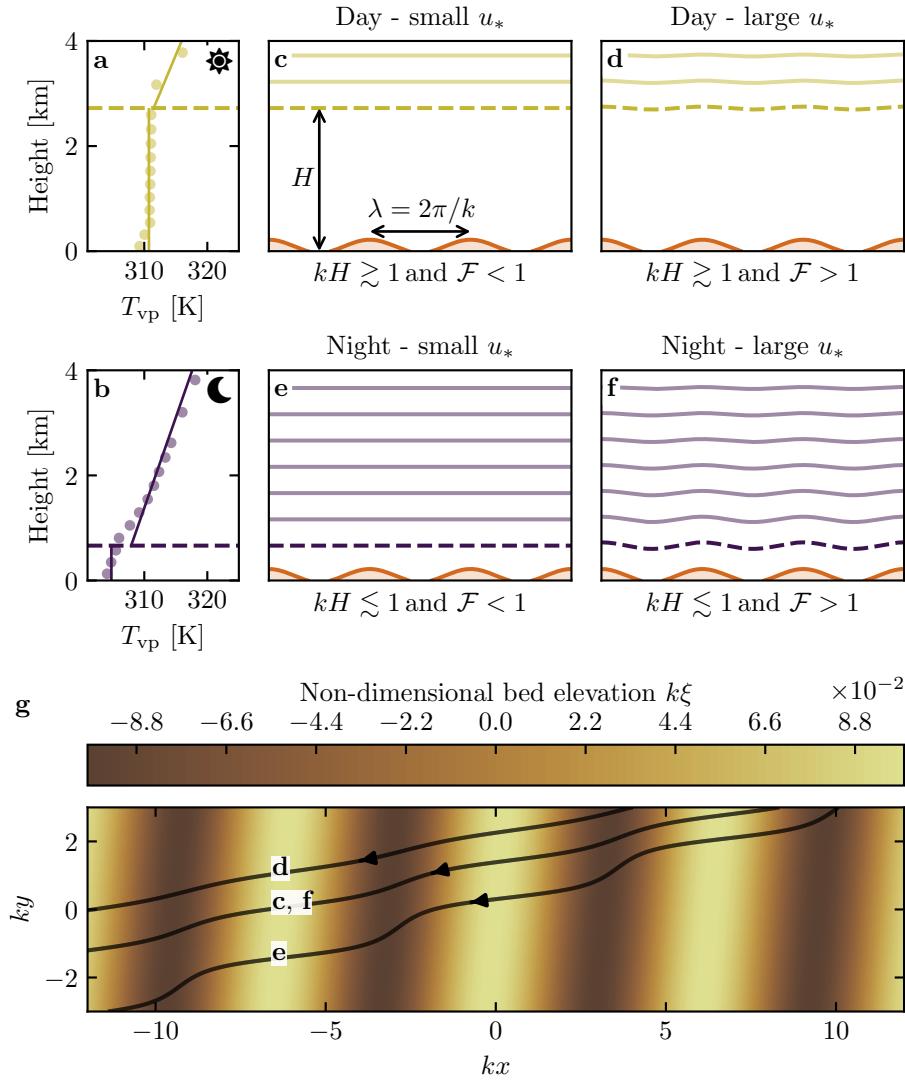
In the following, we discuss the relevant parameters leading to different hydrodynamical interactions with topographical obstacles, and interpret the data with respect to the corresponding physical mechanisms.



**Fig. 3** Distributions of wind direction at the Deep Sea Station for the Era5Land climate reanalysis (orange) and the local measurements (blue). In each subplot, both distributions are plotted from the same time steps, selected with constraints on the wind direction (columns) and/or wind velocity (rows) of the Era5Land dataset. The gray vertical dashed lines indicate the dune orientation. The numbers at the top right give the percentage of time steps selected in each sub-range, as well as the percentage of them corresponding to the day (between 1000 UTC and 2200 UTC). The purple frame highlights the regime (small wind velocities, nocturnal summer wind) in which the data from both datasets differs. A similar figure can be obtained for the Deep Sea station (see Fig. S6).

173 3.1 Relevant non-dimensional parameters and physical considerations

174 Flow deflection over ridges can be understood from the Bernoulli principle (Hesp  
 175 et al. 2015). As the flow approaches the ridge crest, the compression of the  
 176 streamlines results in larger flow velocities, and thus lower pressures (Rubin  
 177 and Hunter 1987). An incident flow oblique to the ridge is then deflected to-  
 178 wards lower pressure zones, i.e towards the crest. Turbulent dissipation at the  
 179 bottom and non-linearities tends to increase this effect downstream, resulting  
 180 in along the crest wind deflection in the lee side (Hesp et al. 2015; Gadal et al.  
 181 2019).



**Fig. 4** **a–b:** Vertical profiles of the virtual potential temperature at 2 different time steps (day - 31/03/2017 - 1200 UTC, night - 21/03/2017 - 2200 UTC) at the Deep Sea station. Dots: data from the ERA5 reanalysis. Transparent dashed lines: boundary layer height given by the ERA5 reanalysis, calculated from the bulk Richardson number (Seidel et al. 2012). Plain lines: vertical (boundary layer) and linear (free atmosphere) fits to estimate the stratification properties. **c–f:** Sketches representing the interaction between the giant dunes and the atmospheric flow for different meteorological conditions. **g:** Streamlines qualitatively representing the effect of low, medium and high flow confinement. For details on the streamline derivation, see Appendix 4.

Another way to increase the flow deflection is its confinement below a capping surface, that results in further streamline compression. This is the case for bedforms forming in open channel flows such as rivers (Fourrière et al. 2010; Unsworth et al. 2018), but also for eolian dunes. These dunes evolve in the turbulent atmospheric boundary layer (ABL), which is capped by a transitional layer separating it from the stratified atmosphere above (see Fig. 4). Two different mechanisms control the possibility of this additional streamline compression.

On one hand, it depends if the flow disturbance induced by the underlying topography reach the surface. As obstacles typically disturb flow over a characteristic height similar to their length, the potential of interaction between the dunes and the overlying surface is well captured by the parameter  $kH$ , where  $k = 2\pi/\lambda$  is the wavenumber and  $H$  the ABL depth. Note that  $H$  is directly related to the radiative fluxes at the Earth surface, and thus varies with the circadian and seasonal cycles. Here, the giant dunes have kilometric wavelengths, such that  $0.02 \lesssim kH \lesssim 5$ , and they interact most of the time with the capping layer and the stratified free atmosphere (FA) above (Andreotti et al. 2009). Interestingly, the limit of no-interactions between the topography and the boundary layer structure ( $kH \gg 1$ ), in which the properties of the overlying atmospheric structure are irrelevant, is never reached here, in the case of giant dunes.

On the other hand, it depends on rigidity of the capping surface, as its deformation releases the confinement effect inducing streamline compression. This is typically quantified using the Froude number (Vosper 2004; Stull 2006; Sheridan and Vosper 2006; Hunt et al. 2006; Jiang 2014):

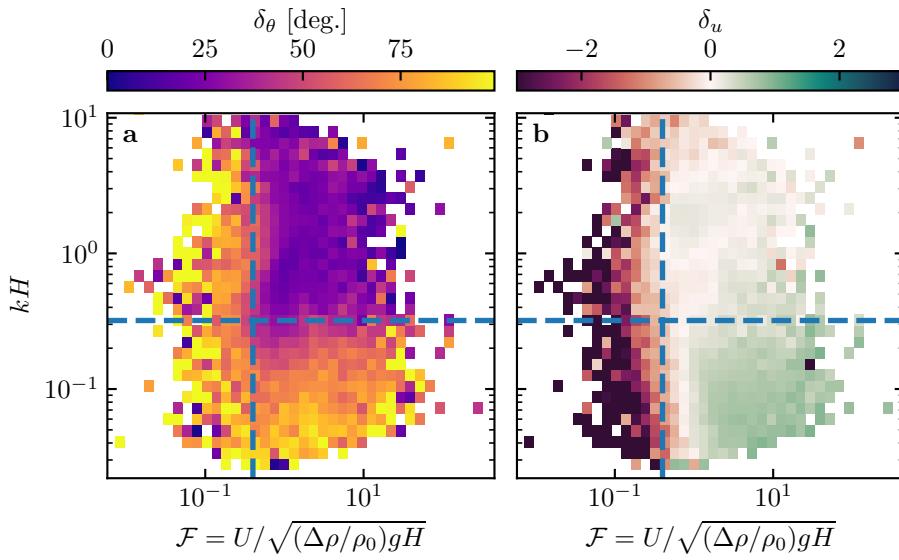
$$\mathcal{F} = \frac{U}{\sqrt{\frac{\Delta\rho}{\rho_0} g H}}, \quad (1)$$

where  $U$  is the wind velocity at the top of the ABL,  $\rho_0$  its average density,  $\Delta\rho$  the density jump between the ABL and the FA.

The smallest wind disturbances are expected during the day, when the ABL depth is comparable to the dune wavelength ( $kH \gtrsim 1$ ) and for large wind velocities, which correspond to a weak confinement situation (Fig. 4d). On the contrary, large wind disturbances are expected to occur during the night, when the confinement is mainly induced by shallow ABL (Fig. 4e–f). Note that this strong confinement can be somewhat reduced in the case of strong winds (corresponding to large Froude numbers, see Fig. 4f), explaining the transition from deflected to non-deflected winds related to low and high velocities observed in the data (see section 2.3).

