

## Discrete and continuous mechanical systems (3)

single particle dynamics  $\rightarrow$  can be inferred from Lagrange's equation of motion.

$$q_i = \text{generalized coordinate} \quad \frac{d}{dt} \left( \frac{\partial L}{\partial \dot{q}_i} \right) - \left( \frac{\partial L}{\partial q_i} \right) = 0 \quad 1)$$

which is derived from Hamilton's variational principle

$$\delta \int_{t_1}^{t_2} L(q_i, \dot{q}_i) dt = 0 \quad 2)$$

$L$  is the Lagrangian (assumed here not to depend explicitly on time)

$$L = T - V \quad 3)$$

$T$  = kinetic energy

$V$  = potential energy

The variation from  $t_1$  to  $t_2$  in Eq. 2) is to be taken over an arbitrary path  $q_i(t)$  such that

$$\delta q_i(t_1) = \delta q_i(t_2) = 0$$

The Hamiltonian of the system can be obtained (4)  
from the Lagrangian by performing a Legendre  
transformation

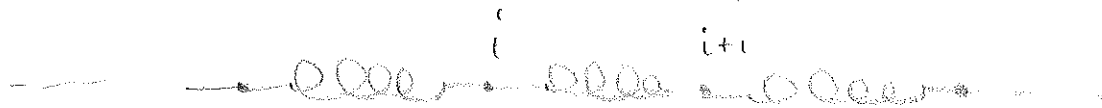
$$H = \sum_i p_i \dot{q}_i - L$$

where  $p_i$  is the "canonical conjugate" of  $q_i$

$$p_i = \frac{\partial L}{\partial \dot{q}_i}$$

All of these considerations can be generalized to  
a system with many particles.

Let's consider  $N$  particles connected with identical  
springs with constant  $k$ , aligned in 1-dim.



$$L = \frac{1}{2} \sum_{i=1}^N \left[ m \dot{\eta}_i^2 - k (\eta_{i+1} - \eta_i)^2 \right] =$$

$\eta_i$  = displacement  
of the  $i$ th  
particle from  
its equilibrium  
position.

$$= \sum_{i=1}^N a \frac{1}{2} \left[ \frac{m}{a} \dot{\eta}_i^2 - ka \left( \frac{\eta_{i+1} - \eta_i}{a} \right)^2 \right] =$$

$$= \sum_i a \mathcal{L}_i$$

$a$  = separation distance between the equilibrium  
positions of two neighboring particles

(5)

$\mathcal{L}_i$  = linear Lagrangian density, that is  
Lagrangian density per unit length.

From the discrete to the continuous:

$n^\circ$  of degrees of freedom  $\rightarrow \infty \Rightarrow$  separation  $\rightarrow$  infinitesimal

$a \rightarrow dx$ ;  $\frac{m}{a} \rightarrow \mu$  = linear mass density

$$\frac{\eta_{i+1} - \eta_i}{a} \rightarrow \frac{\partial \eta}{\partial x}; \quad ka \Rightarrow Y \text{ Young's modulus}$$

Now we have that

$$L = \int \mathcal{L} dx$$

where

$$\mathcal{L} = \frac{1}{2} \left[ \mu \dot{\eta}^2 - Y \left( \frac{\partial \eta}{\partial x} \right)^2 \right]$$

$\eta = \eta(x, t)$  function of continuous parameters  $x, t$

$\eta$  = generalized coordinate  $\Leftrightarrow q_i$  in  $L$

• Variational principle in the continuous case

$$\delta \int_{t_1}^{t_2} L dt = \delta \int_{t_1}^{t_2} dt \int dx \mathcal{L}(\eta, \dot{\eta}, \frac{\partial \eta}{\partial x})$$