

## Homework 6

### Problem 9.1

$$\phi(x) = \frac{i}{\hbar} \int_0^a p(x) dx, \quad \text{with}$$

$$p(x) = \sqrt{2m(E - V(x))}$$

- Based on  $V(x)$ :

$$\begin{aligned} \phi(x) &= \frac{i}{\hbar} \sqrt{2m} \left[ \int_0^{a/2} \sqrt{E - V_0} dx + \int_{a/2}^a \sqrt{E} dx \right] \\ &= \frac{i}{\hbar} \sqrt{2m} \frac{a}{2} (\sqrt{E - V_0} + \sqrt{E}) \end{aligned}$$

- Now,

$$\begin{aligned} \psi(x) &= \frac{1}{\sqrt{p(x)}} e^{i\phi(x)} \\ &= \frac{1}{\sqrt{p(x)}} [C_1 \cos \phi + C_2 \sin \phi] \end{aligned}$$

- For our case, boundary conditions:

$$\begin{array}{cc} \psi(0) = \psi(a) = 0 \\ \textcircled{1} \quad \quad \quad \textcircled{2} \end{array}$$

- ① implies  $C_1 = 0$ ; ② implies:

$$\begin{aligned} \phi(x) &= \frac{a\sqrt{2m}}{2\hbar} (\sqrt{E - V_0} + \sqrt{E}) = n\pi \\ &= \frac{2ma^2}{4\hbar^2} (E - V_0 + E + 2\sqrt{E(E - V_0)}) = n^2\pi^2 \\ &= \frac{1}{4} (2E - V_0 + 2\sqrt{E(E - V_0)}) = E_n^0 \end{aligned}$$

$$\rightarrow 2\sqrt{E(E - V_0)} = 4E_n^0 + V_0 - 2E$$

$$4(E^2 - EV_0) = 16(E_n^0)^2 + V_0^2 + 4E^2 + 8E_n^0 V_0 - 16E_n^0 E - 4V_0 E$$

$$0 = 16(E_n^0)^2 + V_0^2 + 8E_n^0 V_0 - 16E_n^0 E$$

$$16E_n^0 E = 16(E_n^0)^2 + V_0^2 + 8E_n^0 V_0$$

$$\rightarrow \boxed{E = E_n^0 + \frac{V_0}{2} + \frac{V_0^2}{16E_n^0}}$$

- If  $n$  is large or  $V_0$  is small,  $V_0^2$  term vanishes, exactly matching our result from perturbation theory.

## Problem 9.2

- Given  $\psi(x) = e^{i f(x)/\hbar}$ ,

we use  $\psi''(x) = -\frac{p^2}{\hbar^2} \psi$

- First:  $\psi'(x) = \frac{i f'(x)}{\hbar} e^{i f(x)/\hbar}$

$$\rightarrow \psi''(x) = \frac{i}{\hbar} f''(x) e^{i f(x)/\hbar} + \left[ \frac{i}{\hbar} f'(x) \right]^2 e^{i f(x)/\hbar}$$

$$\psi''(x) = \frac{i}{\hbar} f''(x) - \frac{1}{\hbar^2} [f'(x)]^2 \psi = -\frac{p^2}{\hbar^2} \psi$$

$$\rightarrow \underbrace{i\hbar f''(x) - [f'(x)]^2 + p^2 = 0}$$

If we expand  $f$  in a power series:

$$\begin{cases} f(x) = f_0(x) + \hbar f_1(x) + \hbar^2 f_2(x) \\ f'(x) = f'_0(x) + \hbar f'_1(x) + \hbar^2 f'_2(x) \\ \dots = \dots \end{cases}$$

