

Temperature around a nuclear waste rod using
finite difference implicit method

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The governing equation in heat transfer is:

$$\frac{\partial T}{\partial t}(r,t) + \bar{u} * \bar{\nabla} T(r,t) = k \nabla^2 T(r,t) + S1(r,t) \quad K = \text{thermal diffusivity } m^2/s$$

For solids u is zero

$$\frac{1}{k} \frac{\partial (T)}{\partial (t)}(r,t) - (\nabla^2 T(r,t)) = S(r,t)$$

The source term due to radioactive decay of rod is defined as:

$$S(r,t) = \left\{ T_{rod} e^{-t/\tau} / a^2 \text{ for } r \leq a \right\}$$

Where, a=25cm

$T_{rod}=1k$

$\tau=100\text{years}$

$r_c=100\text{cm}$

$T_E=300k$

$k=2 \times 10^7 \text{ cm}^2/\text{year}$

$0 < r < 100\text{cm}$

$0 < t < 100\text{years}$

Because the problem has circular symmetry, 2-D problem in (x,y) can be converted to 1-D problem in r

$$\nabla^2 T = \frac{\partial^2 T}{\partial^2 x} + \frac{\partial^2 T}{\partial^2 y} = \frac{\partial^2 T}{\partial^2 r} + \frac{1}{r} \frac{\partial T}{\partial r} + \frac{1}{r^2} \frac{\partial^2 T}{\partial^2 \theta}$$

T(r,t) has circular symmetry, so

$$\nabla^2 T = \frac{\partial^2 T}{\partial^2 x} + \frac{\partial^2 T}{\partial^2 y} = \frac{\partial^2 T}{\partial^2 r} + \frac{1}{r} \frac{\partial T}{\partial r}$$

Hence 1-D heat equation in polar co-ordinates becomes:

$$\frac{1}{k} \frac{\partial T}{\partial t} - \frac{\partial^2 T}{\partial^2 r} - \frac{1}{r} \frac{\partial T}{\partial r} = S(r, t)$$

- We know that in the steady state solution, $S(r, t) \rightarrow 0$ as $t \rightarrow \infty$,
- Far away from the rod the temperature equals the ambient temperature, $T(r = r_c, t) = 300K$.
- $T(r, t) \rightarrow 300K$ once rod has finished radioactive decaying.

Finite Difference Method – Implicit method:

$$\frac{1}{k} \frac{\partial T}{\partial t} - \frac{\partial^2 T}{\partial r^2} - \frac{1}{r} \frac{\partial T}{\partial r} = S(r, t)$$

Initial conditions : $T(r, 0) = 300k$

There is singularity at $r=0$, thereby we take another boundary condition

Neumann boundary condition at $r=0$ as temperature cannot flow into $r=0$ region

$$\frac{\partial T}{\partial r}(r=0, t) = 0$$

Dirichlet boundary condition at $r=r_c$

$$T(r=r_c, t) = 300k$$

Discretization of space and time:

$$dr = r_c/(n+1) \quad dt = T_f/m \quad r_j = j(dr) \quad t_k = k(dt)$$

$$0 \leq j \leq (n+1) \quad 0 \leq k \leq m \quad T(r_j, t_k) = T_j^k \quad S(r_j, t_k) = S_j^k$$

Neumann boundary condition at $r=0$ becomes:

$$\frac{\partial T_j^k}{\partial t}(r=0, t) = \frac{\partial T_0^k}{\partial t} = 0 \approx \frac{T_1^k - T_0^k}{\Delta t} \Rightarrow T_1^k \approx T_0^k$$