### 3.2 Flow regime diagrams

To highlight these different regimes from our data, we compute wind disturbance diagrams in the space defined by the two relevant non-dimensional



**Fig. 5** Regime diagrams of the wind deviation  $\delta_\theta$  and relative attenuation/amplification  $\delta_u$  in the space  $(\mathcal{F}, kH)$ , containing the data from both the Deep Sea and South Namib stations. Blue dashed lines empirically delimit the different regimes. The point density in each bin of the diagrams is shown in Fig. S11. The regime diagrams in the spaces  $(\mathcal{F}_1, kH)$  and  $(\mathcal{F}_1, \mathcal{F})$  are shown in Fig. S12.

numbers presented above,  $(kH, \mathcal{F})$ . Those are calculated from the time series of the geopotential, temperature and specific humidity vertical profiles available in the ERA5 climate reanalysis (see SI section 2). Flow deviation is computed as the minimal angle between the wind orientations from the local measurements, and the regional predictions. The relative velocity modulation is computed as

$$\delta_u = \frac{u_*^{\text{ERA}} - u_*^{\text{station}}}{u_*^{\text{ERA}}}. \quad (2)$$

When representing the two variables  $\delta_\theta$  and  $\delta_u$  in this space, different regime emerges (Fig. 5). Small wind disturbances ( $\delta_\theta \rightarrow 0, \delta_u \rightarrow 0$ ) are located in the top-right part of the diagrams, corresponding to a regime mixing low-interaction and low-confinement ( $kH$  and  $\mathcal{F}$  large enough, Fig. 4d). Lower values of  $kH$  (stronger interaction) or Froude number (stronger confinement) then both lead to an increase in wind disturbances, both in terms of orientation and velocity. Below a threshold value of  $kH \simeq 0.3$ , wind disturbance occurs independently of the Froude number value, probably due to enhanced non-linear effects linked to strong flow modulation by the obstacle in this part of the diagram. The Froude number also controls a transition from damped to amplified wind velocities in the interdune, with a transition at  $\mathcal{F} \simeq 0.4$  (Fig. 5b). This may be linked to a transition in the flow regime in the lee side of the obstacle (lee waves, hydraulic jumps, rotors) but further measurements are needed in order to assess this (Baines 1995; Vosper 2004).

241 3.3 On the influence of the stratification of the free atmosphere

242 The presence of a stratification in the free atmosphere can also impact the  
 243 flow confinement, depending on its ability to deform under the presence of an  
 244 underlying obstacle. This can be quantified using the internal Froude num-  
 245 ber (Vosper 2004; Stull 2006; Sheridan and Vosper 2006; Hunt et al. 2006;  
 246 Jiang 2014):

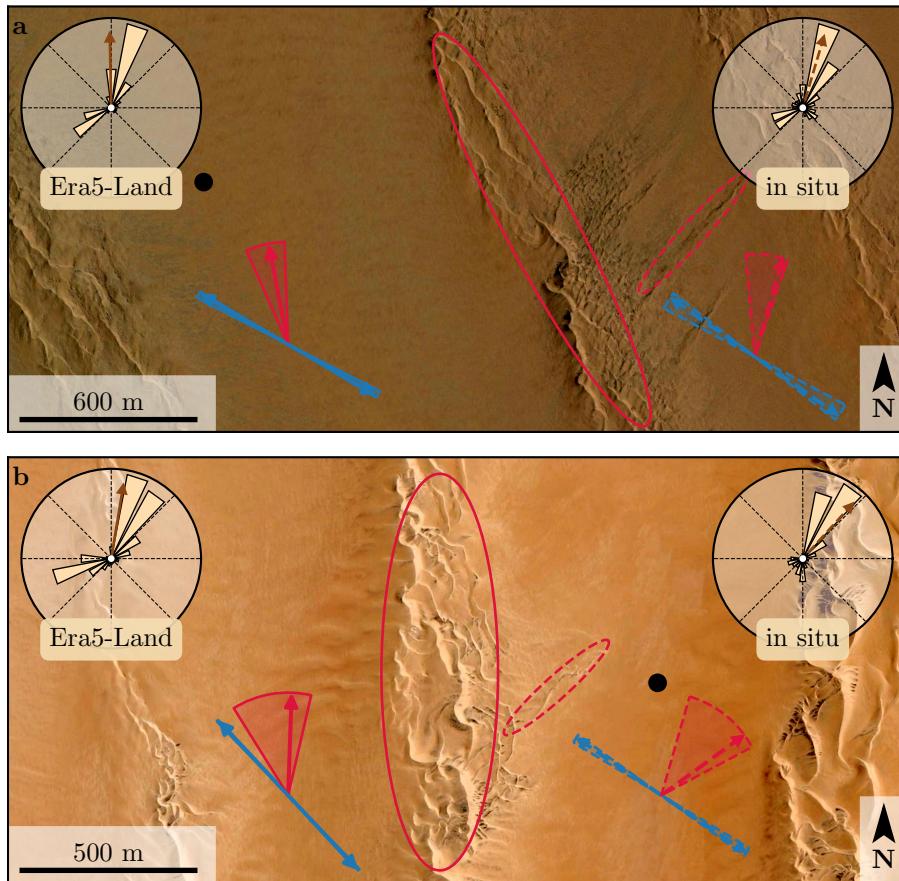
$$\mathcal{F}_I = \frac{kU}{N}, \quad (3)$$

247 where  $N = \sqrt{-(g/\rho_0)(\partial\rho/\partial z)}$  is characteristic of the stratification. Both  
 248 Froude numbers have the same qualitative effect on flow confinement, as they  
 249 quantify the rigidity of the overlying layers. This is confirmed by figure S12,  
 250 where we can also find the different regimes related to wind disturbances de-  
 251 scribed previously for the Froude number  $\mathcal{F}$ .

252 **4 Discussion**

253 The comparison of local (direct measurements) and regional (climate reanal-  
 254 ysis) wind data reveals the giant dunes feedback on the flow. In flat areas,  
 255 the matching between measurements and prediction confirms the ability of the  
 256 ERA5Land climate reanalysis to predict the wind flow down to scales  $\sim 10$  km,  
 257 i.e the grid model. When smaller scale topographies are present (giant dunes  
 258 in our case), locally measured wind regimes can significantly differ from the  
 259 regionally predicted ones. Furthermore, we link these disturbances induced by  
 260 the dunes to their interaction with the lower part of the atmospheric vertical  
 261 structure, and more specifically to its circadian variability. During the night,  
 262 the presence of a shallow atmospheric boundary layer (ABL) induces a strong  
 263 confinement of the flow, associated with large wind deviation and acceleration  
 264 or deceleration. During the day, the capping layer is high enough to prevent  
 265 its interaction with the giant dunes, resulting in a low confinement of the flow,  
 266 and thus smaller wind disturbances. Interestingly, we also found that this ef-  
 267 fect could be counterbalanced by the presence of large wind velocities, capable  
 268 of deforming the capping layer and/or the FA stratification, thus decreasing  
 269 the confinement effect.

270 Simple linear models also suggest that larger wind disturbances occur un-  
 271 der strong flow confinement such as described above Andreotti et al. (2009,  
 272 2012). However, they are unable to reproduce the magnitude of the observed  
 273 deviations, probably due to the presence of hydrodynamical non-linear effects,  
 274 all the more present in high confinement situations linked to strong flow mod-  
 275 ulations (see Fig. S12 and Appendix 1). They also predict different spatial flow  
 276 structures such as lee waves and rotors (Baines 1995; Vosper 2004), which are  
 277 likely to be complicated by these non-linearities, and which cannot be observed  
 278 by our single point measurements. Measurements in different places along and  
 279 across the ridge are then needed in order to properly map these flow structures,  
 280 and allow further comparisons with models.



**Fig. 6** Implications for smaller scale patterns in (a) the South Namib and (b) Deep Sea. The ellipses indicates the different types of elongating dunes, at large (plain) and small scale (dashed). The dune orientations are calculated using the model of Courrech du Pont et al. (2014) from the sand flux angular distributions, shown here for typical sand quartz grains of  $180\ \mu\text{m}$ . The double blue and single red arrows correspond to the two possible dune growth mechanisms, bed instability and elongation, respectively. Likewise, plain arrows are calculated from the ERA5-Land data, and dashed arrows from the local measurements. Wedges show the uncertainty on the orientation calculation, and the arrows correspond to typical parameters found in the literature, i.e. a grain diameter of  $180\ \mu\text{m}$  and a flux-up ratio of 1.6. The black dots indicate the position of the measurement stations. See Appendix 2 for additional details.

This study highlights the interaction between giant dunes and the atmospheric boundary layer. It then supports the debated idea that the capping layer acts as a bounding surface limiting dune growth (Andreotti et al. 2009), as opposed to an unconstrained growth ever-slower with size (Eastwood et al. 2011; Gunn et al. 2021). Once validated, this mechanism would then allow inference of the ABL depth from the giant bedforms spacing where measure-

287 ments are not feasible or available, as performed by Lorenz et al. (2010) on  
 288 Titan.

