

Homework 4

Problem 7.6

$$V = \frac{1}{2}m\omega^2 x^2 - \underbrace{qEx}$$

a) we calculate E_1 :

$$\begin{aligned} E_1 &= \langle \psi_n^0 | H' | \psi_n^0 \rangle \\ &= -qE \cdot \frac{2}{a} \langle \psi_n^0 | \hat{x} | \psi_n^0 \rangle \end{aligned}$$

$$\begin{aligned} \text{w/ } \hat{x} &= \sqrt{\frac{\hbar}{2m\omega}} (\hat{a}_+ + \hat{a}_-) : \\ &= -qE \cdot \frac{2}{a} \sqrt{\frac{\hbar}{2m\omega}} \langle \psi_n^0 | (\hat{a}_+ + \hat{a}_-) | \psi_n^0 \rangle \end{aligned}$$

both terms will be different n , so $E_1 = 0$.

- for 2nd order case:

$$E_n^{(2)} = \sum_{m \neq n} \frac{|\langle \psi_m^0 | H' | \psi_n^0 \rangle|^2}{E_n^0 - E_m^0}$$

$$\text{in general, } \langle \psi_m^0 | H' | \psi_n^0 \rangle = -qE \sqrt{\frac{\hbar}{2m\omega}} \langle \psi_m^0 | \hat{a}_+ + \hat{a}_- | \psi_n^0 \rangle$$

- only non-zero for $m-1=n$, $m+1=n$

- denominator is $(n+\frac{1}{2})\hbar\omega - (m+\frac{1}{2})\hbar\omega = (n-m)\hbar\omega$

$$\Rightarrow \frac{q^2 E^2 \hbar}{2m\omega} \left\{ \frac{[\langle \psi_m^0 | \hat{a}_+ + \hat{a}_- | \psi_n^0 \rangle]^2}{\hbar\omega(n-m)} \delta_{n,m-1} + \frac{[\langle \psi_m^0 | \hat{a}_+ + \hat{a}_- | \psi_n^0 \rangle]^2}{\hbar\omega(n-m)} \delta_{n,m+1} \right\}$$

$$= \frac{q^2 E^2}{2m\omega^2} \left\{ \frac{n}{n-(n-1)} + \frac{n+1}{n-(n+1)} \right\}$$

$$= \frac{q^2 E^2}{2m\omega^2} (n - n - 1) = \boxed{-\frac{q^2 E^2}{2m\omega^2} = E_n^{(2)}}$$

$$\text{b) } \hat{H} = \frac{\hat{p}^2}{2m} + \frac{1}{2}m\omega^2 x^2 - qEx$$

taking $x \rightarrow x' \equiv x - \left(\frac{qE}{m\omega^2}\right)$, we find

$$\rightarrow \frac{\partial f}{\partial x} = \frac{\partial f}{\partial x'} \frac{dx'}{dx} = \frac{\partial f}{\partial x'}$$

$$\rightarrow \frac{1}{2}m\omega^2 \left[x' + \frac{qE}{m\omega^2} \right]^2 - qE \left(x' + \frac{qE}{m\omega^2} \right)$$

$$\frac{1}{2}m\omega^2 \left[x'^2 + \frac{2qE}{m\omega^2} x' + \frac{q^2 E^2}{m^2 \omega^4} \right] - qEx' - \frac{q^2 E^2}{m\omega^2}$$

$$= \frac{1}{2}m\omega^2 x'^2 + qEx' + \frac{1}{2} \frac{q^2 E^2}{m\omega^2} - qEx' - \frac{q^2 E^2}{m\omega^2}$$

$$= \frac{1}{2}m\omega^2 x'^2 - \frac{1}{2} \frac{(qE)^2}{m\omega^2}$$

$$\rightarrow \hat{H}' = \frac{\hat{p}'^2}{2m} + \frac{1}{2}m\omega^2 x'^2 - \frac{1}{2} \frac{(qE)^2}{m\omega^2}$$

- This is normal harmonic oscillator w/ a constant term, so we have normal sols plus this constant term:

$$E_n = \left(n + \frac{1}{2}\right) \hbar\omega - \frac{1}{2} \frac{(qE)^2}{m\omega^2}$$

Problem 7.10

- EQ. (7.34) $\rightarrow \bar{E}' = \pm \epsilon \frac{\hbar \omega}{2}$

- we are considering exact solution for

$$\hat{H} = \frac{\hat{p}^2}{2m} + \frac{1}{2} m \omega^2 (x^2 + y^2) + \epsilon m \omega^2 xy$$

which, in rotated/normal coords, have energies:

$$E_{mn} = \left(m + \frac{1}{2}\right) \hbar \omega_+ + \left(n + \frac{1}{2}\right) \hbar \omega_-$$

$$\rightarrow \omega_{\pm} = \sqrt{1 \pm \epsilon} \omega$$

- Plugging in:

$$\begin{aligned} \bar{E}_{mn} &= \left(m + \frac{1}{2}\right) \hbar \sqrt{1 + \epsilon} \omega + \left(n + \frac{1}{2}\right) \hbar \sqrt{1 - \epsilon} \omega \\ &= E_m \sqrt{1 + \epsilon} + E_n \sqrt{1 - \epsilon} \end{aligned}$$

$$\sqrt{1 + \epsilon} \rightarrow 1 + \frac{\epsilon}{2} + \dots, \quad \sqrt{1 - \epsilon} \rightarrow 1 - \frac{\epsilon}{2} + \dots$$

$$\begin{aligned} E_{mn} &= E_m \left(1 + \frac{\epsilon}{2}\right) + E_n \left(1 - \frac{\epsilon}{2}\right) \\ &= E_m + \frac{\epsilon E_m}{2} + E_n - \frac{\epsilon E_n}{2} \\ &= \bar{E}_m + \bar{E}_n + \frac{\epsilon}{2} (E_m - E_n) \end{aligned}$$

- As in Ex 7.2, we consider $m=1, n=0$ and $m=0, n=1$:

$$\textcircled{1} \quad E_1 + E_0 + \frac{\epsilon}{2} (E_1 - E_0), \quad \bar{E}_1 = \frac{3\hbar\omega}{2}, \quad \bar{E}_0 = \frac{\hbar\omega}{2}$$
$$2\hbar\omega + \frac{\epsilon}{2} \hbar\omega$$

$$\textcircled{2} \quad \bar{E}_0 + \bar{E}_1 + \frac{\epsilon}{2} (E_0 - E_1)$$
$$2\hbar\omega - \frac{\epsilon}{2} \hbar\omega$$

- the 1st-order corrections are

$$\pm \epsilon \frac{\hbar\omega}{2},$$

just like the approximation!

Problem 7.11

$$H' = a^3 V_0 \delta(x - \frac{a}{4}) \delta(y - \frac{a}{2}) \delta(z - \frac{3a}{4})$$

$$\rightarrow \psi_n^0 = \left(\frac{2}{a}\right)^{3/2} \sin\left(\frac{\pi x}{a}\right) \sin\left(\frac{\pi y}{a}\right) \sin\left(\frac{\pi z}{a}\right)$$

$$E'_n = \langle \psi_n^0 | H' | \psi_n^0 \rangle$$

- delta fns kill all integrals!

$$= \left(\frac{2}{a}\right)^3 a^3 V_0 \sin^2\left(\frac{\pi}{a} \cdot \frac{a}{4}\right) \sin^2\left(\frac{\pi}{a} \cdot \frac{a}{2}\right) \sin^2\left(\frac{\pi}{a} \cdot \frac{3a}{4}\right)$$

$$= 8V_0 \sin^2\left(\frac{\pi}{4}\right) \sin^2\left(\frac{\pi}{2}\right) \sin^2\left(\frac{3\pi}{4}\right)$$

$$= 8V_0 \cdot \left(\frac{1}{2}\right) \cdot 1 \cdot \left(\frac{1}{2}\right) = \underline{\underline{2V_0}}$$

- three first excited states: $\{2, 1, 1\}_{\psi_1}$, $\{1, 2, 1\}_{\psi_2}$, $\{1, 1, 2\}_{\psi_3}$

- looking to diagonalize W :

$$W_{aa} = \langle \psi_1^0 | H' | \psi_1^0 \rangle =$$

$$= 8V_0 \sin^2\left(\frac{2\pi}{a} \cdot \frac{a}{4}\right) \sin^2\left(\frac{\pi}{a} \cdot \frac{a}{2}\right) \sin^2\left(\frac{\pi}{a} \cdot \frac{3a}{4}\right)$$

$$= 8V_0 \sin^4\left(\frac{\pi}{2}\right) \sin^2\left(\frac{3\pi}{4}\right) = \frac{8V_0}{2} = \underline{\underline{4V_0}}$$

$$W_{bb} = \langle \psi_2^0 | H' | \psi_2^0 \rangle =$$

$$= 8V_0 \sin^2\left(\frac{\pi}{a} \cdot \frac{a}{4}\right) \sin^2\left(\frac{2\pi}{a} \cdot \frac{a}{2}\right) \sin^2\left(\frac{\pi}{a} \cdot \frac{3a}{4}\right) = \underline{\underline{0}}$$

$$W_{cc} = \langle \psi_3^0 | H' | \psi_3^0 \rangle$$

$$= 8V_0 \sin^2\left(\frac{\pi}{a} \cdot \frac{a}{4}\right) \sin^2\left(\frac{\pi}{a} \cdot \frac{a}{2}\right) \sin^2\left(\frac{2\pi}{a} \cdot \frac{3a}{4}\right)$$

$$= 8V_0 \left(\frac{1}{2}\right) \cdot 1 \cdot (-1)^2 = \underline{\underline{4V_0}}$$

