

HW12

PHYS4500: Quantum Field Theory

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Problem 1.

We are to show that the Dirac Lagrangian

$$\mathcal{L} = i\bar{\psi}\gamma^\mu\partial_\mu\psi - \bar{\psi}M\psi \quad (1.1)$$

with

$$\psi = \begin{pmatrix} \psi_a \\ \psi_b \end{pmatrix} \quad \text{and} \quad M = \begin{pmatrix} m_a & 0 \\ 0 & m_b \end{pmatrix} \quad (1.2)$$

simplifies to the sum of the individual Dirac Lagrangians for the individual spinors ψ_a and ψ_b . First, we know that the quantity $\gamma^\mu\partial_\mu$ is a matrix in spin-space, call it Γ_{ij} , where the i and j indices run from 1 – 4 in spin space. Then, we can explicitly write out the indices in order to manipulate them more easily:

$$\bar{\psi}\gamma^\mu\partial_\mu\psi \rightarrow (\bar{\psi}_a \quad \bar{\psi}_b)_i \Gamma_{ij} \begin{pmatrix} \psi_a \\ \psi_b \end{pmatrix}_j \quad (1.3)$$

$$= (\bar{\psi}_{a,i} \Gamma_{ij} \quad \bar{\psi}_{b,i} \Gamma_{ij}) \begin{pmatrix} \psi_a \\ \psi_b \end{pmatrix}_j \quad (1.4)$$

$$= \bar{\psi}_{a,i} \Gamma_{ij} \psi_{a,j} + i\bar{\psi}_{b,i} \Gamma_{ij} \psi_{b,j}, \quad (1.5)$$

or, since our indices are in valid matrix multiplication order, we can remove them to say:

$$\bar{\psi}\gamma^\mu\partial_\mu\psi = \bar{\psi}_a\gamma^\mu\partial_\mu\psi_a + \bar{\psi}_b\gamma^\mu\partial_\mu\psi_b. \quad (1.6)$$

For the mass part:

$$\bar{\psi}M\psi = (\bar{\psi}_a \quad \bar{\psi}_b) \begin{pmatrix} m_a & 0 \\ 0 & m_b \end{pmatrix} \begin{pmatrix} \psi_a \\ \psi_b \end{pmatrix} = (m_a\bar{\psi}_a \quad m_b\bar{\psi}_b) \begin{pmatrix} \psi_a \\ \psi_b \end{pmatrix} = m_a\bar{\psi}_a\psi_a + m_b\bar{\psi}_b\psi_b. \quad (1.7)$$

Thus, both parts have nicely separated out, so our new Lagrangian is:

$$\mathcal{L} = i\bar{\psi}_a\gamma^\mu\partial_\mu\psi_a + i\bar{\psi}_b\gamma^\mu\partial_\mu\psi_b - m_a\bar{\psi}_a\psi_a - m_b\bar{\psi}_b\psi_b, \quad (1.8)$$

as desired.

Problem 2. (23.1)

We start by considering a doublet ψ and a generic 2×2 matrix M that acts on the doublet.

- a) First, we consider M such that we have $\bar{\psi}'\psi' = \bar{\psi}\psi$. In order for this to be satisfied, M must be unitary. This can be easily shown by simply writing out the transformations explicitly:

$$\bar{\psi}'\psi' = \bar{\psi}M^\dagger M\psi. \quad (2.1)$$

For this to be equal to $\bar{\psi}\psi$, we must have $M^\dagger M = 1$, which is exactly what the unitary condition requires. Thus, M must be unitary.

- b) Now we are to show the determinant of these M matrices must be 1. First, we know the matrix identity $\det AB = (\det A)(\det B)$. For us, we can choose $A = M^\dagger$ and $B = M$ so $\det M^\dagger M = \det 1 = 1 = (\det M^\dagger)(\det M)$. Now, we must show that $\det M^\dagger = (\det M)^*$. Consider an arbitrary complex 2×2 matrix

$$U = \begin{pmatrix} a & b \\ c & d \end{pmatrix}, \quad (2.2)$$

where a, b, c and d are complex numbers. Now,

$$U^\dagger = \begin{pmatrix} a^* & c^* \\ b^* & d^* \end{pmatrix}. \quad (2.3)$$

It is quite easy to see that $\det U = ad - bc$ and that $\det U^\dagger = a^*d^* - b^*c^* = (ad - bc)^* = (\det U)^*$. Thus,

$$1 = (\det M)^*(\det M) \rightarrow |\det M|^2 = 1 \rightarrow |\det M| = 1. \quad (2.4)$$

- c) For a matrix M with $\det M = e^{i\alpha}$, we can define a new matrix $M_{\text{new}} = e^{-i\alpha/2}M$. Using the same matrix identity as in the previous part, we have that $\det M_{\text{new}} = (\det e^{-i\alpha/2})(\det M)$. Now,

$$\det e^{-i\alpha/2} = \det \begin{pmatrix} e^{-i\alpha/2} & 0 \\ 0 & e^{-i\alpha/2} \end{pmatrix} = e^{-i\alpha}. \quad (2.5)$$

Therefore, we have that $\det M_{\text{new}} = e^{-i\alpha}e^{i\alpha} = 1$.

- f) We have used before that elements of $SU(2)$ can be expressed as the exponential of the algebra of the group, something like (in our notation from class)

$$M = e^{-ig\boldsymbol{\sigma} \cdot \boldsymbol{\lambda}/2}. \quad (2.6)$$

We can easily show that this satisfies the $\det M = 1$ constraint meaning M is an element of $SU(2)$ (obviously it is unitary). With the matrix identity that $\det e^A = e^{\text{Tr}[A]}$, we have that

$$\det M = e^{-\frac{i}{2}g\text{Tr}[\boldsymbol{\sigma} \cdot \boldsymbol{\lambda}]}. \quad (2.7)$$

The trace deals with elements down the diagonal. Therefore, the trace is effectively $\text{Tr}[\sigma_z \lambda_z] \rightarrow \lambda_z \text{Tr}[\sigma_z]$ since only the third/ z Pauli matrix contains elements that are on the main diagonal. Now

$$\sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}, \quad (2.8)$$

so it's trace is obviously zero, meaning

$$\det M = e^0 = 1, \quad (2.9)$$

as expected of an element of $SU(2)$.