289 This interaction also have strong implications for smaller scales bedforms,  
 290 as illustrated in Fig. 6. In the Namib Sand Sea, small linear dunes ( $\sim 50$  m  
 291 -wide) are present in the interdune between giant linear dunes ( $\sim 2$  km -wide).  
 292 While differences between larger and smaller scale dune patterns are observed  
 293 ubiquitously, they are now largely attributed to the presence of two different  
 294 dune growth mechanisms, leading to two different dune patterns (orientations  
 295 and/or morphologies) for the same wind regime (Courrech du Pont et al.  
 296 2014; Runyon et al. 2017; Lü et al. 2017; Song et al. 2019; Hu et al. 2021).  
 297 Here, coupling sediment transport and dune growth models, we show that the  
 298 orientations of the small and giant linear dunes can be predicted from the same  
 299 dune growth mechanism, using the locally measured and regionally predicted  
 300 winds, respectively (red arrows in Fig. 6). The giant dune feedback on the  
 301 flow described in this study then provides a mechanism for the existence of  
 302 these small linear dunes elongating across the interdune, as yet unresolved.  
 303 While further studies are needed, these dune type could provide additional  
 304 strong constraints for the inference of local winds from bedforms, as currently  
 305 performed on Mars using ripple orientations (Liu and Zimbelman 2015; Hood  
 306 et al. 2021).

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 308 source python librairies, Matplotlib (Hunter 2007), Numpy (Harris et al. 2020) and Scipy  
 309 (Virtanen et al. 2020), which provide an incredibly efficient ecosystem allowing scientific  
 310 research in Python.

311 Era5 and Era5Land datasets are publicly available at the Copernicus Climate Change  
 312 Service (C3S) Climate Data Store. The locally measured wind data can be found at [up-](#)  
 313 [load on public data repository](#). The digital elevation models from the Shuttle Radar To-  
 314 polography Mission are publicly available from Nasa servers, and can be downloaded at  
 315 <https://dwtkns.com/srtm30m/>. Fully documented codes used to analyze this study are  
 316 available at <https://github.com/Cgadal/GiantDunes> (will be made public upon acceptance  
 317 of this manuscript for publication).

318 [citing all grants ...]

319

## 320 Appendix 1: ABL turbulent wind model

321 Following the work of Fourrière et al. (2010), Andreotti et al. (2012) and Andreotti  
 322 et al. (2009), we briefly expose in this section the linear response of a  
 323 turbulent flow to a small aspect ratio perturbation of the underlying topogra-  
 324 phy. As this topography can be decomposed into several sinusoidal modes, we  
 325 focus on a sinusoidal topography:

$$\xi = \xi_0 \cos [k (\cos(\alpha)x + \sin(\alpha)y)], \quad (4)$$

326 which is also a good approximation for the giant dunes observed in the Deep  
 327 Sea and South Namib Station (see Fig 1 and Fig S4). Here,  $x$  and  $y$  are  
 328 the streamwise and spanwise coordinates,  $k = 2\pi/\lambda$  the wavenumber of the

329 sinusoidal perturbation, and  $\alpha$  its crest orientation, calculated with respect to  
 330 the  $y$ -direction.

331 The two components of the basal shear stress  $\tau = \rho_0 u_* \mathbf{u}_*$ , constant in a  
 332 flat bottom situation, can then be written without loss of generality as:

$$\tau_x = \tau_0 \left( 1 + k\xi_0 \sqrt{\mathcal{A}_x^2 + \mathcal{B}_x^2} \cos [k(\cos(\alpha)x + \sin(\alpha)y) + \phi_x] \right), \quad (5)$$

$$\tau_y = \tau_0 k\xi_0 \sqrt{\mathcal{A}_y^2 + \mathcal{B}_y^2} \cos [k(\cos(\alpha)x + \sin(\alpha)y) + \phi_y], \quad (6)$$

333 where  $\tau_0$  is the basal shear stress on a flat bed, and  $\phi_{x,y} = \tan^{-1}(\mathcal{B}_{x,y}/\mathcal{A}_{x,y})$ .  
 334 The in-phase and in-quadrature hydrodynamical coefficients  $\mathcal{A}_{x,y}$  and  $\mathcal{B}_{x,y}$   
 335 are functions of the flow conditions, i.e the bottom roughness, the vertical  
 336 flow structure or the incident flow direction (Fourrière et al. 2010; Andreotti  
 337 et al. 2009, 2012).

338 Following Andreotti et al. (2012), the impact of the incident wind direction  
 339 can be approximated by the following expressions:

$$\mathcal{A}_x = \mathcal{A}_0 \cos^2 \alpha, \quad (7)$$

$$\mathcal{B}_x = \mathcal{B}_0 \cos^2 \alpha, \quad (8)$$

$$\mathcal{A}_y = \frac{1}{2} \mathcal{A}_0 \cos \alpha \sin \alpha, \quad (9)$$

$$\mathcal{B}_y = \frac{1}{2} \mathcal{B}_0 \cos \alpha \sin \alpha, \quad (10)$$

340 where  $\mathcal{A}_0$  and  $\mathcal{B}_0$  are now two coefficients independent of the dune orientation  
 341  $\alpha$ . In the case of a fully turbulent boundary layer capped by a stratified atmo-  
 342 sphere, they now only depend on  $kH$ ,  $kz_0$ ,  $\mathcal{F}$  and  $\mathcal{F}_I$  Andreotti et al. (2009).  
 343 In this study, we assume a constant hydrodynamic roughness  $z_0 \sim 1$  mm,  
 344 leading to a constant value of  $kz_0 \sim 10^{-6}$ . Measured values of  $z_0$  in the field  
 345 indeed reports a variation of  $z_0$  between 0.1 mm and 10 mm (Sherman and  
 346 Farrell 2008; Field and Pelletier 2018), but  $\mathcal{A}_0$  and  $\mathcal{B}_0$  does not vary much in  
 347 the corresponding range of  $kz_0$  (Fourrière et al. 2010). Note that the linearity  
 348 assumption of this theoretical framework requires  $(|\tau| - \tau_0)/\tau_0 \ll 1$ , which is  
 349 satisfied by  $k\xi \sqrt{\mathcal{A}_0^2 + \mathcal{B}_0^2} \ll 1$ . In our case, the giant dune morphology gives  
 350  $k\xi \simeq 0.1$ , setting the upper bound of the coefficient modulus  $\sqrt{\mathcal{A}_0^2 + \mathcal{B}_0^2}$  to 10.

351 Additionally, we also calculate the time series of the hydrodynamical coeffi-  
 352 cients from the time series of the non-dimensional numbers used in this study.  
 353 The results, shown Fig. S12 under the similar form of the regime diagrams  
 354 presented in Fig. 5 and Fig. S12, exhibit a qualitative matching with the ob-  
 355 servations presented in this study. Small values of these coefficients ( $\mathcal{A}_0 \simeq 3.4$   
 356 and  $\mathcal{B}_0 \simeq 1$ ) are found in low confinement cases, i.e for  $kH \gg 1$  (no inter-  
 357 action between the dunes and the capping layer) or for  $kH \gg 1$  but large  
 358 enough Froude numbers (reduced flow confinement due to the deformation of  
 359 the overlying capping layer and stratification). In contrast, larger values are  
 360 obtained in high confinement cases (small  $kH$  and Froude numbers). How-  
 361 ever, not that most of this part of the diagrams are outside of the linear limit

362 discussed above ( $\sqrt{A_0^2 + B_0^2} \gtrsim 10$ ), which does not allow further quantitative  
 363 comparison with the data.

## 364 Appendix 2: Sediment transport and dune morphodynamics

365 Here, we briefly describe the sediment transport and dune morphodynamics  
 366 theoretical framework leading to the prediction of sand fluxes and dune orient-  
 367 tations from wind data.

368 *Sediment transport* The prediction of sand fluxes from wind data has been  
 369 a long standing issue in geomorphological studies (Fryberger and Dean 1979;  
 370 Pearce and Walker 2005; Sherman and Li 2012; Shen et al. 2019). Based on lab-  
 371 oratory studies in wind tunnels (Rasmussen et al. 1996; Iversen and Rasmussen  
 372 1999; Creyssels et al. 2009; Ho et al. 2011), as well as physical considerations  
 373 (Ungar and Haff 1987; Andreotti 2004; Durán et al. 2011; Páhtz and Durán  
 374 2020), it has been shown that the steady saturated sand flux over a flat sand  
 375 bed depends linearly on the shear stress:

$$\frac{q_{\text{sat}, t}}{Q} = \Omega \sqrt{\Theta_{\text{th}}} (\Theta_t - \Theta_{\text{th}}), \quad (11)$$

376 where  $\Omega$  is a proportionality constant,  $Q = d\sqrt{(\rho_s - \rho_0)gd/\rho_0}$  is a character-  
 377 istic flux,  $\Theta = \rho_0 u_{*,t}^2 / (\rho_s - \rho_0)gd$  the Shields number, and  $\Theta_{\text{th}}$  its threshold  
 378 value for incipient sediment transport. Here,  $\rho_s = 2.6 \text{ g cm}^{-3}$  and  $d = 180 \mu\text{m}$   
 379 are the grain density and diameter, and  $g$  is the gravitational acceleration.

380 Recently, Páhtz and Durán (2020) suggested a quadratic dependency on  
 381 the shear stress by taking into account grain–grain interactions within the  
 382 transport layer, performing better at reproducing laboratory data at high wind  
 383 velocities:

$$\frac{q_{\text{sat}, t}}{Q} = \frac{2\sqrt{\Theta_{\text{th}}}}{\kappa\mu} (\Theta_t - \Theta_{\text{th}}) \left( 1 + \frac{C_M}{\mu} [\Theta_t - \Theta_{\text{th}}] \right), \quad (12)$$

384 where  $\kappa = 0.4$  is the von Kármán constant,  $C_M = 1.7$  a constant and  $\mu$  a  
 385 friction coefficient, taken to be the avalanche slope of the granular material,  
 386 i.e.  $\sim 0.6$ . The fit of this law to the experimental data of Creyssels et al. (2009)  
 387 and Ho et al. (2011) gives  $\Theta_{\text{th}} = 0.0035$ . The sand flux angular distributions  
 388 and the dune orientations in Fig. 6 are calculated using this quartic law (12).  
 389 However, we verified that using the quadratic law (11) instead did not change  
 390 the predicted dune orientations by more than a few degrees.

