

# Electric Fields

## 22.1 THE ELECTRIC FIELD

### Learning Objectives

After reading this module, you should be able to . . .

- 22.1.1** Identify that at every point in the space surrounding a charged particle, the particle sets up an electric field  $\vec{E}$ , which is a vector quantity and thus has both magnitude and direction.
- 22.1.2** Identify how an electric field  $\vec{E}$  can be used to explain how a charged particle can exert an electrostatic force  $\vec{F}$  on a second charged particle

### Key Ideas

- A charged particle sets up an electric field (a vector quantity) in the surrounding space. If a second charged particle is located in that space, an electrostatic force acts on it due to the magnitude and direction of the field at its location.
- The electric field  $\vec{E}$  at any point is defined in terms of the electrostatic force  $\vec{F}$  that would be exerted on a positive test charge  $q_0$  placed there:

$$\vec{E} = \frac{\vec{F}}{q_0}$$

even though there is no contact between the particles.

- 22.1.3** Explain how a small positive test charge is used (in principle) to measure the electric field at any given point.
- 22.1.4** Explain electric field lines, including where they originate and terminate and what their spacing represents.

- Electric field lines help us visualize the direction and magnitude of electric fields. The electric field vector at any point is tangent to the field line through that point. The density of field lines in that region is proportional to the magnitude of the electric field there. Thus, closer field lines represent a stronger field.
- Electric field lines originate on positive charges and terminate on negative charges. So, a field line extending from a positive charge must end on a negative charge.

### What Is Physics?

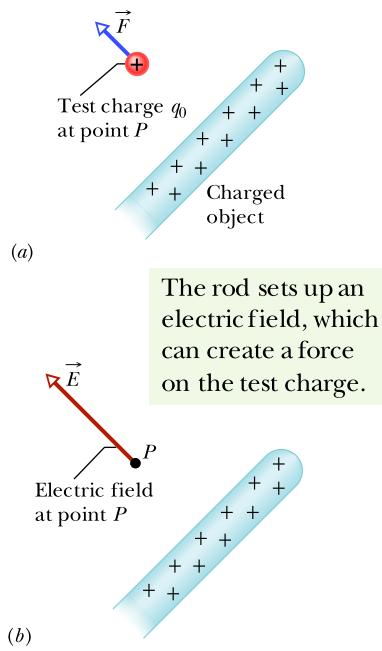
Figure 22.1.1 shows two positively charged particles. From the preceding chapter we know that an electrostatic force acts on particle 1 due to the presence of particle 2. We also know the force direction and, given some data, we can calculate the force magnitude. However, here is a leftover nagging question. How does particle 1 “know” of the presence of particle 2? That is, since the particles do not touch, how can particle 2 push on particle 1—how can there be such an *action at a distance*?

One purpose of physics is to record observations about our world, such as the magnitude and direction of the push on particle 1. Another purpose is to provide an explanation of what is recorded. Our purpose in this chapter is to provide such an explanation to this nagging question about electric force at a distance.

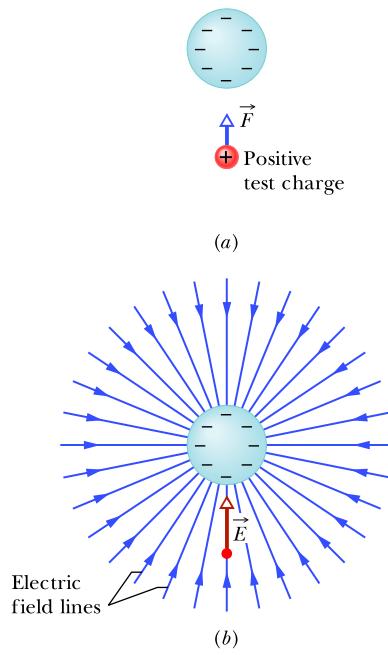
The explanation that we shall examine here is this: Particle 2 sets up an **electric field** at all points in the surrounding space, even if the space is a vacuum. If we place particle 1 at any point in that space, particle 1 knows of the presence of particle 2 because it is affected by the electric field particle 2 has already set up



**Figure 22.1.1** How does charged particle 2 push on charged particle 1 when they have no contact?



**Figure 22.1.2** (a) A positive test charge  $q_0$  placed at point  $P$  near a charged object. An electrostatic force  $\vec{F}$  acts on the test charge. (b) The electric field  $\vec{E}$  at point  $P$  produced by the charged object.



**Figure 22.1.3** (a) The electrostatic force  $\vec{F}$  acting on a positive test charge near a sphere of uniform negative charge. (b) The electric field vector  $\vec{E}$  at the location of the test charge, and the electric field lines in the space near the sphere. The field lines extend toward the negatively charged sphere. (They originate on distant positive charges.)

at that point. Thus, particle 2 pushes on particle 1 not by touching it as you would push on a coffee mug by making contact. Instead, particle 2 pushes by means of the electric field it has set up.

Our goals in this chapter are to (1) define electric field, (2) discuss how to calculate it for various arrangements of charged particles and objects, and (3) discuss how an electric field can affect a charged particle (as in making it move).

## The Electric Field

A lot of different fields are used in science and engineering. For example, a *temperature field* for an auditorium is the distribution of temperatures we would find by measuring the temperature at many points within the auditorium. Similarly, we could define a *pressure field* in a swimming pool. Such fields are examples of *scalar fields* because temperature and pressure are scalar quantities, having only magnitudes and not directions.

In contrast, an electric field is a *vector field* because it is responsible for conveying the information for a force, which involves both magnitude and direction. This field consists of a distribution of electric field vectors  $\vec{E}$ , one for each point in the space around a charged object. In principle, we can define  $\vec{E}$  at some point near the charged object, such as point  $P$  in Fig. 22.1.2a, with this procedure: At  $P$ , we place a particle with a small positive charge  $q_0$ , called a *test charge* because we use it to test the field. (We want the charge to be small so that it does not disturb the object's charge distribution.) We then measure the electrostatic force  $\vec{F}$  that acts on the test charge. The electric field at that point is then

$$\vec{E} = \frac{\vec{F}}{q_0} \quad (\text{electric field}). \quad (22.1.1)$$

Because the test charge is positive, the two vectors in Eq. 22.1.1 are in the same direction, so the direction of  $\vec{E}$  is the direction we measure for  $\vec{F}$ . The magnitude of  $\vec{E}$  at point  $P$  is  $F/q_0$ . As shown in Fig. 22.1.2b, we always represent an electric field with an arrow with its tail anchored on the point where the measurement is made. (This may sound trivial, but drawing the vectors any other way usually results in errors. Also, another common error is to mix up the terms *force* and *field* because they both start with the letter f. Electric force is a push or pull. Electric field is an abstract property set up by a charged object.) From Eq. 22.1.1, we see that the SI unit for the electric field is the newton per coulomb (N/C).

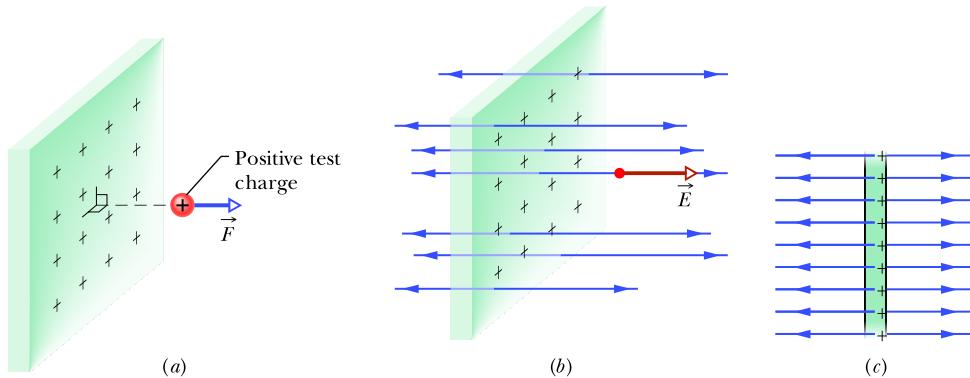
We can shift the test charge around to various other points, to measure the electric fields there, so that we can figure out the distribution of the electric field set up by the charged object. That field exists independent of the test charge. It is something that a charged object sets up in the surrounding space (even vacuum), independent of whether we happen to come along to measure it.

For the next several modules, we determine the field around charged particles and various charged objects. First, however, let's examine a way of visualizing electric fields.

## Electric Field Lines

Look at the space in the room around you. Can you visualize a field of vectors throughout that space—vectors with different magnitudes and directions? As impossible as that seems, Michael Faraday, who introduced the idea of electric fields in the 19th century, found a way. He envisioned lines, now called **electric field lines**, in the space around any given charged particle or object.

Figure 22.1.3 gives an example in which a sphere is uniformly covered with negative charge. If we place a positive test charge at any point near the sphere



**Figure 22.1.4** (a) The force on a positive test charge near a very large, nonconducting sheet with uniform positive charge on one side. (b) The electric field vector  $\vec{E}$  at the test charge's location, and the nearby electric field lines, extending away from the sheet. (c) Side view.

(Fig. 22.1.3a), we find that an electrostatic force pulls on it toward the center of the sphere. Thus at every point around the sphere, an electric field vector points radially inward toward the sphere. We can represent this electric field with electric field lines as in Fig. 22.1.3b. At any point, such as the one shown, the direction of the field line through the point matches the direction of the electric vector at that point.

The rules for drawing electric fields lines are these: (1) At any point, the electric field vector must be tangent to the electric field line through that point and in the same direction. (This is easy to see in Fig. 22.1.3 where the lines are straight, but we'll see some curved lines soon.) (2) In a plane perpendicular to the field lines, the relative density of the lines represents the relative magnitude of the field there, with greater density for greater magnitude.

If the sphere in Fig. 22.1.3 were uniformly covered with positive charge, the electric field vectors at all points around it would be radially outward and thus so would the electric field lines. So, we have the following rule:



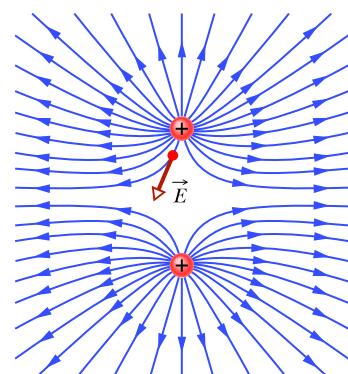
Electric field lines extend away from positive charge (where they originate) and toward negative charge (where they terminate).

In Fig. 22.1.3b, they originate on distant positive charges that are not shown.

For another example, Fig. 22.1.4a shows part of an infinitely large, nonconducting *sheet* (or plane) with a uniform distribution of positive charge on one side. If we place a positive test charge at any point near the sheet (on either side), we find that the electrostatic force on the particle is outward and perpendicular to the sheet. The perpendicular orientation is reasonable because any force component that is, say, upward is balanced out by an equal component that is downward. That leaves only outward, and thus the electric field vectors and the electric field lines must also be outward and perpendicular to the sheet, as shown in Figs. 22.1.4b and c.

Because the charge on the sheet is uniform, the field vectors and the field lines are also. Such a field is a *uniform electric field*, meaning that the electric field has the same magnitude and direction at every point within the field. (This is a lot easier to work with than a *nonuniform field*, where there is variation from point to point.) Of course, there is no such thing as an infinitely large sheet. That is just a way of saying that we are measuring the field at points close to the sheet relative to the size of the sheet and that we are not near an edge.

Figure 22.1.5 shows the field lines for two particles with equal positive charge. Now the field lines are curved, but the rules still hold: (1) The electric



**Figure 22.1.5** Field lines for two particles with equal positive charge. Doesn't the pattern itself suggest that the particles repel each other?

field vector at any given point must be tangent to the field line at that point and in the same direction, as shown for one vector, and (2) a closer spacing means a larger field magnitude. To imagine the full three-dimensional pattern of field lines around the particles, mentally rotate the pattern in Fig. 22.1.5 around the *axis of symmetry*, which is a vertical line through both particles.

### Checkpoint 22.1.1

Electric field lines extend across a lab experiment, from a charged plate on the right to a charged plate on the left. Is the left plate positively charged or negatively charged?

## 22.2 THE ELECTRIC FIELD DUE TO A CHARGED PARTICLE

### Learning Objectives

After reading this module, you should be able to . . .

**22.2.1** In a sketch, draw a charged particle, indicate its sign, pick a nearby point, and then draw the electric field vector  $\vec{E}$  at that point, with its tail anchored on the point.

**22.2.2** For a given point in the electric field of a charged particle, identify the direction of the field vector  $\vec{E}$  when the particle is positively charged and when it is negatively charged.

**22.2.3** For a given point in the electric field of a charged particle, apply the relationship between

the field magnitude  $E$ , the charge magnitude  $|q|$ , and the distance  $r$  between the point and the particle.

**22.2.4** Identify that the equation given here for the magnitude of an electric field applies only to a particle, not an extended object.

**22.2.5** If more than one electric field is set up at a point, draw each electric field vector and then find the net electric field by adding the individual electric fields as vectors (not as scalars).

### Key Ideas

- The magnitude of the electric field  $\vec{E}$  set up by a particle with charge  $q$  at distance  $r$  from the particle is

$$E = \frac{1}{4\pi\epsilon_0} \frac{|q|}{r^2}$$

- The electric field vectors set up by a positively charged particle all point directly away from the

particle. Those set up by a negatively charged particle all point directly toward the particle.

- If more than one charged particle sets up an electric field at a point, the net electric field is the *vector sum* of the individual electric fields—electric fields obey the superposition principle.

## The Electric Field Due to a Point Charge

To find the electric field due to a charged particle (often called a *point charge*), we place a positive test charge  $q_0$  at any point near the particle, at distance  $r$ . From Coulomb's law (Eq. 21.1.4), the force on the test charge due to the particle with charge  $q$  is

$$\vec{F} = \frac{1}{4\pi\epsilon_0} \frac{qq_0}{r^2} \hat{r}$$

As previously, the direction of  $\vec{F}$  is directly away from the particle if  $q$  is positive (because  $q_0$  is positive) and directly toward it if  $q$  is negative. From Eq. 22.1.1, we can now write the electric field set up by the particle (at the location of the test charge) as

$$\vec{E} = \frac{\vec{F}}{q_0} = \frac{1}{4\pi\epsilon_0 r^2} \frac{q}{r^2} \hat{r} \quad (\text{charged particle}). \quad (22.2.1)$$

Let's think through the directions again. The direction of  $\vec{E}$  matches that of the force on the positive test charge: directly away from the point charge if  $q$  is positive and directly toward it if  $q$  is negative.

So, if given another charged particle, we can immediately determine the directions of the electric field vectors near it by just looking at the sign of the charge  $q$ . We can find the magnitude at any given distance  $r$  by converting Eq. 22.2.1 to a magnitude form:

$$E = \frac{1}{4\pi\epsilon_0} \frac{|q|}{r^2} \quad (\text{charged particle}). \quad (22.2.2)$$

We write  $|q|$  to avoid the danger of getting a negative  $E$  when  $q$  is negative, and then thinking the negative sign has something to do with direction. Equation 22.2.2 gives magnitude  $E$  only. We must think about the direction separately.

Figure 22.2.1 gives a number of electric field vectors at points around a positively charged particle, but be careful. Each vector represents the vector quantity at the point where the tail of the arrow is anchored. The vector is not something that stretches from a "here" to a "there" as with a displacement vector.

In general, if several electric fields are set up at a given point by several charged particles, we can find the net field by placing a positive test particle at the point and then writing out the force acting on it due to each particle, such as  $\vec{F}_{01}$  due to particle 1. Forces obey the principle of superposition, so we just add the forces as vectors:

$$\vec{F}_0 = \vec{F}_{01} + \vec{F}_{02} + \cdots + \vec{F}_{0n}.$$

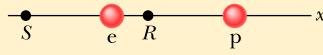
To change over to electric field, we repeatedly use Eq. 22.1.1 for each of the individual forces:

$$\begin{aligned} \vec{E} &= \frac{\vec{F}_0}{q_0} = \frac{\vec{F}_{01}}{q_0} + \frac{\vec{F}_{02}}{q_0} + \cdots + \frac{\vec{F}_{0n}}{q_0} \\ &= \vec{E}_1 + \vec{E}_2 + \cdots + \vec{E}_n. \end{aligned} \quad (22.2.3)$$

This tells us that electric fields also obey the principle of superposition. If you want the net electric field at a given point due to several particles, find the electric field due to each particle (such as  $\vec{E}_1$  due to particle 1) and then sum the fields as vectors. (As with electrostatic forces, you cannot just willy-nilly add up the magnitudes.) This addition of fields is the subject of many of the homework problems.

### Checkpoint 22.2.1

The figure here shows a proton  $p$  and an electron  $e$  on an  $x$  axis. What is the direction of the electric field due to the electron at (a) point  $S$  and (b) point  $R$ ? What is the direction of the net electric field at (c) point  $R$  and (d) point  $S$ ?



### Sample Problem 22.2.1 Net electric field due to three charged particles

Figure 22.2.2a shows three particles with charges  $q_1 = +2Q$ ,  $q_2 = -2Q$ , and  $q_3 = -4Q$ , each a distance  $d$  from the origin. What net electric field  $\vec{E}$  is produced at the origin?

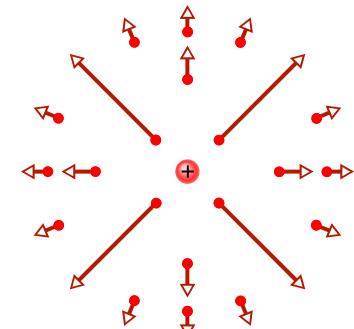
#### KEY IDEA

Charges  $q_1$ ,  $q_2$ , and  $q_3$  produce electric field vectors  $\vec{E}_1$ ,  $\vec{E}_2$ , and  $\vec{E}_3$ , respectively, at the origin, and the net electric field is the vector sum  $\vec{E} = \vec{E}_1 + \vec{E}_2 + \vec{E}_3$ . To find this sum,

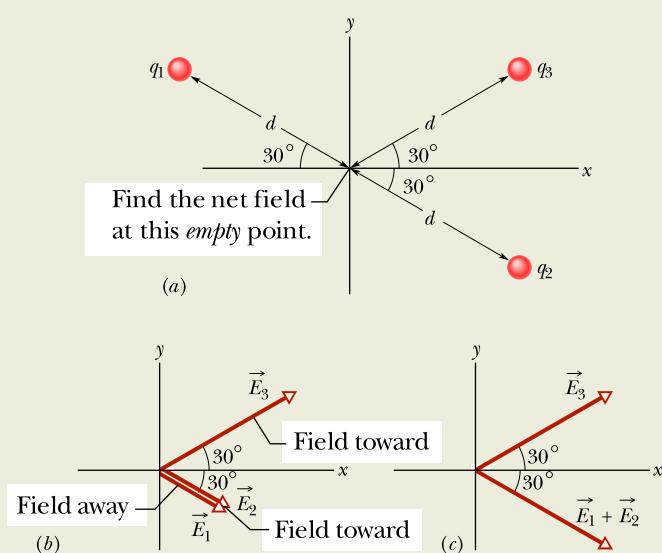
we first must find the magnitudes and orientations of the three field vectors.

