

# Homework 9

## Problem 11.1

a) if  $\kappa$  is just a constant, we simply:

$$\int \frac{1}{f} df = \kappa \int dt$$

$$\ln f = \kappa t + c$$

$$\rightarrow f(t) = Ae^{\kappa t}, \text{ } A \text{ is constant.}$$

b) if  $\kappa = \kappa(t)$ :

$$\int \frac{1}{f} df = \int \kappa dt$$

$$\rightarrow f(t) = \exp(\int \kappa dt),$$

which is not so simple; we need the functional form of  $\kappa$ .

c) We cannot, in general, go from

$$\Psi(t) = \Psi(0) \exp\left\{-\frac{i}{\hbar} [H_1 \tau + H_2 (t - \tau)]\right\}$$

$$\rightarrow \Psi(t) = \Psi(0) \exp\left(-\frac{i}{\hbar} H_1 \tau\right) \exp\left[-\frac{i}{\hbar} H_2 (t - \tau)\right]$$

because this would require that  $H_1$  and  $H_2$  commute, something that doesn't have to be true, in general.  $\kappa = \kappa(t)$  is annoying!

## Problem 11.2

- first, from Laporte's rule,

$$\langle n'l'm' | z | nlm \rangle = 0 \text{ if } l' + l = \text{even}$$

Thus,  $H'_{ii}$  implies  $l' = l$ , and  $l' + l = 2l = \text{even}$ ,

so  $H'_{ii} = 0$ . So too is  $\langle 200 | z | 100 \rangle$ , so

we only care about:

$$\begin{aligned}\psi_{100} &= 2a^{-3/2} e^{-r/a} \sqrt{\frac{1}{4\pi}} \\ \psi_{210} &= \underbrace{\frac{1}{2\sqrt{6}} a^{-3/2} \left(\frac{r}{a}\right)} e^{-r/2a} \sqrt{\frac{3}{4\pi}} \cos\theta \\ \psi_{21\pm 1} &= \leftarrow \cdot \mp \sqrt{\frac{3}{8\pi}} \sin\theta e^{\pm i\phi}\end{aligned}$$

but in  $\langle \psi_{21\pm 1} | z | \psi_{100} \rangle$ , we get

$$\int d\phi e^{\pm i\phi}, \text{ which is zero, so}$$

only integral is

$$\langle 210 | z | 100 \rangle = eE \left( 2a^{-3/2} \sqrt{\frac{1}{4\pi}} \right) \left( \frac{1}{2\sqrt{6}} a^{-5/2} \sqrt{\frac{3}{4\pi}} \right)$$

$$\times \int d^3r z e^{-r/a} r e^{-r/2a} \cos\theta$$

$$= \frac{eE}{4\pi a^4 \sqrt{2}} \int a^3 r r^2 \cos^2\theta e^{-3r/2a}$$

$$= \downarrow \underbrace{\left( \int_0^\infty dr r^4 e^{-3r/2a} \right)}_{(1)} \underbrace{\left( \int_0^\pi d\theta \cos^2\theta \sin\theta \right) \cdot 2\pi}_{(2)}$$

$$\begin{aligned}(1) \quad x \equiv \frac{2a}{3} &\rightarrow \int_0^\infty dr r^4 e^{-r/a} = 4! \left( \frac{2a}{3} \right)^5 \\ &= 24 \left( \frac{2^5}{3^5} \right) a^5\end{aligned}$$

$$(2) \quad u \equiv \cos\theta \quad du = -\sin\theta d\theta \quad \pi \rightarrow -1, \quad 0 \rightarrow 1$$

$$\Rightarrow \int_{-1}^1 du u^2 = \frac{1}{3} [u^3]_{-1}^1 = \frac{2}{3}$$

$$\rightarrow \frac{eE}{4\pi a^4 \sqrt{2}} \cdot 24 \left( \frac{2^5}{3^5} \right) a^5 \cdot \frac{2}{3} \cdot 2\pi$$

$$= \left( \frac{2^5}{3^5} \right) \frac{8eEa}{\sqrt{2}} = \boxed{\left( \frac{2^8}{3^5} \right) \frac{eEa}{\sqrt{2}}}$$

This is the only non-zero matrix element.