391 *Dune orientations* The dune orientations are predicted from the computed  
 392 sand flux time series, using the dimensional model of Courrech du Pont et al.  
 393 (2014). Two orientations are possible depending on the mechanism dominating  
 394 the dune growth: elongation or bed instability (the latter is also known as the  
 395 rule of maximum gross bedform-normal transport from Rubin and Hunter  
 396 (1987)).

<sup>397</sup> The orientation  $\alpha$  corresponding the bed instability is then the one that  
<sup>398</sup> maximizes the following growth rate:

$$\sigma \propto \frac{1}{H_d W_d T} \int_t q_{\text{crest},t} |\sin(\theta_t - \alpha)|, \quad (13)$$

<sup>399</sup> where  $H_d$  and  $W_d$  are dimensional constants representing the dune height and  
<sup>400</sup> width, respectively. The flux at the crest is expressed as:

$$q_{\text{crest},t} = q_{\text{sat},t} [1 + \gamma |\sin(\theta_t - \alpha)|], \quad (14)$$

<sup>401</sup> where the flux-up ratio  $\gamma$  has been calibrated to 1.6 using field studies, under-  
<sup>402</sup> water laboratory experiments and numerical simulations. Similarly, the dune  
<sup>403</sup> orientation corresponding to the elongation mechanism is the one that verifies:

$$\tan(\alpha) = \frac{\langle q_{\text{crest},t}(\alpha) \mathbf{e}_{\theta_t} \rangle \cdot \mathbf{e}_{WE}}{\langle q_{\text{crest},t}(\alpha) \mathbf{e}_{\theta_t} \rangle \cdot \mathbf{e}_{SN}}, \quad (15)$$

<sup>404</sup> where  $\langle \cdot \rangle$  denotes a vectorial time average. The unitary vectors  $\mathbf{e}_{WE}$ ,  $\mathbf{e}_{SN}$  and  
<sup>405</sup>  $\mathbf{e}_{\theta_t}$  are in the West–East, South–North and wind direction, respectively.

<sup>406</sup> The resulting computed dune orientations, blue and red arrows in figure 6,  
<sup>407</sup> are then depending on a certain number of parameters (grain properties, flux-  
<sup>408</sup> up ratio), for which we took typical values for eolian desert on Earth. Due to  
<sup>409</sup> the lack of measurements in the studied places, significant uncertainties can  
<sup>410</sup> however be expected. We therefore run a sensibility test by calculating the  
<sup>411</sup> dune orientations for grain diameters ranging from 100  $\mu\text{m}$  to 400  $\mu\text{m}$  and the  
<sup>412</sup> speed-up ratio from 0.1 to 10 (wedges on figure 6).

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413 **References**

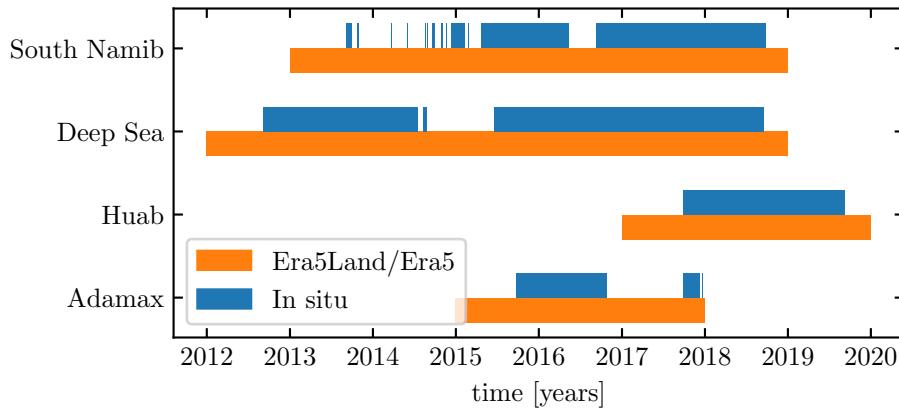
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**Fig. S1** Gant chart representing the valid time steps for the two data sets, for all stations.

630 **Supplementary Material for *Boundary-Layer Meteorology* Sample  
631 Paper: Instructions for Authors**

632 **First Author\* · Second Author · Third Author**

633  
634 \*Affiliation and email address for the corresponding author only (note that  
635 the corresponding author does not need to be the first author).

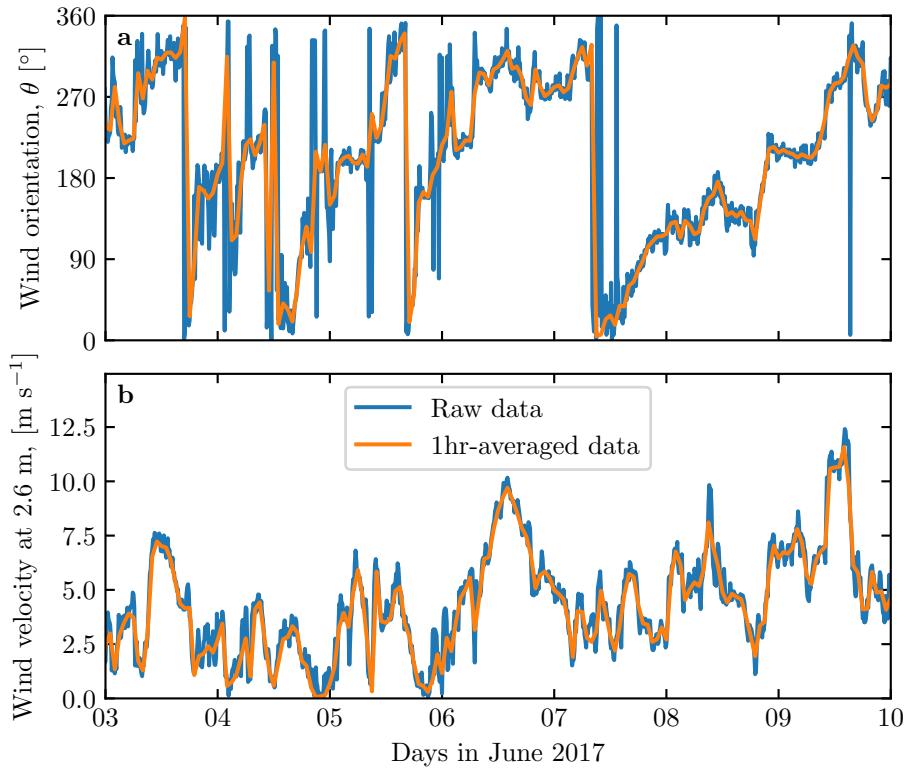
636 **1. Shear velocity and calibration of the hydrodynamical roughness**

637 As the regionally predicted and locally measured velocities are available at  
638 different heights, we can not compare them directly. We then convert all ve-  
639 locities into shear velocities  $u_*$ , characteristic of the turbulent velocity profile  
640 (Spalding 1961; Stull 1988):

$$\frac{u(z)}{u_*} = \frac{1}{\kappa} \ln \left( \frac{z}{z_0} \right), \quad (16)$$

641 where  $z$  is the vertical coordinate,  $\kappa \approx 0.4$  the von Kármán constant and  
642  $z_0$  the hydrodynamic roughness. Several field measurements of hydrodynamic  
643 roughnesses are available. In the absence of sediment transport, it scales with  
644 the geometric roughness of the topography (Pelletier and Field 2016). When  
645 transport occurs, it then scales with the thickness of the transport layer ( $\sim$   
646 1 mm), depending on the wind velocity and grain properties (Sherman and  
647 Farrell 2008; Zhang et al. 2016; Field and Pelletier 2018).

648 The different environments of our stations (vegetated, arid, sandy) makes  
649 it difficult to develop a precise model for the computation of the hydrodynamic  
650 roughness. Furthermore, we do not have precise enough topographical mea-  
651 surements that would allow to compute the geometric roughness. We choose



**Fig. S2** Comparison between raw local wind measurements, and centered averaged data over one hour for the South Namib station. **a:** wind direction. **b:** wind velocity at the measurement height, 2.6 m.