- Plugging on:

$$i\hbar [f''_0(x) + \hbar f''_1(x)] - [f'_0(x) + \hbar f'_1(x) + \hbar^2 f'_2(x)]^2 + p^2 = 0$$

$$i\hbar f''_0(x) + \hbar^2 f''_1(x) - [f'_0(x)]^2 - \hbar^2 [f'_1(x)]^2$$

$$- 2\hbar f'_0(x) f'_1(x) - 2\hbar^2 f'_0(x) f'_2(x) + p^2 = 0$$

$$\hbar^0 [p^2 - (f'_0)^2] + \hbar [i f''_0 - 2 f'_0 f'_1] + \hbar^2 [f''_1 - (f'_1)^2 - 2 f'_0 f'_2] = 0$$

$$\Rightarrow \begin{cases} p^2 - (f'_0)^2 = 0 \\ i f''_0 - 2 f'_0 f'_1 = 0 \\ f''_1 - (f'_1)^2 - 2 f'_0 f'_2 = 0 \end{cases}$$

- Obviously,  $f'_0 = p$ , since  $p$  can run negative.

so,  $f_0 = \int p dx$  and  $f''_0 = p'$

- from 2<sup>nd</sup> EQ:

$$i p' = 2 p f'_1 \rightarrow f'_1 = \frac{i}{2} \frac{p'}{p}$$

$$\Rightarrow f_1 = \frac{i}{2} \int \frac{p'}{p} dx$$

$$= \frac{i}{2} \int \frac{d}{dx} [\ln p] dx$$

$$= \frac{i}{2} [p(x) + c]$$

- w/ this, to order  $\hbar$ :

$$f(x) = \int p dx + \frac{i\hbar}{2} p(x) + c$$

- Thus:

$$\psi(x) = \exp \left\{ \frac{i}{\hbar} \left[ \int p dx + \frac{i\hbar}{2} p(x) + c \right] \right\}$$

$$= C \exp \left[ \frac{i}{\hbar} \int p dx - \frac{1}{2} p(x) \right]$$

$$\psi(x) = \frac{C}{\sqrt{p(x)}} \exp \left[ \frac{i}{\hbar} \int p dx \right]$$

### Problem 9.3

- we know

$$T = e^{-2\gamma}, \quad \gamma = \frac{1}{\hbar} \int_0^{2a} |p(x)| dx$$

- since our barrier  $= V_0$  for  $[0, 2a]$

$$\Rightarrow \gamma = \frac{1}{\hbar} \int_0^{2a} \sqrt{2m(V_0 - E)} dx$$

$$= \frac{2a}{\hbar} \sqrt{2m(V_0 - E)}$$

$$\Rightarrow T = \exp \left[ \frac{-4a}{\hbar} \sqrt{2m(V_0 - E)} \right]$$

- The answer to 2.33 is:

$$T^{-1} = 1 + \frac{V_0^2}{4E(V_0 - E)} \sinh^2 \left( \frac{2a}{\hbar} \sqrt{2m(V_0 - E)} \right)$$

- If  $T \ll 1$ ,  $T^{-1}$  is very large; the  $+1$  can be ignored

- Also:  $\sinh x = \frac{1}{2}(e^x - e^{-x})$  - for us,  $x > 0$ ,  
so the  $e^{-x}$  term would be small, meaning

$$T^{-1} = \frac{V_0^2}{16E(V_0 - E)} \exp \left[ \frac{4a}{\hbar} \sqrt{2m(V_0 - E)} \right]$$

$$\rightarrow T = \frac{16E(V_0 - E)}{V_0^2} \exp \left[ \frac{-4a}{\hbar} \sqrt{2m(V_0 - E)} \right]$$

- The coefficient is  $\sim 1$ , so the behavior is dominated by the exponential; thus

$$T \approx \exp \left[ \frac{-4a}{\hbar} \sqrt{2m(V_0 - E)} \right]$$

which matches the WKB Approx. result.