# Problem 12.13

- From selection rules, if  $l=l'$ , the matrix element

$$\langle \psi_{100} | \vec{r} | \psi_{200} \rangle$$

vanishes, so  $A=0$  and  $\tau = \frac{1}{A} \rightarrow \text{infinite}$ .

- In the other case,

$$\langle \psi_{100} | \vec{r} | \psi_{210} \rangle$$

has  $x$  and  $y$  components  $=0$ ; only  $z$  remains: we calculated it already:

$$\langle \psi_{100} | z | \psi_{210} \rangle = \frac{a}{\sqrt{2}} \left( \frac{2^8}{3^5} \right) = 2a \left( \frac{2^7}{3^5} \right)$$

- For  $\langle \psi_{100} | \vec{r} | \psi_{21\pm 1} \rangle$ , the  $z$  component is 0, and  $x, y$  are related.

$$\psi_{100} = \frac{1}{\sqrt{\pi a^3}} e^{-r/a}$$

$$\begin{aligned} \psi_{21-1} &= \frac{1}{2\sqrt{6a^5}} \sqrt{\frac{3}{8\pi}} r e^{-r/2a} \sin\theta e^{-i\phi} \\ &= \frac{1}{8\sqrt{\pi a^5}} r e^{-r/2a} \sin\theta e^{-i\phi} \end{aligned}$$

$$\begin{aligned} \rightarrow \langle \psi_{100} | x | \psi_{21-1} \rangle &= \frac{1}{8\pi a^4} \int a^3 r e^{-3r/2a} \sin\theta e^{-i\phi} x \\ &= \frac{1}{8\pi a^4} \int_0^\infty dr r^4 e^{-3r/2a} \int_0^\pi d\theta \sin^3\theta \int_0^{2\pi} d\phi e^{-i\phi} \cos\phi \end{aligned}$$

$$\textcircled{1} = 4! \left( \frac{2a}{3} \right)^5 = 24a^5 \left( \frac{2^5}{3^5} \right)$$

$$\textcircled{2} = \frac{1}{4} \int_0^\pi (3\sin\theta - \sin 3\theta) d\theta$$

$$= \frac{1}{4} \left[ -3\cos\theta + \frac{1}{3}\cos 3\theta \right]_0^\pi = \frac{1}{2} \left( 3 - \frac{1}{3} \right) = \frac{4}{3}$$

$$\textcircled{3} = \frac{1}{2} \int_0^{2\pi} e^{-i\phi} (e^{i\phi} + e^{i\phi}) d\phi$$

$$= \frac{1}{2} \int_0^{2\pi} (1 + e^{-2i\phi}) d\phi = \pi$$

$$\rightarrow \frac{1}{8\pi a^4} \cdot 24a^5 \left( \frac{2^5}{3^5} \right) \cdot \frac{4}{3} \cdot \pi = a \left( \frac{2^7}{3^5} \right)$$

- Now,  $\langle 100 | x | 21-1 \rangle = -i \langle 100 | y | 21-1 \rangle$ , so

$$|\langle 100 | \vec{r} | 21-1 \rangle|^2 = 2a^2 \left( \frac{2^7}{3^5} \right)^2$$

- Since  $\langle 100 | \vec{r} | 211 \rangle$  differs only by a minus,

$$|\langle 100 | \vec{r} | 211 \rangle|^2 = 2a^2 \left( \frac{2^7}{3^5} \right)^2$$

- Similarly:  $|\langle 100 | \vec{r} | 210 \rangle| = \uparrow$ ; all equal!