$$W_{ab} = W_{ba} = \langle \psi_1^0 | H' | \psi_2^0 \rangle$$

$$= 8V_0 \sin\left(\frac{2\pi}{a} \cdot \frac{a}{4}\right) \sin\left(\frac{\pi}{a} \cdot \frac{a}{4}\right) \sin\left(\frac{2\pi}{a} \cdot \frac{a}{2}\right) \sin\left(\frac{\pi}{a} \cdot \frac{a}{2}\right) \times \sin^2\left(\frac{\pi}{a} \cdot \frac{3a}{4}\right) = 0$$

$$W_{ac} = W_{ca} = \langle \psi_1^0 | H' | \psi_3^0 \rangle$$

$$= 8V_0 \sin\left(\frac{2\pi}{a} \cdot \frac{a}{4}\right) \sin\left(\frac{\pi}{a} \cdot \frac{a}{4}\right) \sin^2\left(\frac{\pi}{a} \cdot \frac{a}{2}\right) \sin\left(\frac{2\pi}{a} \cdot \frac{3a}{4}\right)$$

$$= 8V_0 (1) \left(\frac{1}{\sqrt{2}}\right) (1) (-1) \left(\frac{1}{\sqrt{2}}\right) = \underline{\underline{-4V_0}} \quad \sin\left(\frac{\pi}{a} \cdot \frac{3a}{4}\right)$$

$$W_{bc} = W_{cb} = \langle \psi_2^0 | H' | \psi_3^0 \rangle$$

$$= 8V_0 \sin^2\left(\frac{\pi}{a} \cdot \frac{a}{4}\right) \sin\left(\frac{2\pi}{a} \cdot \frac{a}{2}\right) \sin\left(\frac{2\pi}{a} \cdot \frac{a}{4}\right) \sin\left(\frac{\pi}{a} \cdot \frac{3a}{4}\right) = 0$$

$$\rightarrow W = 4V_0 \begin{pmatrix} 1 & 0 & -1 \\ 0 & 0 & 0 \\ -1 & 0 & 1 \end{pmatrix}$$

$$\rightarrow \begin{vmatrix} 1-\lambda & 0 & -1 \\ 0 & -\lambda & 0 \\ -1 & 0 & 1-\lambda \end{vmatrix} = 0$$

$$\rightarrow (1-\lambda)(-\lambda)(1-\lambda) - 1(-1)(-\lambda)(-1) = 0$$

$$-\lambda(1-\lambda)^2 + \lambda = 0$$

$$\lambda(-(1-\lambda)^2 + 1) = 0$$

$$\lambda(-\lambda^2 + 2\lambda) = 0 \rightarrow \lambda^2(\lambda - 2) = 0 \rightarrow \lambda = 0, 0, 2$$

$$\Rightarrow \underline{\underline{E' = 0, 0, 8V_0}}$$

Problem 7.13

$$\psi^0 = \sum_{j=1}^n \alpha_j \psi_j^0$$

- Following Section 7.2.1, we first have that:

$$\hat{H} \psi_i^0 = E_i^0 \psi_i^0, \quad \langle \psi_i^0 | \psi_j^0 \rangle = \delta_{ij}$$

$$\rightarrow \hat{H} \psi^0 = E^0 \psi^0$$

- Everything is general up until EQ 7.26:

$$H^0 \psi' + H' \psi^0 = E^0 \psi' + E' \psi^0$$

- Inner product w/ $\langle \psi_i^0 |$:

$$\langle \psi_i^0 | H^0 | \psi' \rangle + \langle \psi_i^0 | H' | \psi^0 \rangle = E^0 \langle \psi_i^0 | \psi' \rangle + E' \langle \psi_i^0 | \psi^0 \rangle$$

- Again, $H^0 = (H^0)^\dagger$, so $\overbrace{\langle \psi_i^0 | H^0 | \psi' \rangle}^{\text{cancel}}$ and $\overbrace{E^0 \langle \psi_i^0 | \psi' \rangle}^{\text{cancel}}$:

$$\langle \psi_i^0 | H' | \psi^0 \rangle = E' \langle \psi_i^0 | \psi^0 \rangle$$

$$\rightarrow \langle \psi_i^0 | H' \left(\sum_{j=0}^n \alpha_j \psi_j^0 \right) \rangle = E' \alpha_i$$

$$\rightarrow \sum_{j=0}^n \alpha_j \langle \psi_i^0 | H' | \psi_j^0 \rangle = E' \alpha_i$$

$$\rightarrow \underline{\sum_{j=0}^n \alpha_j W_{ij} = E' \alpha_i}$$

- In the $n=2$ case, we had

$$\begin{pmatrix} W_{aa} & W_{ab} \\ W_{ba} & W_{bb} \end{pmatrix} \begin{pmatrix} \alpha \\ \beta \end{pmatrix} = E' \begin{pmatrix} \alpha \\ \beta \end{pmatrix},$$

from EQ 7.27 & 7.29, our EQ is the generalization of this, meaning the same conclusion applies, that we have an eigenval equation which yields 1st-order energy corrections.