Dirichlet boundary condition at $r=r_c$ becomes:

$$T_j^k(r=r_c, t) = T_{(n+1)}^k = 300k$$

Writing spacial derivatives in r at future time step k+1 : implicit method

$$T_t(t_{k+1}, r_j) = \frac{T_j^{k+1} - T_j^k}{\Delta t}$$

$$T_{rr}(t_{k+1}, r_j) = \frac{T_{j+1}^{k+1} - 2T_j^{k+1} + T_{j-1}^{k+1}}{\Delta r^2}$$

$$T_r(t_{k+1}, r_j) = \frac{T_{j+1}^{k+1} - T_{j-1}^{k+1}}{2 \Delta r}$$

Using $r = j^* dr$, the discretized pde becomes:

$$\frac{1}{k} \frac{\partial T}{\partial t} - \frac{\partial^2 T}{\partial^2 r} - \frac{1}{r} \frac{\partial T}{\partial r} = S(r, t)$$

$$\frac{1}{k} \frac{T_j^{k+1} - T_j^k}{\Delta t} - \frac{T_{j+1}^{k+1} - 2T_j^{k+1} + T_{j-1}^{k+1}}{\Delta r^2} - \frac{1}{j \Delta r} \frac{T_{j+1}^{k+1} - T_{j-1}^{k+1}}{2 \Delta r} = S_j^k$$

$$s = k^* dt / dr^2$$

$$T_{j+1}^{k+1} \left[-s - \frac{s}{2j} \right] + T_{j+1}^{k+1} \left[-s + \frac{s}{2j} \right] + T_j^{k+1} [1 + 2s] = T_j^k + S_j^k k \Delta t \quad \text{Tridiagonal matrix for } 1 \leq j \leq n$$

Numerical solution for n=3

- $T_0^k \approx T_1^k$, Neumann boundary condition
- $T_4^k = 300K$ from $T(r = r_c, t) = 300$
- The initial conditions are $T_j^0 = 300K$.
- We solve the equation for T_1^k, T_2^k, T_3^k at each time step (t^k):

$$\begin{pmatrix} 1 + 2s & (-s - \frac{s}{2j}) & 0 \\ (-s + \frac{s}{2j}) & 1 + 2s & (-s - \frac{s}{2j}) \\ 0 & (-s + \frac{s}{2j}) & 1 + 2s \end{pmatrix} \begin{pmatrix} T_1^{k+1} \\ T_2^{k+1} \\ T_3^{k+1} \end{pmatrix} + \begin{pmatrix} (-s + \frac{s}{2j})T_0^{k+1} \\ 0 \\ (-s - \frac{s}{2j})T_4^{k+1} \end{pmatrix} = \begin{pmatrix} T_1^k \\ T_2^k \\ T_3^k \end{pmatrix} + \kappa \Delta t \begin{pmatrix} S_1^k \\ S_2^k \\ S_3^k \end{pmatrix}$$

Using the boundary conditions: $T_0^{k+1} \approx T_1^{k+1}$, $T_4^{k+1} = 300K$

$$\begin{pmatrix} (1 + s + \frac{s}{2j}) & (-s - \frac{s}{2j}) & 0 \\ (-s + \frac{s}{2j}) & (1 + 2s) & (-s - \frac{s}{2j}) \\ 0 & (-s + \frac{s}{2j}) & (1 + 2s) \end{pmatrix} \begin{pmatrix} T_1^{k+1} \\ T_2^{k+1} \\ T_3^{k+1} \end{pmatrix} = \begin{pmatrix} T_1^k \\ T_2^k \\ T_3^k \end{pmatrix} + \kappa \Delta t \begin{pmatrix} S_1^k \\ S_2^k \\ S_3^k \end{pmatrix} - \begin{pmatrix} 0 \\ 0 \\ (-s - \frac{s}{2j})T_4^{k+1} \end{pmatrix}$$