$$\rightarrow |\vec{P}|^2 = 2e^2 a^2 \left( \frac{2^7}{3^5} \right)^2 = \frac{(ea)^2 2^{15}}{3^{10}}$$

$$\rightarrow A = \frac{\omega_0^3 |\vec{P}|^2}{3\pi\epsilon_0 \hbar c^3}$$

$$\text{- Here, } \omega_0 = \frac{1}{\hbar} (E_2 - E_1) = -\frac{1}{\hbar} \left( \frac{E_1}{4} - E_1 \right) = \frac{-3E_1}{4\hbar}$$

$$\begin{aligned} \Rightarrow A &= -\frac{3^3}{4^3} \frac{E_1^3}{\hbar^3} \frac{e^2 a^2 2^{15}}{3^{10}} \frac{1}{3\pi\epsilon_0 \hbar c^3} \\ &= -\frac{2^9}{3^8} \frac{E_1^3 e^2 a^2}{\pi\epsilon_0 \hbar^4 c^3} \end{aligned}$$

- Plugging everything in, we get a unit of  $s^{-1}$ :

$$A \approx 6.244 \times 10^8 \text{ s}^{-1}$$

$$\rightarrow \tau = \frac{1}{A} \approx 1.6 \times 10^9 \text{ s}$$

## Problem 11.14

$$\langle n'l'm' | [L_z, x] | nlm \rangle = i\hbar \langle n'l'm' | y | nlm \rangle$$

$$\langle n'l'm' | L_z x | nlm \rangle$$

$$- \langle n'l'm' | x L_z | nlm \rangle = \quad \text{---}$$

-  $L_z = L_z^\dagger$  so

$$(m'-m) \langle n'l'm' | x | nlm \rangle = i \langle n'l'm' | y | nlm \rangle \quad (1)$$

- Similarly

$$(m'-m) \langle n'l'm' | y | nlm \rangle = -i \langle n'l'm' | x | nlm \rangle \quad (2)$$

- Lastly:  $(m'-m) \langle n'l'm' | z | nlm \rangle = 0 \quad (3)$

- If  $m=m'$ , from (1) and (2) we can obviously tell the first line of 11.76 (but can't say anything about  $z$ )

- If  $m'=m\pm 1$ , from (3) the  $z$ -component must be zero, and we can easily derive the 2<sup>nd</sup> line of 11.76 from (2).

- Lastly, if  $\Delta m = \pm k$  w/  $k > 1$ , then from (2):

$$\langle n'l'm' | x | nlm \rangle = \pm k i \langle n'l'm' | y | nlm \rangle,$$

but from (1):

$$\langle n'l'm' | x | nlm \rangle = \pm \frac{i}{k} \langle n'l'm' | y | nlm \rangle$$

which is only true if the matrix element vanishes; thus,  $\Delta m = 0$  or  $\pm 1$  is required.

# Problem 11.15

- First:  $[L^2, z] = [L_x^2, z] + [L_y^2, z] + [L_z^2, z]$

$$= \underbrace{L_x [L_x, z] + [L_x, z] L_x}_{(x \rightarrow y)} + (x \rightarrow y)$$

$$\rightarrow L_x (-i\hbar y) + (-i\hbar y) L_x = -i\hbar (L_x y + y L_x)$$

- For  $(x \rightarrow y)$ , add a minus since  $[L_y, z] = i\hbar x$ :

$$[L^2, z] = i\hbar (L_y x + x L_y - L_x y - y L_x)$$

- Now,  $[L_y, x] = L_y x - x L_y$

$$\rightarrow L_y x = [L_y, x] + x L_y = -i\hbar z + x L_y$$

$$L_x y = i\hbar z + y L_x$$

$$\Rightarrow \underline{[L^2, z] = 2i\hbar (x L_y - y L_x - i\hbar z)}$$

- Next,

$$[L^2, [L^2, z]] = 2i\hbar ([L^2, x L_y] - [L^2, y L_x] - i\hbar [L^2, z])$$

- Generically:

$$\begin{aligned} [A, BC] &= ABC - BCA \\ &= ABC - BAC + BAC - BCA \\ &= [A, B]C + B[A, C] \end{aligned}$$

