

HW1

PHYS4500: Quantum Field Theory

Casey Hampson

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### Question 1. (2.2)

We are given  $p^\mu = [5m \quad -4m \quad 0 \quad 0]^\top$ .

a) We know that the 4-momentum has the following property:

$$p_\mu p^\mu = M^2, \quad (1)$$

so

$$p_\mu p^\mu = (5m)^2 - (-4m)^2 = 25m^2 - 16m^2 = 9m^2 = M^2 \quad (2)$$

$$\rightarrow \boxed{M = 3m}. \quad (3)$$

c) This simple, since  $p^0 = E$  and we just solved for  $M$ :

$$K = E - M = 5m - 3m = 2m. \quad (4)$$

$$\rightarrow \boxed{K = 2m}. \quad (5)$$

d) We apply the Lorentz transformation  $p^{\mu'} = \Lambda^{\mu'}_\mu p^\mu$ :

$$\begin{bmatrix} p^{0'} \\ p^{1'} \\ p^{2'} \\ p^{3'} \end{bmatrix} = \begin{bmatrix} \gamma & -\beta\gamma & 0 & 0 \\ -\beta\gamma & \gamma & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} p^0 \\ p^1 \\ p^2 \\ p^3 \end{bmatrix}. \quad (6)$$

We must determine the value of  $\gamma$ :

$$\gamma = \frac{1}{\sqrt{1 - \left(\frac{4}{5}\right)^2}} = \frac{1}{\sqrt{1 - \frac{16}{25}}} = \frac{1}{\sqrt{\frac{9}{25}}} = \frac{5}{3}. \quad (7)$$

Now,

$$p^{0'} = \gamma p^0 - \beta\gamma p^1 = \frac{5}{3}(5m) - \frac{4}{5} \frac{5}{3}(-4m) = \frac{25}{3}m + \frac{16}{3}m = \frac{41}{3}m, \quad (8)$$

$$p^{1'} = -\beta\gamma p^0 + \gamma p^1 = -\frac{4}{5} \frac{5}{3}(5m) + \frac{5}{3}(-4m) = -\frac{20}{m} - \frac{20}{m} = -\frac{40}{3}m, \quad (9)$$

$$p^{2'} = p^2 = 0, \quad (10)$$

$$p^{3'} = p^3 = 0. \quad (11)$$

So,

$$\boxed{p^{\mu'} = \begin{bmatrix} 41m/3 \\ -40m/3 \\ 0 \\ 0 \end{bmatrix}}. \quad (12)$$

### Question 2. (3.2)

We are given  $A^\mu = [a(t^2 - x^2) \quad bx^2 \quad cx \quad 0]^\top$ .

a) Using the metric tensor:

$$A_\mu = g_{\mu\nu} A^\nu. \quad (13)$$

This really just involves a negating of the sign on all the spatial terms, so we have:

$$A_\mu = [a(t^2 - x^2) \quad -bx^2 \quad -cx \quad 0]. \quad (14)$$

b) Since the indicies are different, this quantity is a rank-2 tensor:

$$\partial_\mu A_\nu = T_{\mu\nu} = \begin{bmatrix} \frac{1}{c} \frac{\partial}{\partial t} [a(t^2 - x^2)] & \frac{1}{c} \frac{\partial}{\partial t} [bx^2] & \frac{1}{c} \frac{\partial}{\partial t} [cx] & 0 \\ \frac{\partial}{\partial x} [a(t^2 - x^2)] & \frac{\partial}{\partial x} [bx^2] & \frac{\partial}{\partial x} [cx] & 0 \\ \frac{\partial}{\partial y} [a(t^2 - x^2)] & \frac{\partial}{\partial y} [bx^2] & \frac{\partial}{\partial y} [cx] & 0 \\ \frac{\partial}{\partial z} [a(t^2 - x^2)] & \frac{\partial}{\partial z} [bx^2] & \frac{\partial}{\partial z} [cx] & 0 \end{bmatrix} = \begin{bmatrix} 2at/c & 0 & 0 & 0 \\ -2ax & 2bx & c & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix}. \quad (15)$$

As a note about the problem's construction, the constants in  $A^\mu$  should have been chosen such that there were not 2  $c$ 's. The  $c$  in the top left matrix element is the speed of light, while the one in the second row and third column is just a constant.

c) This time we have a contraction, so it will just be a scalar:

$$\partial_\mu A^\mu = \frac{1}{c} \frac{\partial}{\partial t} [a(t^2 - c^2)] + \frac{\partial}{\partial x} [bx^2] + \frac{\partial}{\partial y} [cx] \quad (16)$$

$$= \boxed{\frac{2at}{c} + 2bx} \quad (17)$$

### Question 3. (5.4)

We are given a plane-wave solution  $\phi(x^\mu) = A \exp(-ik_\mu x^\mu)$  to the Klein-Gordon Equation, where  $A$  is a constant.

a) In order to show eventually that the phase speed is greater than  $c$ , I'll reintroduce  $c$  into the Klein-Gordon equation, but leave  $\hbar = 1$ :

$$\partial_\mu \partial^\mu \phi + m^2 c^2 \phi = 0. \quad (18)$$

We can simply plug in our plane wave solution into the Klein-Gordon equation and determine what the components of  $k^\mu$  must be, then solve for the phase speed. First,

$$\partial_\mu \partial^\mu (A \exp(-ik_\mu x^\mu)) = -A k_\mu k^\mu \exp(-ik_\mu x^\mu), \quad (19)$$

so we have that

$$-A k_\mu k^\mu \exp(-ik_\mu x^\mu) + m^2 c^2 A \exp(-ik_\mu x^\mu) = 0, \quad (20)$$

$$A \exp(-ik_\mu x^\mu) (-k_\mu k^\mu + m^2 c^2) = 0. \quad (21)$$

Since we don't want the trivial solution where  $A = 0$  and since the exponential is never zero,

$$-k_\mu k^\mu + m^2 c^2 = 0, \quad (22)$$

$$-\frac{\omega^2}{c^2} + |\mathbf{k}|^2 + m^2 c^2 = 0, \quad (23)$$

$$\omega^2 = |\mathbf{k}|^2 c^2 + m^2 c^4, \quad (24)$$

$$\omega = \sqrt{|\mathbf{k}|^2 c^2 + m^2 c^4}. \quad (25)$$

Now, by the definition of the phase speed  $v_p = \omega/|\mathbf{k}|$ , we have

$$\omega/|\mathbf{k}| \equiv \omega/k = \sqrt{c^2 + \frac{m^2 c^4}{k^2}} = c \sqrt{1 + \frac{m^2 c^2}{k^2}}. \quad (26)$$

Since the fraction in the exponent will always be positive, the quantity in the square root must always be greater than one, meaning we can safely say that  $v_p > c$ , as expected.

b) Here, we will take our original expression for  $\omega$  and simply differentiate with respect to  $|\mathbf{k}| \equiv k$ :

$$v_g \equiv \frac{d\omega}{dk} = \frac{d}{dk} \left[ \sqrt{k^2 c^2 + m^2 c^4} \right], \quad (27)$$

$$= \frac{1}{2} (k^2 c^2 + m^2 c^4)^{-1/2} \cdot 2kc^2, \quad (28)$$

$$= \frac{kc^2}{\sqrt{k^2 c^2 + m^2 c^4}}, \quad (29)$$

$$= \frac{kc}{\sqrt{k^2 + m^2}}. \quad (30)$$

Here, the coefficient of  $c$  is always less than 1 due to the additive factor of  $m^2$  in the square root in the denominator. This time, then,  $d\omega/dk < c$ , as expected.