### Problem 9.4

$U^{238}$  has  $Z = 92$  and  $A = 238$ . First:

$$m_p = 238.05078826 \text{ u},$$

$$m_d = m_{Th^{234}} = 234.0436 \text{ u},$$

$$m_\alpha = 4.001506 \text{ u}, \text{ with}$$

$$1 \text{ u} = 931.4941 \text{ MeV}/c^2$$

$$\rightarrow E = m_p c^2 - m_d c^2 - m_\alpha c^2 = 5.2930$$

$$\rightarrow r_1 = 1.07 \sqrt[3]{238} \text{ fm} = 6.63 \text{ fm}$$

$$\begin{aligned} \rightarrow \gamma &= 1.980 \cdot \frac{92}{\sqrt{5.2930}} - 1.485 \sqrt{(92)(6.63)} \\ &= 42.502 \end{aligned}$$

$$\begin{aligned} \rightarrow V &= \sqrt{\frac{2E}{m_\alpha}} = \sqrt{\frac{(2)(5.2930)}{(4.001506)(931.4941)}} = 0.05329 c \\ &= 1.599 \times 10^7 \text{ m/s} = 1.599 \times 10^{22} \text{ fm/s} \end{aligned}$$

$$\begin{aligned} \Rightarrow \tau &= \frac{2 \cdot 6.63}{1.599 \times 10^{22}} \exp[85.004] \text{ s} \\ &\approx 6.846 \times 10^{15} \text{ s} \cdot \frac{1 \text{ yr}}{3.154 \times 10^7 \text{ s}} \end{aligned}$$

$$\boxed{\tau \approx 2.171 \times 10^8 \text{ yr}}$$

For  $Po^{212}$ ,  $Z = 84$ ,  $A = 212$ ,  $m = 211.9889 \text{ u}$

daughter is  $Pb^{208}$   $m = 207.9767$

$$\begin{aligned} \rightarrow E &= 211.9889 - 207.9767 - 4.001506 \text{ u} \\ &\approx 9.96 \text{ MeV} \end{aligned}$$

$$r_1 = 1.07 \sqrt[3]{212} \approx 6.38 \text{ fm}$$

$$\begin{aligned} \gamma &= 1.98 \frac{84}{\sqrt{9.96}} - 1.485 \sqrt{(84)(6.38)} \\ &\approx 18.323 \end{aligned}$$

$$\begin{aligned} V &= \sqrt{\frac{2E}{m_\alpha}} = \sqrt{\frac{(2)(9.96)}{(4.001506)(931.4941)}} \\ &\approx 0.0731 c = 2.193 \times 10^7 \text{ m/s} \\ &= 2.193 \times 10^{22} \text{ fm/s} \end{aligned}$$

$$\begin{aligned} \tau &= \frac{2 \cdot 6.38}{2.193 \times 10^{22}} \exp[36.646] \\ &\approx 4.786 \times 10^{-6} \text{ s} \end{aligned}$$

$$\boxed{\tau \approx 4.786 \text{ } \mu\text{s}}$$

### Problem 9.5

- need functional form of  $V(x)$
- at first, it's just a wall of  $V(x) = E_g$ , but with the decreasing factor from electric field:
- if we also consider the electron w/ energy  $E$ , then the potential looks like:

$$V(x) = E + E_g - eE_0 x$$

$$\begin{aligned}\rightarrow \gamma &= \frac{1}{\hbar} \int_0^a \left| \sqrt{2m[E - (E + E_g - eE_0 x)]} \right| dx \\&= \frac{1}{\hbar} \int_0^a \sqrt{2m(E_g - eE_0 x)} dx \quad \begin{array}{l} u = E_g - eE_0 x \\ du = -eE_0 dx \end{array} \\&= \frac{\sqrt{2m}}{\hbar} \int -\frac{1}{eE_0} u^{1/2} du \\&= -\frac{\sqrt{2m}}{eE_0} \cdot \frac{2}{3} \left[ (E_g - eE_0 x)^{3/2} \right]_0^a \\&= \frac{-2\sqrt{2m}}{3eE_0} \left[ (E_g - eE_0 a)^{3/2} - E_g^{3/2} \right]\end{aligned}$$

- But, as this is the classical region, we know

$$V(a) = E$$

$$\rightarrow E + E_g - eE_0 a = E \Rightarrow E_g = eE_0 a,$$

so

$$\gamma = \frac{2\sqrt{2m}}{3eE_0} E_g^{3/2},$$

meaning:

$$T \approx \exp \left[ \frac{-4\sqrt{2m}}{3eE_0} E_g^{3/2} \right]$$