$$\rightarrow [L^2, x L_y] = [L^2, x] L_y + x [L^2, L_y]$$

$$= [L^2, x] L_y$$

$$\rightarrow 2i\hbar ([L^2, x] L_y - [L^2, y] L_x - i\hbar (L^2 z - z L^2))$$

$\Rightarrow x, y, z$  are cyclic; from a)

$$[L^2, x] = 2i\hbar (y L_z - z L_y - i\hbar x)$$

$$[L^2, y] = 2i\hbar (z L_x - x L_z - i\hbar y)$$

$$\rightarrow 2i\hbar [2i\hbar (y L_z - z L_y - i\hbar x) L_y - 2i\hbar (z L_x - x L_z - i\hbar y) L_x - i\hbar (L^2 z + z L^2) + 2i\hbar z L^2]$$

$$= 2\hbar^2 (z L^2 + L^2 z) - 4\hbar^2 [y L_z L_y - z L_y^2 - i\hbar x L_y - z L_x^2 + x L_z L_x + i\hbar y L_x + z L^2]$$

$$= \sim [(y L_z - i\hbar x) L_y + (x L_z + i\hbar y) L_x + z L^2]$$

- Now,  $[L_z, y] = L_z y - y L_z = -i\hbar x$

$$\rightarrow y L_z - i\hbar x = L_z y \quad \text{and}$$

$$x L_z + i\hbar y = L_z x$$

$$\rightarrow \sim [L_z (y L_y + x L_x + z L_z)]$$

$$= L_z \vec{r} \cdot \vec{L} = 0$$

$$\Rightarrow \underline{[L^2, [L^2, z]] = 2\hbar^2 (z L^2 + L^2 z)} \quad \checkmark$$

$\rightarrow$  obviously, in 3d we will have

$$= \underline{2\hbar^2 (\vec{r} L^2 + L^2 \vec{r})}$$

- Plugging in:

$$\begin{aligned} \langle n'l'm' | [L^2 [L^2, \vec{r}]] | nlm \rangle &= \\ &= \langle n'l'm' | L^2 [L^2, \vec{r}] | nlm \rangle - \langle n'l'm' | [L^2, \vec{r}] L^2 | nlm \rangle \\ &= \hbar^2 [l'(l'+1) - l(l+1)] \langle n'l'm' | [L^2, \vec{r}] | nlm \rangle \\ &= \hbar^4 [l'(l'+1) - l(l+1)]^2 \langle n'l'm' | \vec{r} | nlm \rangle \end{aligned}$$

$$\text{also} = 2\hbar^2 \langle n'l'm' | \vec{r} L^2 + L^2 \vec{r} | nlm \rangle$$

$$= 2\hbar^4 [l'(l'+1) + l(l+1)] \langle n'l'm' | \vec{r} | nlm \rangle$$

- So, we must have

$$2[l'(l'+1) + l(l+1)] = [l'(l'+1) - l(l+1)]^2$$

or else the matrix element is zero.

- Rearranging first term:

$$2l'^2 + 2l' + 2l^2 + 2l$$

$$l'^2 + 2l' + 2l'l' + 2l + l^2 + 1 + l'^2 - 2l'l' + l^2 - 1$$

$$\rightarrow (l' + l + 1)^2 + (l' - l)^2 - 1$$

- 2<sup>nd</sup>:  $l'^2 + l' - l^2 - l$

$$(l'^2 - l^2) + (l' - l)$$

$$(l' - l)(l' + l) + (l' - l)$$

$$(l' - l)(l' + l + 1)$$

- Now, let  $a \equiv l' + l + 1$ ,  $b \equiv l' - l$

$$\rightarrow a^2 + b^2 - 1 = a^2 b^2$$

$$a^2 - a^2 b^2 + b^2 - 1 = 0$$

$$a^2(1 - b^2) - 1(1 - b^2) = 0$$

$$(a^2 - 1)(1 - b^2) = 0$$

- Thus,  $(l' - l)^2 = 1$

$$l' - l = \pm 1$$

$$\underline{\underline{l' = l \pm 1}}$$