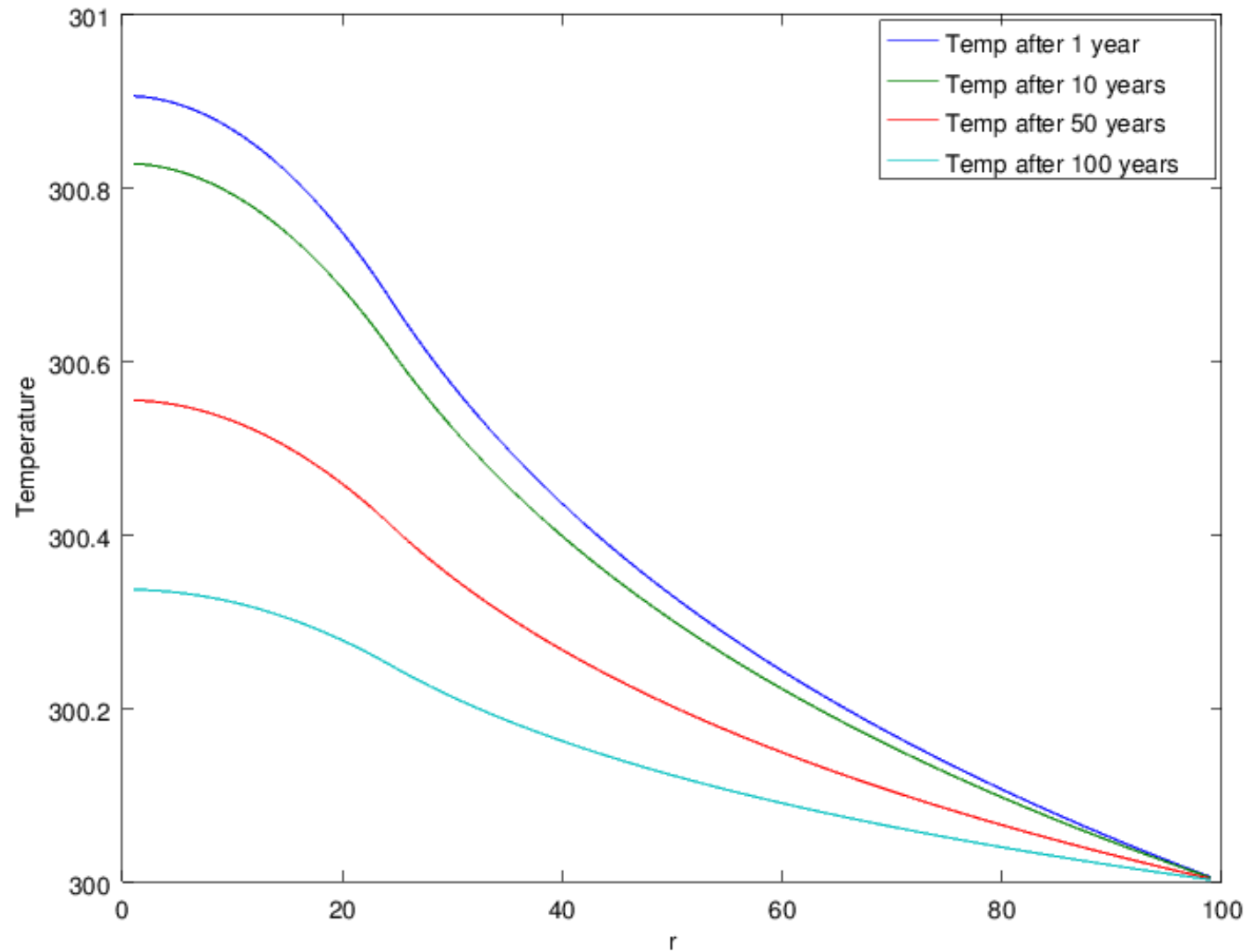
$$\Rightarrow \begin{pmatrix} (1 + s + \frac{s}{2}) & (-s - \frac{s}{2}) & 0 \\ (-s + \frac{s}{4}) & (1 + 2s) & (-s - \frac{s}{4}) \\ 0 & (-s + \frac{s}{6}) & (1 + 2s) \end{pmatrix} \begin{pmatrix} T_1^{k+1} \\ T_2^{k+1} \\ T_3^{k+1} \end{pmatrix}$$

$$= \begin{pmatrix} T_1^k \\ T_2^k \\ T_3^k \end{pmatrix} + \kappa \Delta t \begin{pmatrix} S_1^k \\ S_2^k \\ S_3^k \end{pmatrix} - \begin{pmatrix} 0 \\ 0 \\ (-s - \frac{s}{6})300 \end{pmatrix}$$

Fortran 90 code for the problem:

</home/chanakya/Fortran/sid/cylinder.f90>

Documented output



http://espace.library.uq.edu.au/view/UQ:239427/Lectures_Book.pdf

Output:

