Calibration of the ATLAS electromagnetic calorimeter and measurement of the couplings of the (Brout-Englert-)Higgs boson in the diphoton channel

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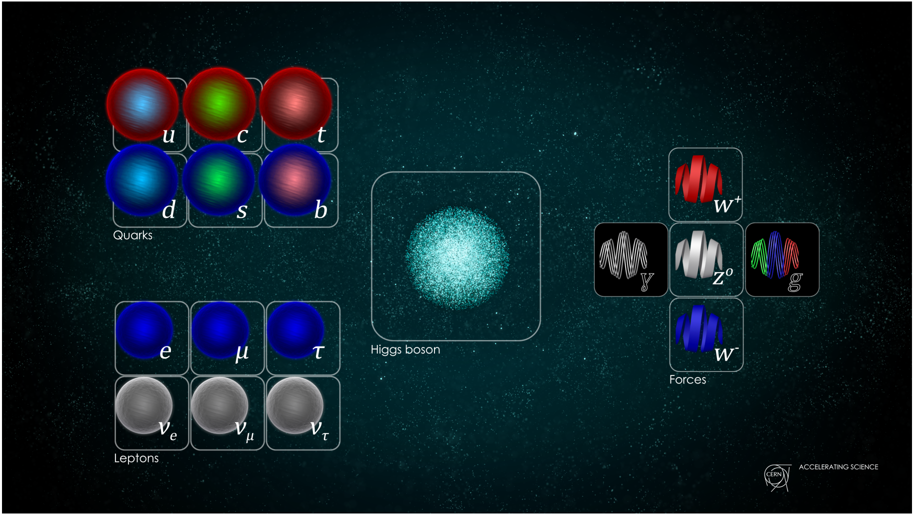
Introduction

1. The Standard Model of matter
2. Experimental conditions and data processing
3. Calibration of electromagnetic objects
4. Measurement of Higgs boson couplings

Particle content of matter

Over the XXth century, elementary particles have been organised into a well structured model.



A mathematical framework

Matter knowledge is embedded into a well defined mathematical framework based on a Lagrangian L.

$$L = \frac{m\dot{\vec{q}}^2}{2} - V(\vec{q}) \tag{1}$$

The dirac lagrangian describes a massive fermion field :

$$L = \bar{\psi}(i\not{\partial} - m)\psi \tag{2}$$

Imposing least action principle (similar to classical mechanic) lead to equations of motion :

$$\frac{\partial L}{\partial q} - \frac{d}{dt}\frac{\partial L}{\partial \dot{q}_i} = 0 \tag{3}$$

Gauge invariance

Symmetries are transformations which leave a system unchanged.
Imposing symmetries on a Lagrangian changes the theory it describes.

$$\psi(x) \rightarrow e^{i\alpha(x)}\psi(x) \tag{4}$$

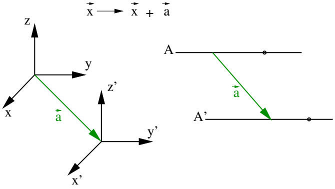
Derivative affects $e^{i\alpha}$

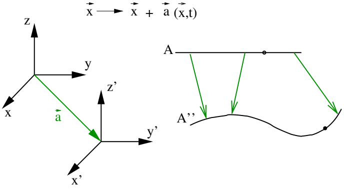
$\rightarrow$ **Invariance achieved by adding a field A_μ and changing L.**

$\rightarrow$ mass term for A ($m^2 A_\mu A^\mu$) is forbidden by symmetries

$$\partial_\mu \rightarrow D_\mu = \partial_\mu + ieA_\mu(x) \tag{5}$$

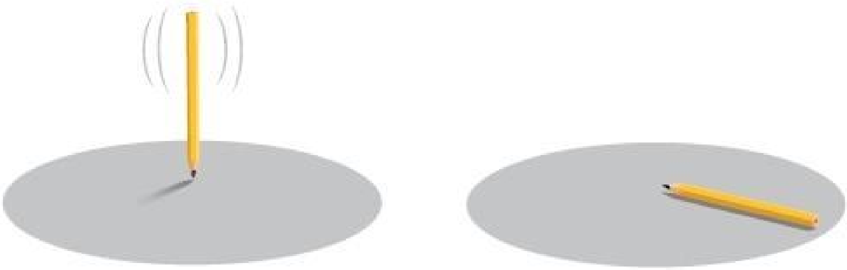
$$A_\mu \rightarrow A_\mu(x) - \frac{1}{e}\partial_\mu \alpha(x) \tag{6}$$





Spontaneous Symmetry Breaking

SSB describes a system for which its ground state has less symmetry than its Lagrangian.



- Unstable equilibrium has cylindrical symmetry
- Ground state (fallen pen) "has chosen" a direction. The cylindrical symmetry has been broken.

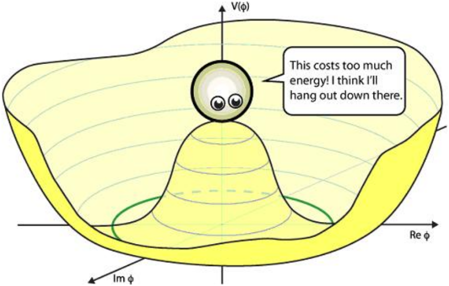
SSB in field theory

SSB is created by imposing a “mexican hat” potential on a field.

$$V(\phi) = \frac{1}{2}\mu^2\phi^*\phi + \frac{1}{4}\lambda(\phi^*\phi)^2 \quad (7)$$

with $\lambda > 0$ and $\mu^2 < 0$.

- Potential has rotational symmetry
- Ground state $|\Phi| = \sqrt{-\frac{\mu^2}{2\lambda^2}} = \frac{v}{\sqrt{2}}$ breaks symmetry.
- Describe a massless and a massive ($m^2 = v^2\lambda$) bosons.



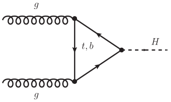
The Standard Model

The SM is composed of :

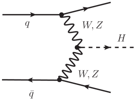
- Local gauge symmetries
  - $SU(3)_c$ for strong interaction. 8 gluons couple to quarks.
  - $SU(2)_L \times U(1)_Y$ for electroweak sector. Bosons $W^{\pm}$, Z and photon couple to quarks and leptons.
- SSB of $SU(2)_L$ by introduction of scalar field Φ
  - $(D_\mu\Phi)^\dagger(D^\mu\Phi)$ gives mass to $W^{\pm}$ and Z.
  - A physical and massive degree of freedom : the (Brout-Englert)-Higgs boson H.
  - Yukawa coupling gives mass to fermions.

Higgs boson production

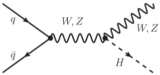
H boson predictions are function of its mass.



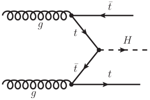
(a) Gluon fusion (ggH)



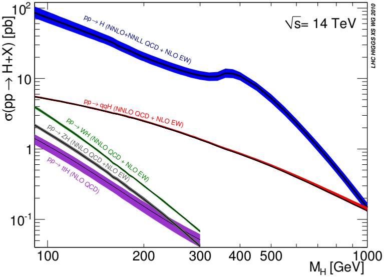
(b) Weak vector boson fusion (VBF)



(c) Higgsstrahung



(d) Associated production with $t\bar{t}$ pair (ttH)

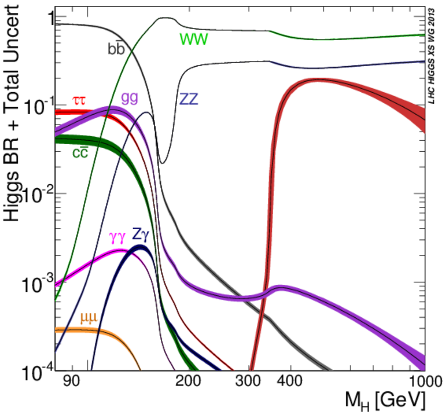


4 dominant production modes

- Gluon fusion ($ggH \simeq 86\%$) probes coupling to gluons through loop.
- Vector Boson Fusion probes direct coupling to electroweak bosons.
- Higgsstrahlung also probes $W\pm$ and Z couplings.
- Associated top production probes couplings to heaviest quark.

Higgs boson decays

A Higgs boson with a mass around 125 GeV opens a wide range of decay channels.



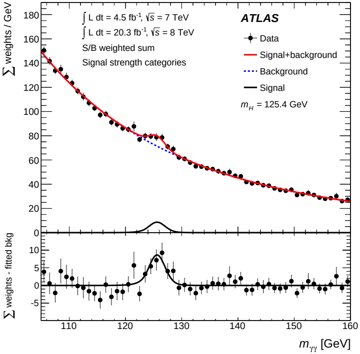
- $H \to bb$ (57 %) probes couplings to b quark. Difficult due to large hadronic background.
- $H \to \tau\tau$ probes couplings to heaviest lepton.
- $H \to VV$ $(V = W, Z)$ probes H boson couplings to EW bosons. Clean signature in leptonic decays of V but low statistics.
- $H \to \gamma\gamma$ probes H boson couplings to photon through loop. Large but smooth background.

H boson Status

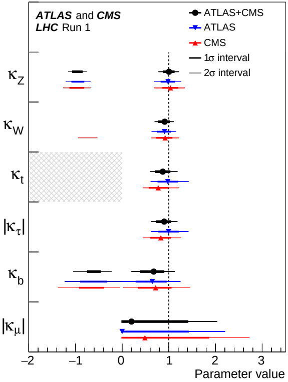
Run 1 of the LHC (2011/2012) allowed the observation of a Higgs like particles and its properties have been measured combining ATLAS+CMS.

mass measurement
PhysRevLett.114.191803

$m_H =$
$125.09 \pm 0.21(\text{stat}) \pm 0.11(\text{syst})$ GeV



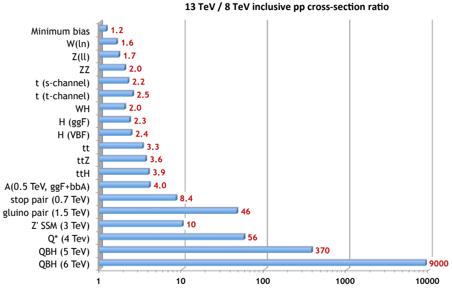
couplings $\kappa_i = \frac{\sigma_i^{exp}}{\sigma_i^{SM}}$
CERN-EP-2016-100



The measured properties are in agreement with the SM H boson.

Run 2 objectives

- LHC energy and luminosity increase
  $\rightarrow$ **30 times more Higgses are expected**
- With reduced statistical uncertainties
  $\rightarrow$ **need to reduce systematic uncertainties**
- Theory uncertainty reduced with ggH N^{3}LO calculation
- Resolution uncertainty dominant at run 1 for couplings
  $\rightarrow$ **Need to improve calibration resolution uncertainty**



Uncertainty group	$\sigma_\mu^{\text{syst.}}$
Theory (yield)	0.09
Experimental (yield)	0.02
Luminosity	0.03
MC statistics	< 0.01
Theory (migrations)	0.03
Experimental (migrations)	0.02
Resolution	0.07
Mass scale	0.02
Background shape	0.02

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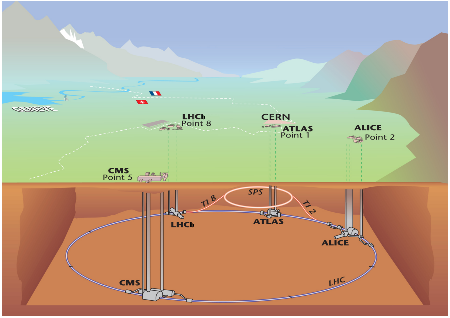
The Large Hadron Collider

The LHC aims at accelerating and colliding protons. Analysing debris of collisions allows to probe SM and/or beyond.

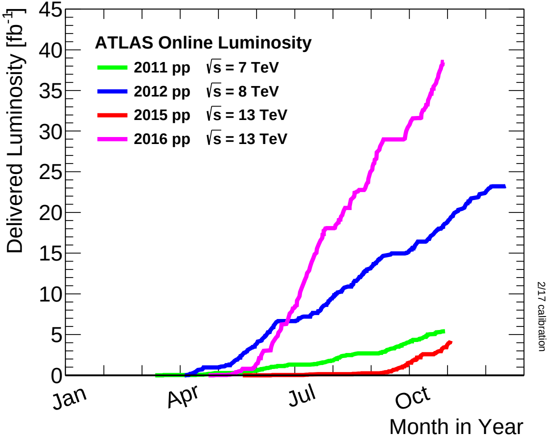
- located at Geneva.
- 27 km long.
- 100 m underground
- collision every 25 ns.
- $\sqrt{s} = 13$ TeV
- 4 collision points equipped with detectors : ALICE, ATLAS, CMS and LHCb.



LHC data taking condition

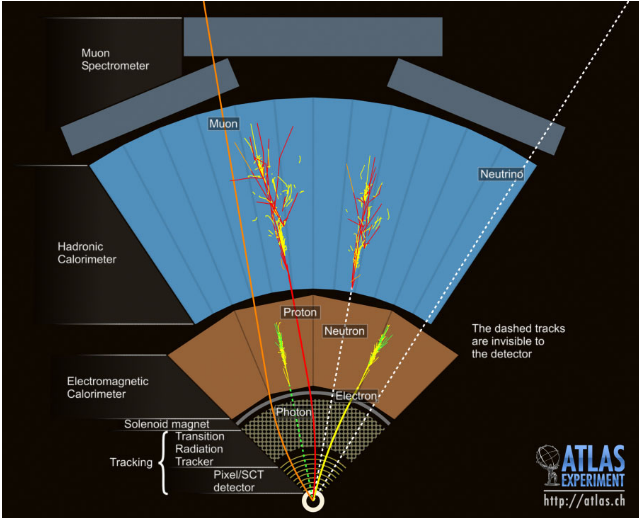
The collision conditions at LHC have significantly changed since its contruction.

- Major increase of integrated luminosity per year.
- Large increase in colliisions per bunch crossing.

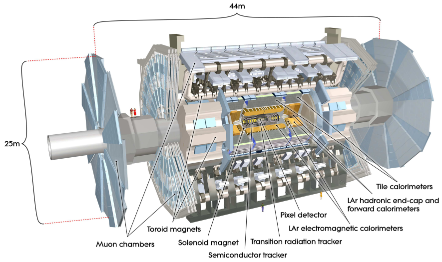


General purpose apparatus

A general purpose particle detector (ATLAS and CMS) is designed as layers with specific measurement goal.



ATLAS experiment



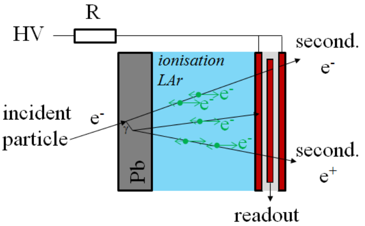
Performance goals of the ATLAS detector

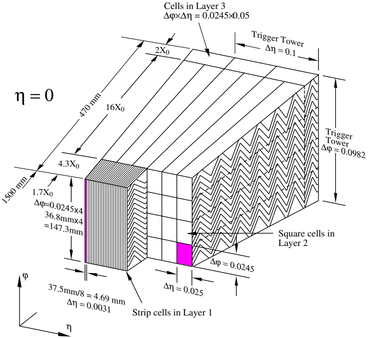
Detector component	Required resolution	η coverage	
		Measurement	Trigger
Tracking	$\sigma_{p_T}/p_T = 0.05\%\ p_T \oplus 1\%$	± 2.5	
EM calorimetry	$\sigma_E/E = 10\%/\sqrt{E} \oplus 0.7\%$	± 3.2	± 2.5
Hadronic calorimetry (jets)			
barrel and end-cap	$\sigma_E/E = 50\%/\sqrt{E} \oplus 3\%$	± 3.2	± 3.2
forward	$\sigma_E/E = 100\%/\sqrt{E} \oplus 10\%$	$3.1 < \lvert\eta\rvert < 4.9$	$3.1 < \lvert\eta\rvert < 4.9$
Muon spectrometer	σ_{p_T}/p_T=10% at p_T = 1 TeV	± 2.7	± 2.4

- **Large acceptance**
- **Radiation hard**

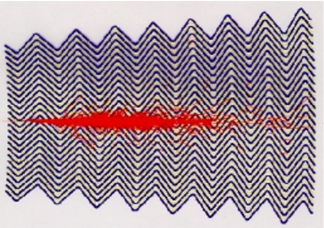
- **Silicon and TRT tracker in 2T magnetic field**
  Measure position and momentum of charged particles
- **Liquid argon electromagnetic calorimeter (LAr)**
  Measure energy of electrons and photons.
- **Scintillating tiles hadronic calorimeter**
  Measure energy of jets
- **Muon chambers**

Electromagnetic calorimeter (LAr)

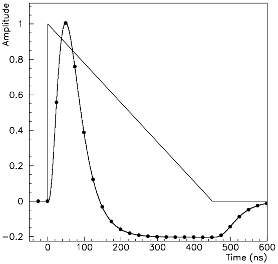




- 1.4m $< r <$ 2m
- Sampling calorimeter :
  - absorber : lead
  - active material : **Liquid Argon** (88K)
- **Accordion geometry** gives uniformity and hermeticity along ϕ.
- **Longitudinally segmented** for pion discrimination



Energy measurement in LAr



- **Signal drift time** ($\sim$ 600ns) **too long** for collisions every 25ns (pile-up).
- Analog signal pass through an **bipolar filter** to reduce signal time. Shape optimize signal over pileup and electronic noise.
- ADC sampling every 25ns (4 points are kept).
- Energy computed using **calibration constants and optimal filtering of the samples**.

$$E_{cell} = \underbrace{\sum_{i=1}^{n_{samples}} a_i(s_i - ped)}_{ADC} \cdot G_{ADC \to DAC} \cdot \left(\frac{M_{phys}}{M_{calib}} \right)^{-1} \cdot F_{DAC \to \mu A} \cdot F_{\mu A \to MeV}$$

Reconstruction & Identification

Reconstruction links the energy deposit in detector cells to a **physical particle and its properties**.

- Divide the central part ($|\eta| = |ln(tan(\theta/2))| < 2.47$) into towers of size $\Delta\eta \times \Delta\phi = 0.25 \times 0.25$
- Sum energies from all cells and all layers of the tower
- Sliding window (3×5 towers) algorithm look for 2.5 GeV of transverse energy
- **Track matching and clustering** :
  - no track $\rightarrow$ photon $\rightarrow 3 \times 7$ cluster
  - track $\rightarrow$ electron $\rightarrow 3 \times 7$ cluster
  - conversion vertex $\rightarrow$ converted photon $\rightarrow 3 \times 7$ cluster

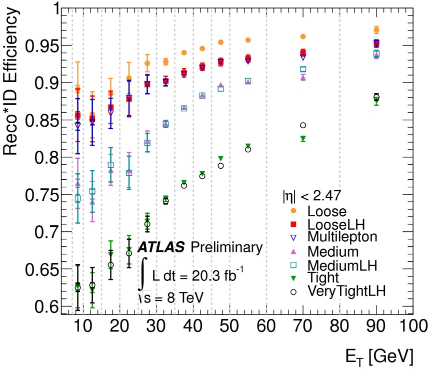
Identification is to separate prompt electrons from both jets and other electrons from either hadron decay or photon conversion.

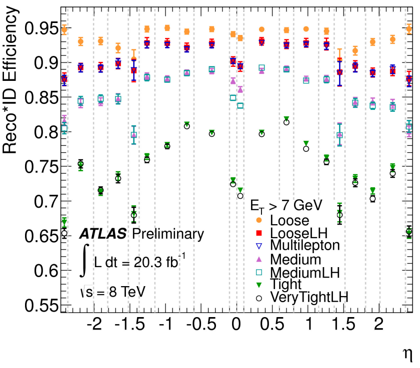
A multivariate likelihood method using 23 variables

of energy deposit and tracking is used.

Reconstruction & Identification efficiencies

Not all electrons pass the reconstruction and identification criteria.
3 menus with increasing purity (but deceasing efficiencies) are defined : loose, medium, tight. The efficiency of these procedures is given as a function of the p_T and $\eta = -ln(tan(\theta/2))$.





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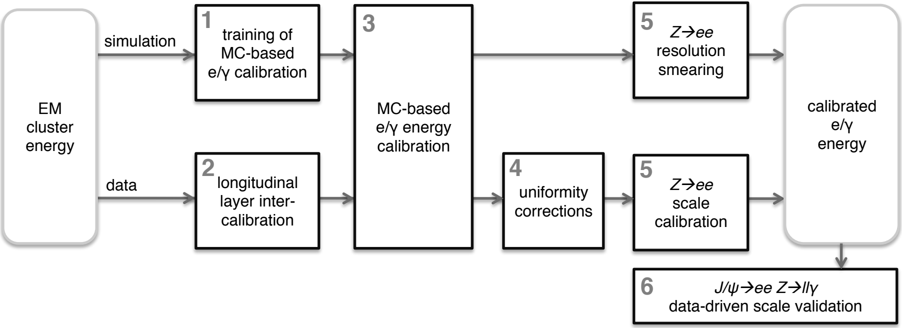
Full calibration

To reach the physics analyses, data and simulated reconstructed events must pass a calibration procedure. This procedures aim to correct the measured energy to **retrieve the true energy of the particle at the interaction point**.



Electrons and photons follow the same steps but with dedicated analyses.

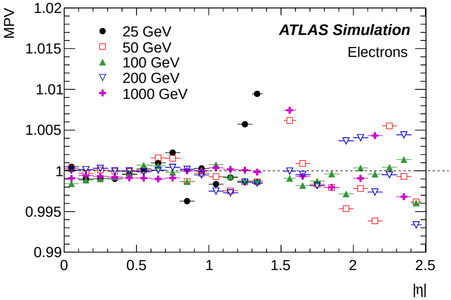
MVA calibration

- Simulated events are passed through a full GEANT4 simulation of the ATLAS detector.
- Events are then categorized in η and p_T bins, separately for electrons and photons.
- **A multivariate analysis (MVA) is performed to compute the true energy from detector observables**.

Plot shows most probable value (MVP) of E^{corr}/E^{true}.

MVA uses :

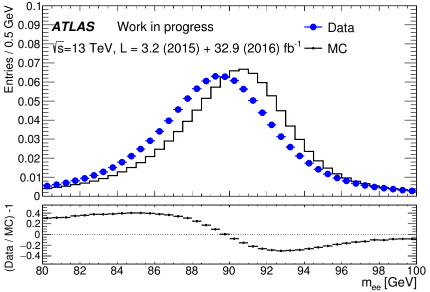
- Energies in all layers of the ECAL
- EM shower shape variables
- Barycenters of energy deposits



Energy scale factors

After MVA calibration, mass distribution of $Z \rightarrow ee$ for data and MC still have **discrepancy**.

A data-driven analysis is performed to match data to MC distribution (relative matching).



A correction, applied to both electrons of Z decay, is computed to shift the central value of data distribution :

energy scale factor (α)

$$E^{corr} = E^{meas}(1 + \alpha)$$

Resolution constant term

$$\frac{\sigma}{E} = \frac{a}{\sqrt{E}} \oplus \frac{b}{E} \oplus c$$

- a : sampling term (10%). Linked to the fluctuations of electromagnetic showers.
  Can be simulated.
- b/E : noise term ($350 cosh(\eta)$ MeV). Measured in dedicated runs.
- **c : constant term (0.7%)**. Must be measured on data.

We observe that data distribution is larger than MC. An **additional constant term (C)** is measure to enlarge MC up to the data width. Both MC electrons undergo the correction :

Resolution constant term (C)

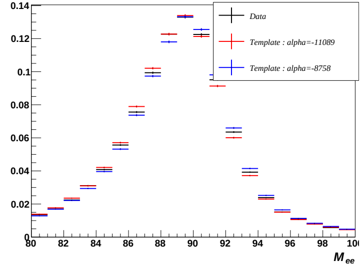
$$E^{corr} = E^{meas}(1 + N(0,1) * C)$$

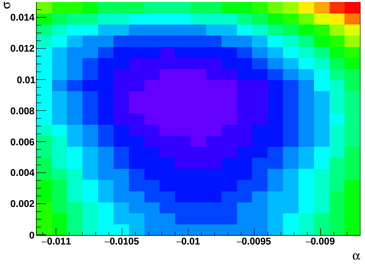
$N(0,1)$: a Gaussian distributed random number

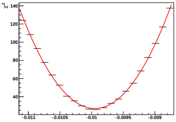
Template method

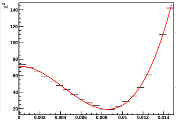
The template method is used to measure α and C simultaneously.

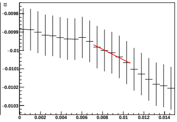
- Create distorded MC (templates) with test values of α and C.
- **Compute χ^2 between Z mass distribution of data and template**.
- **Fit the minimum of the χ^2 distribution** in the (α, C) plane.
- Fit performed in 2 steps of 1D fits :
  - fit $\chi^2 = f(\alpha)$ at constant C (lines) $\rightarrow (\alpha_{min}, \chi^2_{min})$.
  - fit $\chi^2_{min} = f(C) \rightarrow (C, \Delta C)$
  - project C in $\alpha_{min} = f(C)$, corresponding bin gives $(\alpha, \Delta\alpha)$.



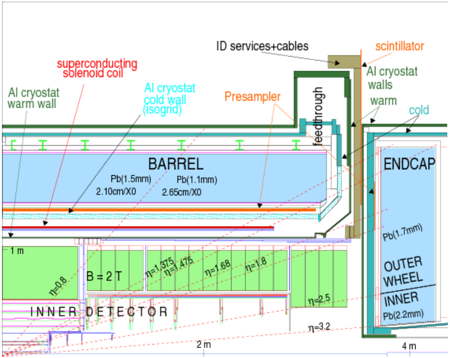








Detector splitting



- Detector is not uniform along η.
- To improve resolution, **calibration is performed in bin of η_{calo}.**
- 68 and 24 bins are used respectively for α and C.