**Magnitudes and directions:** To find the magnitude of  $\vec{E}_1$ , which is due to  $q_1$ , we use Eq. 22.2.2, substituting  $d$  for  $r$  and  $2Q$  for  $q$  and obtaining

$$E_1 = \frac{1}{4\pi\epsilon_0} \frac{2Q}{d^2}.$$



**Figure 22.2.1** The electric field vectors at various points around a positive point charge.



**Figure 22.2.2** (a) Three particles with charges  $q_1$ ,  $q_2$ , and  $q_3$  are at the same distance  $d$  from the origin. (b) The electric field vectors  $\vec{E}_1$ ,  $\vec{E}_2$ , and  $\vec{E}_3$ , at the origin due to the three particles. (c) The electric field vector  $\vec{E}_3$  and the vector sum  $\vec{E}_1 + \vec{E}_2$  at the origin.

Similarly, we find the magnitudes of  $\vec{E}_2$  and  $\vec{E}_3$  to be

$$E_2 = \frac{1}{4\pi\epsilon_0} \frac{2Q}{d^2} \quad \text{and} \quad E_3 = \frac{1}{4\pi\epsilon_0} \frac{4Q}{d^2}.$$

We next must find the orientations of the three electric field vectors at the origin. Because  $q_1$  is a positive charge, the field vector it produces points directly *away* from it, and because  $q_2$  and  $q_3$  are both negative, the field vectors they produce point directly *toward* each

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## 22.3 THE ELECTRIC FIELD DUE TO A DIPOLE

### Learning Objectives

After reading this module, you should be able to . . .

**22.3.1** Draw an electric dipole, identifying the charges (sizes and signs), dipole axis, and direction of the electric dipole moment.

**22.3.2** Identify the direction of the electric field at any given point along the dipole axis, including between the charges.

**22.3.3** Outline how the equation for the electric field due to an electric dipole is derived from the equations for the electric field due to the individual charged particles that form the dipole.

**22.3.4** For a single charged particle and an electric dipole, compare the rate at which the electric field

of them. Thus, the three electric fields produced at the origin by the three charged particles are oriented as in Fig. 22.2.2b. (*Caution:* Note that we have placed the tails of the vectors at the point where the fields are to be evaluated; doing so decreases the chance of error. Error becomes very probable if the tails of the field vectors are placed on the particles creating the fields.)

**Adding the fields:** We can now add the fields vectorially just as we added force vectors in Chapter 21. However, here we can use symmetry to simplify the procedure. From Fig. 22.2.2b, we see that electric fields  $\vec{E}_1$  and  $\vec{E}_2$  have the same direction. Hence, their vector sum has that direction and has the magnitude

$$\begin{aligned} E_1 + E_2 &= \frac{1}{4\pi\epsilon_0} \frac{2Q}{d^2} + \frac{1}{4\pi\epsilon_0} \frac{2Q}{d^2} \\ &= \frac{1}{4\pi\epsilon_0} \frac{4Q}{d^2}, \end{aligned}$$

which happens to equal the magnitude of field  $\vec{E}_3$ .

We must now combine two vectors,  $\vec{E}_3$  and the vector sum  $\vec{E}_1 + \vec{E}_2$ , that have the same magnitude and that are oriented symmetrically about the  $x$  axis, as shown in Fig. 22.2.2c. From the symmetry of Fig. 22.2.2c, we realize that the equal  $y$  components of our two vectors cancel (one is upward and the other is downward) and the equal  $x$  components add (both are rightward). Thus, the net electric field  $\vec{E}$  at the origin is in the positive direction of the  $x$  axis and has the magnitude

$$\begin{aligned} E &= 2E_{3x} = 2E_3 \cos 30^\circ \\ &= (2) \frac{1}{4\pi\epsilon_0} \frac{4Q}{d^2} (0.866) = \frac{6.93Q}{4\pi\epsilon_0 d^2}. \quad (\text{Answer}) \end{aligned}$$

magnitude decreases with increase in distance. That is, identify which drops off faster.

**22.3.5** For an electric dipole, apply the relationship between the magnitude  $p$  of the dipole moment, the separation  $d$  between the charges, and the magnitude  $q$  of either of the charges.

**22.3.6** For any distant point along a dipole axis, apply the relationship between the electric field magnitude  $E$ , the distance  $z$  from the center of the dipole, and either the dipole moment magnitude  $p$  or the product of charge magnitude  $q$  and charge separation  $d$ .

## Key Ideas

- An electric dipole consists of two particles with charges of equal magnitude  $q$  but opposite signs, separated by a small distance  $d$ .
- The electric dipole moment  $\vec{p}$  has magnitude  $qd$  and points from the negative charge to the positive charge.
- The magnitude of the electric field set up by an electric dipole at a distant point on the dipole axis (which runs through both particles) can be written in terms of either the product  $qd$  or the magnitude  $p$  of the dipole moment:

$$E = \frac{1}{2\pi\epsilon_0} \frac{qd}{z^3} = \frac{1}{2\pi\epsilon_0} \frac{p}{z^3},$$

where  $z$  is the distance between the point and the center of the dipole.

- Because of the  $1/z^3$  dependence, the field magnitude of an electric dipole decreases more rapidly with distance than the field magnitude of either of the individual charges forming the dipole, which depends on  $1/r^2$ .

## The Electric Field Due to an Electric Dipole

Figure 22.3.1 shows the pattern of electric field lines for two particles that have the same charge magnitude  $q$  but opposite signs, a very common and important arrangement known as an **electric dipole**. The particles are separated by distance  $d$  and lie along the *dipole axis*, an axis of symmetry around which you can imagine rotating the pattern in Fig. 22.3.1. Let's label that axis as a  $z$  axis. Here we restrict our interest to the magnitude and direction of the electric field  $\vec{E}$  at an arbitrary point  $P$  along the dipole axis, at distance  $z$  from the dipole's midpoint.

Figure 22.3.2a shows the electric fields set up at  $P$  by each particle. The nearer particle with charge  $+q$  sets up field  $E_{(+)}$  in the positive direction of the  $z$  axis (directly away from the particle). The farther particle with charge  $-q$  sets up a smaller field  $E_{(-)}$  in the negative direction (directly toward the particle). We want the net field at  $P$ , as given by Eq. 22.2.3. However, because the field vectors are along the same axis, let's simply indicate the vector directions with plus and minus signs, as we commonly do with forces along a single axis. Then we can write the magnitude of the net field at  $P$  as

$$\begin{aligned} E &= E_{(+)} - E_{(-)} \\ &= \frac{1}{4\pi\epsilon_0 r_{(+)}^2} q - \frac{1}{4\pi\epsilon_0 r_{(-)}^2} q \\ &= \frac{q}{4\pi\epsilon_0(z - \frac{1}{2}d)^2} - \frac{q}{4\pi\epsilon_0(z + \frac{1}{2}d)^2} \end{aligned} \quad (22.3.1)$$

After a little algebra, we can rewrite this equation as

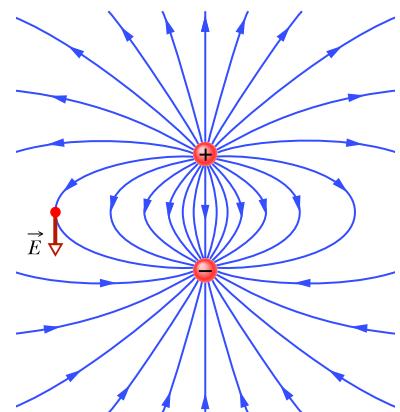
$$E = \frac{q}{4\pi\epsilon_0 z^2} \left( \frac{1}{\left(1 - \frac{d}{2z}\right)^2} - \frac{1}{\left(1 + \frac{d}{2z}\right)^2} \right). \quad (22.3.2)$$

After forming a common denominator and multiplying its terms, we come to

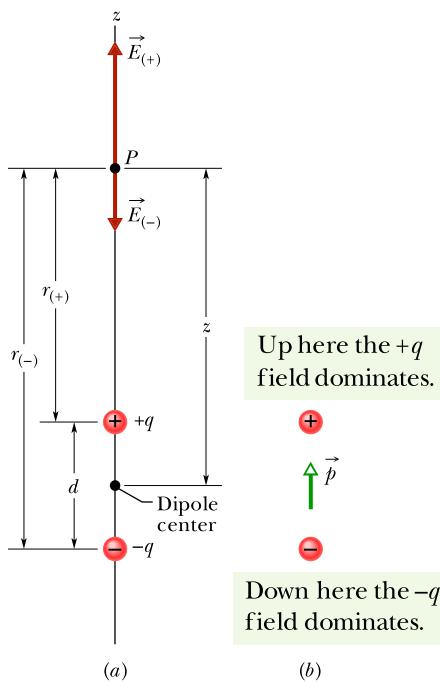
$$E = \frac{q}{4\pi\epsilon_0 z^2} \left( \frac{2d/z}{\left(1 - \left(\frac{d}{2z}\right)^2\right)^2} \right) = \frac{q}{2\pi\epsilon_0 z^3} \left( \frac{d}{\left(1 - \left(\frac{d}{2z}\right)^2\right)^2} \right) \quad (22.3.3)$$

We are usually interested in the electrical effect of a dipole only at distances that are large compared with the dimensions of the dipole—that is, at distances such that  $z \gg d$ . At such large distances, we have  $d/2z \ll 1$  in Eq. 22.3.3. Thus, in our approximation, we can neglect the  $d/2z$  term in the denominator, which leaves us with

$$E = \frac{1}{2\pi\epsilon_0} \frac{qd}{z^3}. \quad (22.3.4)$$



**Figure 22.3.1** The pattern of electric field lines around an electric dipole, with an electric field vector  $\vec{E}$  shown at one point (tangent to the field line through that point).



**Figure 22.3.2** (a) An electric dipole. The electric field vectors  $\vec{E}_{(+)}$  and  $\vec{E}_{(-)}$  at point  $P$  on the dipole axis result from the dipole's two charges. Point  $P$  is at distances  $r_{(+)}$  and  $r_{(-)}$  from the individual charges that make up the dipole. (b) The dipole moment  $\vec{p}$  of the dipole points from the negative charge to the positive charge.

The product  $qd$ , which involves the two intrinsic properties  $q$  and  $d$  of the dipole, is the magnitude  $p$  of a vector quantity known as the **electric dipole moment**  $\vec{p}$  of the dipole. (The unit of  $\vec{p}$  is the coulomb-meter.) Thus, we can write Eq. 22.3.4 as

$$E = \frac{1}{2\pi\epsilon_0} \frac{p}{z^3} \quad (\text{electric dipole}). \quad (22.3.5)$$

The direction of  $\vec{p}$  is taken to be from the negative to the positive end of the dipole, as indicated in Fig. 22.3.2b. We can use the direction of  $\vec{p}$  to specify the orientation of a dipole.

Equation 22.3.5 shows that, if we measure the electric field of a dipole only at distant points, we can never find  $q$  and  $d$  separately; instead, we can find only their product. The field at distant points would be unchanged if, for example,  $q$  were doubled and  $d$  simultaneously halved. Although Eq. 22.3.5 holds only for distant points along the dipole axis, it turns out that  $E$  for a dipole varies as  $1/r^3$  for all distant points, regardless of whether they lie on the dipole axis; here  $r$  is the distance between the point in question and the dipole center.

Inspection of Fig. 22.3.2 and of the field lines in Fig. 22.3.1 shows that the direction of  $\vec{E}$  for distant points on the dipole axis is always the direction of the dipole moment vector  $\vec{p}$ . This is true whether point  $P$  in Fig. 22.3.2a is on the upper or the lower part of the dipole axis.

Inspection of Eq. 22.3.2 shows that if you double the distance of a point from a dipole, the electric field at the point drops by a factor of 8. If you double the distance from a single point charge, however (see Eq. 22.2.2), the electric field drops only by a factor of 4. Thus the electric field of a dipole decreases more rapidly with distance than does the electric field of a single charge. The physical reason for this rapid decrease in electric field for a dipole is that from distant points a dipole looks like two particles that almost—but not quite—coincide. Thus, because they have charges of equal magnitude but opposite signs, their electric fields at distant points almost—but not quite—cancel each other.

### Checkpoint 22.3.1

At a distant point on the dipole axis, how does the direction of the field vector  $\vec{E}$  compare with the direction for the dipole moment vector  $\vec{p}$  for a point (a) above the dipole and (b) below the dipole?

### Sample Problem 22.3.1 Electric dipole and atmospheric sprites

Sprites (Fig. 22.3.3a) are huge flashes that occur far above a large thunderstorm. They were seen for decades by pilots flying at night, but they were so brief and dim that most pilots figured they were just illusions. Then in the 1990s sprites were captured on video. They are still not well understood but are believed to be produced when especially powerful lightning occurs between the ground and storm clouds, particularly when the lightning transfers a huge amount of negative charge  $-q$  from the ground to the base of the clouds (Fig. 22.3.3b).

Just after such a transfer, the ground has a complicated distribution of positive charge. However, we can model the electric field due to the charges in the clouds and the ground by assuming a vertical electric dipole that has charge  $-q$  at

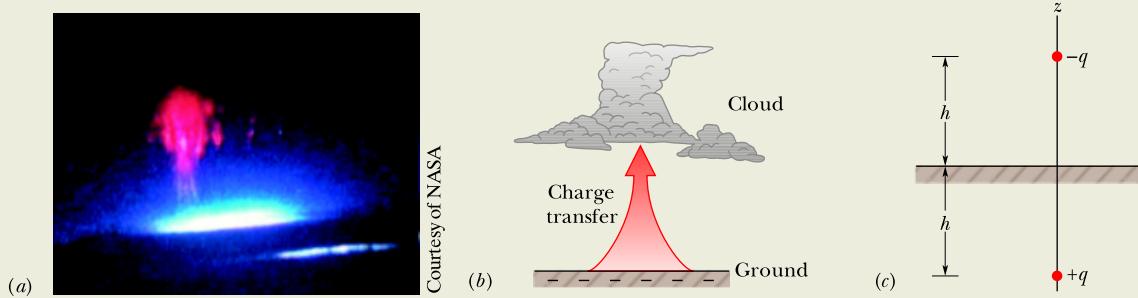
cloud height  $h$  and charge  $+q$  at below-ground depth  $h$  (Fig. 22.3.3c). If  $q = 200\text{C}$  and  $h = 6.0\text{ km}$ , what is the magnitude of the dipole's electric field at altitude  $z_1 = 30\text{ km}$  somewhat above the clouds and altitude  $z_2 = 60\text{ km}$  somewhat above the stratosphere? FCP

### KEY IDEA

We can approximate the magnitude  $E$  of an electric dipole's electric field on the dipole axis with Eq. 22.3.4.

**Calculations:** We write that equation as

$$E = \frac{1}{2\pi\epsilon_0} \frac{q(2h)}{z^3},$$



**Figure 22.3.3** (a) Photograph of a sprite. (b) Lightning in which a large amount of negative charge is transferred from ground to cloud base. (c) The cloud–ground system modeled as a vertical electric dipole.

where  $2h$  is the separation between  $-q$  and  $+q$  in Fig. 22.3.3c. For the electric field at altitude  $z_1 = 30$  km, we find

$$\begin{aligned} E &= \frac{1}{2\pi\epsilon_0} \frac{(200 \text{ C})(2)(6.0 \times 10^3 \text{ m})}{(30 \times 10^3 \text{ m})^3} \\ &= 1.6 \times 10^3 \text{ N/C.} \end{aligned} \quad (\text{Answer})$$

Similarly, for altitude  $z_2 = 60$  km, we find

$$E = 2.0 \times 10^2 \text{ N/C.} \quad (\text{Answer})$$

As we discuss in Module 22.6, when the magnitude of an electric field exceeds a certain critical value

$E_c$ , the field can pull electrons out of atoms (ionize the atoms), and then the freed electrons can run into other atoms, causing those atoms to emit light. The value of  $E_c$  depends on the density of the air in which the electric field exists. At altitude  $z_2 = 60$  km the density of the air is so low that  $E = 2.0 \times 10^2 \text{ N/C}$  exceeds  $E_c$ , and thus light is emitted by the atoms in the air. That light forms sprites. Lower down, just above the clouds at  $z_1 = 30$  km, the density of the air is much higher,  $E = 1.6 \times 10^3 \text{ N/C}$  does not exceed  $E_c$ , and no light is emitted. Hence, sprites occur only far above storm clouds.

**WileyPLUS** Additional examples, video, and practice available at *WileyPLUS*

## 22.4 THE ELECTRIC FIELD DUE TO A LINE OF CHARGE

### Learning Objectives

After reading this module, you should be able to . . .

**22.4.1** For a uniform distribution of charge, find the linear charge density  $\lambda$  for charge along a line, the surface charge density  $\sigma$  for charge on a surface, and the volume charge density  $\rho$  for charge in a volume.

**22.4.2** For charge that is distributed uniformly along a line, find the net electric field at a given point near

the line by splitting the distribution up into charge elements  $dq$  and then summing (by integration) the electric field vectors  $d\vec{E}$  set up at the point by each element.

**22.4.3** Explain how symmetry can be used to simplify the calculation of the electric field at a point near a line of uniformly distributed charge.

### Key Ideas

- The equation for the electric field set up by a particle does not apply to an extended object with charge (said to have a continuous charge distribution).

- To find the electric field of an extended object at a point, we first consider the electric field set up by a charge element  $dq$  in the object, where the element is small enough for us to apply the equation for a particle.

Then we sum, via integration, components of the electric fields  $d\vec{E}$  from all the charge elements.

- Because the individual electric fields  $d\vec{E}$  have different magnitudes and point in different directions, we first see if symmetry allows us to cancel out any of the components of the fields, to simplify the integration.

## The Electric Field Due to a Line of Charge

So far we have dealt with only charged particles, a single particle or a simple collection of them. We now turn to a much more challenging situation in which a thin (approximately one-dimensional) object such as a rod or ring is charged with a huge number of particles, more than we could ever even count. In the next module, we consider two-dimensional objects, such as a disk with charge spread over a surface. In the next chapter we tackle three-dimensional objects, such as a sphere with charge spread through a volume.

