

A Laurent expansion around $\epsilon=0$ gives

(21)

$$\Sigma(p) = \frac{g_0^2}{(4\pi)^2} G \left(\frac{1}{\epsilon} - \gamma + 1 - \ln\left(\frac{-p^2}{4\pi\mu^2}\right) \right) + O(\epsilon)$$

where we used

$$\Gamma(\epsilon) = \frac{1}{\epsilon} - \gamma_\epsilon + O(\epsilon)$$

$$(1-\epsilon) B(1-\epsilon, 1-\epsilon) = 1 + \epsilon + O(\epsilon^2)$$

γ_ϵ = Euler - Mascheroni constant = 0.57721

Note: In this calculation we haven't encountered γ^5 . γ^5 in D-dim requires a special treatment as it cannot be defined explicitly for arbitrary dimensions.

DR conventions

1. D-dim space-time metric $g^{\mu\nu} = (+, -, \dots, -)$
2. $\text{Tr}[1] = 4$ in the space of gamma matrices
3. $\int \frac{d^D k}{(2\pi)^D}$ defines the integral measure
4. γ_5 is an object that satisfies $\{\gamma_5, \gamma^\mu\} = 0$

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Self energy correction in an arbitrary covariant gauge

$$\Sigma(\not{p}) = g_s^2 C_F \int \frac{d^D k}{(2\pi)^D} \frac{1}{i k^2 (k-p)^2} \left\{ \gamma_\mu (\not{k} - \not{p}) \gamma^\mu (1-\alpha) \frac{\not{k} (\not{k} - \not{p}) \not{k}}{k^2} \right\}$$

we need to calculate the term prop to $(1-\alpha)$. we write

$$\Sigma(\not{p}) = \Sigma_1(\not{p}) - (1-\alpha) \Sigma_2(\not{p})$$

$$\Sigma_1(\not{p}) = \frac{g_s^2}{(4\pi)^2} C_F \not{p} \left(\frac{1}{\epsilon} - \gamma + 1 - \ln \left(\frac{-p^2}{4\pi\mu^2} \right) \right) + O(\epsilon)$$

calculated at tag 16.

$$\Sigma_2(\not{p}) = g_s^2 C_F \int \frac{d^D k}{(2\pi)^D} \frac{\not{k} (\not{k} - \not{p}) \not{k}}{i (k^2)^2 (k-p)^2}$$

Now we use the following Feynman parameter,

$$\frac{1}{AB^2} = 2 \int_0^1 dx \frac{(1-x)}{\{xA + (1-x)B\}^3}$$

$$\text{Again } L = -x(1-x)p^2$$

At this point we need to "massage" the denominator in such a way that we can use a shift in the k momentum.

$$k \rightarrow k - xp + xp$$

The $\Sigma_2(p)$ expression then becomes

(23)

$$\Sigma_2(p) = -2g_s^2 G \int_0^1 dx (1-x) \int \frac{d^D k}{(2\pi)^D i} \frac{\cancel{k}(\cancel{k}-\cancel{p})\cancel{k}}{\{-(k-xp)^2 + L\}^3}$$

At this point we shift the k -momentum

$$k' \Rightarrow k - xp$$

and discard all the terms odd in k'

$$\Sigma_2(p) = 2g_s^2 G \int_0^1 dx (1-x) \int \frac{d^D k'}{(2\pi)^D} \frac{(1-x)\cancel{k}'\cancel{k}' - 2x\cancel{k}'^2 - xL\cancel{k}'}{(-k'^2 + L)^3}$$

For any integrable function $f(k'^2)$ we use the following formula

$$\int d^D k' k'_\mu k'_\nu f(k'^2) = \frac{1}{D} g_{\mu\nu} \int d^D k' k'^2 f(k'^2)$$

In fact, remember that $g^{\mu\nu} g_{\mu\nu} = D \Rightarrow$

$$\frac{D}{D} \int d^D k' k'_\mu k'_\nu f(k'^2) = \frac{g_{\mu\nu}}{D} \int d^D k' k'^2 f(k'^2)$$

We obtain

$$\Sigma_2(p) = 2g_s^2 G \int_0^1 dx (1-x) \int \frac{d^D k'}{(2\pi)^D i} \frac{1}{[-k'^2 + L]^3} \left\{ \left(\frac{2(1-x)}{D} - 1 - x \right) k'^2 - xL \right\}$$

which becomes, after a Wick rotation,

$$\Sigma_2(p) = 2g_s^2 G_F \not{p} \int_0^1 dx (1-x) \int \frac{d^D k}{(2\pi)^D} \left\{ \frac{1}{(k^2 + L)^3} \left[\left(\frac{2(1-x)}{D} - 1 - 2x \right) L \right. \right. \quad (24)$$

$$\left. \left. - \frac{1}{(k^2 + L)^2} \left[\frac{2(1-x)}{D} - 1 - x \right] \right] \right\}$$

We use the generalized result for the k integration

$$\int \frac{d^D k}{(2\pi)^D} \frac{1}{(k^2 + L)^a} = \frac{\Gamma(a - D/2)}{(4\pi)^{D/2} \Gamma(a)} L^{D/2 - a} \quad a \in \mathbb{C} \ni \text{Re}(a) > 0$$

We use the following relations for the x integration

$$B(p, q) = \int_0^1 dx x^{p-1} (1-x)^{q-1}$$

$$B(p, q) = \frac{\Gamma(p) \Gamma(q)}{\Gamma(p+q)}$$

$$\Gamma(z+1) = z \Gamma(z)$$

With these results the expression for $\Sigma_2(p)$ is given by

$$\Sigma_2(p) = 2g_s^2 G_F \not{p} (-p^2)^{D/2-2} (D-1) B(D/2, D/2) \Gamma(2-D/2)$$

which is exactly equal to $\Sigma(p)$ calculated at pag 16!