0 0.1 0.2 0.3 0.4 0.5 0.6 0.7 0.8 0.9 1 1.1 1.2 1.285

1.37 1.42 1.47 1.51 1.55 1.59 1.63 1.6775 1.725

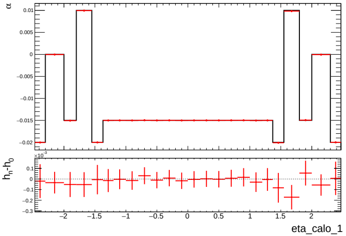
1.7625 1.8 1.9 2 2.05 2.1 2.2 2.3 2.35 2.4 2.435 2.47

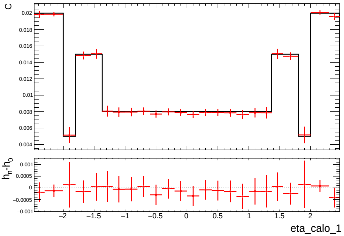
Electrons are labelled by their η bin, hence Z are labeled by the combination of electrons bins. **Scales are computed for each combination.**

Inversion Procedure

Obtaining **electron scales from Z scales** need the minimizations of the following χ^2's

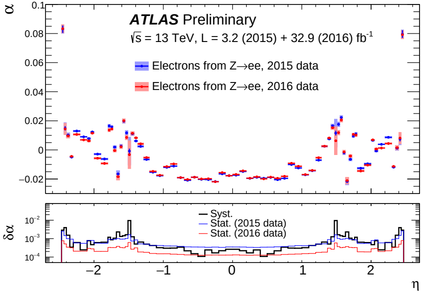
$$\chi^2 = \sum_{i,j\leq i} \frac{(\alpha_i+\alpha_j-2\alpha_{ij})^2}{(\Delta\alpha_{ij})^2}$$
$$\chi^2 = \sum_{i,j\leq i} \frac{\left(\sqrt{\frac{c_i^2+c_j^2}{2}}-c_{ij}\right)^2}{\Delta^2 c_{ij}} \tag{8}$$

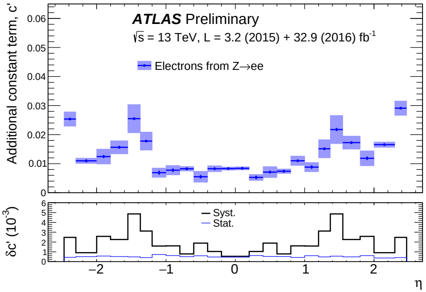




Run 2 results

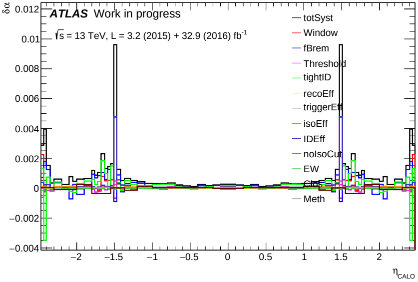
- α measured independently for each year.
- c measured on combined data.

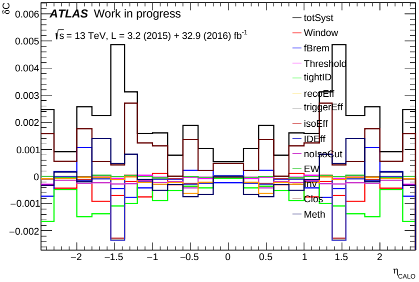




Uncertainties

12 (13) sources of uncertainties have been evaluated for α (c).



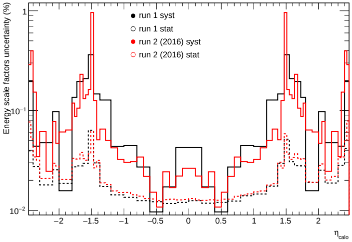


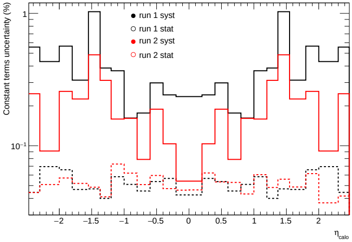
Main sources

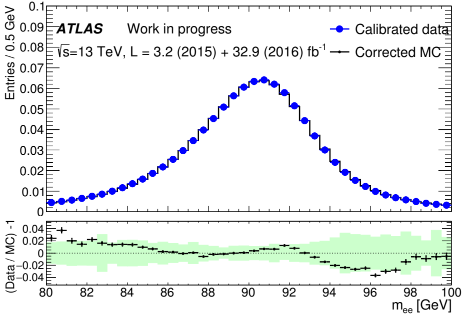
- Diff. btw tight and medium electrons
- Closure
- Bremsstrahlung impact on electron momentum

Runs comparison

Performances of run 2 in-situ calibration surpass run 1 results.



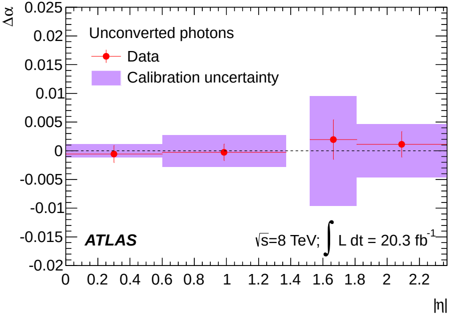


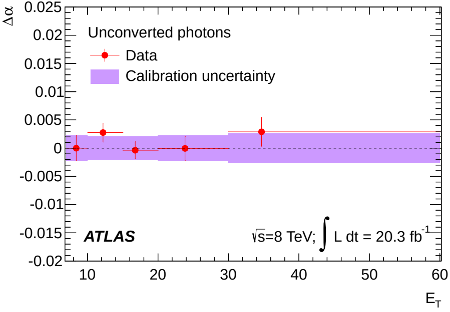


Photon correction

Electrons scale factors are also applied to photons. A residual scale factor ($\Delta\alpha$) is measured from $Z \to ll\gamma$.

No significant deviation observed.





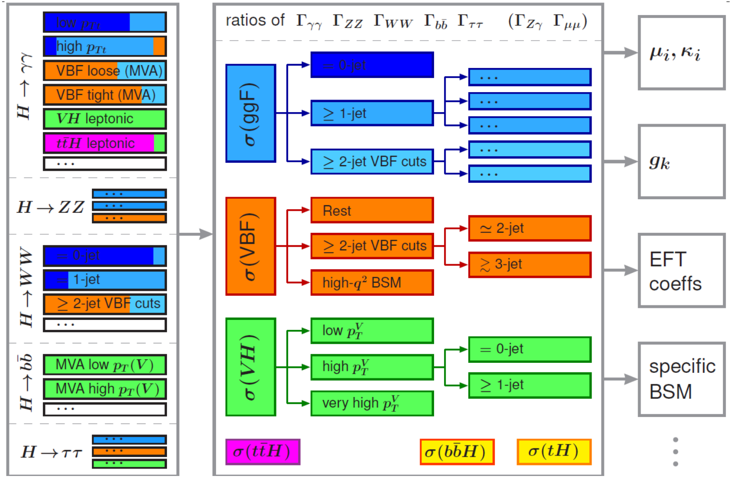
1 The Standard Model of matter

2 Experimental conditions and data processing

3 Calibration of electromagnetic objects

4 Measurement of Higgs boson couplings

Simplified Template Cross Section (STXS) framework



Replace signal strengths ($\mu = \frac{\sigma^{exp}}{\sigma^{th}}$) by cross-section in exclusive phase space regions.

Couplings measurement strategy

Inclusive selection

- 2 tight photons
- $\frac{p_T^{\gamma}}{m_{\gamma\gamma}} > 0.35(0.25)$
- $|\eta| \in [0, 1.137] \cup [1.52, 2.37]$
- $m_{\gamma\gamma} \in [105, 160]$ GeV

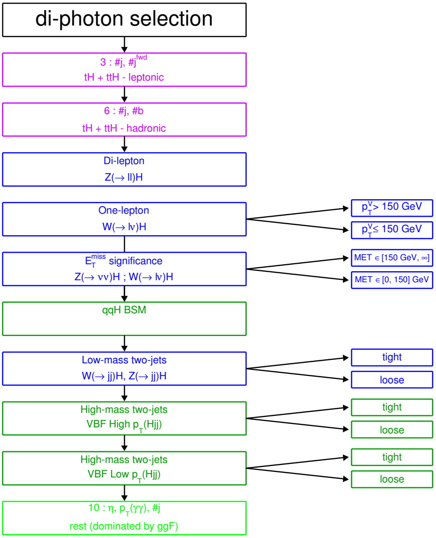
Dataset properties

- $\sim$ 330k events
- 42% signal efficiency
- $\simeq$ 1730 expected signal yield

Analysis strategy

- Define reconstructed categories targetting specific truth bin.
- Measure acceptance of each category wrt truth bins.
- Evaluate systematics effects on signal model.
- Combined fit of $m_{\gamma\gamma}$ distribution with signal+bkg model.

Reconstructed categories

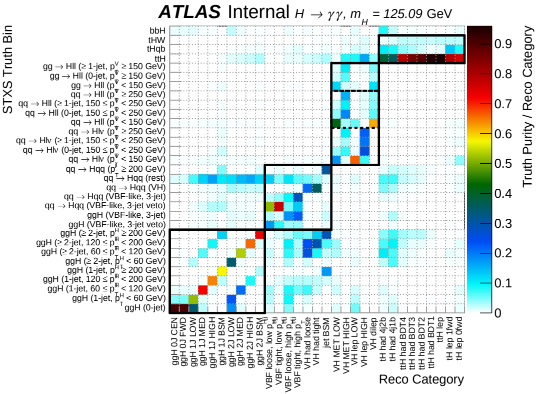


Optimised sensitivity to :

- rare processes
- truth bins

Category	Selection
tH lep 0fwd	$N_{\mathrm{lep}} = 1$, $N_{\mathrm{jets}}^{\mathrm{cen}} \leq 3$, $N_{\mathrm{b-tag}} \geq 1$, $N_{\mathrm{jets}}^{\mathrm{fwd}} = 0$ ($p_{\mathrm{T}}^{\mathrm{jet}} > 25$ GeV)
tH lep 1fwd	$N_{\mathrm{lep}} = 1$, $N_{\mathrm{jets}}^{\mathrm{cen}} \leq 4$, $N_{\mathrm{b-tag}} \geq 1$, $N_{\mathrm{jets}}^{\mathrm{fwd}} \geq 1$ ($p_{\mathrm{T}}^{\mathrm{jet}} > 25$ GeV)
ttH lep	$N_{\mathrm{lep}} \geq 1$, $N_{\mathrm{jets}}^{\mathrm{cen}} \geq 2$, $N_{\mathrm{b-tag}} \geq 1$, $Z_{\ell\ell}$ veto ($p_{\mathrm{T}}^{\mathrm{jet}} > 25$ GeV)
ttH had BDT1	$N_{\mathrm{lep}} = 0$, $N_{\mathrm{jets}} \geq 3$, $N_{\mathrm{b-tag}} \geq 1$, $\mathrm{BDT}_{\mathrm{ttH}} > 0.92$
ttH had BDT2	$N_{\mathrm{lep}} = 0$, $N_{\mathrm{jets}} \geq 3$, $N_{\mathrm{b-tag}} \geq 1$, $0.83 < \mathrm{BDT}_{\mathrm{ttH}} < 0.92$
ttH had BDT3	$N_{\mathrm{lep}} = 0$, $N_{\mathrm{jets}} \geq 3$, $N_{\mathrm{b-tag}} \geq 1$, $0.79 < \mathrm{BDT}_{\mathrm{ttH}} < 0.83$
ttH had BDT4	$N_{\mathrm{lep}} = 0$, $N_{\mathrm{jets}} \geq 3$, $N_{\mathrm{b-tag}} \geq 1$, $0.52 < \mathrm{BDT}_{\mathrm{ttH}} < 0.79$
tH had 4j1b	$N_{\mathrm{lep}} = 0$, $N_{\mathrm{jets}}^{\mathrm{cen}} = 4$, $N_{\mathrm{b-tag}} = 1$ ($p_{\mathrm{T}}^{\mathrm{jet}} > 25$ GeV)
tH had 4j2b	$N_{\mathrm{lep}} = 0$, $N_{\mathrm{jets}}^{\mathrm{cen}} = 4$, $N_{\mathrm{b-tag}} \geq 2$ ($p_{\mathrm{T}}^{\mathrm{jet}} > 25$ GeV)
VH dilep	$N_{\mathrm{lep}} \geq 2$, $70\,\mathrm{GeV} \leq m_{\ell\ell} \leq 110\,\mathrm{GeV}$
VH lep HIGH	$N_{\mathrm{lep}} = 1$, $\lvert m_{e\gamma} - 89\,\mathrm{GeV}\rvert > 5\,\mathrm{GeV}$, $p_{\mathrm{T}}^{l+E_{\mathrm{T}}^{\mathrm{miss}}} > 150\,\mathrm{GeV}$
VH lep LOW	$N_{\mathrm{lep}} = 1$, $\lvert m_{e\gamma} - 89\,\mathrm{GeV}\rvert > 5\,\mathrm{GeV}$, $p_{\mathrm{T}}^{l+E_{\mathrm{T}}^{\mathrm{miss}}} < 150\,\mathrm{GeV}$, $E_{\mathrm{T}}^{\mathrm{miss}}$ significance > 1
VH MET HIGH	$150\,\mathrm{GeV} < E_{\mathrm{T}}^{\mathrm{miss}} < 250\,\mathrm{GeV}$, $E_{\mathrm{T}}^{\mathrm{miss}}$ significance > 9 or $E_{\mathrm{T}}^{\mathrm{miss}} > 250\,\mathrm{GeV}$
VH MET LOW	$80\,\mathrm{GeV} < E_{\mathrm{T}}^{\mathrm{miss}} < 150\,\mathrm{GeV}$, $E_{\mathrm{T}}^{\mathrm{miss}}$ significance > 8
jet BSM	$p_{\mathrm{T,j1}} > 200\,\mathrm{GeV}$
VH had tight	$60\,\mathrm{GeV} < m_{\mathrm{jj}} < 120\,\mathrm{GeV}$, $\mathrm{BDT}_{\mathrm{VH}} > 0.78$
VH had loose	$60\,\mathrm{GeV} < m_{\mathrm{jj}} < 120\,\mathrm{GeV}$, $0.35 < \mathrm{BDT}_{\mathrm{VH}} < 0.78$
VBF tight, high p_T^{Hjj}	$\Delta\eta_{\mathrm{jj}} > 2$, $\lvert\eta_{\gamma\gamma} - 0.5(\eta_{\mathrm{j1}} + \eta_{\mathrm{j2}})\rvert < 5$, $p_T^{Hjj} > 25\,\mathrm{GeV}$, $\mathrm{BDT}_{\mathrm{VBF}} > 0.47$
VBF loose, high p_T^{Hjj}	$\Delta\eta_{\mathrm{jj}} > 2$, $\lvert\eta_{\gamma\gamma} - 0.5(\eta_{\mathrm{j1}} + \eta_{\mathrm{j2}})\rvert < 5$, $p_T^{Hjj} > 25\,\mathrm{GeV}$, $-0.32 < \mathrm{BDT}_{\mathrm{VBF}} < 0.47$
VBF tight, low p_T^{Hjj}	$\Delta\eta_{\mathrm{jj}} > 2$, $\lvert\eta_{\gamma\gamma} - 0.5(\eta_{\mathrm{j1}} + \eta_{\mathrm{j2}})\rvert < 5$, $p_T^{Hjj} < 25\,\mathrm{GeV}$, $\mathrm{BDT}_{\mathrm{VBF}} > 0.87$
VBF loose, low p_T^{Hjj}	$\Delta\eta_{\mathrm{jj}} > 2$, $\lvert\eta_{\gamma\gamma} - 0.5(\eta_{\mathrm{j1}} + \eta_{\mathrm{j2}})\rvert < 5$, $p_T^{Hjj} < 25\,\mathrm{GeV}$, $0.26 < \mathrm{BDT}_{\mathrm{VBF}} < 0.87$
ggH 2J BSM	≥ 2 jets, $p_{\mathrm{T}}^{\gamma\gamma} \geq 200\,\mathrm{GeV}$
ggH 2J HIGH	≥ 2 jets, $p_{\mathrm{T}}^{\gamma\gamma} \in [120, 200]$ GeV
ggH 2J MED	≥ 2 jets, $p_{\mathrm{T}}^{\gamma\gamma} \in [60, 120]$ GeV
ggH 2J LOW	≥ 2 jets, $p_{\mathrm{T}}^{\gamma\gamma} \in [0, 60]$ GeV
ggH 1J BSM	$= 1$ jet, $p_{\mathrm{T}}^{\gamma\gamma} \geq 200\,\mathrm{GeV}$
ggH 1J HIGH	$= 1$ jet, $p_{\mathrm{T}}^{\gamma\gamma} \in [120, 200]$ GeV
ggH 1J MED	$= 1$ jet, $p_{\mathrm{T}}^{\gamma\gamma} \in [60, 120]$ GeV
ggH 1J LOW	$= 1$ jet, $p_{\mathrm{T}}^{\gamma\gamma} \in [0, 60]$ GeV
ggH 0J FWD	$= 0$ jets, one photon with $\lvert\eta\rvert > 0.95$
ggH 0J CEN	$= 0$ jets, two photons with $\lvert\eta\rvert \leq 0.95$

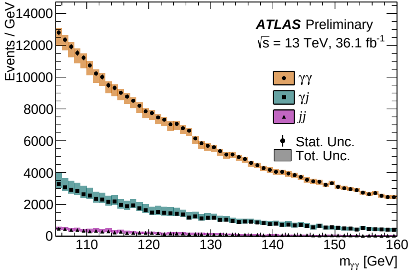
Truth bin distribution



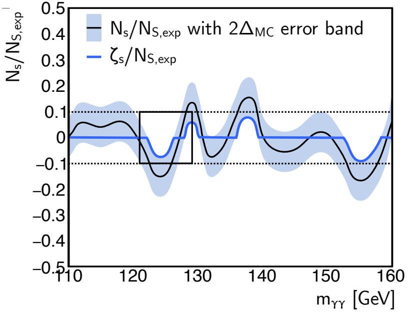
- Columns : distribution of events of a given category over the truth bins.
- Rectangles : process optimised categories

Good performances of process targetting

Background parametrization



- MC not reliable for MC description
- Shape fitted on data
- Spurious signal evaluated for selection of functional form.

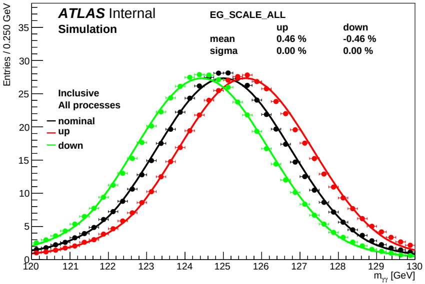


$$\zeta_s = \begin{cases} (N_s+2\Delta_{MC}), & N_s+2\Delta_{MC} < 0 \\ (N_s-2\Delta_{MC}), & N_s-2\Delta_{MC} > 0 \\ 0, & \text{otherwise} \end{cases}$$

Calibration uncertainties methodology

For a given systematic source :

- Create distributions of $m_{\gamma\gamma}^{nom}$, $m_{\gamma\gamma}^{up}$, $m_{\gamma\gamma}^{down}$
- Fit main parameter of the systematic with DSCB :
  - Fit $m_{\gamma\gamma} \in [105, 160]$GeV
  - Fixing $n_{high} = 5$ and $n_{low} = 9$
  - Fixing $\alpha_{high} = \hat{\alpha}_{high}^{nom}$, $\alpha_{low}^{nom} = \hat{\alpha}_{low}^{nom}$, $X = \hat{X}^{nom}$
- Systematic variation : $\frac{X^{fluct}}{X^{nom}} - 1$, $X \in \{\mu, \sigma\}$



$$CB(m_{\gamma\gamma}) = \begin{cases} e^{-t^2/2} & \text{if } -\alpha_{low} \leq t \leq \alpha_{high} \\ \dfrac{e^{-\frac{1}{2}\alpha_{low}^2}}{\left[\frac{1}{R_{low}}(R_{low} - \alpha_{low} - t)\right]^{n_{low}}} & \text{if } t < -\alpha_{low} \\ \dfrac{e^{-\frac{1}{2}\alpha_{high}^2}}{\left[\frac{1}{R_{high}}\left(R_{high} - \alpha_{high} + t\right)\right]^{n_{high}}} & \text{if } t > \alpha_{high} \\ t = (m_{\gamma\gamma} - \mu)/\sigma,\ R_{low} = \frac{\alpha_{low}}{n_{low}},\ R_{high} = \frac{\alpha_{high}}{n_{high}} \end{cases} \quad (9)$$

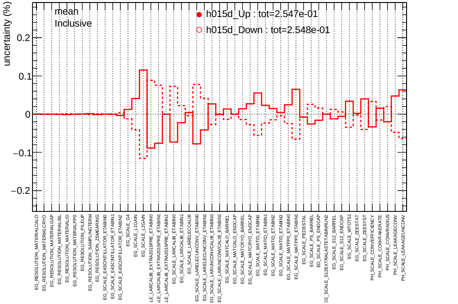
Calibration uncertainties results

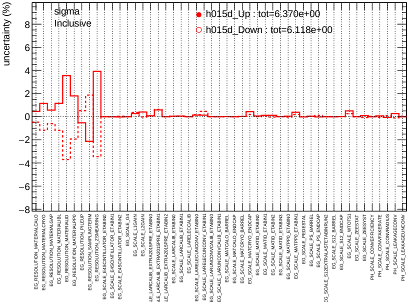
49 Nuisance Parameters

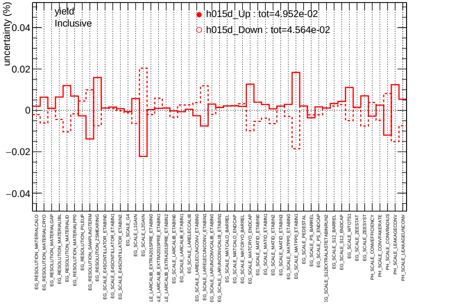
- 9 for resolution
- 40 for scale

affecting

- resolution
- mass
- yield per category

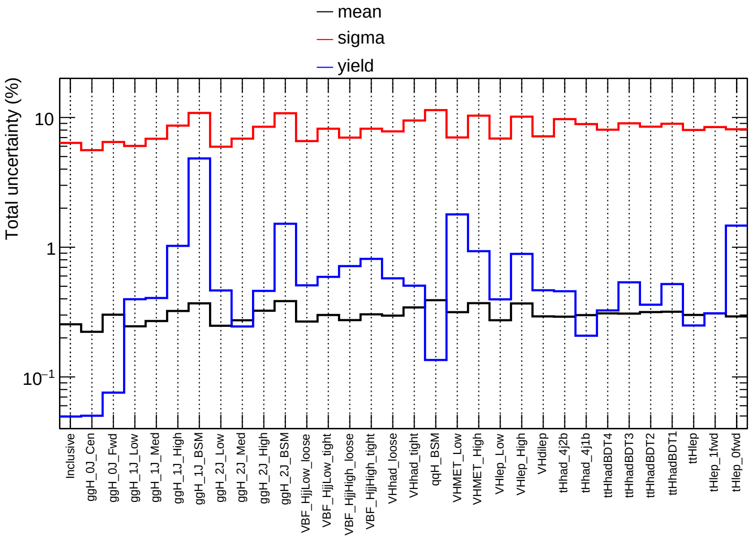






Total uncertainty

Total calibration uncertainty as a function of reconstructed category.



Likelihood Method

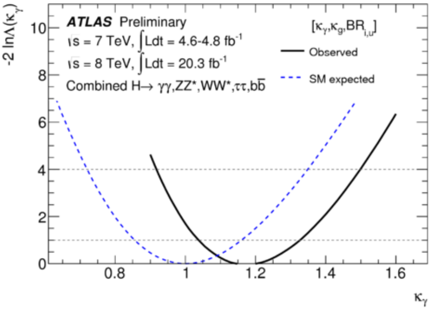
A function (**likelihood**) is built to **evaluate the best set of parameters** $(\vec{\mu},\vec{\theta})$ of a model to agree the best with a dataset in a category.

$$\mathcal{L} = \underbrace{\frac{(n_s(\vec{\mu},\vec{\theta})+b)^{n_{obs}}}{n_{obs}!} e^{-(n_s(\vec{\mu},\vec{\theta})+b)}}_{(1)} \overbrace{\prod_j^{n_{obs}} \psi(\vec{x_j};\vec{\mu},\vec{\theta})}^{(2)} \underbrace{e^{-\frac{\theta^2}{2}}}_{(3)}$$

(1) **Poissonian law** to evaluate the probability to observe n_{obs}($\equiv$ signal + background) events when $(n_s + b)$ are expected.

(2) **Probability density function** of the observables $\vec{x}$ (diphoton invariant mass for example) for the j^{th} event.

(3) Constraint on the nuisance parameter θ. See next slide.



Nuisance parameters

There are some **external measurements** that contribute to the likelihood and have some **uncertainties**. A **free nuisance parameter** is added for each of these measurements. In order to take into account these external measurements, a **constraint is put on these nuisance parameters**.

For example, the luminosity is re-defined as $L(1+\delta_L\theta_L)$, with θ_L the nuisance parameter and δ_L the uncertainty on the luminosity (assumed to be Gaussian). In this case, a Gaussian constraint is chosen.
The contribution from luminosity will hence be :

$$L(1+\delta_L\theta_L)e^{-\theta_L^2/2}$$

Error Estimation

A test statistic is defined as : $t_\mu = -2ln\frac{\mathcal{L}(\mu,\hat{\hat{\theta}})}{\mathcal{L}(\hat{\mu},\hat{\theta})}$, with $\hat{\theta}$ and $\hat{\mu}$ the best (fitted) parameters, and $\hat{\hat{\theta}}$ the fitted nuisance parameters for a fixed μ.
Uncertainty are given by : $\mathbf{t_{\hat{\mu}\pm 1\sigma} = 1}$ and $\mathbf{t_{\hat{\mu}\pm 2\sigma} = 4}$ in 1D Gaussian limit.

Run 2 couplings results

Due to lack of statistics, some truth bins have been merged.

$$\sigma(ggH, 0\text{ jet}) \times B(H \to \gamma\gamma) = 63\,^{+17}_{-16}\ \text{fb} = 63\,^{+15}_{-15}\,(\text{stat.})\,^{+8}_{-6}\,(\text{syst.})\ \text{fb}$$

$$\sigma(ggH, 1\text{ jet}, p_T^H < 60\text{ GeV}) \times B(H \to \gamma\gamma) = 15\,^{+13}_{-12}\ \text{fb} = 15\,^{+12}_{-12}\,(\text{stat.})\,^{+6}_{-4}\,(\text{syst.})\ \text{fb}$$

$$\sigma(ggH, 1\text{ jet}, 60 \leq p_T^H < 120\text{ GeV}) \times B(H \to \gamma\gamma) = 10\,^{+7}_{-6}\ \text{fb} = 10\,^{+6}_{-6}\,(\text{stat.})\,^{+2}_{-1}\,(\text{syst.})\ \text{fb}$$

$$\sigma(ggH, 1\text{ jet}, 120 \leq p_T^H < 200\text{ GeV}) \times B(H \to \gamma\gamma) = 1.7\,^{+1.7}_{-1.6}\ \text{fb} = 1.7\,^{+1.6}_{-1.6}\,(\text{stat.})\,^{+0.6}_{-0.4}\,(\text{syst.})\ \text{fb}$$

$$\sigma(ggH, \geq 2\text{ jet}) \times B(H \to \gamma\gamma) = 11\,^{+8}_{-8}\ \text{fb} = 11\,^{+8}_{-8}\,(\text{stat.})\,^{+3}_{-2}\,(\text{syst.})\ \text{fb}$$

$$\sigma(qq \to Hqq, p_T^j < 200\text{ GeV}) \times B(H \to \gamma\gamma) = 10\,^{+6}_{-5}\ \text{fb} = 10\,^{+5}_{-5}\,(\text{stat.})\,^{+2}_{-1}\,(\text{syst.})\ \text{fb}$$

$$\sigma(ggH + qq \to Hqq, \text{BSM}-\text{like}) \times B(H \to \gamma\gamma) = 1.8\,^{+1.4}_{-1.4}\ \text{fb} = 1.8\,^{+1.3}_{-1.3}\,(\text{stat.})\,^{+0.5}_{-0.5}\,(\text{syst.})\ \text{fb}$$

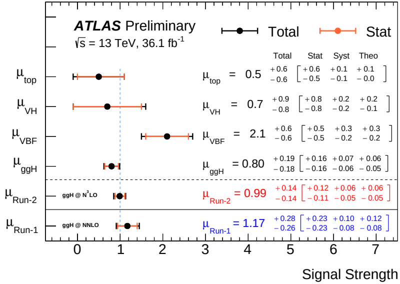
$$\sigma(\text{VH, leptonic}) \times B(H \to \gamma\gamma) = 1.4\,^{+1.4}_{-1.2}\ \text{fb} = 1.4\,^{+1.3}_{-1.2}\,(\text{stat.})\,^{+0.3}_{-0.3}\,(\text{syst.})\ \text{fb}$$

$$\sigma(\text{top}) \times B(H \to \gamma\gamma) = 1.3\,^{+0.9}_{-0.8}\ \text{fb} = 1.3\,^{+0.9}_{-0.8}\,(\text{stat.})\,^{+0.3}_{-0.1}\,(\text{syst.})\ \text{fb}$$

Run 2 signal strength results

STXS difficult to interpret directly.

Measurement of signal strengths $\mu_i = \frac{\sigma_i^{exp}}{\sigma_i^{SM}}$ performed.



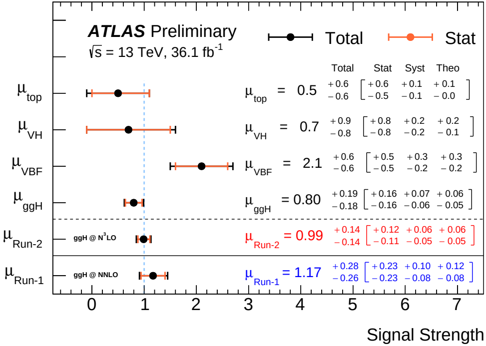
Uncertainty Group	$\sigma_\mu^{\text{syst.}}$
Theory (yield)	0.03
Experimental (yield)	0.02
Luminosity	0.03
Theory (migrations)	0.05
Experimental (migrations)	0.01
Mass resolution	0.03
Mass scale	0.04
Background shape	0.03

- Major theory improvement
- Major resolution improvement
- Increase of mass scale impact
- **No deviation from SM**

Conclusion

- Outstanding LC performance $\rightarrow$ large statistic available
- Major improvement of calibration uncertainties
  resolution uncertainties no longer dominant on μ
- Couplings measurement performed with 36.1 fbat $\sqrt{s} = 13$ TeV
- **No significant deviation from standard model**

- Much more data expected until end of run 2
- **Work on experimental systematics required**



Calibration of the ATLAS electromagnetic calorimeter and measurement of the couplings of the (Brout-Englert-)Higgs boson in the diphoton channel

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PhD defense
Orsay, September 12, 2017

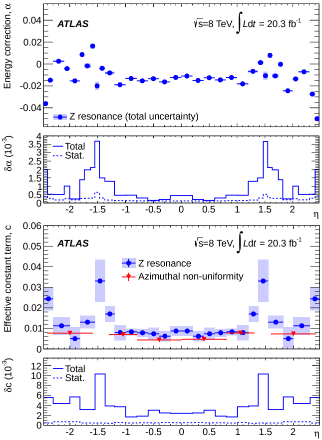
Identification variables

Type	Description	Name
Hadronic leakage	Ratio of E_T in the first layer of the hadronic calorimeter to E_T of the EM cluster (used over the range $\|\eta\| < 0.8$ or $\|\eta\| > 1.37$)	R_{Had1}
	Ratio of E_T in the hadronic calorimeter to E_T of the EM cluster (used over the range $0.8 < \|\eta\| < 1.37$)	R_{Had}
Back layer of EM calorimeter	Ratio of the energy in the back layer to the total energy in the EM accordion calorimeter	f_3
Middle layer of EM calorimeter	Lateral shower width, $\sqrt{(\Sigma E_i \eta_i^2)/(\Sigma E_i) - ((\Sigma E_i \eta_i)/(\Sigma E_i))^2}$, where E_i is the energy and η_i is the pseudorapidity of cell i and the sum is calculated within a window of 3×5 cells	$W_{\eta 2}$
	Ratio of the energy in 3×3 cells over the energy in 3×7 cells centered at the electron cluster position	R_ϕ
	Ratio of the energy in 3×7 cells over the energy in 7×7 cells centered at the electron cluster position	R_η
Strip layer of EM calorimeter	Shower width, $\sqrt{(\Sigma E_i (i - i_{\mathrm{max}})^2)/(\Sigma E_i)}$, where i runs over all strips in a window of $\Delta\eta \times \Delta\phi \approx 0.0625 \times 0.2$, corresponding typically to 20 strips in η, and i_{max} is the index of the highest-energy strip	w_{stot}
	Ratio of the energy difference between the largest and second largest energy deposits in the cluster over the sum of these energies	E_{ratio}
	Ratio of the energy in the strip layer to the total energy in the EM accordion calorimeter	f_1
Track quality	Number of hits in the B-layer (discriminates against photon conversions)	n_{Blayer}
	Number of hits in the pixel detector	n_{Pixel}
	Number of total hits in the pixel and SCT detectors	n_{Si}
	Transverse impact parameter	d_0
	Significance of transverse impact parameter defined as the ratio of d_0 and its uncertainty	σ_{d_0}
	Momentum lost by the track between the perigee and the last measurement point divided by the original momentum	$\Delta p/p$
TRT	Total number of hits in the TRT	n_{TRT}
	Ratio of the number of high-threshold hits to the total number of hits in the TRT	F_{HT}
Track-cluster matching	$\Delta\eta$ between the cluster position in the strip layer and the extrapolated track	$\Delta\eta_1$
	$\Delta\phi$ between the cluster position in the middle layer and the extrapolated track	$\Delta\phi_2$
	Defined as $\Delta\phi_2$, but the track momentum is rescaled to the cluster energy before extrapolating the track to the middle layer of the calorimeter	$\Delta\phi_{\mathrm{res}}$
	Ratio of the cluster energy to the track momentum	E/p
Conversions	Veto electron candidates matched to reconstructed photon conversions	isConv

Calibration in-situ : run 1 results and uncertainties

Uncertainties are evaluated as the difference between official scales and the ones measured with a changed parameter. They include :

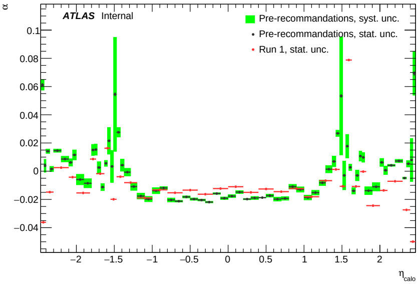
- electron identification quality from medium to tight.
- Z mass window
- electron p_T cut
- uncertainties on efficiencies scale factors
- energy loss through bremshtrahlung
- background
- pile-up
- measurement method

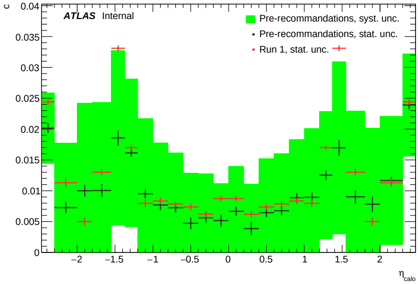


Run 2 prerecommandations

Run 2 early analyses need scales factors for 13TeV but not enough data will be available. Need to **estimate run 2 scales from run 1 data**.
Pre-recommandations are computed using 8 TeV data reprocessed with :

- new detector geometry
- new reconstruction algorithm
- new calibration machine learning



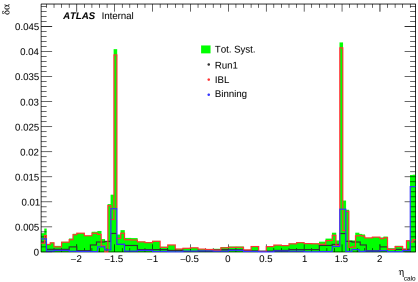


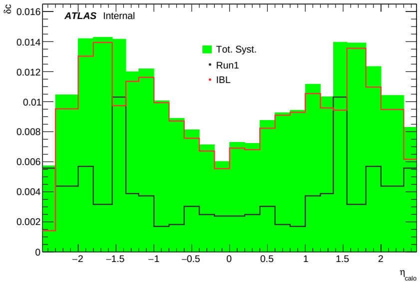
Calibration in-situ : run 2 pre-recommandations systematics

2012 systematics are used for the pre-recommandations.

Two more systematics are added in quadrature :

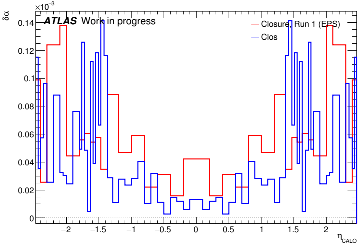
- Increasing the number of bin for α shows sub-patterns. Systematic is defined as difference between a bin value and the average of its sub-bins.
- Pre-recommandations being computed with 8TeV datasets, one needs to evaluate the impact of the center of mass energy. Systematic is defined as the scale measured from 13 TeV MC on 8 TeV templates.

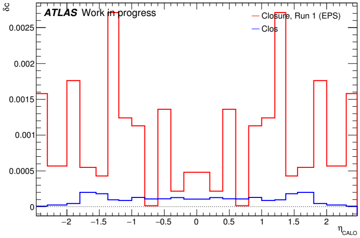


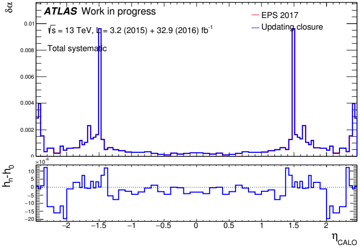


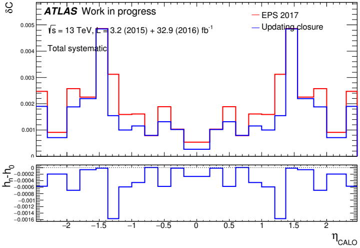
Calibration in-situ : closure improvement

Comparison of run 1 and run 2 closure systematics. A major improvement has been achieved for c.



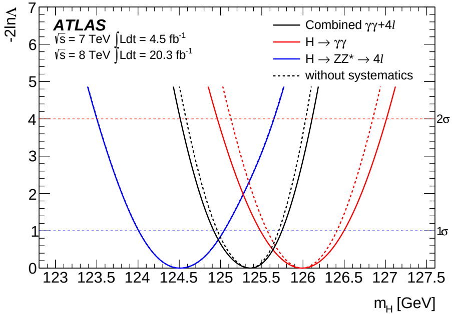






ATLAS run 1 H boson mass measurement

$$m_H = 125.36 \pm 0.37(\text{stat}) \pm 0.18(\text{syst})$$



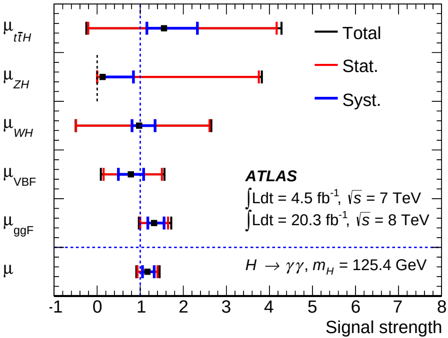
Systematic	Uncertainty on m_H [MeV]
LAr syst on material before presampler (barrel)	70
LAr syst on material after presampler (barrel)	20
LAr cell non-linearity (layer 2)	60
LAr cell non-linearity (layer 1)	30
LAr layer calibration (barrel)	50
Lateral shower shape (conv)	50
Lateral shower shape (unconv)	40
Presampler energy scale (barrel)	20
ID material model ($\lvert\eta\rvert < 1.1$)	50
$H \rightarrow \gamma\gamma$ background model (unconv rest low p_T)	40
$Z \rightarrow ee$ calibration	50
Primary vertex effect on mass scale	20
Muon momentum scale	10
Remaining systematic uncertainties	70
Total	180

Statistical uncertainties highly dominant.

Run 2 will increase sensitivity to systematics.

$\mu_{\gamma\gamma}$ measurement

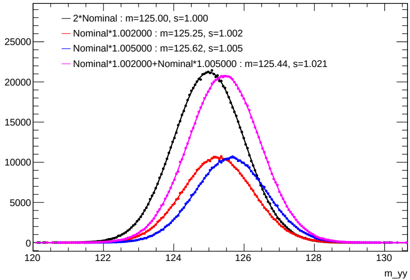
$$\mu_{\gamma\gamma} = \frac{(\sigma \times BR)^{meas}}{(\sigma \times BR)^{SM}} = 1.17 \pm 0.23(\text{stat}) \, {}^{+0.10}_{-0.08}(\text{syst}) \, {}^{+0.12}_{-0.08}(\text{theory})$$



Uncertainty group	$\sigma_{\mu}^{\text{syst.}}$
Theory (yield)	0.09
Experimental (yield)	0.02
Luminosity	0.03
MC statistics	< 0.01
Theory (migrations)	0.03
Experimental (migrations)	0.02
Resolution	0.07
Mass scale	0.02
Background shape	0.02

If no improvements, **calibration uncertainty will be dominant in run 2**.

Scale impact on width

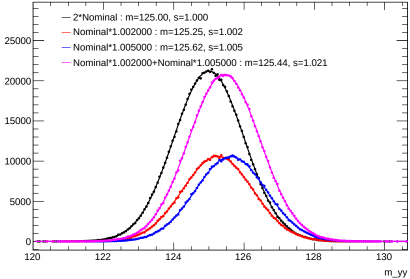


1M random numbers generated on a Gaussian($\mu = 125, \sigma = 1$).

- Initial numbers distribution.
- Half events multiplied by 1.002.
- Remaining events multiplied by 1.005.
- Combined distribution of red and blue.

Mean (m) and RMS (s) of a fitted gaussians are given in the legend. Interpretationof the curve in the next slides.

μ/σ scale correlation



Lets assume a gaussian distributed energy distribution. Applying energy scale correction gives :

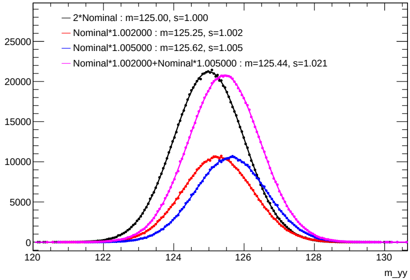
$$E \to E(1+a)$$

Hence the distribution will be changed to :

$$exp(-\frac{(E-\mu)^2}{2\sigma^2} \;\to\; exp(-\frac{(\frac{E}{1+a}-\mu)^2}{2\sigma^2}) = exp(-\frac{(E-\mu(1+a))^2}{2\sigma^2(1+a)^2}) \quad (10)$$

The new distribution is a shifted gaussian with scaled RMS.
Given the medium shift of EG_SCALE_ALL, we expect $^{+0.4}_{-0.4}\%$ change in resolution.

Inhomegenous scale



The RMS of two points separated by d is $d/4$.
If d is the difference between two scale factors,

$$d \sim 3.10^{-3} \cdot E_\gamma = 0.18$$

$$\frac{\text{RMS}}{\text{Resolution}} = \frac{d/4}{1.5\text{GeV}} = 3\%$$

The inhomogeneity of the scale factors uncertainties changes the width of the distribution at the percent level. This effect will always increase the width.

Black and pink distribution show an illustration of this effect.

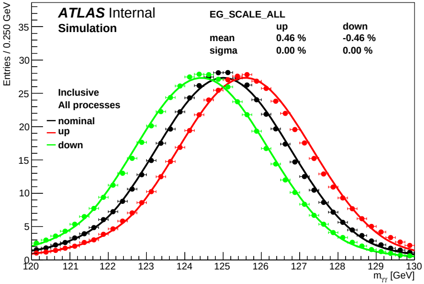
Scale factors interpretation

Assume the up fluctuation (red) as data and nominal distribution (black) as MC in the template method. One has

$$m_H^{up} = m_H^{nom}(1+\alpha)$$

Hence

$$\delta_{m_H} = \alpha$$



Furthermore :

$$\sigma_H^{up} = \sigma_H^{nom} \oplus cE$$

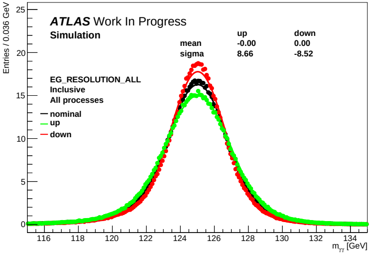
Hence

$$\delta_{\sigma_H} = \sqrt{1 + \frac{c^2E^2}{\sigma_H^2}} - 1$$

One has to be carefull with resolution uncertainty as the template method is weak to measure small differences.

Method comparison

4 different fitting methods are compared : fitting in 3 different ranges and template method cross-check within $[122, 128]$GeV. Methods compared on h013 simplified model (2NP).

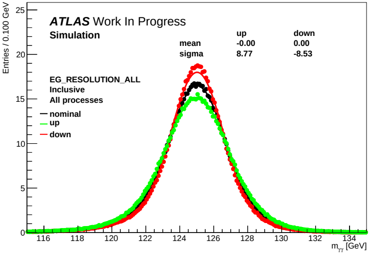


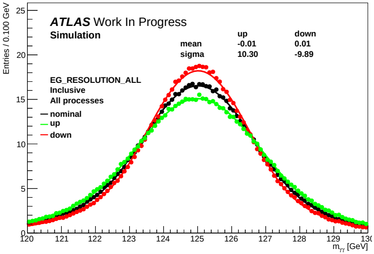
$\leftarrow$

$[105, 160]$

$\rightarrow$

$[115, 135]$



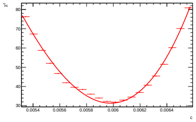


$\leftarrow$

$[120, 130]$

$\rightarrow$

$[122, 128]$



$$c = (0.598 \pm 0.009)\%$$
$$\rightarrow \delta_{\sigma_H} = (8.82 \pm 0.25)\%$$