**Heads Up.** Many students consider this module to be the most difficult in the book for a variety of reasons. There are lots of steps to take, a lot of vector features to keep track of, and after all that, we set up and then solve an integral. The worst part, however, is that the procedure can be different for different arrangements of the charge. Here, as we focus on a particular arrangement (a charged ring), be aware of the general approach, so that you can tackle other arrangements in the homework (such as rods and partial circles).

Figure 22.4.1 shows a thin ring of radius  $R$  with a uniform distribution of positive charge along its circumference. It is made of plastic, which means that the charge is fixed in place. The ring is surrounded by a pattern of electric field lines, but here we restrict our interest to an arbitrary point  $P$  on the central axis (the axis through the ring's center and perpendicular to the plane of the ring), at distance  $z$  from the center point.

The charge of an extended object is often conveyed in terms of a charge density rather than the total charge. For a line of charge, we use the *linear charge density*  $\lambda$  (the charge per unit length), with the SI unit of coulomb per meter. Table 22.4.1 shows the other charge densities that we shall be using for charged surfaces and volumes.

**First Big Problem.** So far, we have an equation for the electric field of a particle. (We can combine the field of several particles as we did for the electric dipole to generate a special equation, but we are still basically using Eq. 22.2.2.) Now take a look at the ring in Fig. 22.4.1. That clearly is not a particle and so Eq. 22.2.2 does not apply. So what do we do?

The answer is to mentally divide the ring into differential elements of charge that are so small that we can treat them as though they *are* particles. Then we *can* apply Eq. 22.2.2.

**Second Big Problem.** We now know to apply Eq. 22.2.2 to each charge element  $dq$  (the front  $d$  emphasizes that the charge is very small) and can write an expression for its contribution of electric field  $d\vec{E}$  (the front  $d$  emphasizes that the contribution is very small). However, each such contributed field vector at  $P$  is in its own direction. How can we add them to get the net field at  $P$ ?

The answer is to split the vectors into components and then separately sum one set of components and then the other set. However, first we check to see if one set simply all cancels out. (Canceling out components saves lots of work.)

**Third Big Problem.** There is a huge number of  $dq$  elements in the ring and thus a huge number of  $d\vec{E}$  components to add up, even if we can cancel out one set of components. How can we add up more components than we could even count? The answer is to add them by means of integration.

**Do It.** Let's do all this (but again, be aware of the general procedure, not just the fine details). We arbitrarily pick the charge element shown in Fig. 22.4.1. Let  $ds$  be the arc length of that (or any other)  $dq$  element. Then in terms of the linear density  $\lambda$  (the charge per unit length), we have

$$dq = \lambda ds. \quad (22.4.1)$$

**An Element's Field.** This charge element sets up the differential electric field  $d\vec{E}$  at  $P$ , at distance  $r$  from the element, as shown in Fig. 22.4.1. (Yes, we are

introducing a new symbol that is not given in the problem statement, but soon we shall replace it with “legal symbols.”) Next we rewrite the field equation for a particle (Eq. 22.2.2) in terms of our new symbols  $dE$  and  $dq$ , but then we replace  $dq$  using Eq. 22.4.1. The field magnitude due to the charge element is

$$dE = \frac{1}{4\pi\epsilon_0} \frac{dq}{r^2} = \frac{1}{4\pi\epsilon_0} \frac{\lambda ds}{r^2}. \quad (22.4.2)$$

Notice that the illegal symbol  $r$  is the hypotenuse of the right triangle displayed in Fig. 22.4.1. Thus, we can replace  $r$  by rewriting Eq. 22.4.2 as

$$dE = \frac{1}{4\pi\epsilon_0} \frac{\lambda ds}{(z^2 + R^2)}. \quad (22.4.3)$$

Because every charge element has the same charge and the same distance from point  $P$ , Eq. 22.4.3 gives the field magnitude contributed by each of them. Figure 22.4.1 also tells us that each contributed  $d\vec{E}$  leans at angle  $\theta$  to the central axis (the  $z$  axis) and thus has components perpendicular and parallel to that axis.

**Cancelling Components.** Now comes the neat part, where we eliminate one set of those components. In Fig. 22.4.1, consider the charge element on the opposite side of the ring. It too contributes the field magnitude  $dE$  but the field vector leans at angle  $\theta$  in the opposite direction from the vector from our first charge element, as indicated in the side view of Fig. 22.4.2. Thus the two perpendicular components cancel. All around the ring, this cancellation occurs for every charge element and its *symmetric partner* on the opposite side of the ring. So we can neglect all the perpendicular components.

**Adding Components.** We have another big win here. All the remaining components are in the positive direction of the  $z$  axis, so we can just add them up as scalars. Thus we can already tell the direction of the net electric field at  $P$ : directly away from the ring. From Fig. 22.4.2, we see that the parallel components each have magnitude  $dE \cos \theta$ , but  $\theta$  is another illegal symbol. We can replace  $\cos \theta$  with legal symbols by again using the right triangle in Fig. 22.4.1 to write

$$\cos \theta = \frac{z}{r} = \frac{z}{(z^2 + R^2)^{1/2}}. \quad (22.4.4)$$

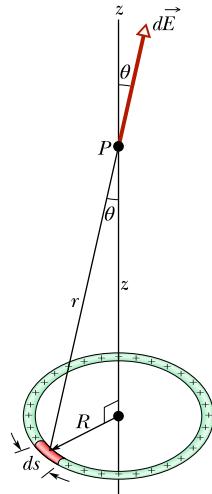
Multiplying Eq. 22.4.3 by Eq. 22.4.4 gives us the parallel field component from each charge element:

$$dE \cos \theta = \frac{1}{4\pi\epsilon_0} \frac{z\lambda}{(z^2 + R^2)^{3/2}} ds. \quad (22.4.5)$$

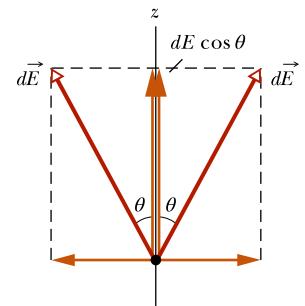
**Integrating.** Because we must sum a huge number of these components, each small, we set up an integral that moves along the ring, from element to element, from a starting point (call it  $s = 0$ ) through the full circumference ( $s = 2\pi R$ ). Only the quantity  $s$  varies as we go through the elements; the other symbols in Eq. 22.4.5 remain the same, so we move them outside the integral. We find

$$\begin{aligned} E &= \int dE \cos \theta = \frac{z\lambda}{4\pi\epsilon_0(z^2 + R^2)^{3/2}} \int_0^{2\pi R} ds \\ &= \frac{z\lambda(2\pi R)}{4\pi\epsilon_0(z^2 + R^2)^{3/2}} \end{aligned} \quad (22.4.6)$$

This is a fine answer, but we can also switch to the total charge by using  $\lambda = q/(2\pi R)$ :



**Figure 22.4.1** A ring of uniform positive charge. A differential element of charge occupies a length  $ds$  (greatly exaggerated for clarity). This element sets up an electric field  $d\vec{E}$  at point  $P$ .



**Figure 22.4.2** The electric fields set up at  $P$  by a charge element and its symmetric partner (on the opposite side of the ring). The components perpendicular to the  $z$  axis cancel; the parallel components add.

$$E = \frac{qz}{4\pi\epsilon_0(z^2 + R^2)^{3/2}} \quad (\text{charged ring}). \quad (22.4.7)$$

If the charge on the ring is negative, instead of positive as we have assumed, the magnitude of the field at  $P$  is still given by Eq. 22.4.7. However, the electric field vector then points toward the ring instead of away from it.

Let us check Eq. 22.4.7 for a point on the central axis that is so far away that  $z \gg R$ . For such a point, the expression  $z^2 + R^2$  in Eq. 22.4.7 can be approximated as  $z^2$ , and Eq. 22.4.7 becomes

$$E = \frac{1}{4\pi\epsilon_0} \frac{q}{z^2} \quad (\text{charged ring at large distance}). \quad (22.4.8)$$

This is a reasonable result because from a large distance, the ring “looks like” a point charge. If we replace  $z$  with  $r$  in Eq. 22.4.8, we indeed do have the magnitude of the electric field due to a point charge, as given by Eq. 22.2.2.

Let us next check Eq. 22.4.7 for a point at the center of the ring—that is, for  $z = 0$ . At that point, Eq. 22.4.7 tells us that  $E = 0$ . This is a reasonable result because if we were to place a test charge at the center of the ring, there would be no net electrostatic force acting on it; the force due to any element of the ring would be canceled by the force due to the element on the opposite side of the ring. By Eq. 22.1.1, if the force at the center of the ring were zero, the electric field there would also have to be zero.

### Sample Problem 22.4.1 Electric field of a charged circular rod

Figure 22.4.3a shows a plastic rod with a uniform charge  $-Q$ . It is bent in a  $120^\circ$  circular arc of radius  $r$  and symmetrically placed across an  $x$  axis with the origin at the center of curvature  $P$  of the rod. In terms of  $Q$  and  $r$ , what is the electric field  $\vec{E}$  due to the rod at point  $P$ ?

#### KEY IDEA

Because the rod has a continuous charge distribution, we must find an expression for the electric fields due to differential elements of the rod and then sum those fields via calculus.

**An element:** Consider a differential element having arc length  $ds$  and located at an angle  $\theta$  above the  $x$  axis (Figs. 22.4.3b and c). If we let  $\lambda$  represent the linear charge density of the rod, our element  $ds$  has a differential charge of magnitude

$$dq = \lambda ds. \quad (22.4.9)$$

**The element's field:** Our element produces a differential electric field  $d\vec{E}$  at point  $P$ , which is a distance  $r$  from the element. Treating the element as a point charge, we can rewrite Eq. 22.2.2 to express the magnitude of  $d\vec{E}$  as

$$dE = \frac{1}{4\pi\epsilon_0} \frac{dq}{r^2} = \frac{1}{4\pi\epsilon_0} \frac{\lambda ds}{r^2}. \quad (22.4.10)$$

The direction of  $d\vec{E}$  is toward  $ds$  because charge  $dq$  is negative.

**Symmetric partner:** Our element has a symmetrically located (mirror image) element  $ds'$  in the bottom half of the rod. The electric field  $d\vec{E}'$  set up at  $P$  by  $ds'$  also has the magnitude given by Eq. 22.4.10, but the field vector points toward  $ds'$  as shown in Fig. 22.4.3d. If we resolve the electric field vectors of  $ds$  and  $ds'$  into  $x$  and  $y$  components as shown in Figs. 22.4.3e and f, we see that their  $y$  components cancel (because they have equal magnitudes and are in opposite directions). We also see that their  $x$  components have equal magnitudes and are in the same direction.

**Summing:** Thus, to find the electric field set up by the rod, we need sum (via integration) only the  $x$  components of the differential electric fields set up by all the differential elements of the rod. From Fig. 22.4.3f and Eq. 22.4.10, we can write the component  $dE_x$  set up by  $ds$  as

$$dE_x = dE \cos \theta = \frac{1}{4\pi\epsilon_0} \frac{\lambda}{r^2} \cos \theta ds. \quad (22.4.11)$$

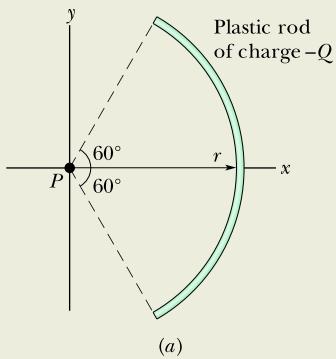
Equation 22.4.11 has two variables,  $\theta$  and  $s$ . Before we can integrate it, we must eliminate one variable. We do so by replacing  $ds$ , using the relation

$$ds = r d\theta,$$

in which  $d\theta$  is the angle at  $P$  that includes arc length  $ds$  (Fig. 22.4.3g). With this replacement, we can integrate Eq. 22.4.11 over the angle made by the rod at  $P$ , from  $\theta = -60^\circ$  to  $\theta = 60^\circ$ ; that will give us the field magnitude at  $P$ :

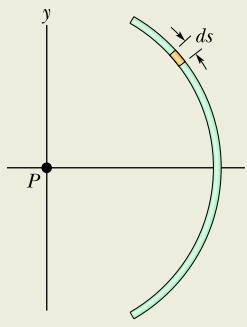


This negatively charged rod is obviously not a particle.



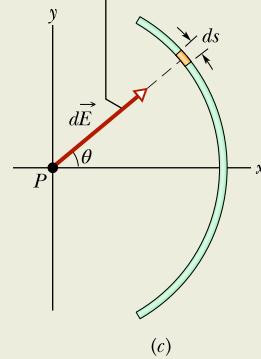
(a)

But we can treat this element as a particle.

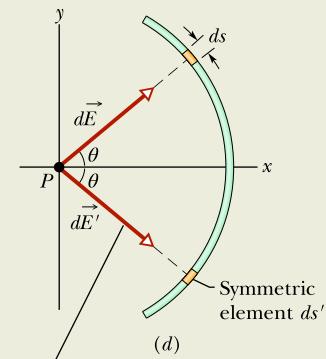


(b)

Here is the field the element creates.

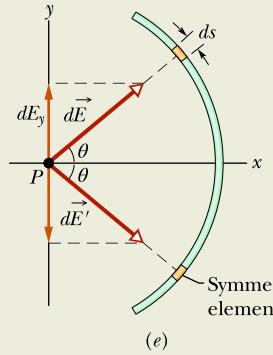


(c)



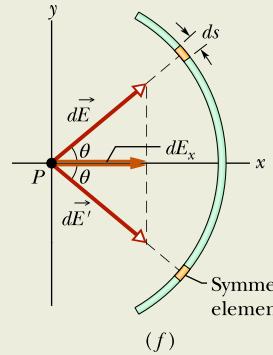
(d)

These y components just cancel, so neglect them.



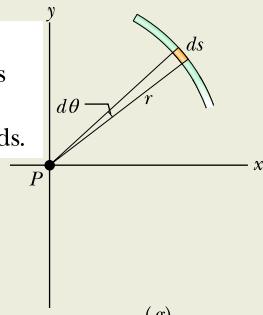
(e)

These x components add. Our job is to add all such components.



(f)

Here is the field created by the symmetric element, same size and angle.



(g)

**Figure 22.4.3** Available in WileyPLUS as an animation with voiceover. (a) A plastic rod of charge  $-Q$  is a circular section of radius  $r$  and central angle  $120^\circ$ ; point  $P$  is the center of curvature of the rod. (b)–(c) A differential element in the top half of the rod, at an angle  $\theta$  to the  $x$  axis and of arc length  $ds$ , sets up a differential electric field  $d\vec{E}$  at  $P$ . (d) An element  $ds'$ , symmetric to  $ds$  about the  $x$  axis, sets up a field  $d\vec{E}'$  at  $P$  with the same magnitude. (e)–(f) The field components. (g) Arc length  $ds$  makes an angle  $d\theta$  about point  $P$ .

$$\begin{aligned}
 E &= \int dE_x = \int_{-60^\circ}^{60^\circ} \frac{1}{4\pi\epsilon_0 r^2} \frac{\lambda}{r^2} \cos \theta r d\theta \\
 &= \frac{\lambda}{4\pi\epsilon_0 r} \int_{-60^\circ}^{60^\circ} \cos \theta d\theta = \frac{\lambda}{4\pi\epsilon_0 r} [\sin \theta]_{-60^\circ}^{60^\circ} \\
 &= \frac{\lambda}{4\pi\epsilon_0 r} [\sin 60^\circ - \sin(-60^\circ)] \\
 &= \frac{1.73\lambda}{4\pi\epsilon_0 r}.
 \end{aligned} \tag{22.4.12}$$

(If we had reversed the limits on the integration, we would have gotten the same result but with a minus sign. Since the integration gives only the magnitude of  $\vec{E}$ , we would then have discarded the minus sign.)

**Charge density:** To evaluate  $\lambda$ , we note that the full rod subtends an angle of  $120^\circ$  and so is one-third of a full

circle. Its arc length is then  $2\pi r/3$ , and its linear charge density must be

$$\lambda = \frac{\text{charge}}{\text{length}} = \frac{Q}{2\pi r/3} = \frac{0.477Q}{r}.$$

Substituting this into Eq. 22.4.12 and simplifying give us

$$\begin{aligned}
 E &= \frac{(1.73)(0.477Q)}{4\pi\epsilon_0 r^2} \\
 &= \frac{0.83Q}{4\pi\epsilon_0 r^2}.
 \end{aligned} \tag{Answer}$$

The direction of  $\vec{E}$  is toward the rod, along the axis of symmetry of the charge distribution. We can write  $\vec{E}$  in unit-vector notation as

$$\vec{E} = \frac{0.83Q\hat{z}}{4\pi\epsilon_0 r^2}.$$

### Problem-Solving Tactics A Field Guide for Lines of Charge

Here is a generic guide for finding the electric field  $\vec{E}$  produced at a point  $P$  by a line of uniform charge, either circular or straight. The general strategy is to pick out an element  $dq$  of the charge, find  $d\vec{E}$  due to that element, and integrate  $d\vec{E}$  over the entire line of charge.

**Step 1.** If the line of charge is circular, let  $ds$  be the arc length of an element of the distribution. If the line is straight, run an  $x$  axis along it and let  $dx$  be the length of an element. Mark the element on a sketch.

**Step 2.** Relate the charge  $dq$  of the element to the length of the element with either  $dq = \lambda ds$  or  $dq = \lambda dx$ . Consider  $dq$  and  $\lambda$  to be positive, even if the charge is actually negative. (The sign of the charge is used in the next step.)

**Step 3.** Express the field  $d\vec{E}$  produced at  $P$  by  $dq$  with Eq. 22.2.2, replacing  $q$  in that equation with either  $\lambda ds$  or  $\lambda dx$ . If the charge on the line is positive, then at  $P$  draw a vector  $d\vec{E}$  that points directly away from  $dq$ . If the charge is negative, draw the vector pointing directly toward  $dq$ .

**Step 4.** Always look for any symmetry in the situation. If  $P$  is on an axis of symmetry of the charge distribution, resolve the field  $d\vec{E}$  produced by  $dq$  into components that are perpendicular and parallel to the axis of symmetry. Then consider a second element  $dq'$  that is located symmetrically to  $dq$  about the line of symmetry. At  $P$  draw the vector  $d\vec{E}'$  that this symmetric element produces and resolve it into components. One of the components produced by  $dq$  is a *cancelling component*; it is canceled by the corresponding component produced by  $dq'$  and needs no further attention. The other component produced by  $dq$  is an *adding component*; it adds to the corresponding component produced by  $dq'$ . Add the adding components of all the elements via integration.