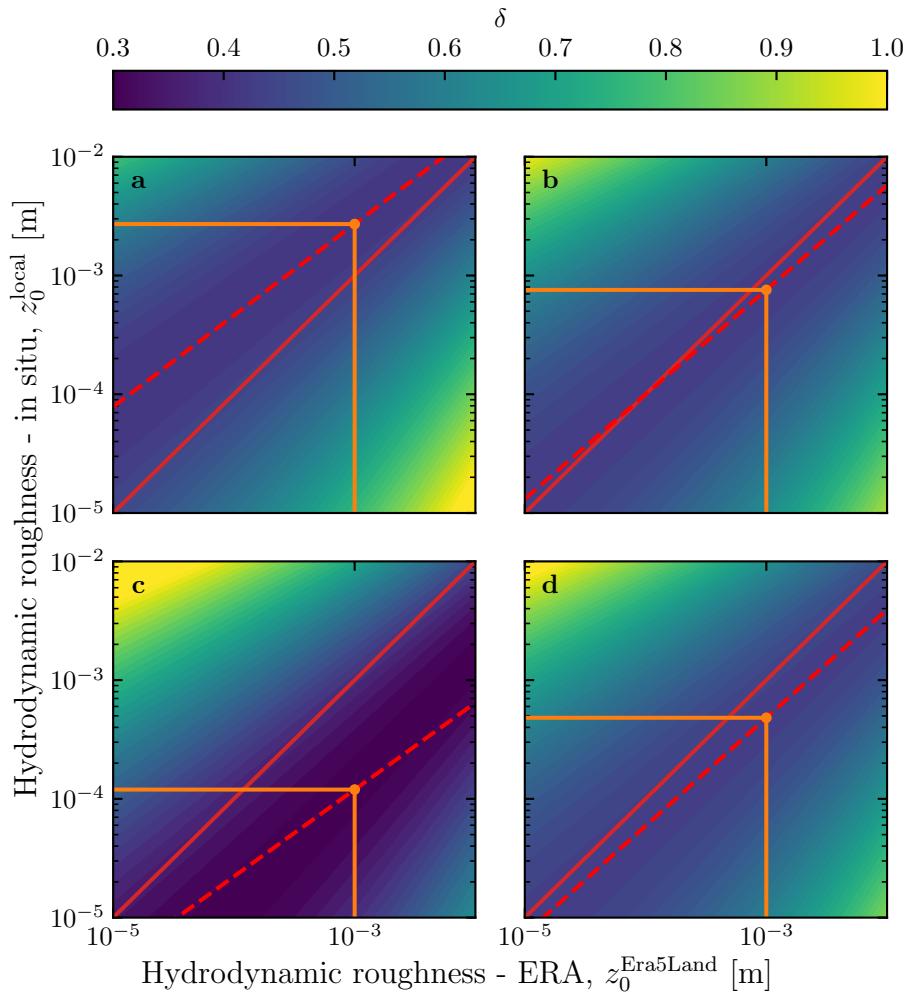
652 to leave aside this complexity, and use a different approach. The selected hydrodynamic roughness is then the one that leads to the best possible matching  
 653 between the regionally predicted and locally measured winds.  
 654

655 For each station, the hydrodynamic roughness is calibrated by finding the  
 656 one that minimizes the relative difference  $\delta$  between the wind vectors of both  
 657 datasets:

$$\delta = \frac{\sqrt{\langle \| \mathbf{u}_{*,\text{era}} - \mathbf{u}_{*,\text{station}} \|^2 \rangle_t}}{\sqrt{\langle \| \mathbf{u}_{*,\text{era}} \|^2 \rangle_t} \sqrt{\langle \| \mathbf{u}_{*,\text{station}} \|^2 \rangle_t}} \quad (17)$$

658 This  $\delta$ -parameter is computed for hydrodynamic roughness values ranging  
 659 from  $10^{-5}$  m to  $10^{-2}$  m for the different stations. Note that for the Deep Sea  
 660 and South Namib stations, where the giant dunes feedback presumably affect  
 661 the wind, we take into account the non-deflected winds only in the calculation  
 662 of  $\delta$  (with a  $15^\circ$  tolerance).

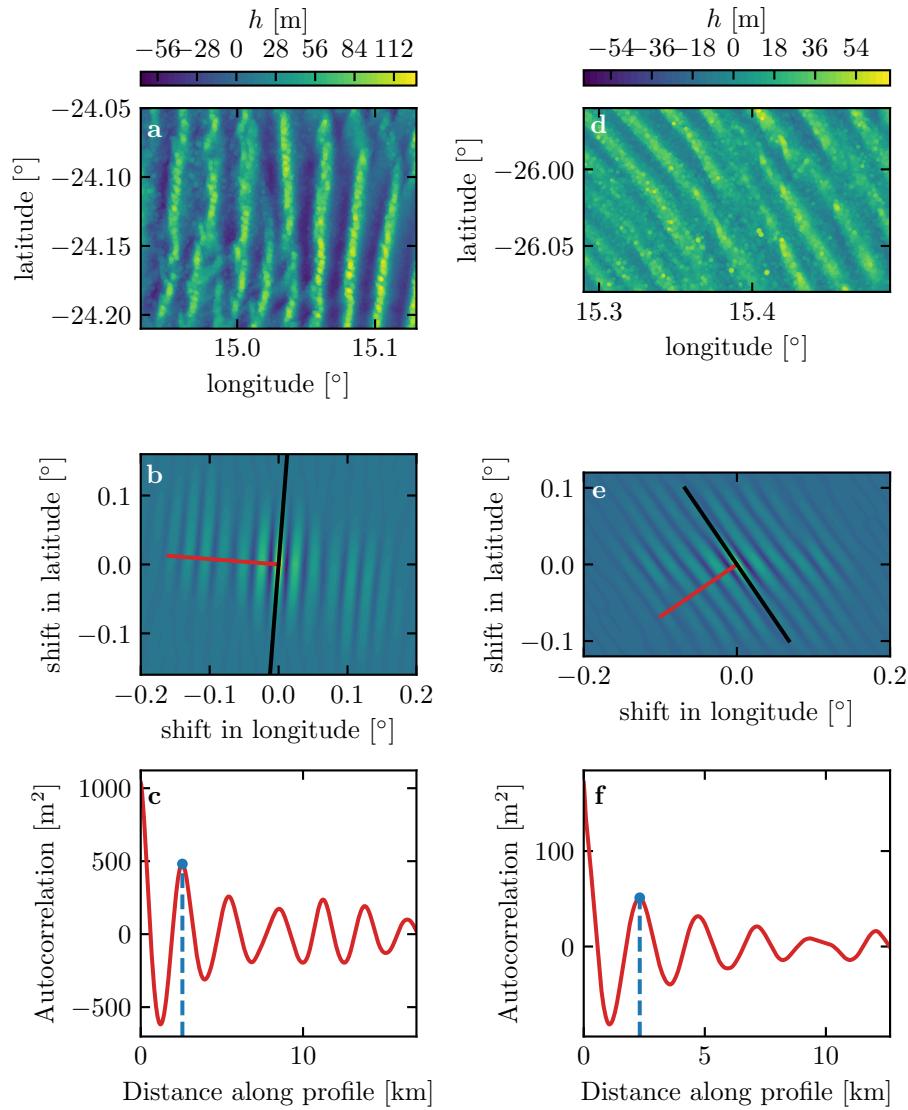
663 A shown by figure S3, the minimum of  $\delta$  in the space ( $z_0^{\text{Era5Land}}$ ,  $z_0^{\text{local}}$ )  
 664 forms a line. We thus take the roughness of the Era5Land dataset as the typical  
 665 value when sediment transport occurs,  $10^{-3}$  m, corresponding to the thickness  
 666 of the transport layer (Durán et al. 2011). It leads for the Adamax, Deep Sea,



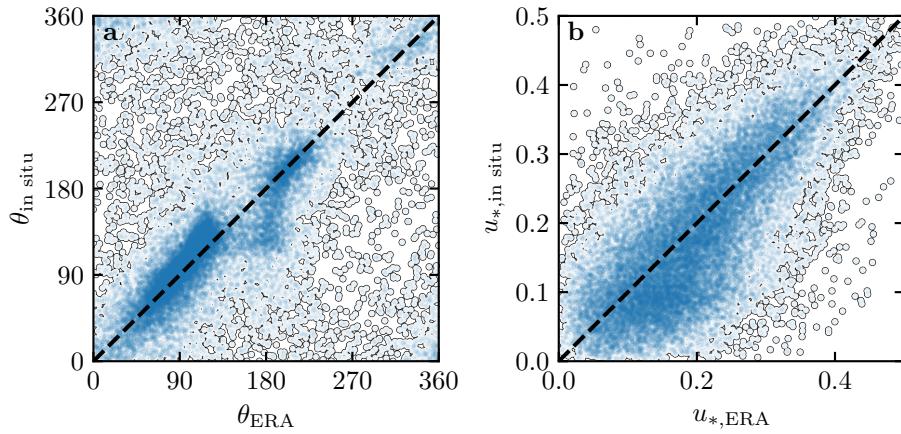
**Fig. S3** Calibration of the hydrodynamic roughnesses. The metric  $\delta$  defined in (17) is represented in colorscale as a function of the hydrodynamic roughnesses chosen for the Era5-Land and local winds, for the (a) Adamax, (b) Deep Sea, (c) Huab and (d) South Namib stations. The red dashed and plain lines shows the minima of  $\delta$  and the identity line, respectively. The orange lines and dots highlight the chosen hydrodynamic roughnesses for the local winds, by imposing  $z_0^{\text{Era5Land}} = 1 \text{ mm}$ . It leads for for each station to 2.7 mm, 0.76 mm, 0.12 mm and 0.48 mm, respectively.

667 Huab and South Namib stations values of 2.7 mm, 0.76 mm, 0.12 mm and  
668 0.48 mm, respectively.

669 The choice of the hydrodynamic roughness values impacts the calculated  
670 shear velocities only, but not the wind directions. As such, most of our conclu-  
671 sions are independent of such a choice. However, it may affect the magnitude  
672 of the wind velocity attenuation/amplification in flow confinement situations.



**Fig. S4** Analysis of the DEMs of the Deep Sea (left column – **a**, **b**, **c**) and South Namib (right column – **d**, **e**, **f**) stations. **a–d**: Detrended topography (a second order polynomial is first fitted and then removed). **b–e**: autocorrelation matrix shown in colorscale. The black line shows the detected orientation, and the red line the profile along which the wavelength is calculated, shown in **c–f**. The blue lines and dots show the first peak of the autocorrelation profile, whose abscissa gives the characteristic wavelength of the dune pattern.



**Fig. S5** Statistical agreement of the wind orientation (a) and velocity (b) between the Era5Land dataset and the local measurements for the Huab and Adamax stations. Note how the points are clustered around identity lines (dashed and black).

## 673 2. Extraction of the ABL properties

674 The estimation of the non-dimensional numbers requires the computation of  
 675 meteorological quantities representative of the current atmospheric boundary  
 676 layer. In arid areas, the vertical structure of the atmosphere can be approxi-  
 677 mated by well mixed convective boundary layer of height  $H$ , topped by the  
 678 stratified free atmosphere (Stull 1988; Shao 2008). In this type of structure,  
 679 the virtual potential temperature is constant inside the boundary layer, and  
 680 then increases linearly in the stratified free atmosphere (Fig. S9a):

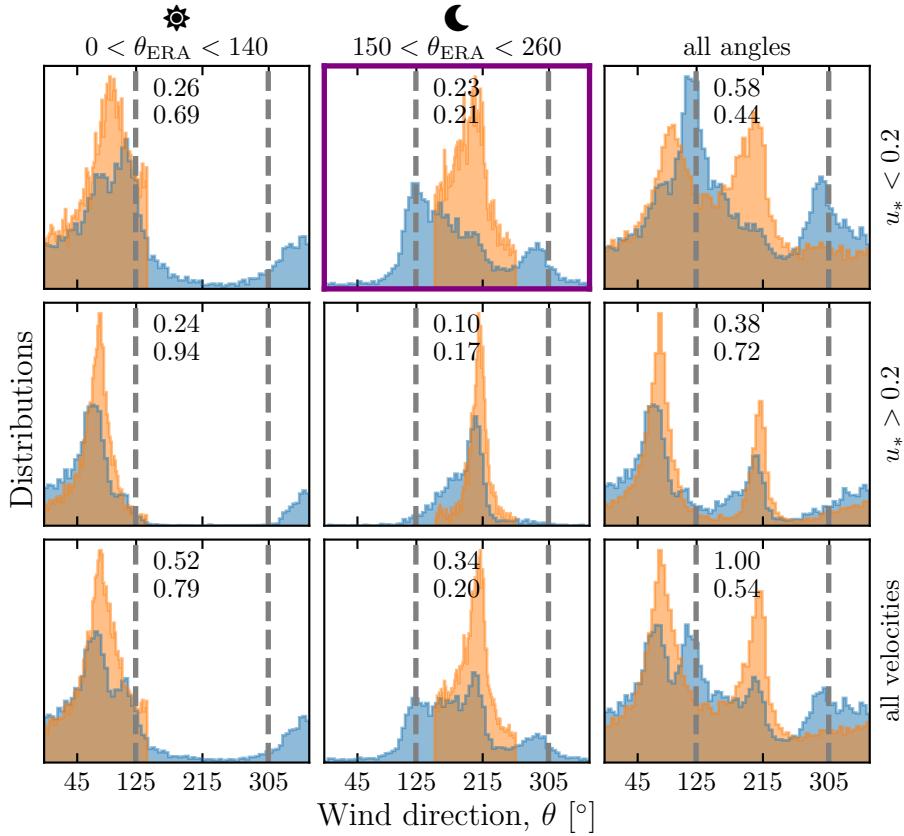
$$T_{\text{vp}}(z) = \begin{cases} T_0 & \text{for } z < H, \\ T_0 \left( 1 + \frac{\Delta T_{\text{vp}}}{T_0} + \frac{N^2}{g}(z - H) \right) & \text{for } z > H, \end{cases} \quad (18)$$

681 where,  $T_0$  is the surface virtual potential temperature,  $\Delta T_{\text{vp}}$  the discontinuity  
 682 at the capping layer and  $N$  the Brunt-Väisälä frequency, characteristic of the  
 683 stratification.

684 The Era5 dataset provides a direct estimate of the ABL depth  $H$  from a  
 685 bulk Richardson number calculation, as well as vertical profiles of the geo-  
 686 potential  $\phi$ , temperature  $T$  and specific humidity  $e_w$  at given pressure levels  $P$ .  
 687 From these quantities, the virtual potential temperature can be calculated as:

$$T_{\text{vp}} = T (1 + [R_M - 1] e_w) \left( \frac{P_0}{P} \right)^{P_c(1-0.24e_w)}, \quad (19)$$

688 where  $P_0 = 10^5$  Pa is the standard pressure,  $P_c = 0.2854$  the Poisson coefficient  
 689 for dry air and  $R_M = 1.61$  is the ratio between the molecular masses of dry



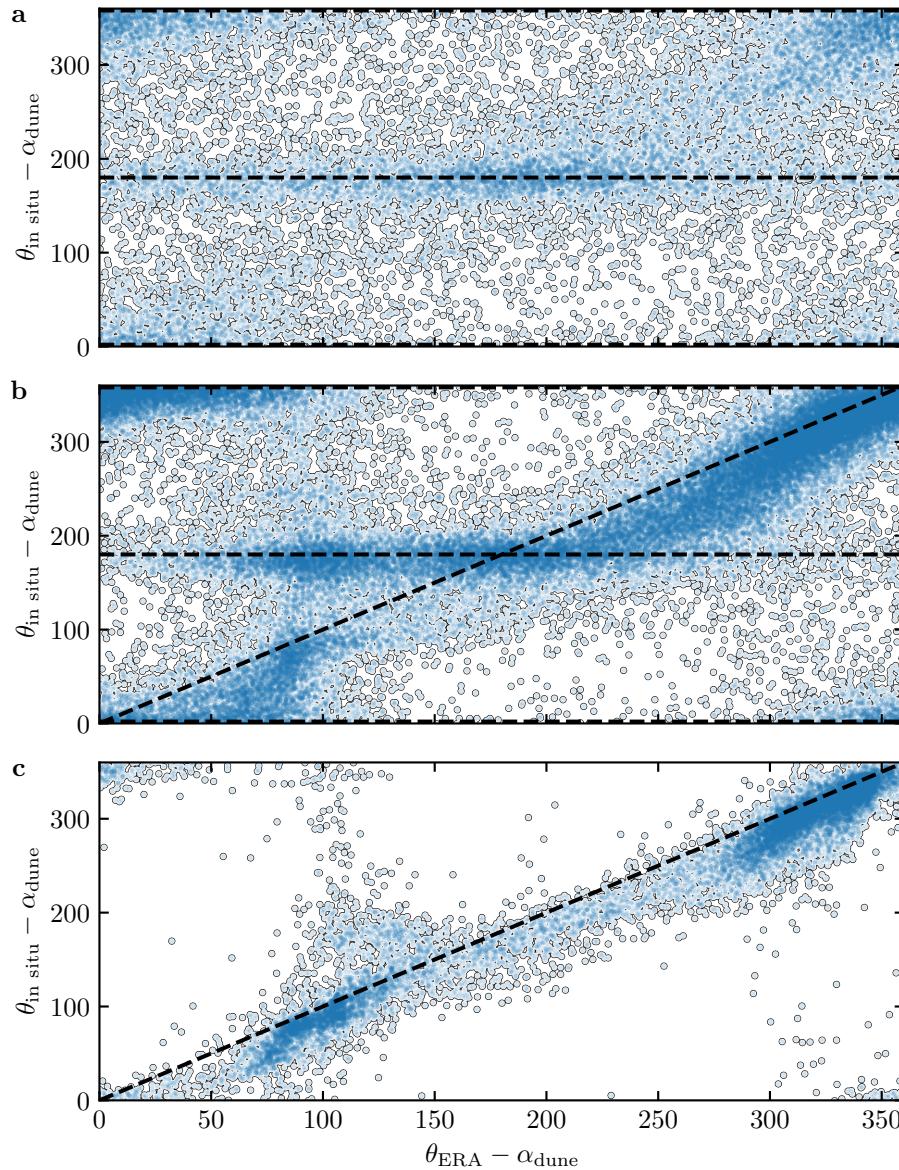
**Fig. S6** Distributions of wind direction at the South Namib Station for the Era5Land climate reanalysis (orange) and the local measurements (blue). In each subplot, both distributions are plotted from the same time steps, selected with constraints on the wind direction (columns) and/or wind velocity (rows) in the Era5Land dataset. The gray vertical dashed lines indicate the dune orientation. The numbers at the top center give the percentage of time steps selected in each sub-range, as well as the percentage of them corresponding to the day (between 1000 UTC and 2200 UTC). The purple frame highlights the regime (small wind velocities, nocturnal summer wind) during which the wind data from both datasets differs. A similar figure can be obtained for the South Namib station (see Fig. 3).

690 air and water. The vertical coordinates are calculated as:

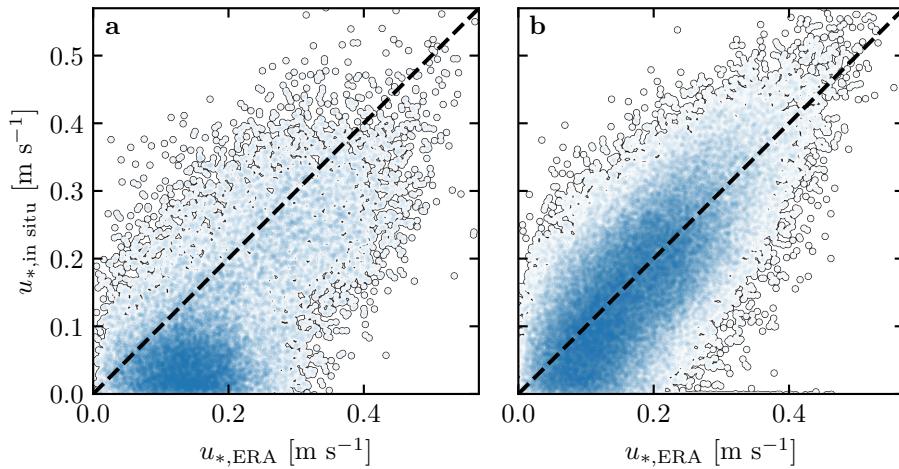
$$z = \frac{\phi R_t}{g R_t - \phi}, \quad (20)$$

691 where  $R_t = 6356766$  m is the average Earth radius, and  $g = 9.81$  m s<sup>-2</sup> the  
692 gravitational acceleration.

693 Example of obtained vertical profiles of the virtual potential temperature  
694 are shown in Fig. S9a. On each of them, an average is computed below the ABL  
695 depth directly given by the Era5 dataset, and a linear function is fitted above.  
696 Note that under the Boussinesq approximation, the temperature variations are

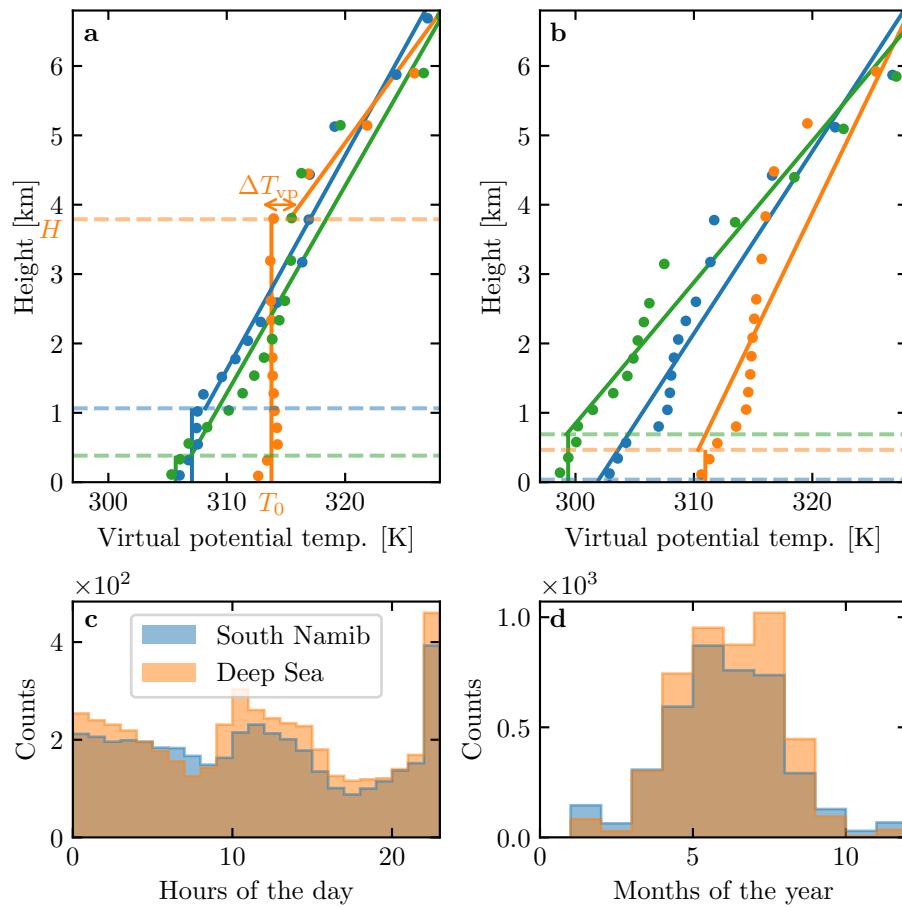


**Fig. S7** Statistical comparison of the wind orientation between the Era5Land dataset and the local measurements for the South Namib and Deep Sea stations, for different velocity ranges. **a:**  $u_{*,\text{ERA}} < 0.1 \text{ m s}^{-1}$ . **b:**  $0.1 < u_{*,\text{ERA}} \leq 0.25 \text{ m s}^{-1}$ . **c:**  $u_{*,\text{ERA}} \geq 0.25 \text{ m s}^{-1}$ . Note that the measured dune orientations are subtracted to the wind orientation, which allows to plot both stations on the same graph. Black dashed lines indicate locally measured orientations aligned with the dune crests (here  $0^\circ$ ,  $180^\circ$  and  $360^\circ$  – **a, b**), as well as the identity lines (**b, c**).

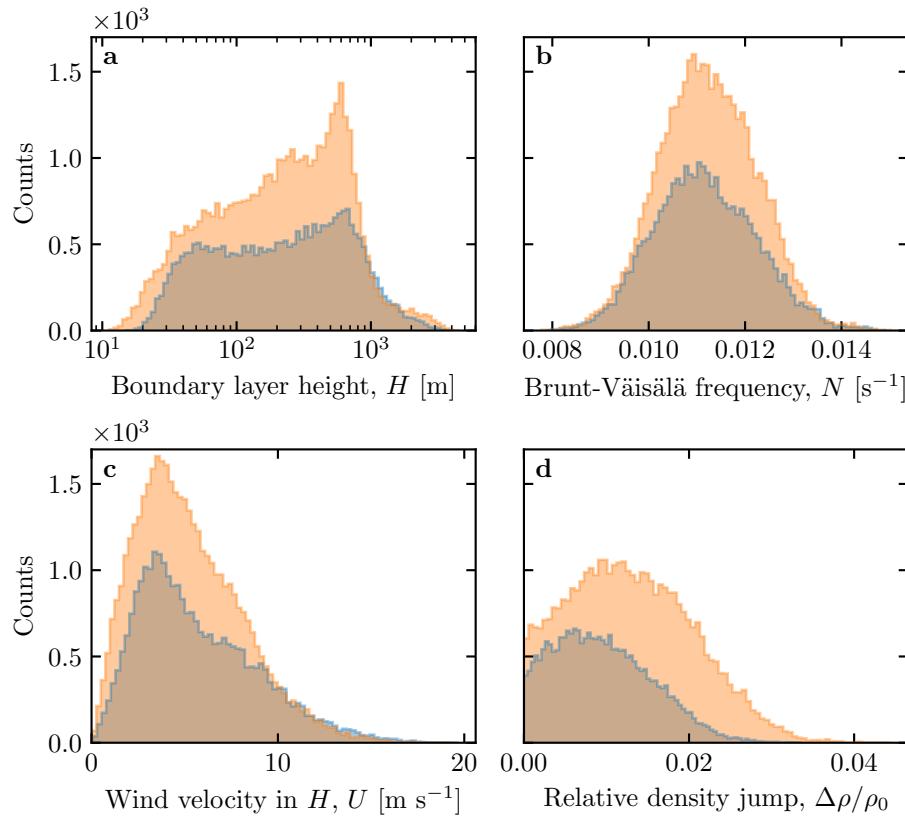


**Fig. S8** Statistical comparison of the wind velocity between the Era5Land dataset and the local measurements for the South Namib and Deep Sea stations. **a:** Nocturnal summer easterly wind. **b:** Diurnal southerly wind. Black dashed lines are identity lines. The angle ranges used to select diurnal and nocturnal summer winds are those taken in Fig. 3 and Fig. S6.

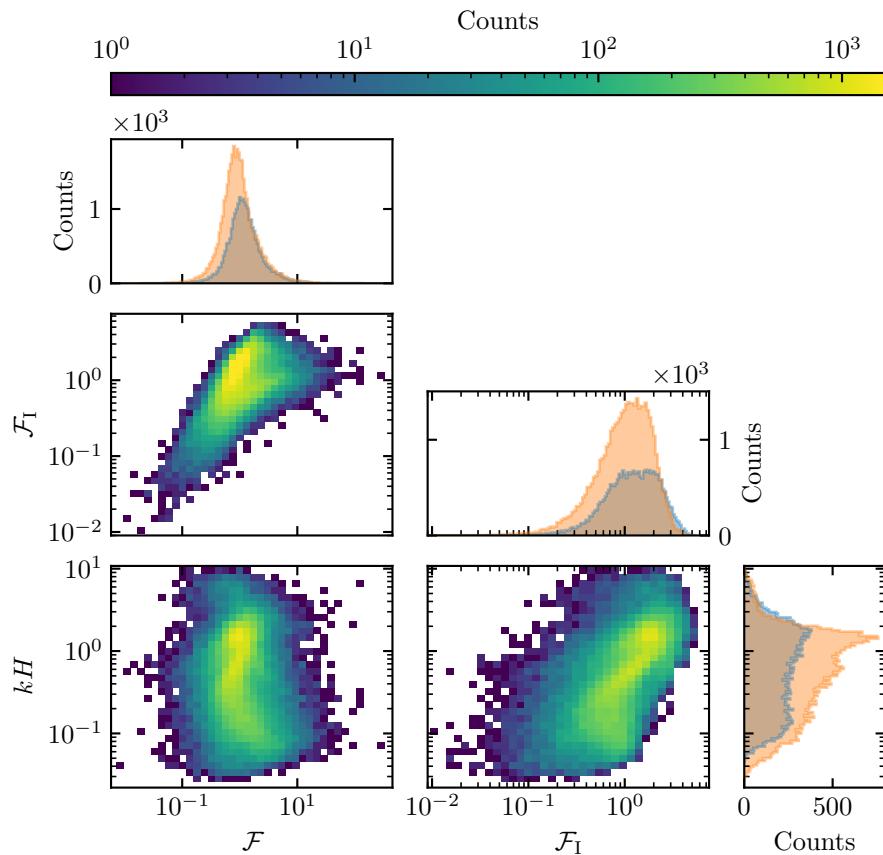
697 assumed to induce most of those of the density, leading to  $\Delta\rho/\rho_0 \simeq \Delta T_{\text{vp}}/T_0$   
 698 (see supplementary material of (Andreotti et al. 2009)). Note that we removed  
 699 some profiles that displayed a vertical structure that could not be approxi-  
 700 mated by the simple model used here, leading for example to negative values  
 701 of  $\Delta T_{\text{vp}}$  (see Fig. S9b). While these profiles dominantly occurs in winter, when  
 702 the two winds blows, they are evenly spread across the hours of the day, and  
 703 represent 12 % of our data only (see Fig. S9c-d). For these reasons, we are  
 704 confident this does not affect pour conclusions.



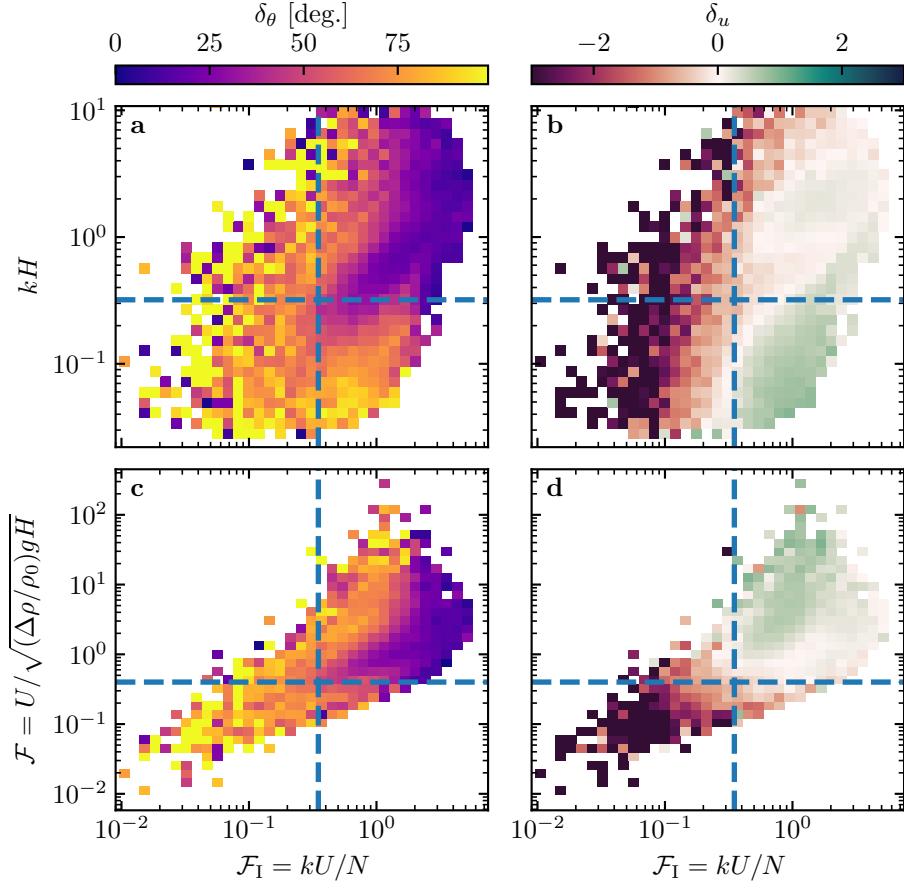
**Fig. S9** **a:** Vertical profiles of the virtual potential temperature at 3 different time steps (blue - 29/11/2012 - 1100 UTC, orange - 21/03/2017 - 1200 UTC, green - 21/03/2017 - 2000 UTC) at the South Namib station. Dots: data from the ERA5 reanalysis. Transparent dashed lines: boundary layer height given by the ERA5 reanalysis, calculated from the bulk Richardson number (Seidel et al. 2012). Plain lines: vertical (ABL) and linear (FA) fits to estimate the quantities in Fig. S10. **b:** Examples of ill-processed vertical profiles at 3 different time steps (blue - 2/12/2013 - 2300 UTC, orange - 20/03/2017 - 0000 UTC, green - 14/07/2017 - 1400 UTC) at the South Namib station. **c:** Hourly distribution of ill-processed vertical profiles. **d:** Monthly distribution of ill-processed vertical profiles.



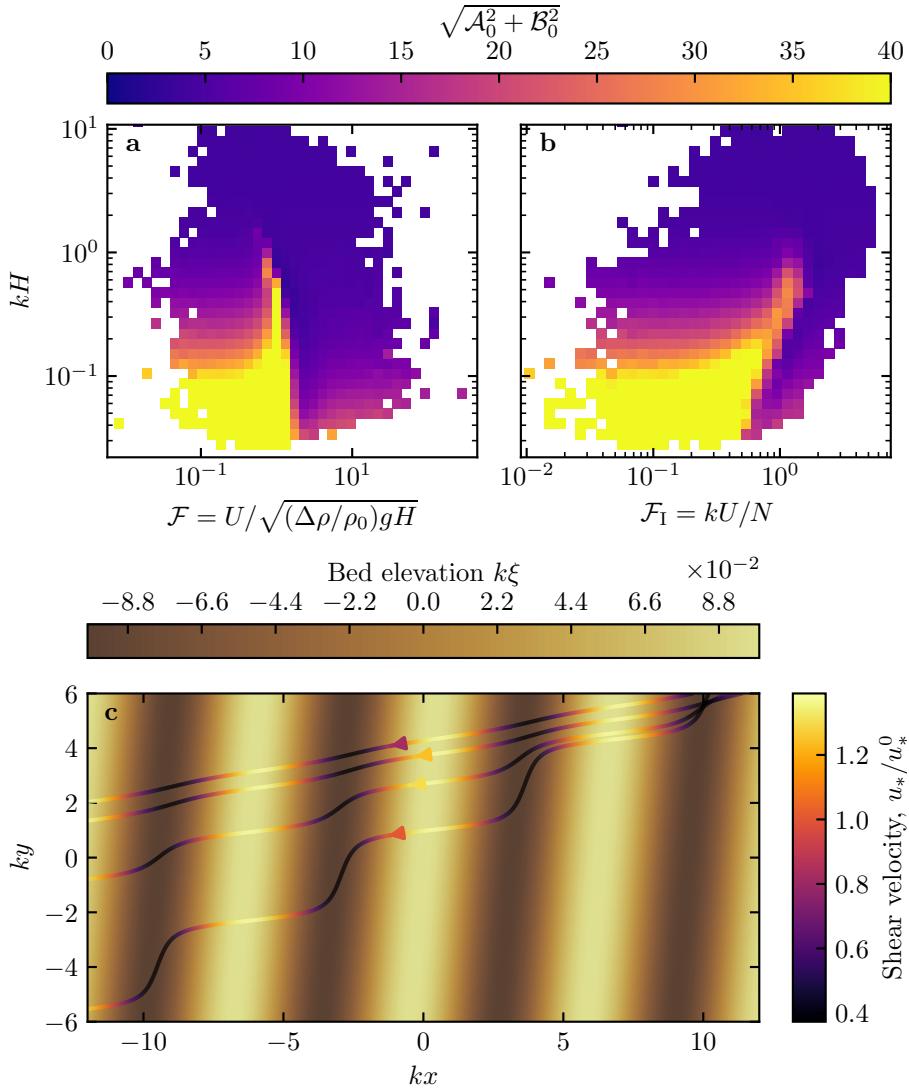
**Fig. S10** Distributions of the meteorological parameters resulting from the processing of the Era5-Land data for the South Namib (blue) and the Deep Sea (orange) stations.



**Fig. S11** Non-dimensional parameters distributions. For the marginal distributions, the orange correspond to the South Namib station, and the blue to the Deep Sea station.



**Fig. S12** Regime diagrams of the wind deviation  $\delta_\theta$  and relative attenuation/amplification  $\delta_u$  in the spaces  $(\mathcal{F}_I, kH)$  and  $(\mathcal{F}_I, \mathcal{F})$ , containing the data from both the Deep Sea and South Namib stations. Blue dashed lines empirically delimit the different regimes. The point density in each bin of the diagrams is shown in Fig. S11. The regime diagrams in the space  $(\mathcal{F}, kH)$  are shown in Fig. 5 of the main article.



**Fig. S13** Physical interpretation of the flow disturbance. (a) and (b) Magnitude of the disturbance induced by a sinusoidal topography calculated from the time series of the non-dimensional numbers presented in Figures 4 and 5 using the linear model of Andreotti et al. (2009). (c) Shear velocity streamlines represented in the case of the Deep Sea station, for increasing values of  $\sqrt{\mathcal{A}_0^2 + \mathcal{B}_0^2}$ . From the upper to the lower streamline, values of  $(kH, Fr_{\text{surface}}, Fr_{\text{internal}}, \mathcal{A}_0, \mathcal{B}_0, \sqrt{\mathcal{A}_0^2 + \mathcal{B}_0^2})$  are  $(1.9, 0.6, 1.5, 3.4, 1.0, 3.5)$ ,  $(1.5, 0.3, 0.4, 4.8, 1.4, 5.0)$ ,  $(0.1, 3.5, 1.0, 8.6, 0.1, 8.6)$ ,  $(0.5, 0.05, 0.04, 9.6, 2.5, 9.9)$ .