**Step 5.** Here are four general types of uniform charge distributions, with strategies for the integral of step 4.

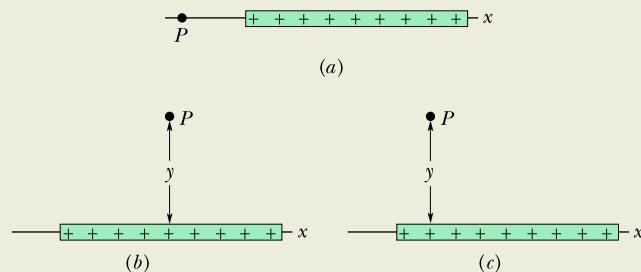
*Ring*, with point  $P$  on (central) axis of symmetry, as in Fig. 22.4.1. In the expression for  $dE$ , replace  $r^2$  with  $z^2 + R^2$ , as in Eq. 22.4.3. Express the adding component of  $d\vec{E}$  in terms of  $\theta$ . That introduces  $\cos \theta$ , but  $\theta$  is identical for all elements and thus is not a variable. Replace  $\cos \theta$  as in Eq. 22.4.4. Integrate over  $s$ , around the circumference of the ring.

*Circular arc*, with point  $P$  at the center of curvature, as in Fig. 22.4.4. Express the adding component of  $d\vec{E}$  in terms of  $\theta$ . That introduces either  $\sin \theta$  or  $\cos \theta$ . Reduce the resulting two variables  $s$  and  $\theta$  to one,  $\theta$ , by replacing  $ds$  with  $r d\theta$ . Integrate over  $\theta$  from one end of the arc to the other end.

*Straight line*, with point  $P$  on an extension of the line, as in Fig. 22.4.4a. In the expression for  $dE$ , replace  $r$  with  $x$ . Integrate over  $x$ , from end to end of the line of charge.

*Straight line*, with point  $P$  at perpendicular distance  $y$  from the line of charge, as in Fig. 22.4.4b. In the expression for  $dE$ , replace  $r$  with an expression involving  $x$  and  $y$ . If  $P$  is on the perpendicular bisector of the line of charge, find an expression for the adding component of  $d\vec{E}$ . That will introduce either  $\sin \theta$  or  $\cos \theta$ . Reduce the resulting two variables  $x$  and  $\theta$  to one,  $x$ , by replacing the trigonometric function with an expression (its definition) involving  $x$  and  $y$ . Integrate over  $x$  from end to end of the line of charge. If  $P$  is not on a line of symmetry, as in Fig. 22.4.4c, set up an integral to sum the components  $dE_x$ , and integrate over  $x$  to find  $E_x$ . Also set up an integral to sum the components  $dE_y$ , and integrate over  $x$  again to find  $E_y$ . Use the components  $E_x$  and  $E_y$  in the usual way to find the magnitude  $E$  and the orientation of  $\vec{E}$ .

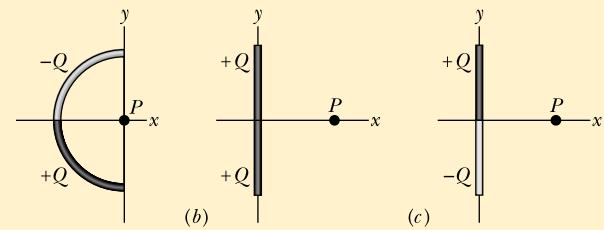
**Step 6.** One arrangement of the integration limits gives a positive result. The reverse gives the same result with a minus sign; discard the minus sign. If the result is to be stated in terms of the total charge  $Q$  of the distribution, replace  $\lambda$  with  $Q/L$ , in which  $L$  is the length of the distribution.



**Figure 22.4.4** (a) Point  $P$  is on an extension of the line of charge. (b)  $P$  is on a line of symmetry of the line of charge, at perpendicular distance  $y$  from that line. (c) Same as (b) except that  $P$  is not on a line of symmetry.

**Checkpoint 22.4.1**

The figure here shows three nonconducting rods, one circular and two straight. Each has a uniform charge of magnitude  $Q$  along its top half and another along its bottom half. For each rod, what is the direction of the net electric field at point  $P$ ?



## 22.5 THE ELECTRIC FIELD DUE TO A CHARGED DISK

**Learning Objectives**

After reading this module, you should be able to . . .

**22.5.1** Sketch a disk with uniform charge and indicate the direction of the electric field at a point on the central axis if the charge is positive and if it is negative.

**22.5.2** Explain how the equation for the electric field on the central axis of a uniformly charged ring can be

used to find the equation for the electric field on the central axis of a uniformly charged disk.

**22.5.3** For a point on the central axis of a uniformly charged disk, apply the relationship between the surface charge density  $\sigma$ , the disk radius  $R$ , and the distance  $z$  to that point.

**Key Idea**

- On the central axis through a uniformly charged disk,

$$E = \frac{\sigma}{2\epsilon_0} \left( 1 - \frac{z}{\sqrt{z^2 + R^2}} \right)$$

gives the electric field magnitude. Here  $z$  is the distance along the axis from the center of the disk,  $R$  is the radius of the disk, and  $\sigma$  is the surface charge density.

### The Electric Field Due to a Charged Disk

Now we switch from a line of charge to a surface of charge by examining the electric field of a circular plastic disk, with a radius  $R$  and a uniform surface charge density  $\sigma$  (charge per unit area, Table 22.4.1) on its top surface. The disk sets up a pattern of electric field lines around it, but here we restrict our attention to the electric field at an arbitrary point  $P$  on the central axis, at distance  $z$  from the center of the disk, as indicated in Fig. 22.5.1.

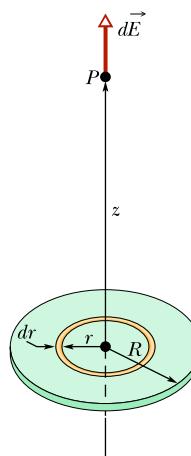
We could proceed as in the preceding module but set up a two-dimensional integral to include all of the field contributions from the two-dimensional distribution of charge on the top surface. However, we can save a lot of work with a neat shortcut using our earlier work with the field on the central axis of a thin ring.

We superimpose a ring on the disk as shown in Fig. 22.5.1, at an arbitrary radius  $r \leq R$ . The ring is so thin that we can treat the charge on it as a charge element  $dq$ . To find its small contribution  $dE$  to the electric field at point  $P$ , we rewrite Eq. 22.4.7 in terms of the ring's charge  $dq$  and radius  $r$ :

$$dE = \frac{dq z}{4\pi\epsilon_0(z^2 + r^2)^{3/2}}. \quad (22.5.1)$$

The ring's field points in the positive direction of the  $z$  axis.

To find the total field at  $P$ , we are going to integrate Eq. 22.5.1 from the center of the disk at  $r = 0$  out to the rim at  $r = R$  so that we sum all the  $dE$  contributions (by sweeping our arbitrary ring over the entire disk surface). However, that means we want to integrate with respect to a variable radius  $r$  of the ring.



**Figure 22.5.1** A disk of radius  $R$  and uniform positive charge. The ring shown has radius  $r$  and radial width  $dr$ . It sets up a differential electric field  $d\vec{E}$  at point  $P$  on its central axis.

We get  $dr$  into the expression by substituting for  $dq$  in Eq. 22.5.1. Because the ring is so thin, call its thickness  $dr$ . Then its surface area  $dA$  is the product of its circumference  $2\pi r$  and thickness  $dr$ . So, in terms of the surface charge density  $\sigma$ , we have

$$dq = \sigma dA = \sigma (2\pi r dr). \quad (22.5.2)$$

After substituting this into Eq. 22.5.1 and simplifying slightly, we can sum all the  $dE$  contributions with

$$E = \int dE = \frac{\sigma z}{4\epsilon_0} \int_0^R (z^2 + r^2)^{-3/2} (2r) dr, \quad (22.5.3)$$

where we have pulled the constants (including  $z$ ) out of the integral. To solve this integral, we cast it in the form  $\int X^m dX$  by setting  $X = (z^2 + r^2)$ ,  $m = -\frac{3}{2}$ , and  $dX = (2r) dr$ . For the recast integral we have

$$\int X^m dX = \frac{X^{m+1}}{m+1},$$

and so Eq. 22.5.3 becomes

$$E = \frac{\sigma z}{4\epsilon_0} \left[ \frac{(z^2 + r^2)^{-1/2}}{-\frac{1}{2}} \right]_0^R. \quad (22.5.4)$$

Taking the limits in Eq. 22.5.4 and rearranging, we find

$$E = \frac{\sigma}{2\epsilon_0} \left( 1 - \frac{z}{\sqrt{z^2 + R^2}} \right) \quad (\text{charged disk}) \quad (22.5.5)$$

as the magnitude of the electric field produced by a flat, circular, charged disk at points on its central axis. (In carrying out the integration, we assumed that  $z \geq 0$ .)

If we let  $R \rightarrow \infty$  while keeping  $z$  finite, the second term in the parentheses in Eq. 22.5.5 approaches zero, and this equation reduces to

$$E = \frac{\sigma}{2\epsilon_0} \quad (\text{infinite sheet}). \quad (22.5.6)$$

This is the electric field produced by an infinite sheet of uniform charge located on one side of a nonconductor such as plastic. The electric field lines for such a situation are shown in Fig. 22.1.4.

We also get Eq. 22.5.6 if we let  $z \rightarrow 0$  in Eq. 22.5.5 while keeping  $R$  finite. This shows that at points very close to the disk, the electric field set up by the disk is the same as if the disk were infinite in extent.

### Checkpoint 22.5.1

In Fig. 22.5.1, as we sweep the ring outward, what happens to the ring's contribution of electric field at point  $P$ ?

## 22.6 A POINT CHARGE IN AN ELECTRIC FIELD

### Learning Objectives

After reading this module, you should be able to . . .

- 22.6.1** For a charged particle placed in an external electric field (a field due to other charged objects), apply the relationship between the electric field  $\vec{E}$  at that point, the particle's charge  $q$ , and the electrostatic force  $\vec{F}$  that acts on the particle, and identify the relative directions of the force and the

field when the particle is positively charged and negatively charged.

- 22.6.2** Explain Millikan's procedure of measuring the elementary charge.

- 22.6.3** Explain the general mechanism of ink-jet printing.

## Key Ideas

- If a particle with charge  $q$  is placed in an external electric field  $\vec{E}$ , an electrostatic force  $\vec{F}$  acts on the particle:

$$\vec{F} = q\vec{E}.$$

- If charge  $q$  is positive, the force vector is in the same direction as the field vector. If charge  $q$  is negative, the force vector is in the opposite direction (the minus sign in the equation reverses the force vector from the field vector).

## A Point Charge in an Electric Field

In the preceding four modules we worked at the first of our two tasks: given a charge distribution, to find the electric field it produces in the surrounding space. Here we begin the second task: to determine what happens to a charged particle when it is in an electric field set up by other stationary or slowly moving charges.

What happens is that an electrostatic force acts on the particle, as given by

$$\vec{F} = q\vec{E}, \quad (22.6.1)$$

in which  $q$  is the charge of the particle (including its sign) and  $\vec{E}$  is the electric field that other charges have produced at the location of the particle. (The field is *not* the field set up by the particle itself; to distinguish the two fields, the field acting on the particle in Eq. 22.6.1 is often called the *external field*. A charged particle or object is not affected by its own electric field.) Equation 22.6.1 tells us



The electrostatic force  $\vec{F}$  acting on a charged particle located in an external electric field  $\vec{E}$  has the direction of  $\vec{E}$  if the charge  $q$  of the particle is positive and has the opposite direction if  $q$  is negative.

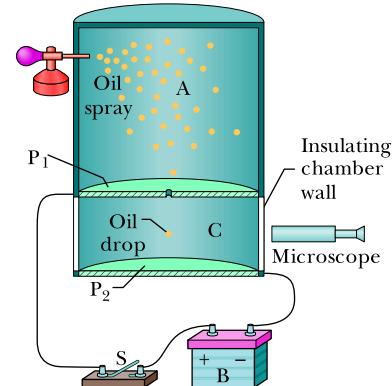
### Measuring the Elementary Charge

Equation 22.6.1 played a role in the measurement of the elementary charge  $e$  by American physicist Robert A. Millikan in 1910–1913. Figure 22.6.1 is a representation of his apparatus. When tiny oil drops are sprayed into chamber A, some of them become charged, either positively or negatively, in the process. Consider a drop that drifts downward through the small hole in plate  $P_1$  and into chamber C. Let us assume that this drop has a negative charge  $q$ .

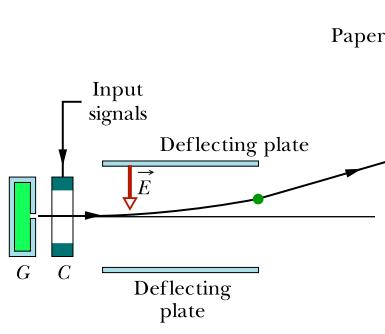
If switch S in Fig. 22.6.1 is open as shown, battery B has no electrical effect on chamber C. If the switch is closed (the connection between chamber C and the positive terminal of the battery is then complete), the battery causes an excess positive charge on conducting plate  $P_1$  and an excess negative charge on conducting plate  $P_2$ . The charged plates set up a downward-directed electric field  $\vec{E}$  in chamber C. According to Eq. 22.6.1, this field exerts an electrostatic force on any charged drop that happens to be in the chamber and affects its motion. In particular, our negatively charged drop will tend to drift upward.

By timing the motion of oil drops with the switch opened and with it closed and thus determining the effect of the charge  $q$ , Millikan discovered that the values of  $q$  were always given by

$$q = ne, \quad \text{for } n = 0, \pm 1, \pm 2, \pm 3, \dots, \quad (22.6.2)$$



**Figure 22.6.1** The Millikan oil-drop apparatus for measuring the elementary charge  $e$ . When a charged oil drop drifted into chamber C through the hole in plate  $P_1$ , its motion could be controlled by closing and opening switch S and thereby setting up or eliminating an electric field in chamber C. The microscope was used to view the drop, to permit timing of its motion.



**Figure 22.6.2** Ink-jet printer. Drops shot from generator  $G$  receive a charge in charging unit  $C$ . An input signal from a computer controls the charge and thus the effect of field  $\vec{E}$  on where the drop lands on the paper.

in which  $e$  turned out to be the fundamental constant we call the *elementary charge*,  $1.60 \times 10^{-19} \text{ C}$ . Millikan's experiment is convincing proof that charge is quantized, and he earned the 1923 Nobel Prize in physics in part for this work. Modern measurements of the elementary charge rely on a variety of interlocking experiments, all more precise than the pioneering experiment of Millikan.

### Ink-Jet Printing

The need for high-quality, high-speed printing has caused a search for an alternative to impact printing, such as occurs in an old typewriter. Building up letters by squirting tiny drops of ink at the paper is one such alternative.

Figure 22.6.2 shows a negatively charged drop moving between two conducting deflecting plates, between which a uniform, downward-directed electric field  $\vec{E}$  has been set up. The drop is deflected upward according to Eq. 22.6.1 and then strikes the paper at a position that is determined by the magnitudes of  $\vec{E}$  and the charge  $q$  of the drop.

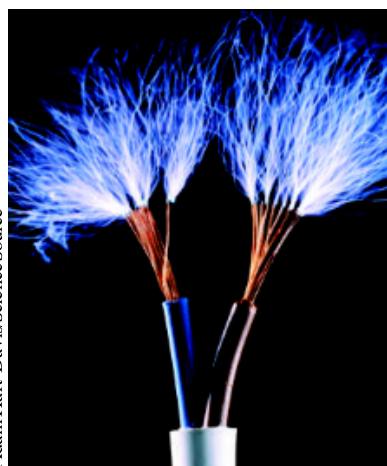
In practice,  $E$  is held constant and the position of the drop is determined by the charge  $q$  delivered to the drop in the charging unit, through which the drop must pass before entering the deflecting system. The charging unit, in turn, is activated by electronic signals that encode the material to be printed.

### Electrical Breakdown and Sparking

If the magnitude of an electric field in air exceeds a certain critical value  $E_c$ , the air undergoes *electrical breakdown*, a process whereby the field removes electrons from the atoms in the air. The air then begins to conduct electric current because the freed electrons are propelled into motion by the field. As they move, they collide with any atoms in their path, causing those atoms to emit light. We can see the paths, commonly called sparks, taken by the freed electrons because of that emitted light. Figure 22.6.3 shows sparks above charged metal wires where the electric fields due to the wires cause electrical breakdown of the air.

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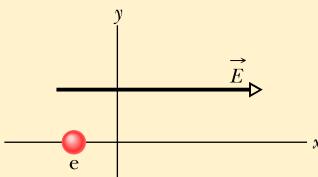
Adam Hart-Davis/Science Source



**Figure 22.6.3** The metal wires are so charged that the electric fields they produce in the surrounding space cause the air there to undergo electrical breakdown.

### Checkpoint 22.6.1

- In the figure, what is the direction of the electrostatic force on the electron due to the external electric field shown?
- In which direction will the electron accelerate if it is moving parallel to the  $y$  axis before it encounters the external field?
- If, instead, the electron is initially moving rightward, will its speed increase, decrease, or remain constant?



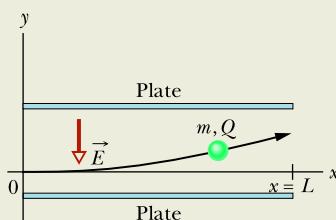
### Sample Problem 22.6.1 Motion of a charged particle in an electric field

Figure 22.6.4 shows the deflecting plates of an ink-jet printer, with superimposed coordinate axes. An ink drop with a mass  $m$  of  $1.3 \times 10^{-10}$  kg and a negative charge of magnitude  $Q = 1.5 \times 10^{-13}$  C enters the region between the plates, initially moving along the  $x$  axis with speed  $v_x = 18$  m/s. The length  $L$  of each plate is

1.6 cm. The plates are charged and thus produce an electric field at all points between them. Assume that field  $\vec{E}$  is downward directed, is uniform, and has a magnitude of  $1.4 \times 10^6$  N/C. What is the vertical deflection of the drop at the far edge of the plates? (The gravitational force on the drop is small relative to the electrostatic force acting on the drop and can be neglected.)

#### KEY IDEA

The drop is negatively charged and the electric field is directed *downward*. From Eq. 22.6.1, a constant



**Figure 22.6.4** An ink drop of mass  $m$  and charge magnitude  $Q$  is deflected in the electric field of an ink-jet printer.

electrostatic force of magnitude  $QE$  acts *upward* on the charged drop. Thus, as the drop travels parallel to the  $x$  axis at constant speed  $v_x$ , it accelerates upward with some constant acceleration  $a_y$ .

**Calculations:** Applying Newton's second law ( $F = ma$ ) for components along the  $y$  axis, we find that

$$a_y = \frac{F}{m} = \frac{QE}{m}. \quad (22.6.3)$$

Let  $t$  represent the time required for the drop to pass through the region between the plates. During  $t$  the vertical and horizontal displacements of the drop are

$$y = \frac{1}{2}a_y t^2 \quad \text{and} \quad L = v_x t, \quad (22.6.4)$$

respectively. Eliminating  $t$  between these two equations and substituting Eq. 22.6.3 for  $a_y$ , we find

$$\begin{aligned} y &= \frac{QEL^2}{2mv_x^2} \\ &= \frac{(1.5 \times 10^{-13} \text{ C})(1.4 \times 10^6 \text{ N/C})(1.6 \times 10^{-2} \text{ m})^2}{(2)(1.3 \times 10^{-10} \text{ kg})(18 \text{ m/s})^2} \\ &= 6.4 \times 10^{-4} \text{ m} \\ &= 0.64 \text{ mm.} \end{aligned} \quad (\text{Answer})$$

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## 22.7 A DIPOLE IN AN ELECTRIC FIELD

### Learning Objectives

After reading this module, you should be able to . . .

**22.7.1** On a sketch of an electric dipole in an external electric field, indicate the direction of the field, the direction of the dipole moment, the direction of the electrostatic forces on the two ends of the dipole, and the direction in which those forces tend to rotate the dipole, and identify the value of the net force on the dipole.

**22.7.2** Calculate the torque on an electric dipole in an external electric field by evaluating a cross product of the dipole moment vector and the electric field vector, in magnitude-angle notation and unit-vector notation.

**22.7.3** For an electric dipole in an external electric field, relate the potential energy of the dipole to the work done by a torque as the dipole rotates in the electric field.

**22.7.4** For an electric dipole in an external electric field, calculate the potential energy by taking a dot product of the dipole moment vector and the electric field vector, in magnitude-angle notation and unit-vector notation.

**22.7.5** For an electric dipole in an external electric field, identify the angles for the minimum and maximum potential energies and the angles for the minimum and maximum torque magnitudes.

## Key Ideas

- The torque on an electric dipole of dipole moment  $\vec{p}$  when placed in an external electric field  $\vec{E}$  is given by a cross product:

$$\vec{\tau} = \vec{p} \times \vec{E}.$$

- A potential energy  $U$  is associated with the orientation of the dipole moment in the field, as given by a dot product:

$$U = -\vec{p} \cdot \vec{E}.$$

- If the dipole orientation changes, the work done by the electric field is

$$W = -\Delta U.$$

If the change in orientation is due to an external agent, the work done by the agent is  $W_a = -W$ .

## A Dipole in an Electric Field

We have defined the electric dipole moment  $\vec{p}$  of an electric dipole to be a vector that points from the negative to the positive end of the dipole. As you will see, the behavior of a dipole in a uniform external electric field  $\vec{E}$  can be described completely in terms of the two vectors  $\vec{E}$  and  $\vec{p}$ , with no need of any details about the dipole's structure.

A molecule of water ( $H_2O$ ) is an electric dipole; Fig. 22.7.1 shows why. There the black dots represent the oxygen nucleus (having eight protons) and the two hydrogen nuclei (having one proton each). The colored enclosed areas represent the regions in which electrons can be located around the nuclei.

In a water molecule, the two hydrogen atoms and the oxygen atom do not lie on a straight line but form an angle of about  $105^\circ$ , as shown in Fig. 22.7.1. As a result, the molecule has a definite “oxygen side” and “hydrogen side.” Moreover, the 10 electrons of the molecule tend to remain closer to the oxygen nucleus than to the hydrogen nuclei. This makes the oxygen side of the molecule slightly more negative than the hydrogen side and creates an electric dipole moment  $\vec{p}$  that points along the symmetry axis of the molecule as shown. If the water molecule is placed in an external electric field, it behaves as would be expected of the more abstract electric dipole of Fig. 22.3.2.

To examine this behavior, we now consider such an abstract dipole in a uniform external electric field  $\vec{E}$ , as shown in Fig. 22.7.2a. We assume that the dipole is a rigid structure that consists of two centers of opposite charge, each of magnitude  $q$ , separated by a distance  $d$ . The dipole moment  $\vec{p}$  makes an angle  $\theta$  with field  $\vec{E}$ .

Electrostatic forces act on the charged ends of the dipole. Because the electric field is uniform, those forces act in opposite directions (as shown in Fig. 22.7.2a) and with the same magnitude  $F = qE$ . Thus, *because the field is uniform*, the net force on the dipole from the field is zero and the center of mass of the dipole does not move. However, the forces on the charged ends do produce a net torque  $\vec{\tau}$  on the dipole about its center of mass. The center of mass lies on the line connecting the charged ends, at some distance  $x$  from one end and thus a distance  $d - x$  from the other end. From Eq. 10.6.1 ( $\tau = rF \sin \phi$ ), we can write the magnitude of the net torque  $\vec{\tau}$  as

$$\tau = Fx \sin \theta + F(d - x) \sin \theta = Fd \sin \theta. \quad (22.7.1)$$

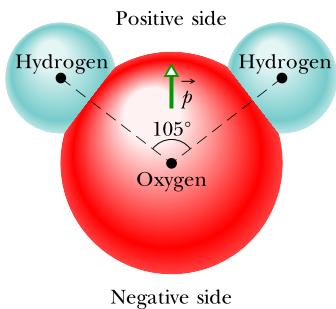
We can also write the magnitude of  $\vec{\tau}$  in terms of the magnitudes of the electric field  $E$  and the dipole moment  $p = qd$ . To do so, we substitute  $qE$  for  $F$  and  $p/q$  for  $d$  in Eq. 22.7.1, finding that the magnitude of  $\vec{\tau}$  is

$$\tau = pE \sin \theta. \quad (22.7.2)$$

We can generalize this equation to vector form as

$$\vec{\tau} = \vec{p} \times \vec{E} \quad (\text{torque on a dipole}). \quad (22.7.3)$$

Vectors  $\vec{p}$  and  $\vec{E}$  are shown in Fig. 22.7.2b. The torque acting on a dipole tends to rotate  $\vec{p}$  (hence the dipole) into the direction of field  $\vec{E}$ , thereby reducing  $\theta$ .



**Figure 22.7.1** A molecule of  $H_2O$ , showing the three nuclei (represented by dots) and the regions in which the electrons can be located. The electric dipole moment  $\vec{p}$  points from the (negative) oxygen side to the (positive) hydrogen side of the molecule.

In Fig. 22.7.2, such rotation is clockwise. As we discussed in Chapter 10, we can represent a torque that gives rise to a clockwise rotation by including a minus sign with the magnitude of the torque. With that notation, the torque of Fig. 22.7.2 is

$$\tau = -pE \sin \theta. \quad (22.7.4)$$

### Potential Energy of an Electric Dipole

Potential energy can be associated with the orientation of an electric dipole in an electric field. The dipole has its least potential energy when it is in its equilibrium orientation, which is when its moment  $\vec{p}$  is lined up with the field  $\vec{E}$  (then  $\vec{\tau} = \vec{p} \times \vec{E} = 0$ ). It has greater potential energy in all other orientations. Thus the dipole is like a pendulum, which has its least gravitational potential energy in its equilibrium orientation—at its lowest point. To rotate the dipole or the pendulum to any other orientation requires work by some external agent.

In any situation involving potential energy, we are free to define the zero-potential-energy configuration in an arbitrary way because only differences in potential energy have physical meaning. The expression for the potential energy of an electric dipole in an external electric field is simplest if we choose the potential energy to be zero when the angle  $\theta$  in Fig. 22.7.2 is  $90^\circ$ . We then can find the potential energy  $U$  of the dipole at any other value of  $\theta$  with Eq. 8.1.1 ( $\Delta U = -W$ ) by calculating the work  $W$  done by the field on the dipole when the dipole is rotated to that value of  $\theta$  from  $90^\circ$ . With the aid of Eq. 10.8.5 ( $W = \int \tau d\theta$ ) and Eq. 22.7.4, we find that the potential energy  $U$  at any angle  $\theta$  is

$$U = -W = - \int_{90^\circ}^{\theta} \tau d\theta = \int_{90^\circ}^{\theta} pE \sin \theta d\theta. \quad (22.7.5)$$

Evaluating the integral leads to

$$U = -pE \cos \theta. \quad (22.7.6)$$

We can generalize this equation to vector form as

$$U = -\vec{p} \cdot \vec{E} \quad (\text{potential energy of a dipole}). \quad (22.7.7)$$

Equations 22.7.6 and 22.7.7 show us that the potential energy of the dipole is least ( $U = -pE$ ) when  $\theta = 0$  ( $\vec{p}$  and  $\vec{E}$  are in the same direction); the potential energy is greatest ( $U = pE$ ) when  $\theta = 180^\circ$  ( $\vec{p}$  and  $\vec{E}$  are in opposite directions).

When a dipole rotates from an initial orientation  $\theta_i$  to another orientation  $\theta_f$ , the work  $W$  done on the dipole by the electric field is

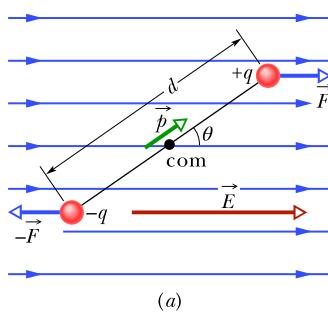
$$W = -\Delta U = -(U_f - U_i), \quad (22.7.8)$$

where  $U_f$  and  $U_i$  are calculated with Eq. 22.7.7. If the change in orientation is caused by an applied torque (commonly said to be due to an external agent), then the work  $W_a$  done on the dipole by the applied torque is the negative of the work done on the dipole by the field; that is,

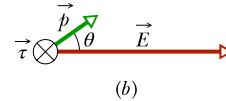
$$W_a = -W = (U_f - U_i). \quad (22.7.9)$$

### Microwave Cooking

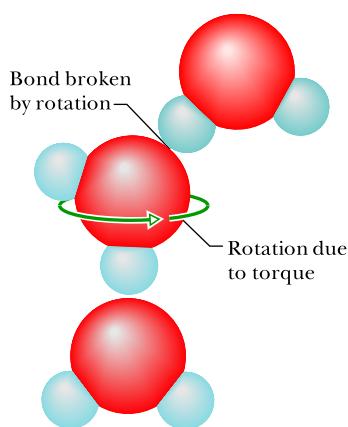
In liquid water, where molecules are relatively free to move around, the electric field produced by each molecular dipole affects the surrounding dipoles. As a result, the molecules bond together in groups of two or three, because the negative (oxygen) end of one dipole and a positive (hydrogen) end of another dipole attract each other. Each time a group forms, electric potential energy is transferred to the random thermal motion of the group and the surrounding molecules. And each time collisions among the molecules break up a group, the transfer is reversed. The temperature of the water (which is associated with the average thermal motion) does not change because, on the average, the net transfer of energy is zero.



The dipole is being torqued into alignment.



**Figure 22.7.2** (a) An electric dipole in a uniform external electric field  $\vec{E}$ . Two centers of equal but opposite charge are separated by distance  $d$ . The line between them represents their rigid connection. (b) Field  $\vec{E}$  causes a torque  $\vec{\tau}$  on the dipole. The direction of  $\vec{\tau}$  is into the page, as represented by the symbol  $\otimes$ .



**Figure 22.7.3** A group of three water molecules. A torque due to an oscillating electric field in a microwave oven breaks one of the bonds between the molecules and thus breaks up the group.

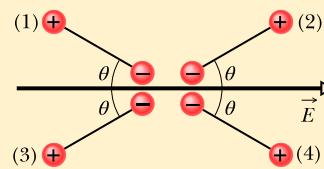
In a microwave oven, the story differs. When the oven is operated, the microwaves produce (in the oven) an electric field that rapidly oscillates back and forth in direction. If there is water in the oven, the oscillating field exerts oscillating torques on the water molecules, continually rotating them back and forth to align their dipole moments with the field direction. Molecules that are bonded as a pair can twist around their common bond to stay aligned, but molecules that are bonded in a group of three must break at least one of their two bonds (Fig. 22.7.3).

The energy to break these bonds comes from the electric field, that is, from the microwaves. Then molecules that have broken away from groups can form new groups, transferring the energy they just gained into thermal energy. Thus, thermal energy is added to the water when the groups form but is not removed when the groups break apart, and the temperature of the water increases. Foods that contain water can be cooked in a microwave oven because of the heating of that water. If water molecules were not electric dipoles, this would not be so and microwave ovens would be useless.

FCP

### Checkpoint 22.7.1

The figure shows four orientations of an electric dipole in an external electric field. Rank the orientations according to (a) the magnitude of the torque on the dipole and (b) the potential energy of the dipole, greatest first.



### Sample Problem 22.7.1 Torque and energy of an electric dipole in an electric field

A neutral water molecule ( $\text{H}_2\text{O}$ ) in its vapor state has an electric dipole moment of magnitude  $6.2 \times 10^{-30} \text{ C} \cdot \text{m}$ .

- (a) How far apart are the molecule's centers of positive and negative charge?

#### KEY IDEA

A molecule's dipole moment depends on the magnitude  $q$  of the molecule's positive or negative charge and the charge separation  $d$ .

**Calculations:** There are 10 electrons and 10 protons in a neutral water molecule; so the magnitude of its dipole moment is

$$p = qd = (10e)(d),$$

in which  $d$  is the separation we are seeking and  $e$  is the elementary charge. Thus,

$$\begin{aligned} d &= \frac{p}{10e} = \frac{6.2 \times 10^{-30} \text{ C} \cdot \text{m}}{(10)(1.60 \times 10^{-19} \text{ C})} \\ &= 3.9 \times 10^{-12} \text{ m} = 3.9 \text{ pm}. \quad (\text{Answer}) \end{aligned}$$

This distance is not only small, but it is also actually smaller than the radius of a hydrogen atom.

- (b) If the molecule is placed in an electric field of  $1.5 \times 10^4 \text{ N/C}$ , what maximum torque can the field exert on it? (Such a field can easily be set up in the laboratory.)

#### KEY IDEA

The torque on a dipole is maximum when the angle  $\theta$  between  $\vec{p}$  and  $\vec{E}$  is  $90^\circ$ .

**Calculation:** Substituting  $\theta = 90^\circ$  in Eq. 22.7.2 yields

$$\begin{aligned} \tau &= pE \sin \theta \\ &= (6.2 \times 10^{-30} \text{ C} \cdot \text{m})(1.5 \times 10^4 \text{ N/C})(\sin 90^\circ) \\ &= 9.3 \times 10^{-26} \text{ N} \cdot \text{m}. \quad (\text{Answer}) \end{aligned}$$

- (c) How much work must an *external agent* do to rotate this molecule by  $180^\circ$  in this field, starting from its fully aligned position, for which  $\theta = 0^\circ$ ?

#### KEY IDEA

The work done by an external agent (by means of a torque applied to the molecule) is equal to the change in the molecule's potential energy due to the change in orientation.

**Calculation:** From Eq. 22.7.9, we find

$$\begin{aligned} W_a &= U_{180^\circ} - U_0 \\ &= (-pE \cos 180^\circ) - (-pE \cos 0^\circ) \\ &= 2pE = (2)(6.2 \times 10^{-30} \text{ C} \cdot \text{m})(1.5 \times 10^4 \text{ N/C}) \\ &= 1.9 \times 10^{-25} \text{ J}. \quad (\text{Answer}) \end{aligned}$$

## Review & Summary

**Electric Field** To explain the electrostatic force between two charges, we assume that each charge sets up an electric field in the space around it. The force acting on each charge is then due to the electric field set up at its location by the other charge.

**Definition of Electric Field** The *electric field*  $\vec{E}$  at any point is defined in terms of the electrostatic force  $\vec{F}$  that would be exerted on a positive test charge  $q_0$  placed there:

$$\vec{E} = \frac{\vec{F}}{q_0}. \quad (22.1.1)$$

**Electric Field Lines** *Electric field lines* provide a means for visualizing the direction and magnitude of electric fields. The electric field vector at any point is tangent to a field line through that point. The density of field lines in any region is proportional to the magnitude of the electric field in that region. Field lines originate on positive charges and terminate on negative charges.

**Field Due to a Point Charge** The magnitude of the electric field  $\vec{E}$  set up by a point charge  $q$  at a distance  $r$  from the charge is

$$E = \frac{1}{4\pi\epsilon_0} \frac{|q|}{r^2}. \quad (22.2.2)$$

The direction of  $\vec{E}$  is away from the point charge if the charge is positive and toward it if the charge is negative.

**Field Due to an Electric Dipole** An *electric dipole* consists of two particles with charges of equal magnitude  $q$  but opposite sign, separated by a small distance  $d$ . Their **electric dipole moment**  $\vec{p}$  has magnitude  $qd$  and points from the negative charge to the positive charge. The magnitude of the electric field set up by the dipole at a distant point on the dipole axis (which runs through both charges) is

$$E = \frac{1}{2\pi\epsilon_0} \frac{p}{z^3} \quad (22.3.5)$$

where  $z$  is the distance between the point and the center of the dipole.

## Questions

- 1 Figure 22.1 shows three arrangements of electric field lines. In each arrangement, a proton is released from rest at point  $A$  and is then accelerated through point  $B$  by the electric field. Points  $A$  and  $B$  have equal separations in the three arrangements. Rank the arrangements according to the linear momentum of the proton at point  $B$ , greatest first.

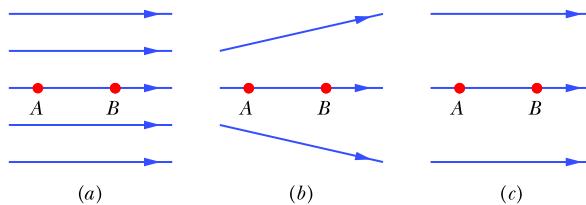


Figure 22.1 Question 1.

**Field Due to a Continuous Charge Distribution** The electric field due to a *continuous charge distribution* is found by treating charge elements as point charges and then summing, via integration, the electric field vectors produced by all the charge elements to find the net vector.

**Field Due to a Charged Disk** The electric field magnitude at a point on the central axis through a uniformly charged disk is given by

$$E = \frac{\sigma}{2\epsilon_0} \left( 1 - \frac{z}{\sqrt{z^2 + R^2}} \right), \quad (22.5.5)$$

where  $z$  is the distance along the axis from the center of the disk,  $R$  is the radius of the disk, and  $\sigma$  is the surface charge density.

**Force on a Point Charge in an Electric Field** When a point charge  $q$  is placed in an external electric field  $\vec{E}$ , the electrostatic force  $\vec{F}$  that acts on the point charge is

$$\vec{F} = q\vec{E}. \quad (22.6.1)$$

Force  $\vec{F}$  has the same direction as  $\vec{E}$  if  $q$  is positive and the opposite direction if  $q$  is negative.

**Dipole in an Electric Field** When an electric dipole of dipole moment  $\vec{p}$  is placed in an electric field  $\vec{E}$ , the field exerts a torque  $\vec{\tau}$  on the dipole:

$$\vec{\tau} = \vec{p} \times \vec{E}. \quad (22.7.3)$$

The dipole has a potential energy  $U$  associated with its orientation in the field:

$$U = -\vec{p} \cdot \vec{E}. \quad (22.7.7)$$

This potential energy is defined to be zero when  $\vec{p}$  is perpendicular to  $\vec{E}$ ; it is least ( $U = -pE$ ) when  $\vec{p}$  is aligned with  $\vec{E}$  and greatest ( $U = pE$ ) when  $\vec{p}$  is directed opposite  $\vec{E}$ .

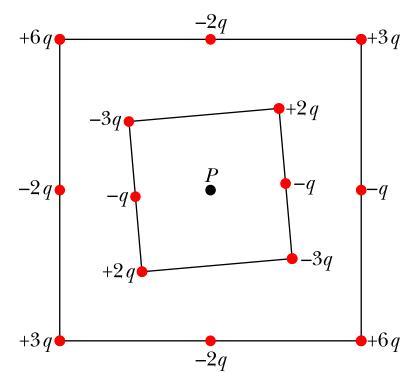


Figure 22.2 Question 2.

- 2 Figure 22.2 shows two square arrays of charged particles. The squares, which are centered on point  $P$ , are misaligned. The particles are separated by either  $d$  or  $d/2$  along the perimeters of the squares. What are the magnitude and direction of the net electric field at  $P$ ?

- 3 In Fig. 22.3, two particles of charge  $-q$  are

arranged symmetrically about the  $y$  axis; each produces an electric field at point  $P$  on that axis. (a) Are the magnitudes of the fields at  $P$  equal? (b) Is each electric field directed toward or away from the charge producing it? (c) Is the magnitude of the net electric field at  $P$  equal to the sum of the magnitudes  $E$  of the two field vectors (is it equal to  $2E$ )? (d) Do the  $x$  components of those two field vectors add or cancel? (e) Do their  $y$  components add or cancel? (f) Is the direction of the net field at  $P$  that of the canceling components or the adding components? (g) What is the direction of the net field?

**4** Figure 22.4 shows four situations in which four charged particles are evenly spaced to the left and right of a central point. The charge values are indicated. Rank the situations according to the magnitude of the net electric field at the central point, greatest first.

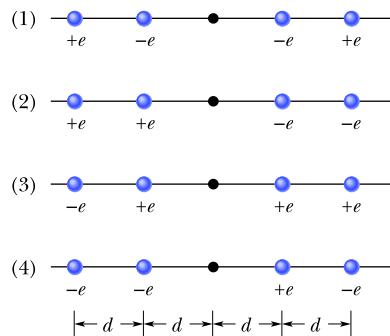


Figure 22.4 Question 4.

**5** Figure 22.5 shows two charged particles fixed in place on an axis. (a) Where on the axis (other than at an infinite distance) is there a point at which their net electric field is zero: between the charges, to their left, or to their right? (b) Is there a point of zero net electric field anywhere off the axis (other than at an infinite distance)?

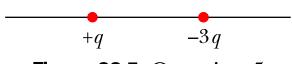


Figure 22.5 Question 5.

**6** In Fig. 22.6, two identical circular nonconducting rings are centered on the same line with their planes perpendicular to the line. Each ring has charge that is uniformly distributed along its circumference. The rings each produce electric fields at points along the line. For three situations, the charges on rings  $A$  and  $B$  are, respectively, (1)  $q_0$  and  $q_0$ , (2)  $-q_0$  and  $-q_0$ , and (3)  $-q_0$  and  $q_0$ . Rank the situations according to the magnitude of the net electric field at (a) point  $P_1$  midway between the rings, (b) point  $P_2$  at the center of ring  $B$ , and (c) point  $P_3$  to the right of ring  $B$ , greatest first.

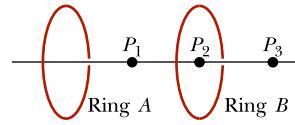


Figure 22.6 Question 6.

**7** The potential energies associated with four orientations of an electric dipole in an electric field are (1)  $-5U_0$ , (2)  $-7U_0$ , (3)  $3U_0$ , and (4)  $5U_0$ , where  $U_0$  is positive. Rank the orientations according to (a) the angle between the electric dipole moment  $\vec{p}$  and the electric field  $\vec{E}$  and (b) the magnitude of the torque on the electric dipole, greatest first.

**8** (a) In Checkpoint 22.7.1, if the dipole rotates from orientation 1 to orientation 2, is the work done on the dipole by the field positive, negative, or zero? (b) If, instead, the dipole rotates from orientation 1 to orientation 4, is the work done by the field more than, less than, or the same as in (a)?

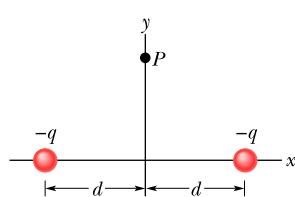


Figure 22.3 Question 3.

**9** Figure 22.7 shows two disks and a flat ring, each with the same uniform charge  $Q$ . Rank the objects according to the magnitude of the electric field they create at points  $P$  (which are at the same vertical heights), greatest first.

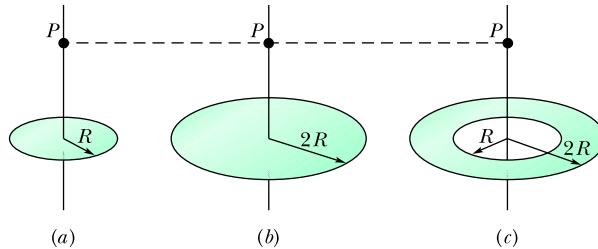


Figure 22.7 Question 9.

**10** In Fig. 22.8, an electron  $e$  travels through a small hole in plate  $A$  and then toward plate  $B$ . A uniform electric field in the region between the plates then slows the electron without deflecting it. (a) What is the direction of the field? (b) Four other particles similarly travel through small holes in either plate  $A$  or plate  $B$  and then into the region between the plates. Three have charges  $+q_1$ ,  $+q_2$ , and  $-q_3$ . The fourth (labeled n) is a neutron, which is electrically neutral. Does the speed of each of those four other particles increase, decrease, or remain the same in the region between the plates?

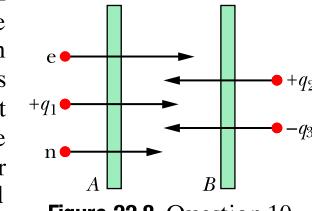


Figure 22.8 Question 10.

**11** In Fig. 22.9a, a circular plastic rod with uniform charge  $+Q$  produces an electric field of magnitude  $E$  at the center of curvature (at the origin). In Figs. 22.9b, c, and d, more circular rods, each with identical uniform charges  $+Q$ , are added until the circle is complete. A fifth arrangement (which would be labeled e) is like that in d except the rod in the fourth quadrant has charge  $-Q$ . Rank the five arrangements according to the magnitude of the electric field at the center of curvature, greatest first.

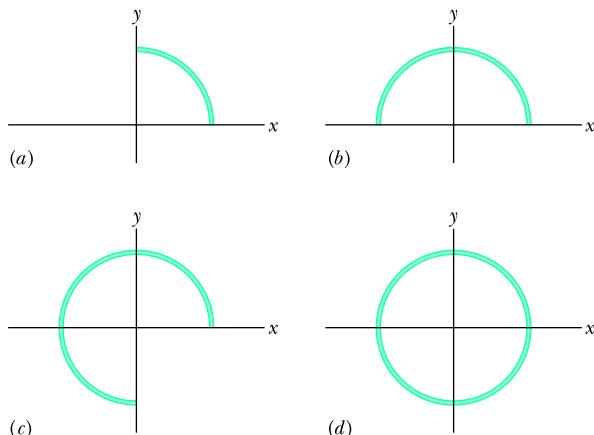


Figure 22.9 Question 11.

**12** When three electric dipoles are near each other, they each experience the electric field of the other two, and the

three-dipole system has a certain potential energy. Figure 22.10 shows two arrangements in which three electric dipoles are side by side. Each dipole has the same magnitude of electric dipole moment, and the spacings between adjacent dipoles are identical. In which arrangement is the potential energy of the three-dipole system greater?

- 13** Figure 22.11 shows three rods, each with the same charge  $Q$  spread uniformly along its length. Rods *a* (of length  $L$ ) and *b*

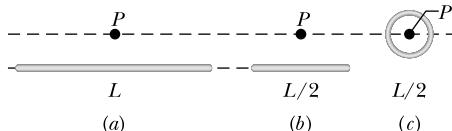


Figure 22.11 Question 13.

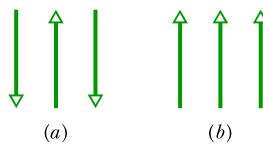


Figure 22.10 Question 12.

(of length  $L/2$ ) are straight, and points *P* are aligned with their midpoints. Rod *c* (of length  $L/2$ ) forms a complete circle about point *P*. Rank the rods according to the magnitude of the electric field they create at points *P*, greatest first.

- 14** Figure 22.12 shows five protons that are launched in a uniform electric field  $\vec{E}$ ; the magnitude and direction of the launch velocities are indicated. Rank the protons according to the magnitude of their accelerations due to the field, greatest first.

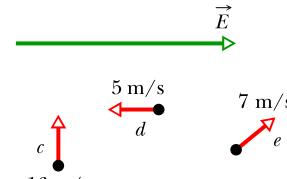


Figure 22.12 Question 14.

## Problems

**GO** Tutoring problem available (at instructor's discretion) in WileyPLUS

**SSM** Worked-out solution available in Student Solutions Manual

**E** Easy **M** Medium **H** Hard

**FCP** Additional information available in *The Flying Circus of Physics* and at [flyingcircusofphysics.com](http://flyingcircusofphysics.com)

**CALC** Requires calculus

**BIO** Biomedical application

### Module 22.1 The Electric Field

- 1 E** Sketch qualitatively the electric field lines both between and outside two concentric conducting spherical shells when a uniform positive charge  $q_1$  is on the inner shell and a uniform negative charge  $-q_2$  is on the outer. Consider the cases  $q_1 > q_2$ ,  $q_1 = q_2$ , and  $q_1 < q_2$ .

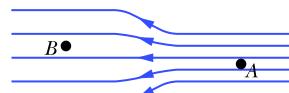


Figure 22.13 Problem 2.

- 2 E** In Fig. 22.13 the electric field lines on the left have twice the separation of those on the right. (a) If the magnitude of the field at *A* is 40 N/C, what is the magnitude of the force on a proton at *B*? (b) What is the magnitude of the field at *B*?

- 7 M SSM** In Fig. 22.14, the four particles form a square of edge length  $a = 5.00 \text{ cm}$  and have charges  $q_1 = +10.0 \text{ nC}$ ,  $q_2 = -20.0 \text{ nC}$ ,  $q_3 = +20.0 \text{ nC}$ , and  $q_4 = -10.0 \text{ nC}$ . In unit-vector notation, what net electric field do the particles produce at the square's center?

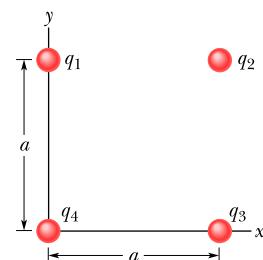


Figure 22.14 Problem 7.

- 8 M GO** In Fig. 22.15, the four particles are fixed in place and have charges  $q_1 = q_2 = +5e$ ,  $q_3 = +3e$ , and  $q_4 = -12e$ . Distance  $d = 5.0 \mu\text{m}$ . What is the magnitude of the net electric field at point *P* due to the particles?

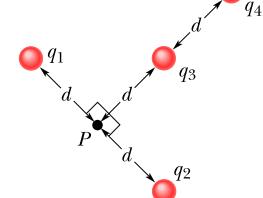


Figure 22.15 Problem 8.

### Module 22.2 The Electric Field Due to a Charged Particle

- 3 E SSM** The nucleus of a plutonium-239 atom contains 94 protons. Assume that the nucleus is a sphere with radius  $6.64 \text{ fm}$  and with the charge of the protons uniformly spread through the sphere. At the surface of the nucleus, what are the (a) magnitude and (b) direction (radially inward or outward) of the electric field produced by the protons?

- 4 E** Two charged particles are attached to an *x* axis: Particle 1 of charge  $-2.00 \times 10^{-7} \text{ C}$  is at position  $x = 6.00 \text{ cm}$  and particle 2 of charge  $+2.00 \times 10^{-7} \text{ C}$  is at position  $x = 21.0 \text{ cm}$ . Midway between the particles, what is their net electric field in unit-vector notation?

- 5 E SSM** A charged particle produces an electric field with a magnitude of  $2.0 \text{ N/C}$  at a point that is  $50 \text{ cm}$  away from the particle. What is the magnitude of the particle's charge?

- 6 E** What is the magnitude of a point charge that would create an electric field of  $1.00 \text{ N/C}$  at points  $1.00 \text{ m}$  away?

- 9 M GO** Figure 22.16 shows two charged particles on an *x* axis:  $-q = -3.20 \times 10^{-19} \text{ C}$  at  $x = -3.00 \text{ m}$  and  $q = 3.20 \times 10^{-19} \text{ C}$  at  $x = +3.00 \text{ m}$ . What are the (a) magnitude and (b) direction (relative to the positive direction of the *x* axis) of the net electric field produced at point *P* at  $y = 4.00 \text{ m}$ ?

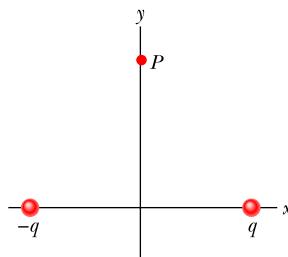


Figure 22.16 Problem 9.

- 10 M GO** Figure 22.17a shows two charged particles fixed in place on an  $x$  axis with separation  $L$ . The ratio  $q_1/q_2$  of their charge magnitudes is 4.00. Figure 22.17b shows the  $x$  component  $E_{\text{net},x}$  of their net electric field along the  $x$  axis just to the right of particle 2. The  $x$  axis scale is set by  $x_s = 30.0 \text{ cm}$ . (a) At what value of  $x > 0$  is  $E_{\text{net},x}$  maximum? (b) If particle 2 has charge  $-q_2 = -3e$ , what is the value of that maximum?

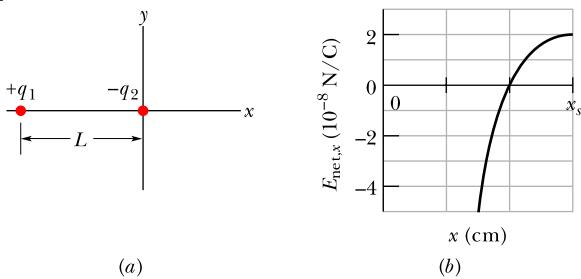


Figure 22.17 Problem 10.

- 11 M SSM** Two charged particles are fixed to an  $x$  axis: Particle 1 of charge  $q_1 = 2.1 \times 10^{-8} \text{ C}$  is at position  $x = 20 \text{ cm}$  and particle 2 of charge  $q_2 = -4.00q_1$  is at position  $x = 70 \text{ cm}$ . At what coordinate on the axis (other than at infinity) is the net electric field produced by the two particles equal to zero?

- 12 M GO** Figure 22.18 shows an uneven arrangement of electrons (e) and protons (p) on a circular arc of radius  $r = 2.00 \text{ cm}$ , with angles  $\theta_1 = 30.0^\circ$ ,  $\theta_2 = 50.0^\circ$ ,  $\theta_3 = 30.0^\circ$ , and  $\theta_4 = 20.0^\circ$ . What are the (a) magnitude and (b) direction (relative to the positive direction of the  $x$  axis) of the net electric field produced at the center of the arc?

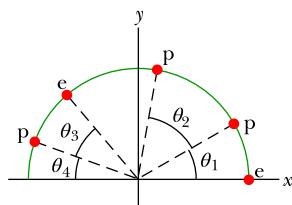


Figure 22.18 Problem 12.

- 13 M GO** Figure 22.19 shows a proton (p) on the central axis through a disk with a uniform charge density due to excess electrons. The disk is seen from an edge-on view. Three of those electrons are shown: electron  $e_c$  at the disk center and electrons  $e_s$  at opposite sides of the disk, at radius  $R$  from the center. The proton is initially at distance  $z = R = 2.00 \text{ cm}$  from the disk. At that location, what are the magnitudes of (a) the electric field  $\vec{E}_c$  due to electron  $e_c$  and (b) the net electric field  $\vec{E}_{s,\text{net}}$  due to electrons  $e_s$ ? The proton is then moved to  $z = R/10.0$ . What then are the magnitudes of (c)  $\vec{E}_c$  and (d)  $\vec{E}_{s,\text{net}}$  at the proton's location? (e) From (a) and (c) we see that as the proton gets nearer to the disk, the magnitude of  $\vec{E}_c$  increases, as expected. Why does the magnitude of  $\vec{E}_{s,\text{net}}$  from the two side electrons decrease, as we see from (b) and (d)?

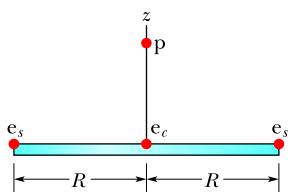


Figure 22.19 Problem 13.

- 14 M** In Fig. 22.20, particle 1 of charge  $q_1 = -5.00q$  and particle 2 of charge  $q_2 = +2.00q$  are fixed to an  $x$  axis. (a) As a multiple of distance  $L$ , at what coordinate on the axis is the net electric field of the particles

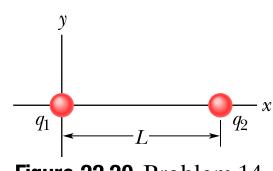


Figure 22.20 Problem 14.

- zero? (b) Sketch the net electric field lines between and around the particles.

- 15 M** In Fig. 22.21, the three particles are fixed in place and have charges  $q_1 = q_2 = +e$  and  $q_3 = +2e$ . Distance  $a = 6.00 \mu\text{m}$ . What are the (a) magnitude and (b) direction of the net electric field at point  $P$  due to the particles?

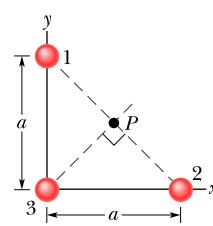


Figure 22.21 Problem 15.

- 16 H** Figure 22.22 shows a plastic ring of radius  $R = 50.0 \text{ cm}$ . Two small charged beads are on the ring: Bead 1 of charge  $+2.00 \mu\text{C}$  is fixed in place at the left side; bead 2 of charge  $+6.00 \mu\text{C}$  can be moved along the ring. The two beads produce a net electric field of magnitude  $E$  at the center of the ring. At what (a) positive and (b) negative value of angle  $\theta$  should bead 2 be positioned such that  $E = 2.00 \times 10^5 \text{ N/C}$ ?

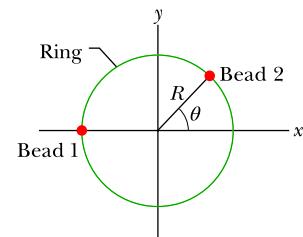


Figure 22.22 Problem 16.

- 17 H** Two charged beads are on the plastic ring in Fig. 22.23a. Bead 2, which is not shown, is fixed in place on the ring, which has radius  $R = 60.0 \text{ cm}$ . Bead 1, which is not fixed in place, is initially on the  $x$  axis at angle  $\theta = 0^\circ$ . It is then moved to the opposite side, at angle  $\theta = 180^\circ$ , through the first and second quadrants of the  $xy$  coordinate system. Figure 22.23b gives the  $x$  component of the net electric field produced at the origin by the two beads as a function of  $\theta$ , and Fig. 22.23c gives the  $y$  component of that net electric field. The vertical axis scales are set by  $E_{xs} = 5.0 \times 10^4 \text{ N/C}$  and  $E_{ys} = -9.0 \times 10^4 \text{ N/C}$ . (a) At what angle  $\theta$  is bead 2 located? What are the charges of (b) bead 1 and (c) bead 2?

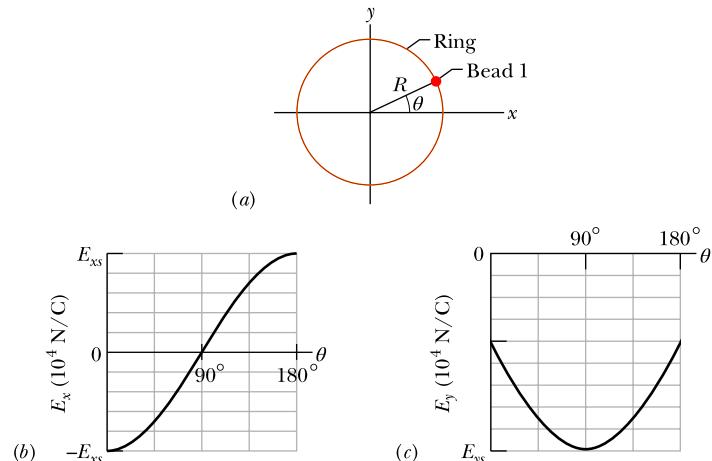


Figure 22.23 Problem 17.

### Module 22.3 The Electric Field Due to a Dipole

- 18 M** The electric field of an electric dipole along the dipole axis is approximated by Eqs. 22.3.4 and 22.3.5. If a binomial expansion is made of Eq. 22.3.3, what is the next term in the expression for the dipole's electric field along the dipole axis? That is, what is  $E_{\text{next}}$  in the expression

$$E = \frac{1}{2\pi\epsilon_0 z^3} q d + E_{\text{next}}$$

- 19 M** Figure 22.24 shows an electric dipole. What are the (a) magnitude and (b) direction (relative to the positive direction of the  $x$  axis) of the dipole's electric field at point  $P$ , located at distance  $r \gg d$ ?

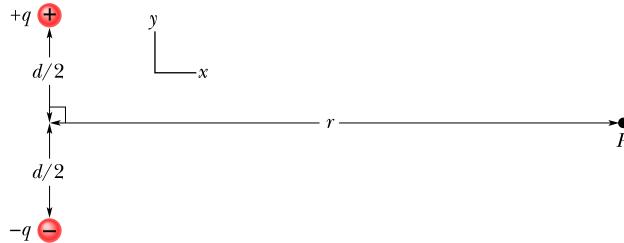


Figure 22.24 Problem 19.

- 20 M** Equations 22.3.4 and 22.3.5 are approximations of the magnitude of the electric field of an electric dipole, at points along the dipole axis. Consider a point  $P$  on that axis at distance  $z = 5.00d$  from the dipole center ( $d$  is the separation distance between the particles of the dipole). Let  $E_{\text{appr}}$  be the magnitude of the field at point  $P$  as approximated by Eqs. 22.3.4 and 22.3.5. Let  $E_{\text{act}}$  be the actual magnitude. What is the ratio  $E_{\text{appr}}/E_{\text{act}}$ ?

- 21 H SSM** *Electric quadrupole.* Figure 22.25 shows a generic electric quadrupole. It consists of two dipoles with dipole moments that are equal in magnitude but opposite in direction. Show that the value of  $E$  on the axis of the quadrupole for a point  $P$  a distance  $z$  from its center (assume  $z \gg d$ ) is given by

$$E = \frac{3Q}{4\pi\epsilon_0 z^4},$$

in which  $Q (= 2qd^2)$  is known as the *quadrupole moment* of the charge distribution.

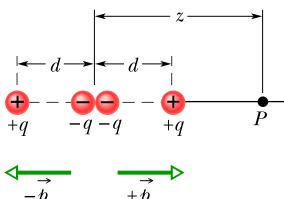


Figure 22.25 Problem 21.

#### Module 22.4 The Electric Field Due to a Line of Charge

- 22 E** *Density, density, density.* (a) A charge  $-300e$  is uniformly distributed along a circular arc of radius 4.00 cm, which subtends an angle of  $40^\circ$ . What is the linear charge density along the arc? (b) A charge  $-300e$  is uniformly distributed over one face of a circular disk of radius 2.00 cm. What is the surface charge density over that face? (c) A charge  $-300e$  is uniformly distributed over the surface of a sphere of radius 2.00 cm. What is the surface charge density over that surface? (d) A charge  $-300e$  is uniformly spread through the volume of a sphere of radius 2.00 cm. What is the volume charge density in that sphere?

- 23 E** Figure 22.26 shows two parallel nonconducting rings with their central axes along a common line. Ring 1 has uniform charge  $q_1$  and radius  $R$ ; ring 2 has uniform charge  $q_2$  and the same radius  $R$ . The rings are separated by distance  $d = 3.00R$ . The net electric

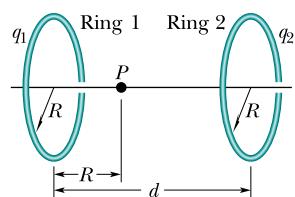


Figure 22.26 Problem 23.

field at point  $P$  on the common line, at distance  $R$  from ring 1, is zero. What is the ratio  $q_1/q_2$ ?

- 24 M CALC** A thin nonconducting rod with a uniform distribution of positive charge  $Q$  is bent into a complete circle of radius  $R$  (Fig. 22.27). The central perpendicular axis through the ring is a  $z$  axis, with the origin at the center of the ring. What is the magnitude of the electric field due to the rod at (a)  $z = 0$  and (b)  $z = \infty$ ? (c) In terms of  $R$ , at what positive value of  $z$  is that magnitude maximum? (d) If  $R = 2.00$  cm and  $Q = 4.00 \mu\text{C}$ , what is the maximum magnitude?

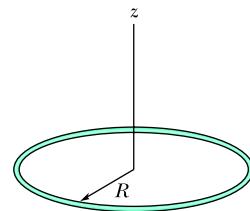


Figure 22.27 Problem 24.

- 25 M** Figure 22.28 shows three circular arcs centered on the origin of a coordinate system. On each arc, the uniformly distributed charge is given in terms of  $Q = 2.00 \mu\text{C}$ . The radii are given in terms of  $R = 10.0$  cm. What are the (a) magnitude and (b) direction (relative to the positive  $x$  direction) of the net electric field at the origin due to the arcs?

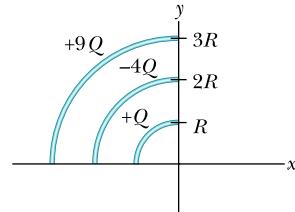


Figure 22.28 Problem 25.

- 26 M GO** In Fig. 22.29, a thin glass rod forms a semicircle of radius  $r = 5.00$  cm. Charge is uniformly distributed along the rod, with  $+q = 4.50 \text{ pC}$  in the upper half and  $-q = -4.50 \text{ pC}$  in the lower half. What are the (a) magnitude and (b) direction (relative to the positive direction of the  $x$  axis) of the electric field  $\vec{E}$  at  $P$ , the center of the semicircle?

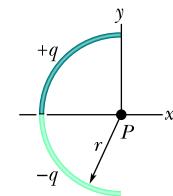


Figure 22.29  
Problem 26.

- 27 M GO** In Fig. 22.30, two curved plastic rods, one of charge  $+q$  and the other of charge  $-q$ , form a circle of radius  $R = 8.50$  cm in an  $xy$  plane. The  $x$  axis passes through both of the connecting points, and the charge is distributed uniformly on both rods. If  $q = 15.0 \text{ pC}$ , what are the (a) magnitude and (b) direction (relative to the positive direction of the  $x$  axis) of the electric field  $\vec{E}$  produced at  $P$ , the center of the circle?

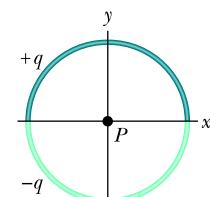


Figure 22.30  
Problem 27.

- 28 M CALC** Charge is uniformly distributed around a ring of radius  $R = 2.40$  cm, and the resulting electric field magnitude  $E$  is measured along the ring's central axis (perpendicular to the plane of the ring). At what distance from the ring's center is  $E$  maximum?

- 29 M GO** Figure 22.31a shows a nonconducting rod with a uniformly distributed charge  $+Q$ . The rod forms a half-circle with radius  $R$  and produces an electric field of magnitude  $E_{\text{arc}}$  at its center of curvature  $P$ . If the arc is collapsed to a point at distance

*R* from *P* (Fig. 22.31*b*), by what factor is the magnitude of the electric field at *P* multiplied?

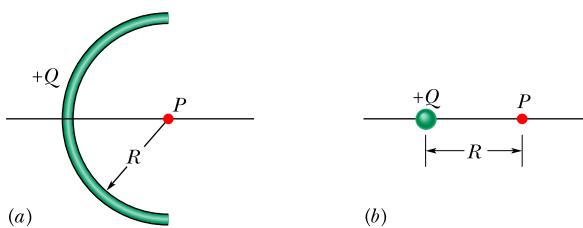


Figure 22.31 Problem 29.

**30 M GO** Figure 22.32 shows two concentric rings, of radii *R* and  $R' = 3.00R$ , that lie on the same plane. Point *P* lies on the central *z* axis, at distance  $D = 2.00R$  from the center of the rings. The smaller ring has uniformly distributed charge  $+Q$ . In terms of  $Q$ , what is the uniformly distributed charge on the larger ring if the net electric field at *P* is zero?

**31 M CALC SSM** In Fig. 22.33, a nonconducting rod of length  $L = 8.15$  cm has a charge  $-q = -4.23 \text{ fC}$  uniformly distributed along its length. (a) What is the linear charge density of the rod? What are the (b) magnitude and (c) direction (relative to the positive direction of the *x* axis) of the electric field produced at point *P*, at distance  $a = 12.0$  cm from the rod? What is the electric field magnitude produced at distance  $a = 50$  m by (d) the rod and (e) a particle of charge  $-q = -4.23 \text{ fC}$  that we use to replace the rod? (At that distance, the rod “looks” like a particle.)

**32 H CALC GO** In Fig. 22.34, positive charge  $q = 7.81 \text{ pC}$  is spread uniformly along a thin nonconducting rod of length  $L = 14.5$  cm. What are the (a) magnitude and (b) direction (relative to the positive direction of the *x* axis) of the electric field produced at point *P*, at distance  $R = 6.00$  cm from the rod along its perpendicular bisector?

**33 H CALC GO** In Fig. 22.35, a “semi-infinite” nonconducting rod (that is, infinite in one direction only) has uniform linear charge density  $\lambda$ . Show that the electric field  $\vec{E}_p$  at point *P* makes an angle of  $45^\circ$  with the rod and that this result is independent of the distance *R*. (*Hint:* Separately find the component of  $\vec{E}_p$  parallel to the rod and the component perpendicular to the rod.)

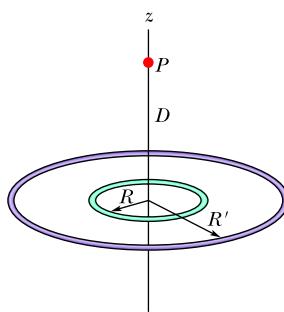


Figure 22.32 Problem 30.

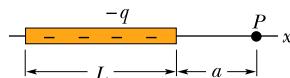


Figure 22.33 Problem 31.

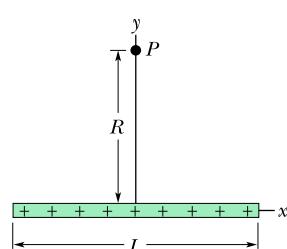


Figure 22.34 Problem 32.

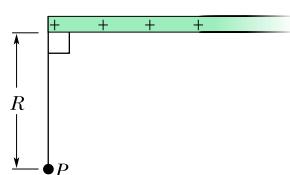


Figure 22.35 Problem 33.

### Module 22.5 The Electric Field Due to a Charged Disk

**34 E** A disk of radius 2.5 cm has a surface charge density of  $5.3 \mu\text{C/m}^2$  on its upper face. What is the magnitude of the electric field produced by the disk at a point on its central axis at distance  $z = 12$  cm from the disk?

**35 E SSM** At what distance along the central perpendicular axis of a uniformly charged plastic disk of radius 0.600 m is the magnitude of the electric field equal to one-half the magnitude of the field at the center of the surface of the disk?

**36 M** A circular plastic disk with radius  $R = 2.00$  cm has a uniformly distributed charge  $Q = +(2.00 \times 10^6)e$  on one face. A circular ring of width  $30 \mu\text{m}$  is centered on that face, with the center of that width at radius  $r = 0.50$  cm. In coulombs, what charge is contained within the width of the ring?

**37 M** Suppose you design an apparatus in which a uniformly charged disk of radius *R* is to produce an electric field. The field magnitude is most important along the central perpendicular axis of the disk, at a point *P* at distance  $2.00R$  from the disk (Fig. 22.36*a*). Cost analysis suggests that you switch to a ring of the same outer radius *R* but with inner radius  $R/2.00$  (Fig.

22.36*b*). Assume that the ring will have the same surface charge density as the original disk. If you switch to the ring, by what percentage will you decrease the electric field magnitude at *P*?

**38 M** Figure 22.37*a* shows a circular disk that is uniformly charged. The central *z* axis is perpendicular to the disk face, with the origin at the disk. Figure 22.37*b* gives the magnitude of the electric field along that axis in terms of the maximum magnitude  $E_m$  at the disk surface. The *z* axis scale is set by  $z_s = 8.0$  cm. What is the radius of the disk?

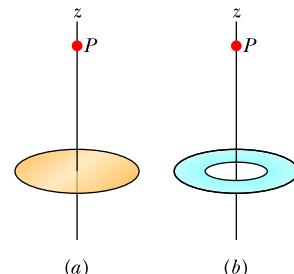


Figure 22.36 Problem 37.

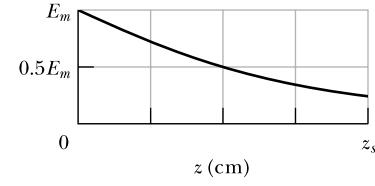
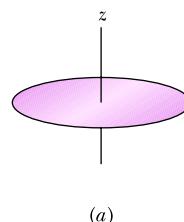


Figure 22.37 Problem 38.

### Module 22.6 A Point Charge in an Electric Field

**39 E** In Millikan’s experiment, an oil drop of radius  $1.64 \mu\text{m}$  and density  $0.851 \text{ g/cm}^3$  is suspended in chamber C (Fig. 22.6.1) when a downward electric field of  $1.92 \times 10^5 \text{ N/C}$  is applied. Find the charge on the drop, in terms of  $e$ .

**40 E GO** An electron with a speed of  $5.00 \times 10^8 \text{ cm/s}$  enters an electric field of magnitude  $1.00 \times 10^3 \text{ N/C}$ , traveling along a field line in the direction that retards its motion. (a) How far will the electron travel in the field before stopping momentarily, and (b) how much time will have elapsed? (c) If the region containing the electric field is 8.00 mm long (too short for the electron to stop within it), what fraction of the electron’s initial kinetic energy will be lost in that region?

**41 E SSM** A charged cloud system produces an electric field in the air near Earth's surface. A particle of charge  $-2.0 \times 10^{-9}$  C is acted on by a downward electrostatic force of  $3.0 \times 10^{-6}$  N when placed in this field. (a) What is the magnitude of the electric field? What are the (b) magnitude and (c) direction of the electrostatic force  $\vec{F}_e$  on a proton placed in this field? (d) What is the magnitude of the gravitational force  $\vec{F}_g$  on the proton? (e) What is the ratio  $F_e/F_g$  in this case?

**42 E** Humid air breaks down (its molecules become ionized) in an electric field of  $3.0 \times 10^6$  N/C. In that field, what is the magnitude of the electrostatic force on (a) an electron and (b) an ion with a single electron missing?

**43 E SSM** An electron is released from rest in a uniform electric field of magnitude  $2.00 \times 10^4$  N/C. Calculate the acceleration of the electron. (Ignore gravitation.)

**44 E** An alpha particle (the nucleus of a helium atom) has a mass of  $6.64 \times 10^{-27}$  kg and a charge of  $+2e$ . What are the (a) magnitude and (b) direction of the electric field that will balance the gravitational force on the particle?

**45 E** An electron on the axis of an electric dipole is 25 nm from the center of the dipole. What is the magnitude of the electrostatic force on the electron if the dipole moment is  $3.6 \times 10^{-29}$  C·m? Assume that 25 nm is much larger than the separation of the charged particles that form the dipole.

**46 E** An electron is accelerated eastward at  $1.80 \times 10^9$  m/s<sup>2</sup> by an electric field. Determine the field (a) magnitude and (b) direction.

**47 E SSM** Beams of high-speed protons can be produced in "guns" using electric fields to accelerate the protons. (a) What acceleration would a proton experience if the gun's electric field were  $2.00 \times 10^4$  N/C? (b) What speed would the proton attain if the field accelerated the proton through a distance of 1.00 cm?

**48 M** In Fig. 22.38, an electron (e) is to be released from rest on the central axis of a uniformly charged disk of radius  $R$ . The surface charge density on the disk is  $+4.00 \mu\text{C}/\text{m}^2$ . What is the magnitude of the electron's initial acceleration if it is released at a distance (a)  $R$ , (b)  $R/100$ , and (c)  $R/1000$  from the center of the disk? (d) Why does the acceleration magnitude increase only slightly as the release point is moved closer to the disk?

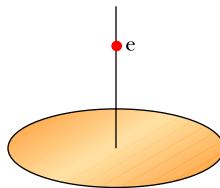


Figure 22.38  
Problem 48.

**49 M** A 10.0 g block with a charge of  $+8.00 \times 10^{-5}$  C is placed in an electric field  $\vec{E} = (3000\hat{i} - 600\hat{j})$  N/C. What are the (a) magnitude and (b) direction (relative to the positive direction of the  $x$  axis) of the electrostatic force on the block? If the block is released from rest at the origin at time  $t = 0$ , what are its (c)  $x$  and (d)  $y$  coordinates at  $t = 3.00$  s?

**50 M** At some instant the velocity components of an electron moving between two charged parallel plates are  $v_x = 1.5 \times 10^5$  m/s and  $v_y = 3.0 \times 10^3$  m/s. Suppose the electric field between the plates is uniform and given by  $\vec{E} = (120 \text{ N/C})\hat{j}$ . In unit-vector notation, what are (a) the electron's acceleration in that field and (b) the electron's velocity when its  $x$  coordinate has changed by 2.0 cm?

**51 M BIO FCP** Assume that a honeybee is a sphere of diameter 1.000 cm with a charge of  $+45.0 \mu\text{C}$  uniformly spread over its

surface. Assume also that a spherical pollen grain of diameter  $40.0 \mu\text{m}$  is electrically held on the surface of the bee because the bee's charge induces a charge of  $-1.00 \mu\text{C}$  on the near side of the grain and a charge of  $+1.00 \mu\text{C}$  on the far side. (a) What is the magnitude of the net electrostatic force on the grain due to the bee? Next, assume that the bee brings the grain to a distance of 1.000 mm from the tip of a flower's stigma and that the tip is a particle of charge  $-45.0 \mu\text{C}$ . (b) What is the magnitude of the net electrostatic force on the grain due to the stigma? (c) Does the grain remain on the bee or does it move to the stigma?

**52 M** An electron enters a region of uniform electric field with an initial velocity of 40 km/s in the same direction as the electric field, which has magnitude  $E = 50$  N/C. (a) What is the speed of the electron 1.5 ns after entering this region? (b) How far does the electron travel during the 1.5 ns interval?

**53 M GO** Two large parallel copper plates are 5.0 cm apart and have a uniform electric field between them as depicted in Fig. 22.39. An electron is released from the negative plate at the same time that a proton is released from the positive plate. Neglect the force of the particles on each other and find their distance from the positive plate when they pass each other. (Does it surprise you that you need not know the electric field to solve this problem?)

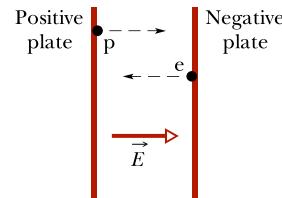


Figure 22.39 Problem 53.

**54 M GO** In Fig. 22.40, an electron is shot at an initial speed of  $v_0 = 2.00 \times 10^6$  m/s, at angle  $\theta_0 = 40.0^\circ$  from an  $x$  axis. It moves through a uniform electric field  $\vec{E} = (5.00 \text{ N/C})\hat{j}$ . A screen for detecting electrons is positioned parallel to the  $y$  axis, at distance  $x = 3.00$  m. In unit-vector notation, what is the velocity of the electron when it hits the screen?

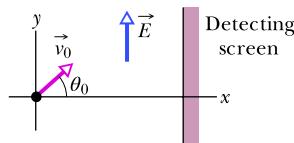


Figure 22.40 Problem 54.

**55 M** A uniform electric field exists in a region between two oppositely charged plates. An electron is released from rest at the surface of the negatively charged plate and strikes the surface of the opposite plate, 2.0 cm away, in a time  $1.5 \times 10^{-8}$  s. (a) What is the speed of the electron as it strikes the second plate? (b) What is the magnitude of the electric field  $\vec{E}$ ?

### Module 22.7 A Dipole in an Electric Field

**56 E** An electric dipole consists of charges  $+2e$  and  $-2e$  separated by 0.78 nm. It is in an electric field of strength  $3.4 \times 10^6$  N/C. Calculate the magnitude of the torque on the dipole when the dipole moment is (a) parallel to, (b) perpendicular to, and (c) antiparallel to the electric field.

**57 E SSM** An electric dipole consisting of charges of magnitude  $1.50 \mu\text{C}$  separated by  $6.20 \mu\text{m}$  is in an electric field of strength 1100 N/C. What are (a) the magnitude of the electric dipole moment and (b) the difference between the potential energies for dipole orientations parallel and antiparallel to  $\vec{E}$ ?

**58 M** A certain electric dipole is placed in a uniform electric field  $\vec{E}$  of magnitude 20 N/C. Figure 22.41 gives the potential

energy  $U$  of the dipole versus the angle  $\theta$  between  $\vec{E}$  and the dipole moment  $\vec{p}$ . The vertical axis scale is set by  $U_s = 100 \times 10^{-28} \text{ J}$ . What is the magnitude of  $\vec{p}$ ?

**59 M** How much work is required to turn an electric dipole  $180^\circ$  in a uniform electric field of magnitude  $E = 46.0 \text{ N/C}$  if the dipole moment has a magnitude of  $p = 3.02 \times 10^{-25} \text{ C}\cdot\text{m}$  and the initial angle is  $64^\circ$ ?

**60 M** A certain electric dipole is placed in a uniform electric field  $\vec{E}$  of magnitude  $40 \text{ N/C}$ . Figure 22.42 gives the magnitude  $\tau$  of the torque on the dipole versus the angle  $\theta$  between field  $\vec{E}$  and the dipole moment  $\vec{p}$ . The vertical axis scale is set by  $\tau_s = 100 \times 10^{-28} \text{ N}\cdot\text{m}$ . What is the magnitude of  $\vec{p}$ ?

**61 M** Find an expression for the oscillation frequency of an electric dipole of dipole moment  $\vec{p}$  and rotational inertia  $I$  for small amplitudes of oscillation about its equilibrium position in a uniform electric field of magnitude  $E$ .

### Additional Problems

**62** (a) What is the magnitude of an electron's acceleration in a uniform electric field of magnitude  $1.40 \times 10^6 \text{ N/C}$ ? (b) How long would the electron take, starting from rest, to attain one-tenth the speed of light? (c) How far would it travel in that time?

**63** A spherical water drop  $1.20 \mu\text{m}$  in diameter is suspended in calm air due to a downward-directed atmospheric electric field of magnitude  $E = 462 \text{ N/C}$ . (a) What is the magnitude of the gravitational force on the drop? (b) How many excess electrons does it have?

**64** Three particles, each with positive charge  $Q$ , form an equilateral triangle, with each side of length  $d$ . What is the magnitude of the electric field produced by the particles at the midpoint of any side?

**65** In Fig. 22.43a, a particle of charge  $+Q$  produces an electric field of magnitude  $E_{\text{part}}$  at point  $P$ , at distance  $R$  from the particle. In Fig. 22.43b, that same amount of charge is spread uniformly along a circular arc that has radius  $R$  and subtends an angle  $\theta$ . The charge on the arc produces an electric field of magnitude  $E_{\text{arc}}$  at its center of curvature  $P$ . For what value of  $\theta$  does  $E_{\text{arc}} = 0.500 E_{\text{part}}$ ? (Hint: You will probably resort to a graphical solution.)

**66** A proton and an electron form two corners of an equilateral triangle of side length  $2.0 \times 10^{-6} \text{ m}$ . What is the magnitude of the net electric field these two particles produce at the third corner?

**67 CALC** A charge (uniform linear density =  $9.0 \text{ nC/m}$ ) lies on a string that is stretched along an  $x$  axis from  $x = 0$  to  $x = 3.0 \text{ m}$ .

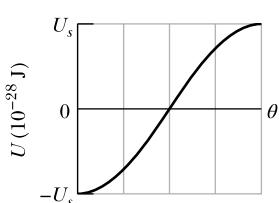


Figure 22.41 Problem 58.

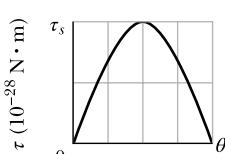


Figure 22.42 Problem 60.

Determine the magnitude of the electric field at  $x = 4.0 \text{ m}$  on the  $x$  axis.

**68** In Fig. 22.44, eight particles form a square in which distance  $d = 2.0 \text{ cm}$ . The charges are  $q_1 = +3e$ ,  $q_2 = +e$ ,  $q_3 = -5e$ ,  $q_4 = -2e$ ,  $q_5 = +3e$ ,  $q_6 = +e$ ,  $q_7 = -5e$ , and  $q_8 = +e$ . In unit-vector notation, what is the net electric field at the square's center?

**69** Two particles, each with a charge of magnitude  $12 \text{ nC}$ , are at two of the vertices of an equilateral triangle with edge length  $2.0 \text{ m}$ . What is the magnitude of the electric field at the third vertex if (a) both charges are positive and (b) one charge is positive and the other is negative?

**70** The following table gives the charge seen by Millikan at different times on a single drop in his experiment. From the data, calculate the elementary charge  $e$ .

$6.563 \times 10^{-19} \text{ C}$	$13.13 \times 10^{-19} \text{ C}$	$19.71 \times 10^{-19} \text{ C}$
$8.204 \times 10^{-19} \text{ C}$	$16.48 \times 10^{-19} \text{ C}$	$22.89 \times 10^{-19} \text{ C}$
$11.50 \times 10^{-19} \text{ C}$	$18.08 \times 10^{-19} \text{ C}$	$26.13 \times 10^{-19} \text{ C}$

**71** A charge of  $20 \text{ nC}$  is uniformly distributed along a straight rod of length  $4.0 \text{ m}$  that is bent into a circular arc with a radius of  $2.0 \text{ m}$ . What is the magnitude of the electric field at the center of curvature of the arc?

**72** An electron is constrained to the central axis of the ring of charge of radius  $R$  in Fig. 22.4.1, with  $z \ll R$ . Show that the electrostatic force on the electron can cause it to oscillate through the ring center with an angular frequency

$$\omega = \sqrt{\frac{eq}{4\pi\epsilon_0 m R^3}}$$

where  $q$  is the ring's charge and  $m$  is the electron's mass.

**73 SSM** The electric field in an  $xy$  plane produced by a positively charged particle is  $7.2(4.0\hat{i} + 3.0\hat{j}) \text{ N/C}$  at the point  $(3.0, 3.0) \text{ cm}$  and  $100\hat{i} \text{ N/C}$  at the point  $(2.0, 0) \text{ cm}$ . What are the (a)  $x$  and (b)  $y$  coordinates of the particle? (c) What is the charge of the particle?

**74** (a) What total (excess) charge  $q$  must the disk in Fig. 22.5.1 have for the electric field on the surface of the disk at its center to have magnitude  $3.0 \times 10^6 \text{ N/C}$ , the  $E$  value at which air breaks down electrically, producing sparks? Take the disk radius as  $2.5 \text{ cm}$ . (b) Suppose each surface atom has an effective cross-sectional area of  $0.015 \text{ nm}^2$ . How many atoms are needed to make up the disk surface? (c) The charge calculated in (a) results from some of the surface atoms having one excess electron. What fraction of these atoms must be so charged?

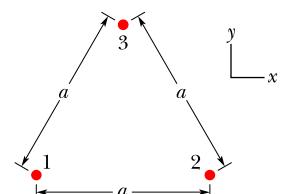
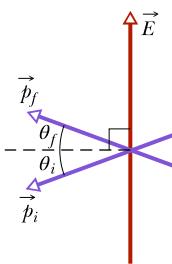


Figure 22.45 Problem 75.

charge  $Q$ ) form an equilateral triangle of edge length  $a$ . For what value of  $Q$  (both sign and magnitude) does the net electric field produced by the particles at the center of the triangle vanish?

**76** In Fig. 22.46, an electric dipole swings from an initial orientation  $i$  ( $\theta_i = 20.0^\circ$ ) to a final orientation  $f$  ( $\theta_f = 20.0^\circ$ ) in a uniform external electric field  $\vec{E}$ . The electric dipole moment is  $1.60 \times 10^{-27} \text{ C} \cdot \text{m}$ ; the field magnitude is  $3.00 \times 10^6 \text{ N/C}$ . What is the change in the dipole's potential energy?



**Figure 22.46**  
Problem 76.

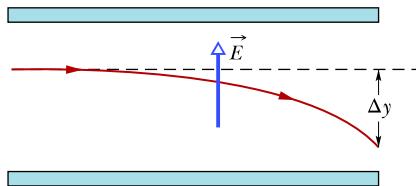
**77** A particle of charge  $-q_1$  is at the origin of an  $x$  axis. (a) At what location on the axis should a particle of charge  $-4q_1$  be placed so that the net electric field is zero at  $x = 2.0 \text{ mm}$  on the axis? (b) If, instead, a particle of charge  $+4q_1$  is placed at that location, what is the direction (relative to the positive direction of the  $x$  axis) of the net electric field at  $x = 2.0 \text{ mm}$ ?

**78** Two particles, each of positive charge  $q$ , are fixed in place on a  $y$  axis, one at  $y = d$  and the other at  $y = -d$ . (a) Write an expression that gives the magnitude  $E$  of the net electric field at points on the  $x$  axis given by  $x = ad$ . (b) Graph  $E$  versus  $\alpha$  for the range  $0 < \alpha < 4$ . From the graph, determine the values of  $\alpha$  that give (c) the maximum value of  $E$  and (d) half the maximum value of  $E$ .

**79** *Water molecule dipole field.* A molecule of water vapor produces an electric field in the surrounding space as if it were an electric dipole. Its dipole moment has a magnitude  $p = 6.2 \times 10^{-30} \text{ C} \cdot \text{m}$ . What is the magnitude of the electric field at a distance  $z = 1.1 \text{ nm}$  from the molecule on its dipole axis, which is far compared to the charge separation in the dipole?

**80** *Dipole oscillation.* Find the frequency of oscillation of an electric dipole, of dipole moment  $p$  and rotational inertia  $I$ , for small amplitudes of oscillation about its equilibrium orientation in a uniform electric field of magnitude  $E$ .

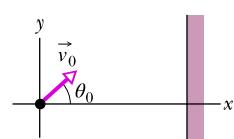
**81** *Conventional TV tube.* Early television sets depended on images being built up on the screen by the deflection of electrons directed toward the screen from the rear of the tube. Figure 22.47 shows such a deflection system. The length of the plates is  $3.0 \text{ cm}$  and the deflecting electric field between the two plates is  $1.0 \times 10^6 \text{ N/C}$  vertically upward. If an electron enters



**Figure 22.47** Problem 81.

the space between the plates with a horizontal speed of  $3.9 \times 10^7 \text{ m/s}$ , what is the vertical displacement of  $\Delta y$  at the end of the plates?

**82** *Electron shot at screen.* In Fig. 22.48, an electron is shot at an initial speed of  $v_0 = 7.00 \times 10^6 \text{ m/s}$ , at angle  $\theta_0 = 30.0^\circ$  from an  $x$  axis. It moves in a region with uniform electric field  $\vec{E} = (8.50 \text{ N/C})\hat{j}$ . A screen for detecting electrons is positioned parallel to the  $y$  axis, at distance  $x = 2.50 \text{ m}$ . (a) In unit-vector notation, what is the velocity of the electron when it hits the screen? (b) When the electron reaches the screen, is it still traveling in the  $+y$  direction and what is its kinetic energy?

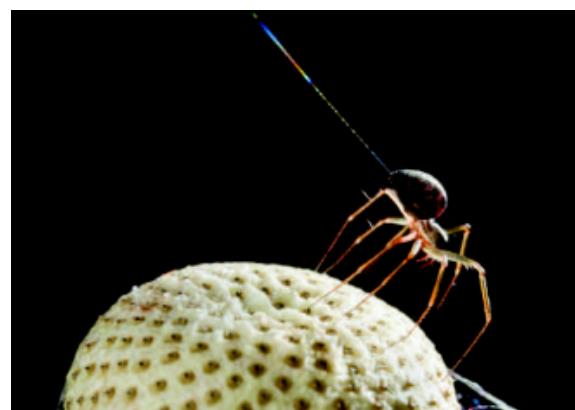


**Figure 22.48** Problem 82.

**83** *Double positive charge.* In Fig. 22.3.2, let both charges be positive with the same magnitude. Assuming  $z \gg d$ , show that the electric field at point  $P$  is given by

$$E = \frac{1}{4\pi\epsilon_0} \frac{2q}{z^2}.$$

**84** **BIO** *Spiders ballooning in electric fields.* Some spiders disperse by a process known as ballooning. When they extrude silk threads (Fig. 22.49), the threads catch on a breeze that can carry a spider far away and up to several kilometers high. However, some spiders can balloon even on a calm day, as was recorded by Charles Darwin during his journey on the *Beagle*. When extruded, the nonconducting silk thread is negatively charged and thus can experience an electric force in the naturally occurring electric field in the atmosphere, especially near the sharp points on leaves, needles, and branch tips. Near those sharp points, the magnitude  $E$  of the field can be  $10 \text{ N/C}$ . (a) What is the minimum charge magnitude  $q$  needed on the silk if a  $0.95 \text{ mg}$  spider is to be lifted by the electric force due to a vertical field with that field magnitude, and (b) what is the corresponding minimum number  $n$  of electrons?



**Figure 22.49** Problem 84.

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