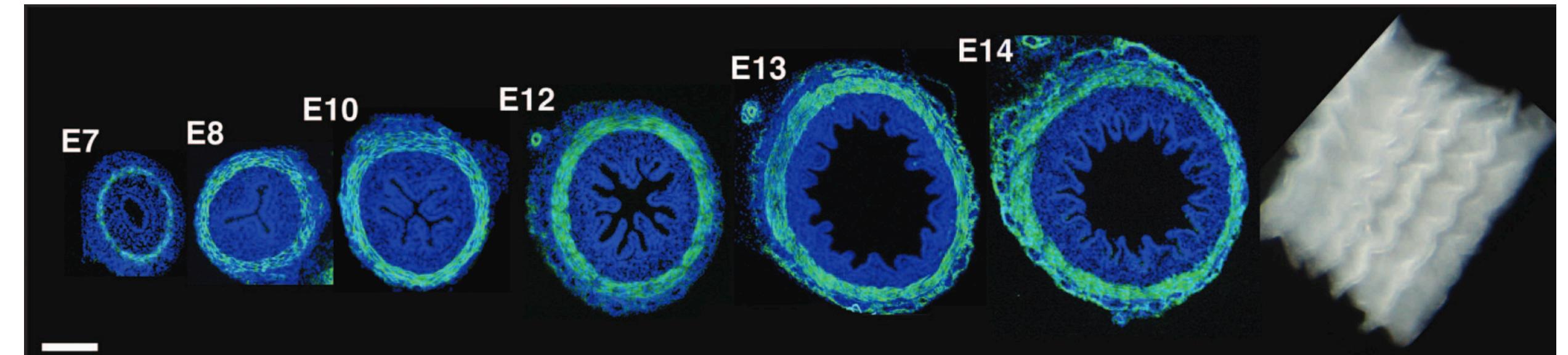
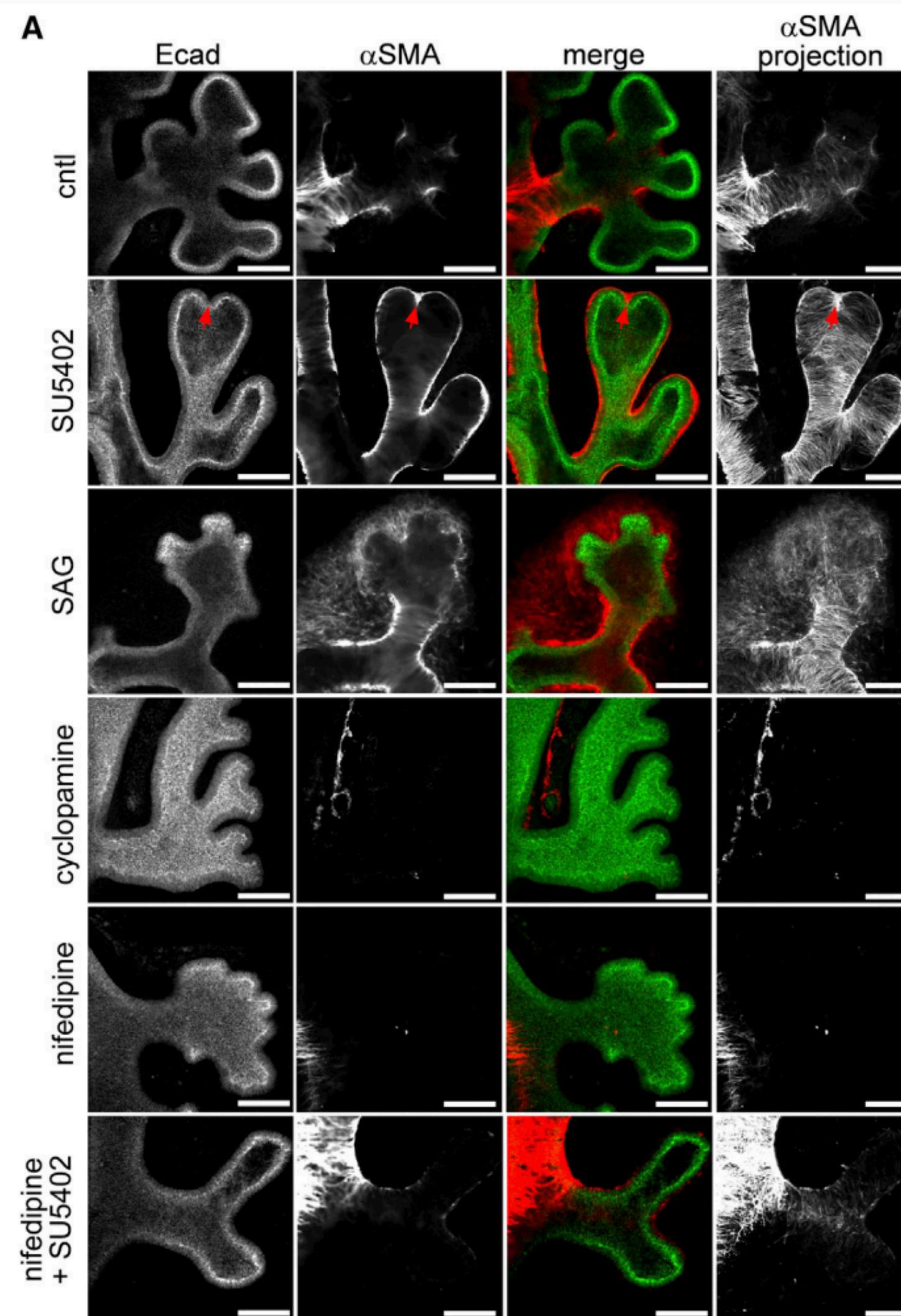


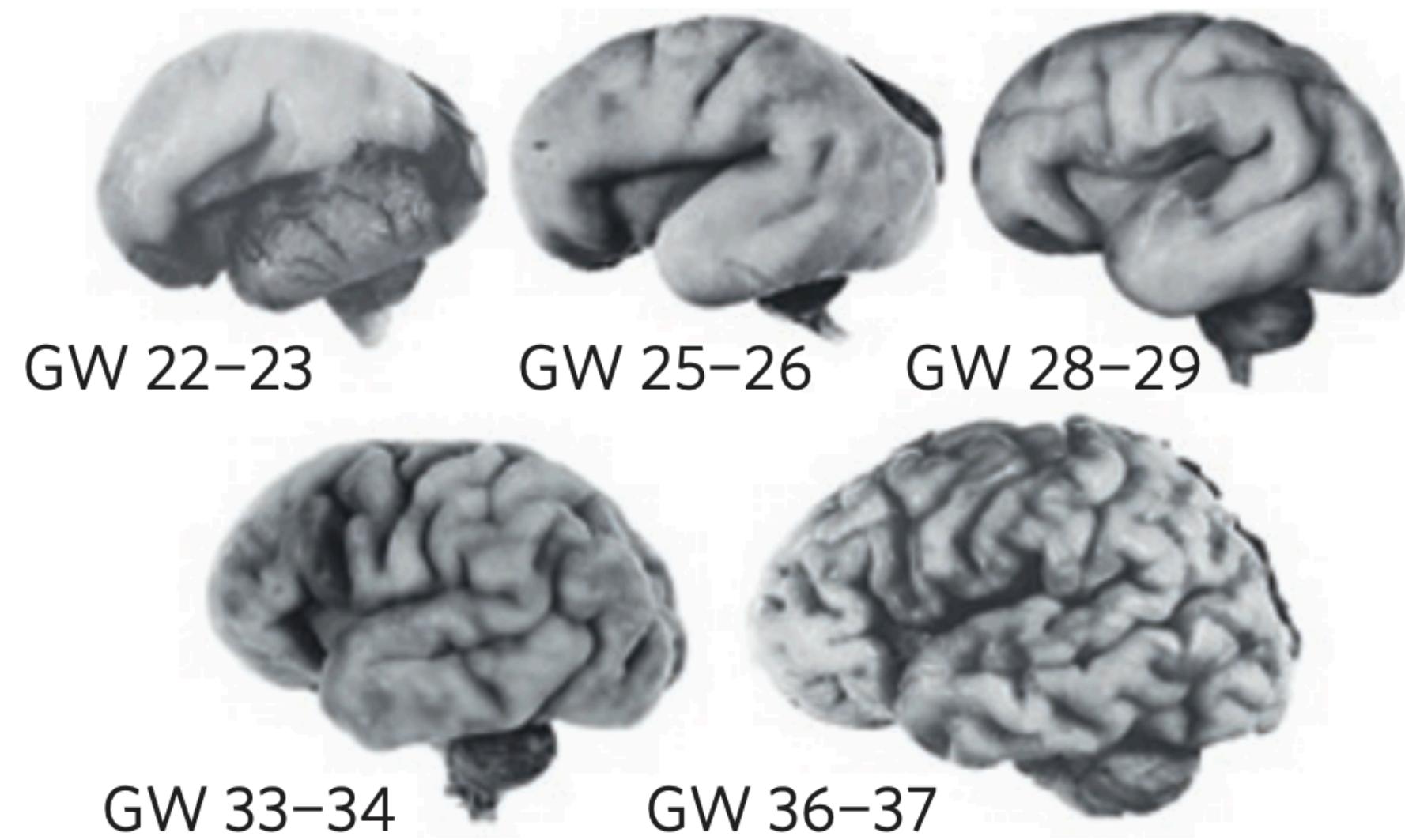
SOME INSTABILITIES IN CONTINUUM MECHANICS OF TISSUES

by:
Chandraniva Guha Ray
Supervisor:
Dr. Pierre Haas

SHAPE DETERMINATION IN BIOLOGY: A FUNDAMENTAL QUESTION



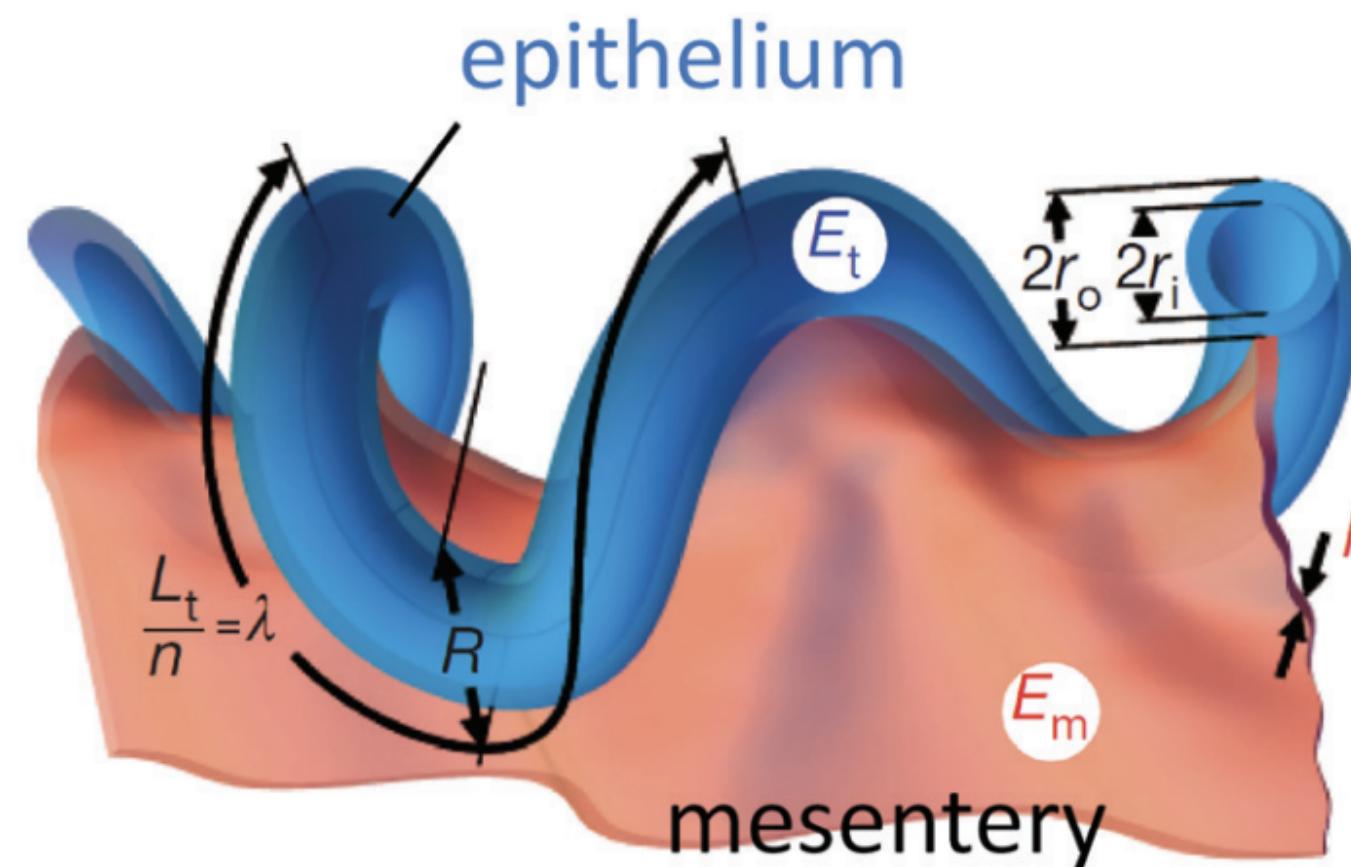
Shyer et al., Villification: How the Gut Gets Its Villi (2013)



Kim et al., Localized Smooth Muscle Differentiation Is Essential for Epithelial Bifurcation during Branching Morphogenesis of the Mammalian Lung (2015)

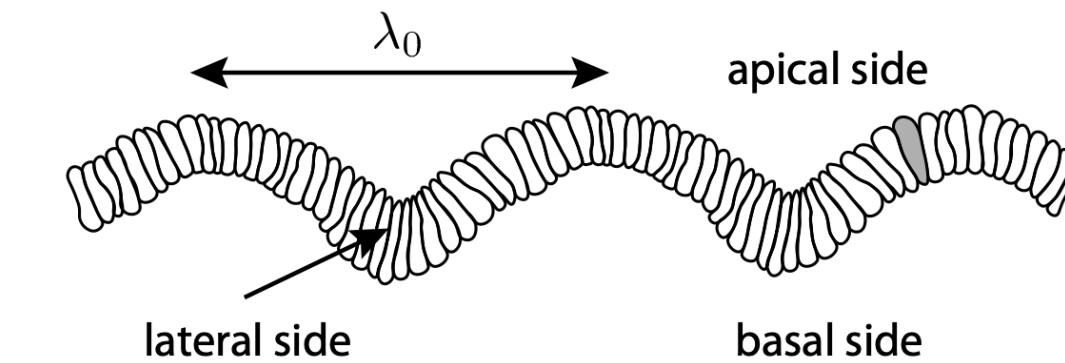
Griffiths et al., Atlas of Fetal and Neonatal Brain MR Imaging (2010)

SHAPES VIA INSTABILITIES

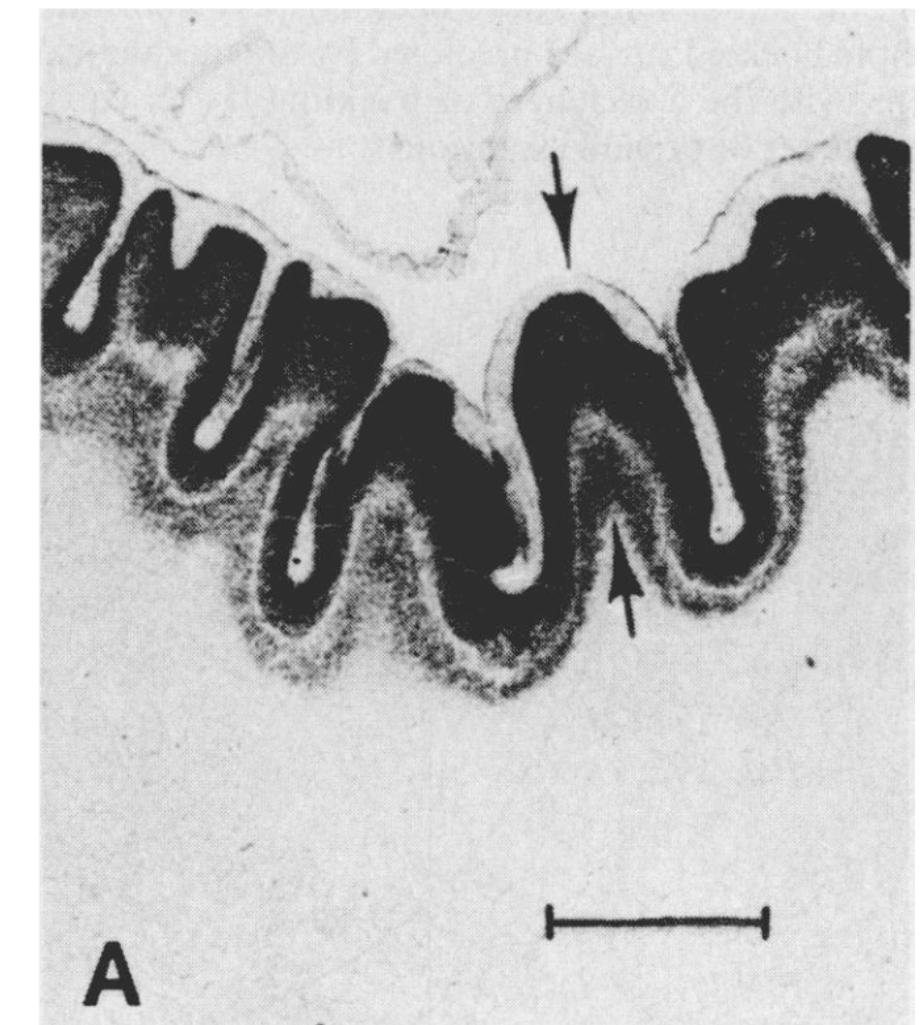


Savin et. al., On the growth and form of the gut (2011)

Wrinkling

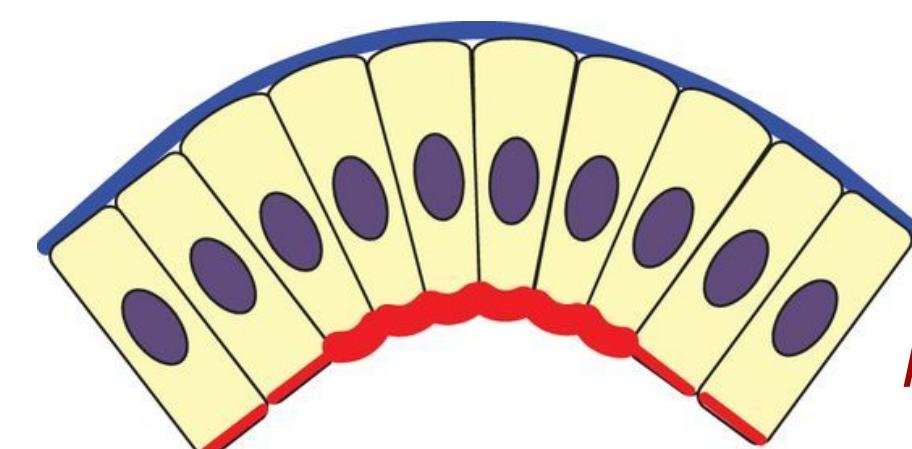
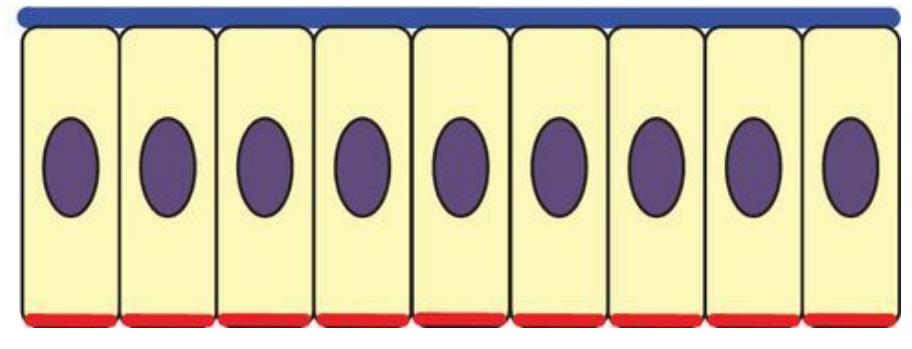


Andrenšek et al., PRL (2023)



Richmond et al., Mechanical model of brain convolution development (1975)

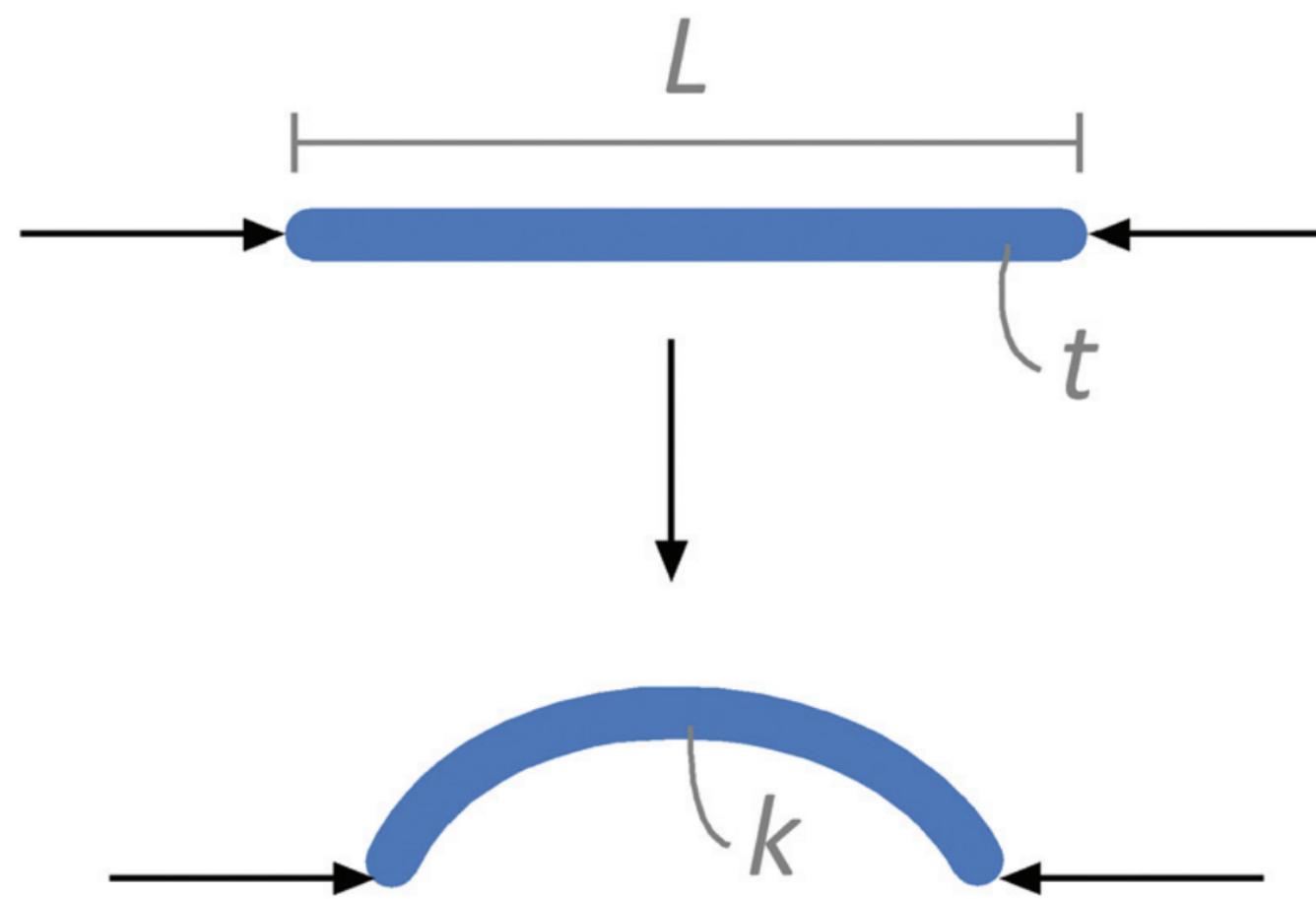
Buckling



Pearl et. al., Cellular systems for epithelial invagination (2016)

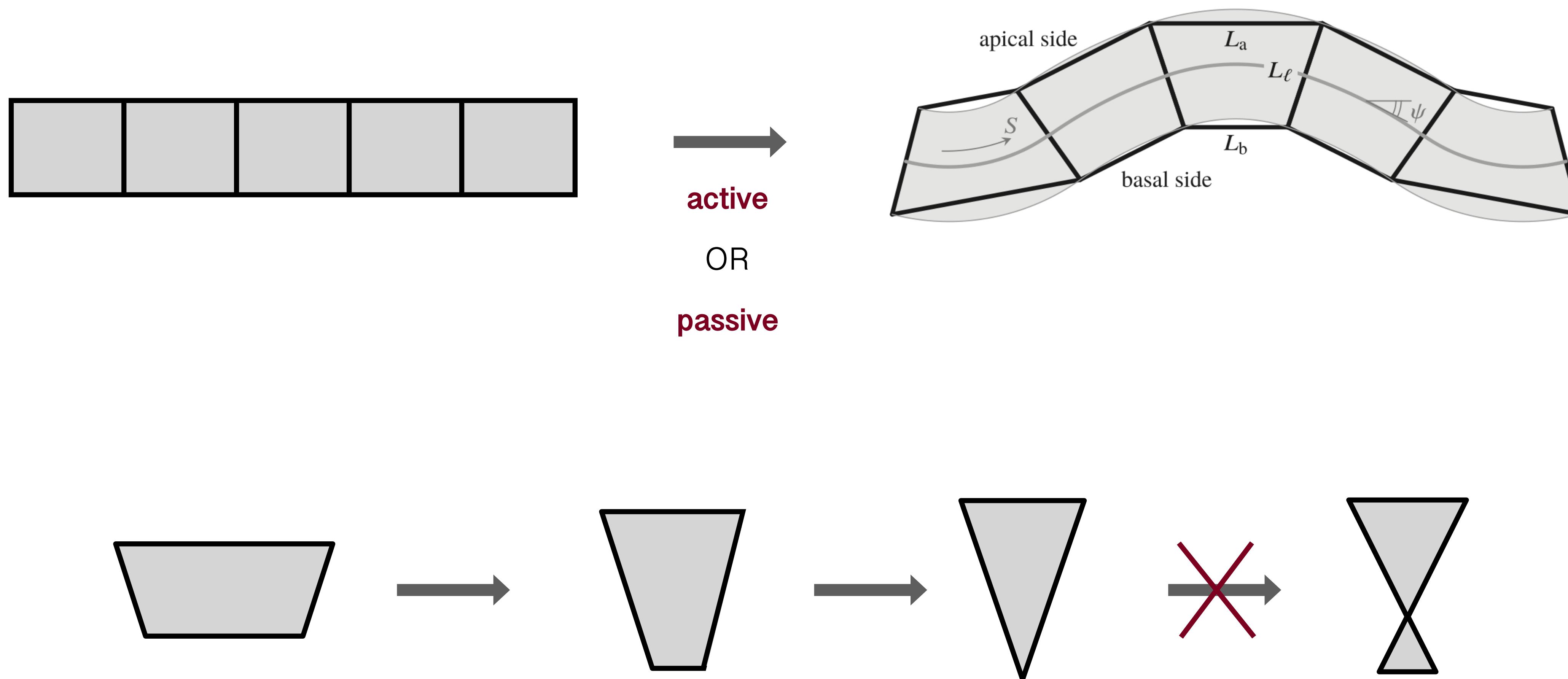
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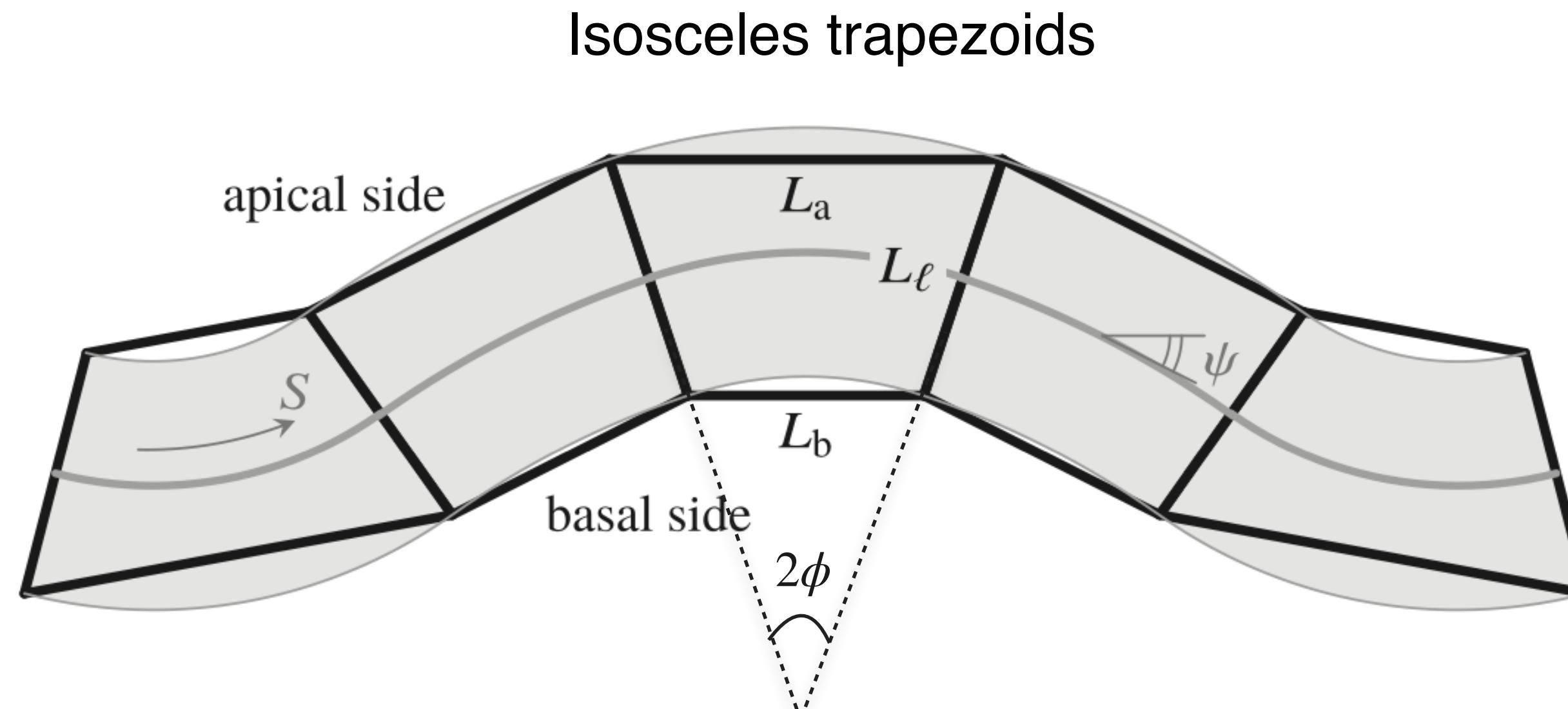
HOW DOES THE DISCRETENESS OF CELLS AFFECT
MECHANICAL INSTABILITIES AT THE TISSUE SCALE?

CELL DISCRETENESS INDUCES MICROSCOPIC CONSTRAINTS



**GOAL: UNDERSTANDING THE TISSUE-MECHANICAL CONSEQUENCES
OF MICROSCOPIC CONSTRAINTS IN A MINIMAL MODEL**

CONTINUUM LIMIT OF A SHEET OF DISCRETE CELLS



Differential surface tensions: $\Lambda_a, \Lambda_b, \Lambda_l$

$$E_{\text{cell}} = \Lambda_a L_a + \Lambda_b L_b + \Lambda_l L_l$$

continuum limit

$$K = (L_a + L_b)/2$$

$$L_a = K + L_\ell \sin \phi,$$

$$L_b = K - L_\ell \sin \phi,$$

$$L = L_\ell \cos \phi.$$

+

Impose incompressibility

$$A_c = KL = \text{constant}$$

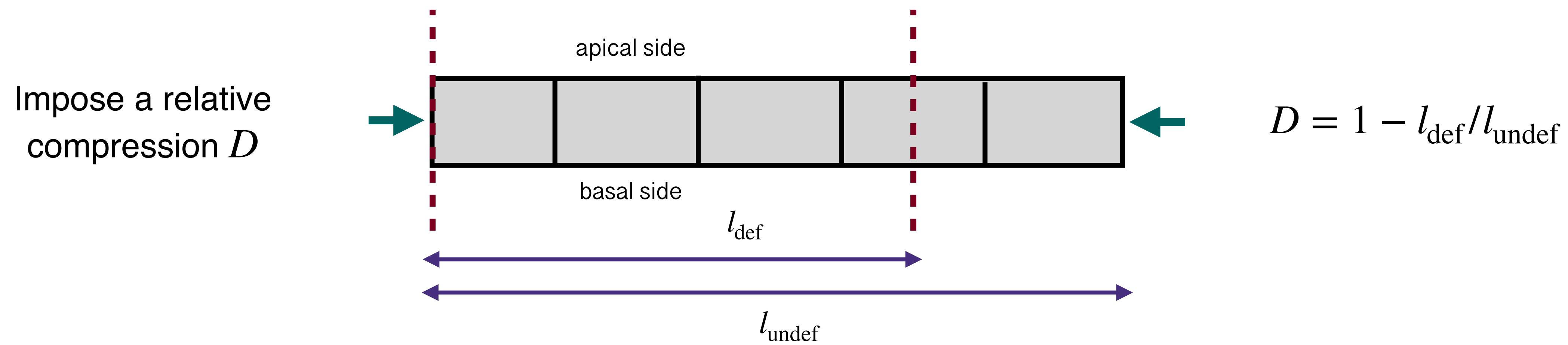
Total energy of monolayer

$$\mathcal{E} = \ell_0^2 \int \left[\frac{\ell_0}{\lambda} \sec \phi + \frac{\lambda}{\ell_0} (1 + \delta \sin \phi) \right] ds,$$

Haas et al., Nonlinear and nonlocal elasticity in coarse-grained differential-tension models of epithelia, PRE (2019)

Energy minimisation problem!

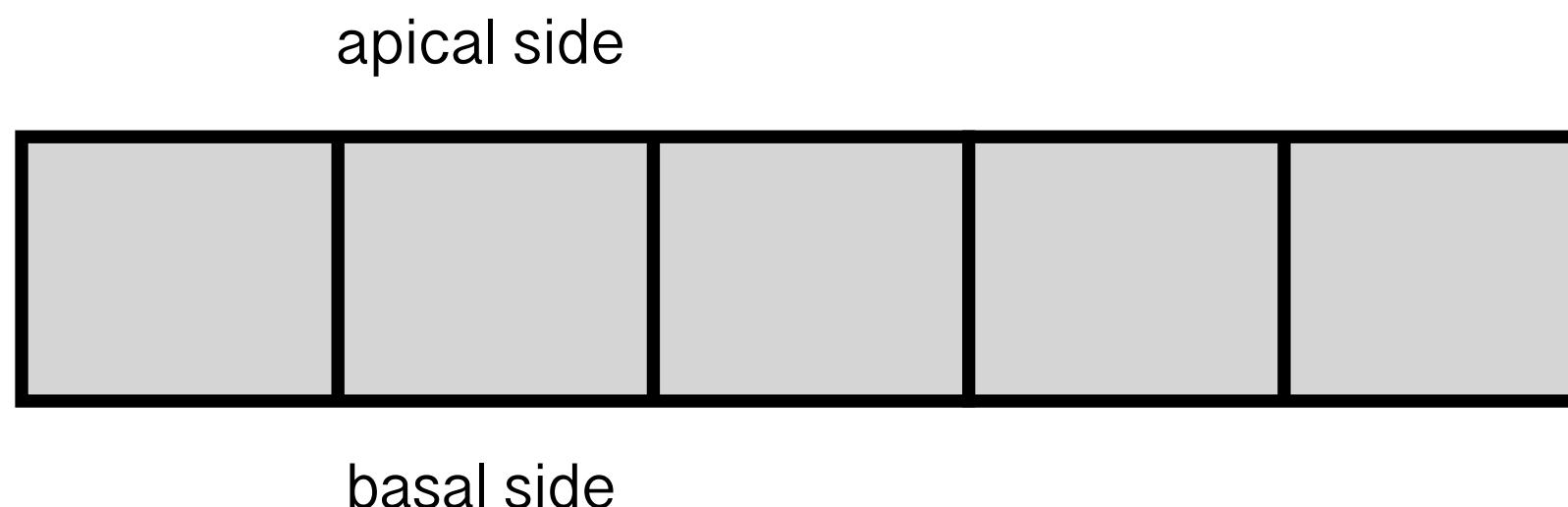
BUCKLING TRANSITION



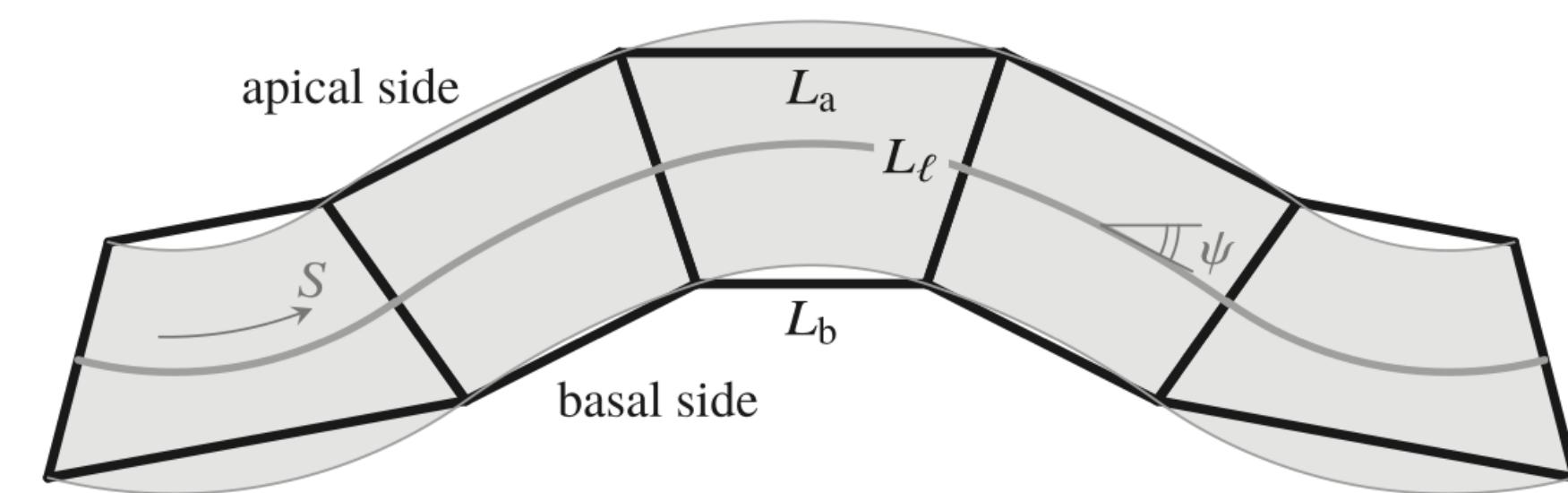
No

Is $E_{\text{buckled}} < E_{\text{flat}}$?

Yes



$D < D_*$: Flat



$D > D_*$: Buckled

VARIATIONAL APPROACH

$$\mathcal{E} = \ell_0^2 \int \left[\frac{\ell_0}{\lambda} \sec \phi + \frac{\lambda}{\ell_0} (1 + \delta \sin \phi) \right] ds.$$

+

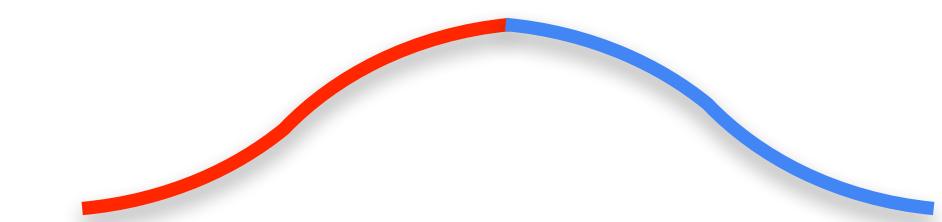
Clamped boundaries:

$$x(2) - x(0) = 2(1 - D)$$

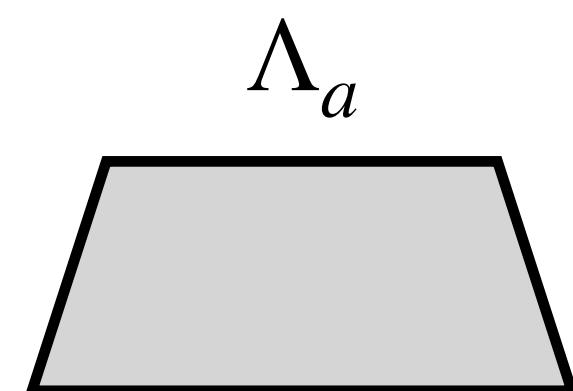
$$y(2) = y(0)$$

$$\psi(0) = \psi(1) = 0$$

Symmetric buckling:



$$\mathcal{L} = \int_0^1 \left[\Lambda + \frac{1}{\Lambda} + \frac{\dot{\psi}^2}{8\Lambda \Xi^2} - \frac{\delta \Lambda \dot{\psi}^3}{48\Xi^3} + \frac{5\dot{\psi}^4}{384\Lambda \Xi^4} + \frac{\ddot{\psi}^2}{48\Lambda \Xi^4} + \frac{\mu \cos \psi}{\Lambda} \left(1 + \frac{\dot{\psi}^2}{8\Xi^2} + \frac{41\dot{\psi}^4}{1920\Xi^4} + \frac{\ddot{\psi}^2}{40\Xi^4} + \frac{\dot{\psi} \ddot{\psi}}{240\Xi^4} \right) + O(\Xi^{-5}) \right] d\sigma.$$

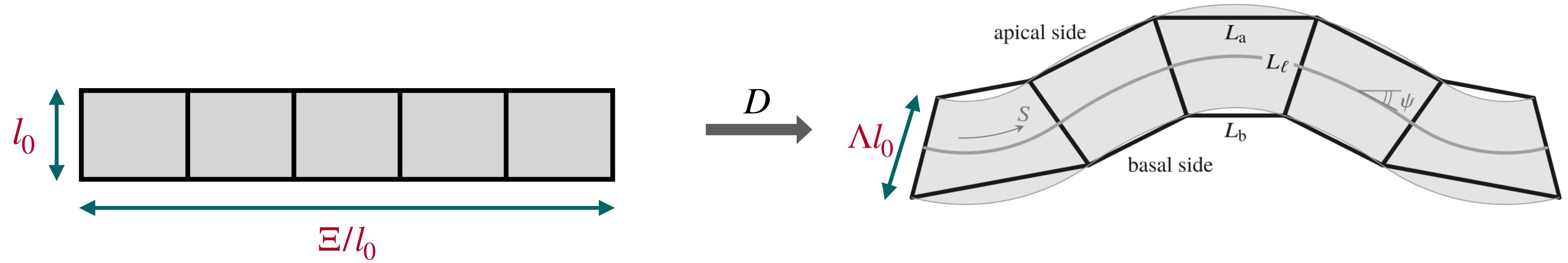


$$\delta \sim |\Lambda_a - \Lambda_b|$$

Horizontal compressive force

$$\delta \mathcal{L} = \delta \mu \frac{d\mathcal{L}}{d\mu} + \delta \Lambda \frac{d\mathcal{L}}{d\Lambda} + \delta \psi \frac{d\mathcal{L}}{d\psi} + \delta \dot{\psi} \frac{d\mathcal{L}}{d\dot{\psi}} = 0$$

GOVERNING EQUATIONS



Exact shape is determined for any D

$$\frac{d\psi}{d\sigma} = \dot{\psi} \quad \frac{d\dot{\psi}}{d\sigma} = \frac{4\mu\Xi^2 \sin\psi}{1 + \mu \cos\psi} \left(\frac{\dot{\psi}^2}{8\Xi^2} - 1 \right) \quad \frac{dI}{d\sigma} = \cos\psi \left(1 + \frac{\dot{\psi}^2}{8\Xi^2} \right) \quad \frac{dJ}{d\sigma} = \frac{\dot{\psi}^2}{8\Xi^2} \quad \frac{d\Lambda}{d\sigma} = 0 \quad \frac{d\mu}{d\sigma} = 0$$

$$\psi(0) = 0$$

$$\psi(1) = 0$$

$$I(0) = 0$$

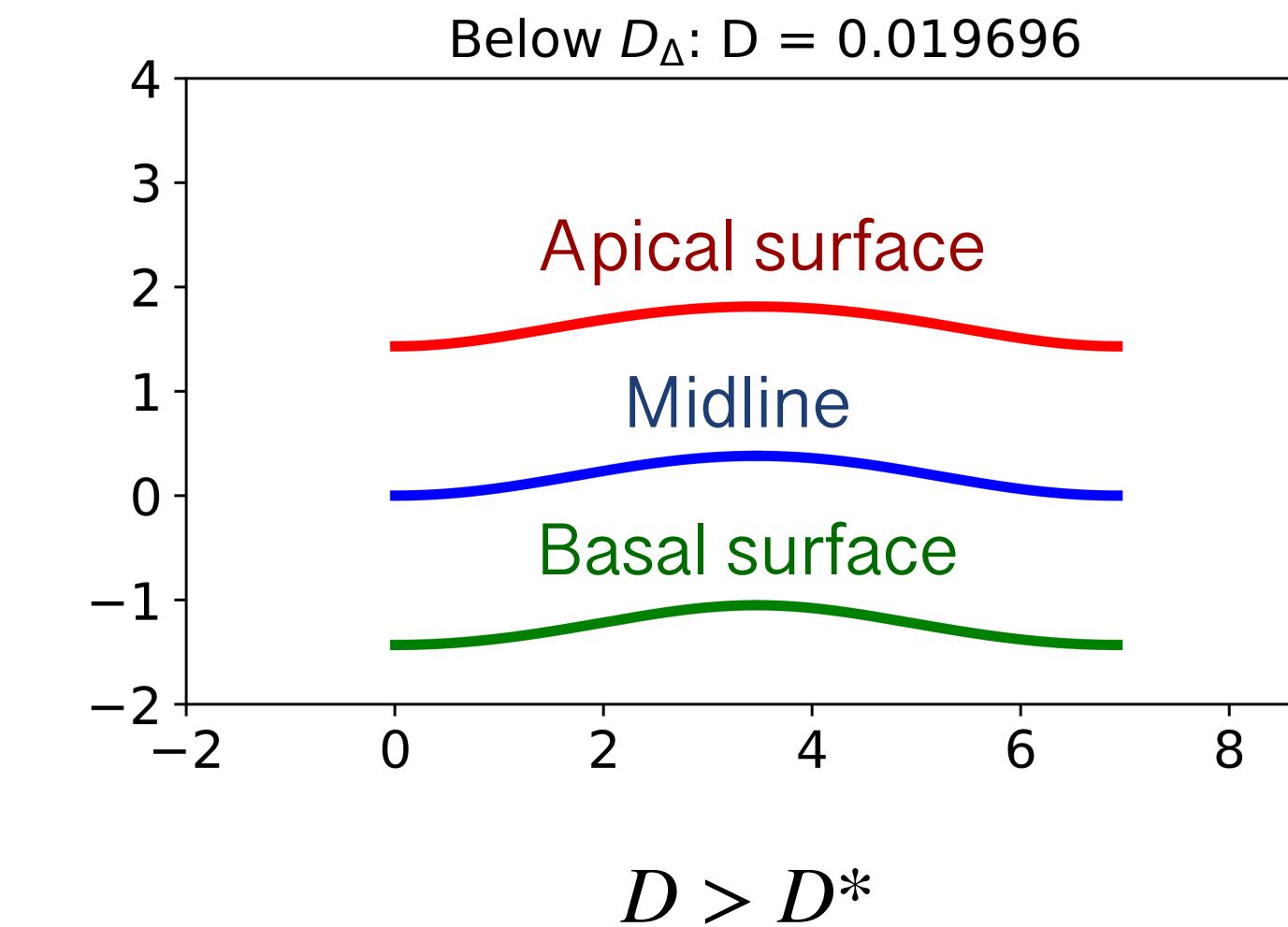
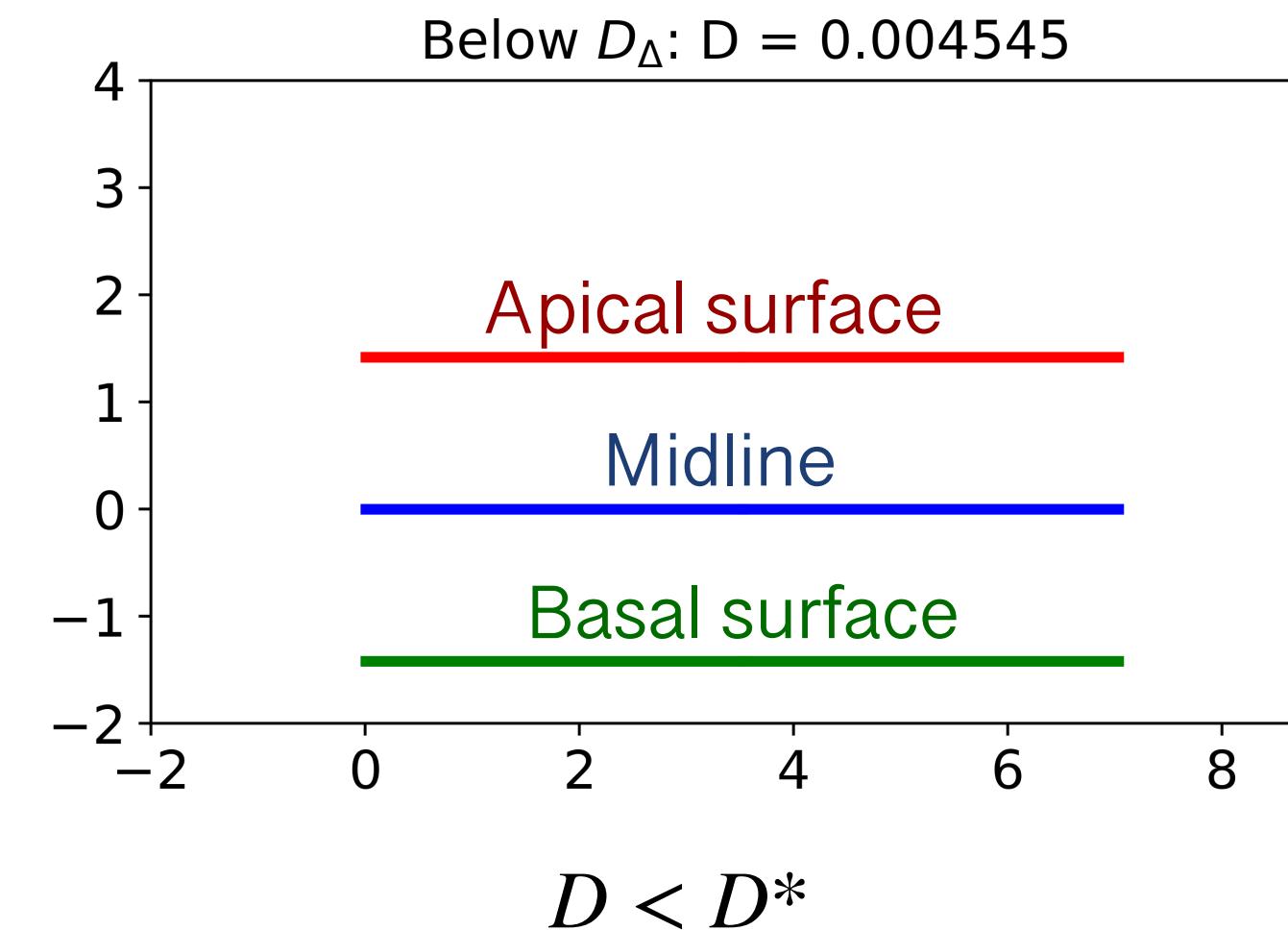
$$I(1) = \Lambda(1 - D)$$

$$J(0) = 0$$

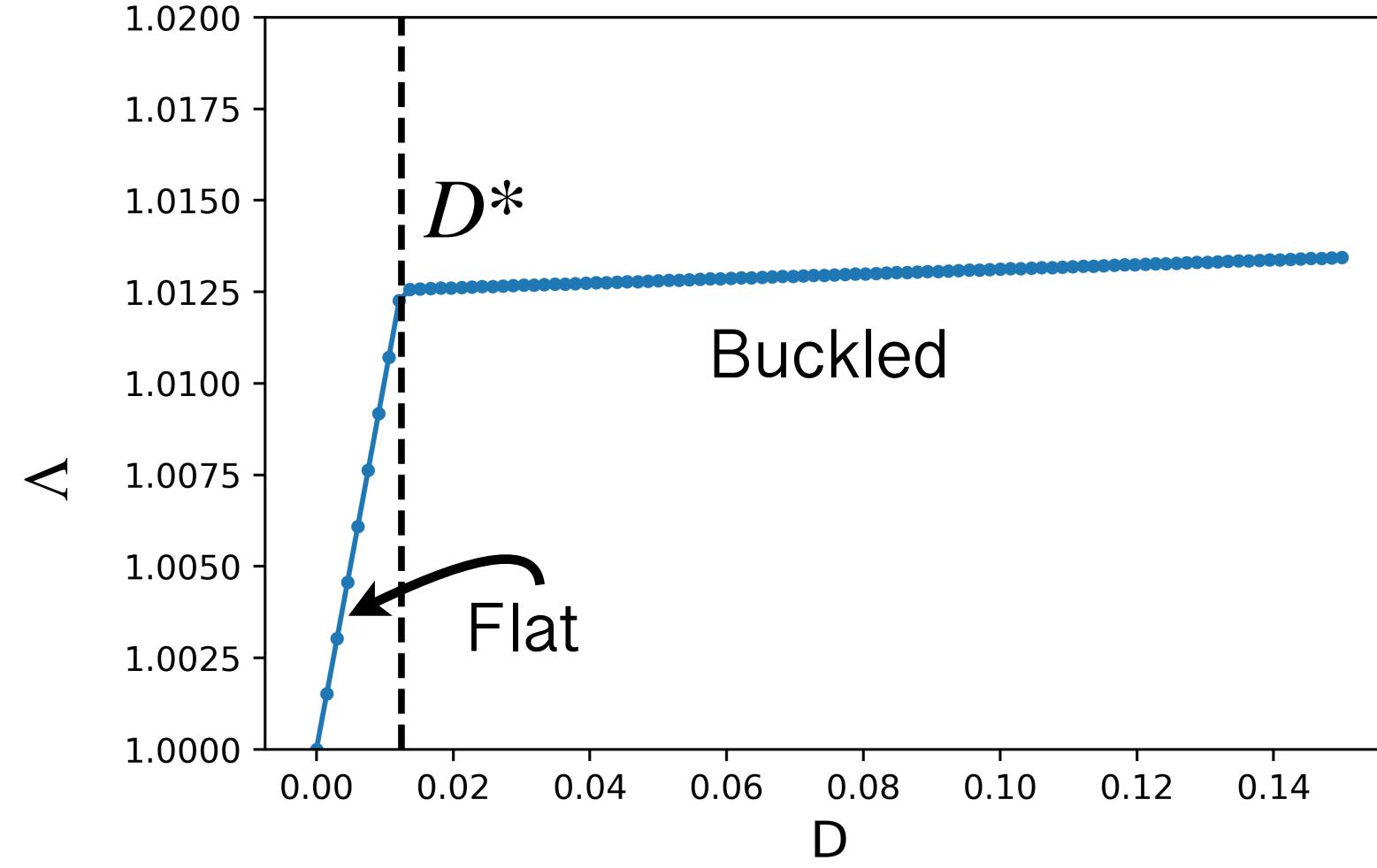
$$J(1) = \Lambda^2 - 1 - \mu\Lambda(1 - D)$$

Boundary value problem!

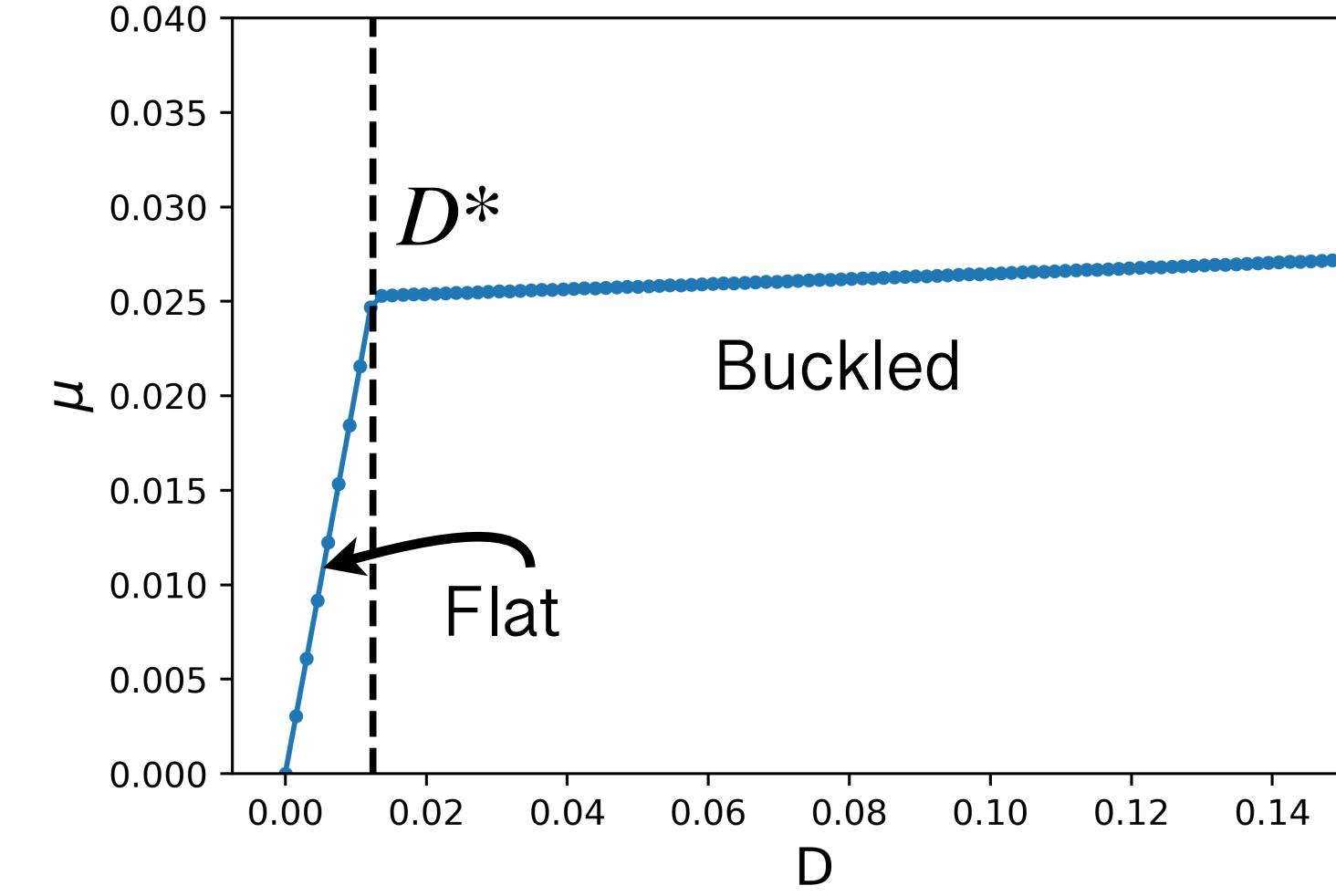
AT BUCKLING TRANSITION: RELEASE OF STIFFNESS



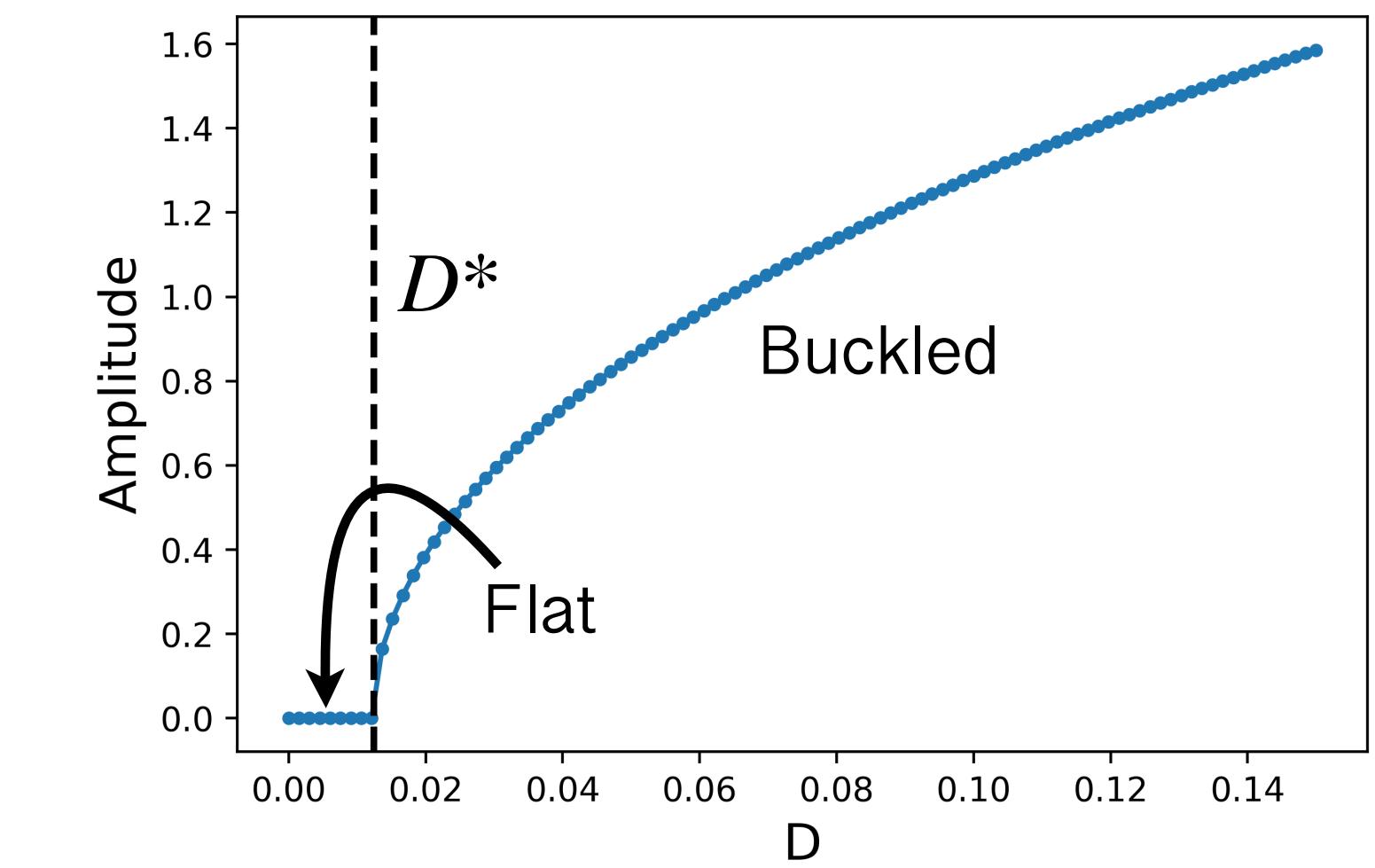
Λ = Deformed relative thickness



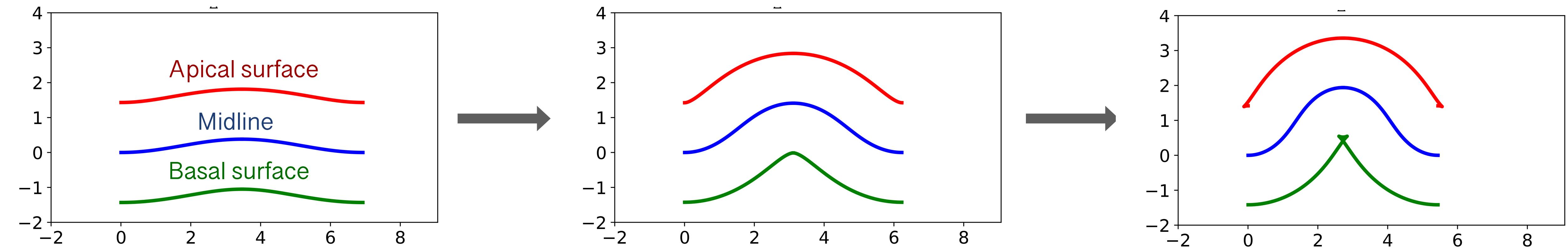
μ = Compressive force



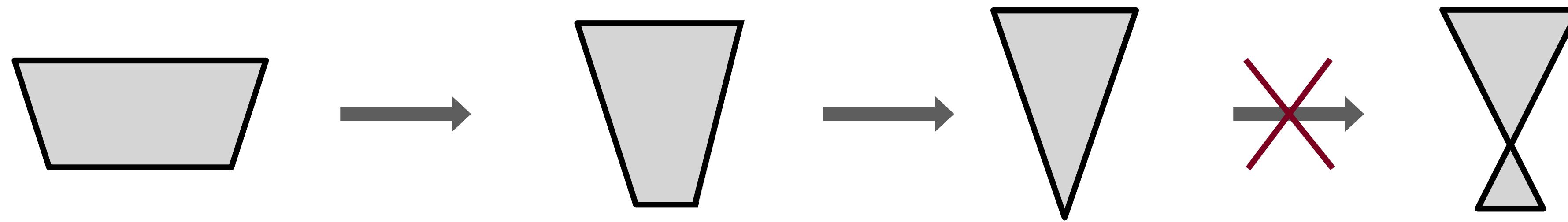
Amplitude = Maximum vertical displacement of midline



WHAT HAPPENS UNDER FURTHER COMPRESSION?

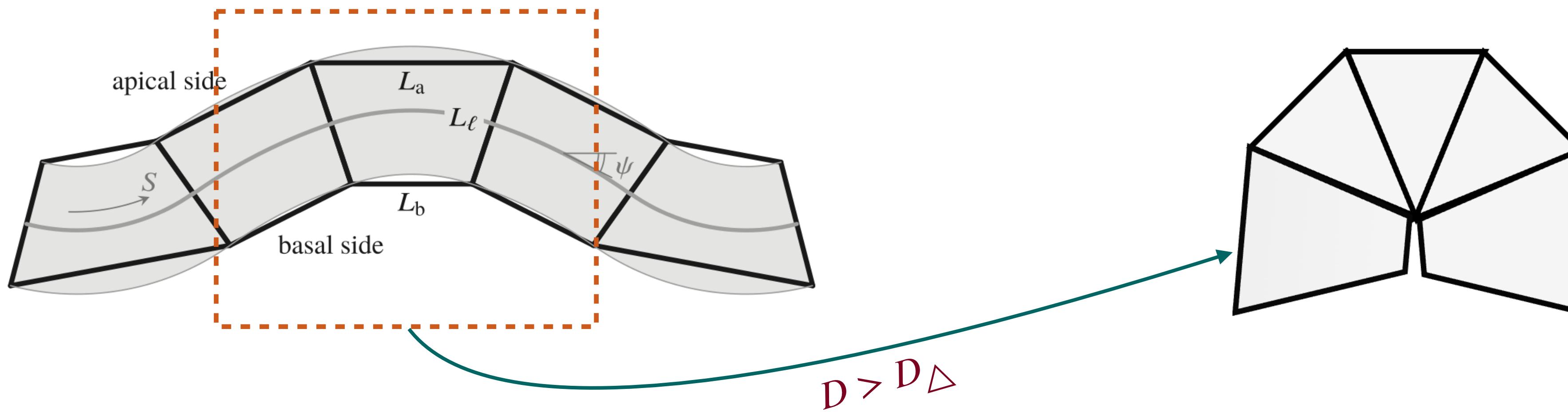


Model breaks down!



Cells self-intersect

NEW GOVERNING EQUATIONS



Introduce $s_* \sim \text{Fraction of } \triangle \text{ cells}$

$\sigma \in [0, s_*]$ \rightarrow \triangle cells that cannot fold further

$\sigma \in \left[s_*, \frac{1}{2}\right]$ \rightarrow Cells that can still fold

$$\sigma = s_* + (1 - 2s_*)\sigma'$$

$$\frac{d\psi}{d\sigma'} = \dot{\psi}$$

$$\frac{dJ}{d\sigma'} = \frac{\dot{\psi}^2}{8\Xi^2}(1 - 2s_*)$$

$$\frac{d\mu}{d\sigma'} = 0$$

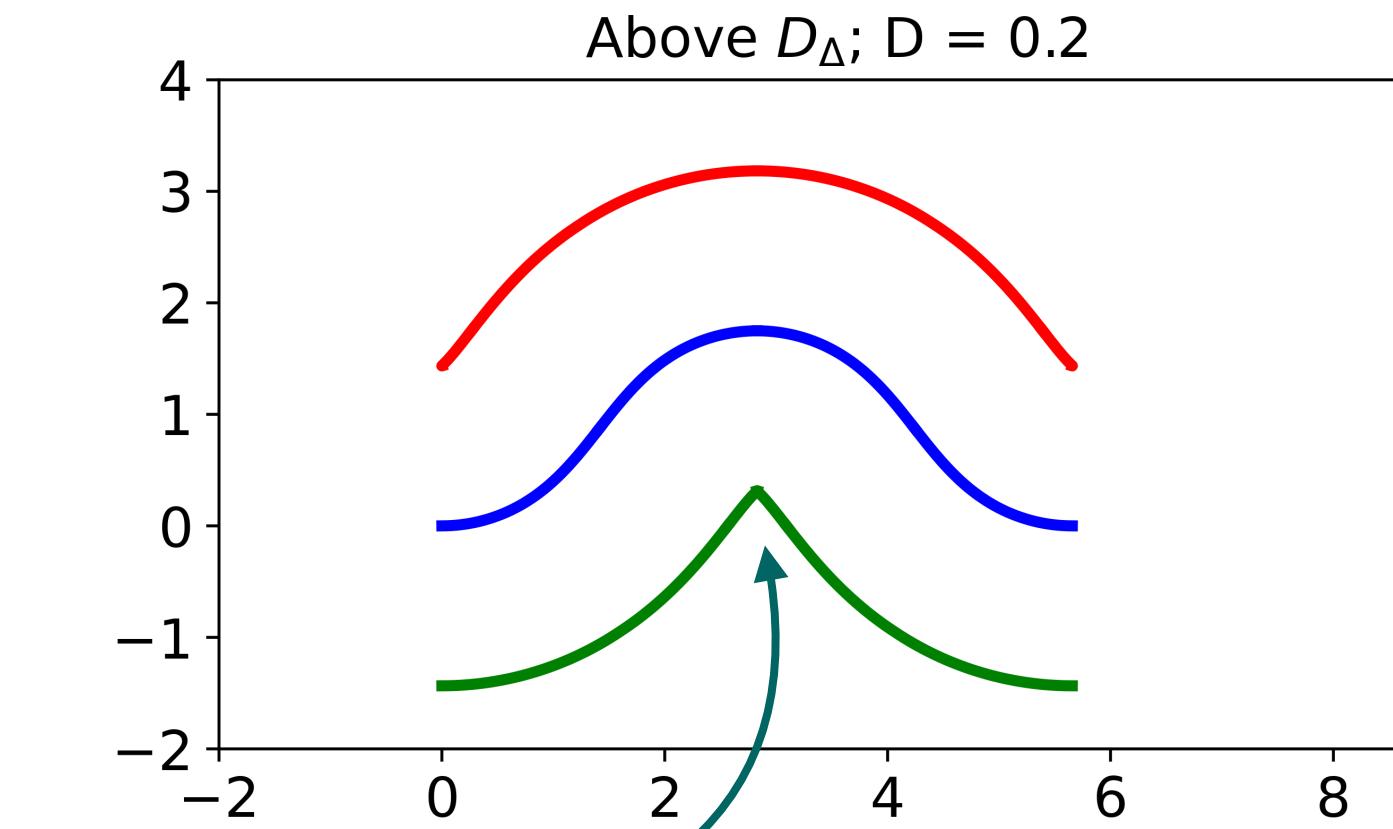
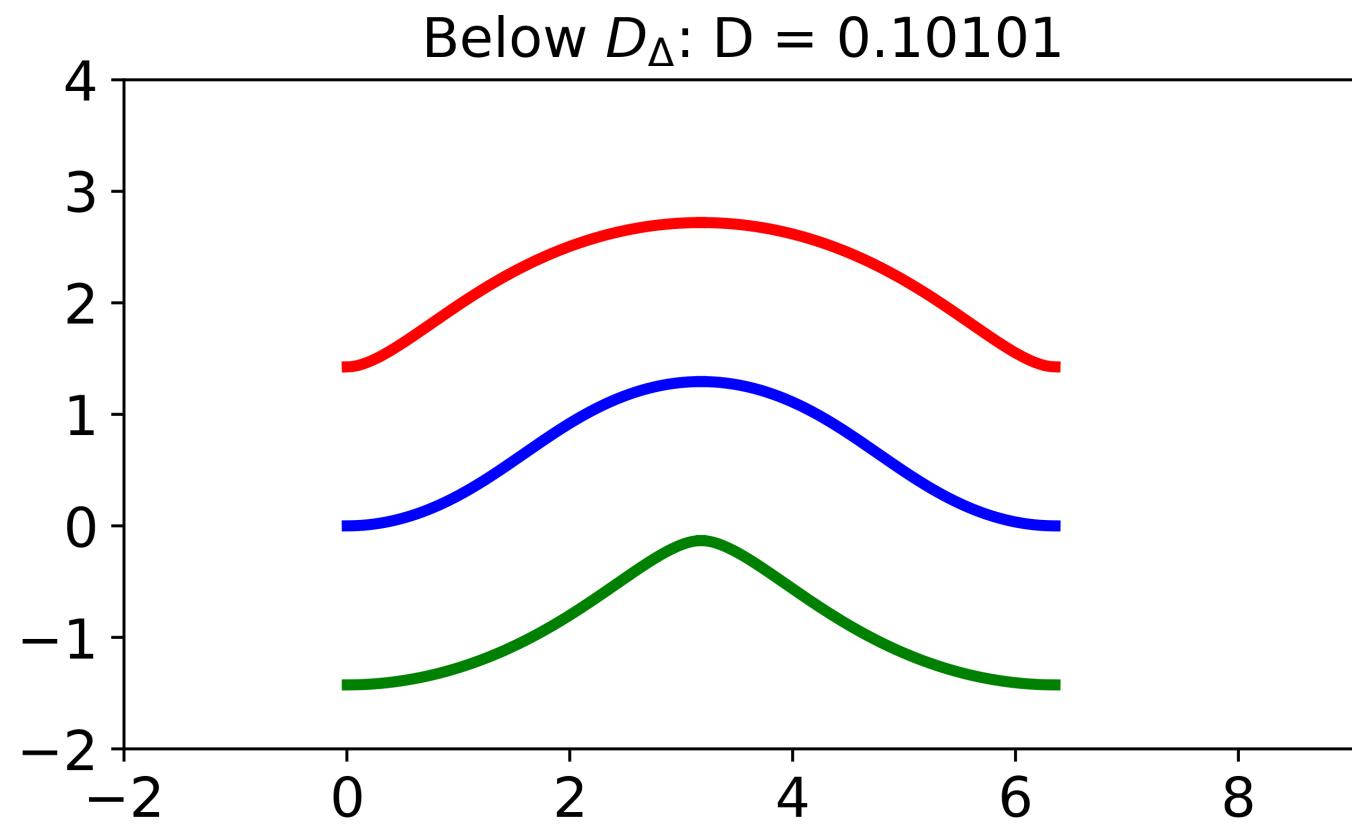
$$\frac{d\dot{\psi}}{d\sigma'} = \frac{4\mu\Xi^2 \sin\psi}{1 + \mu \cos\psi} \left(\frac{\dot{\psi}^2}{8\Xi^2(1 - 2s_*)^2} - 1 \right)(1 - 2s_*)$$

$$\frac{d\Lambda}{d\sigma'} = 0$$

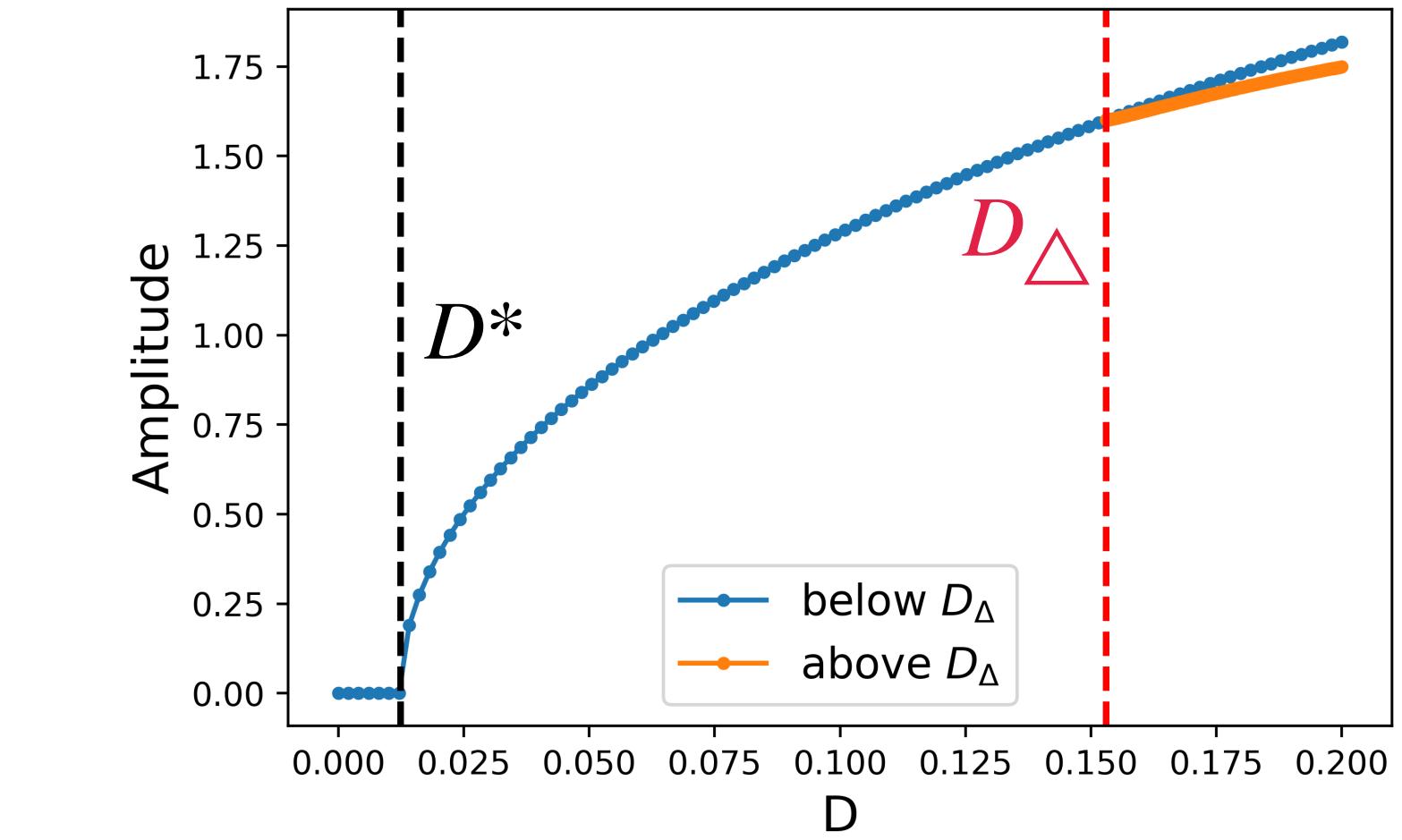
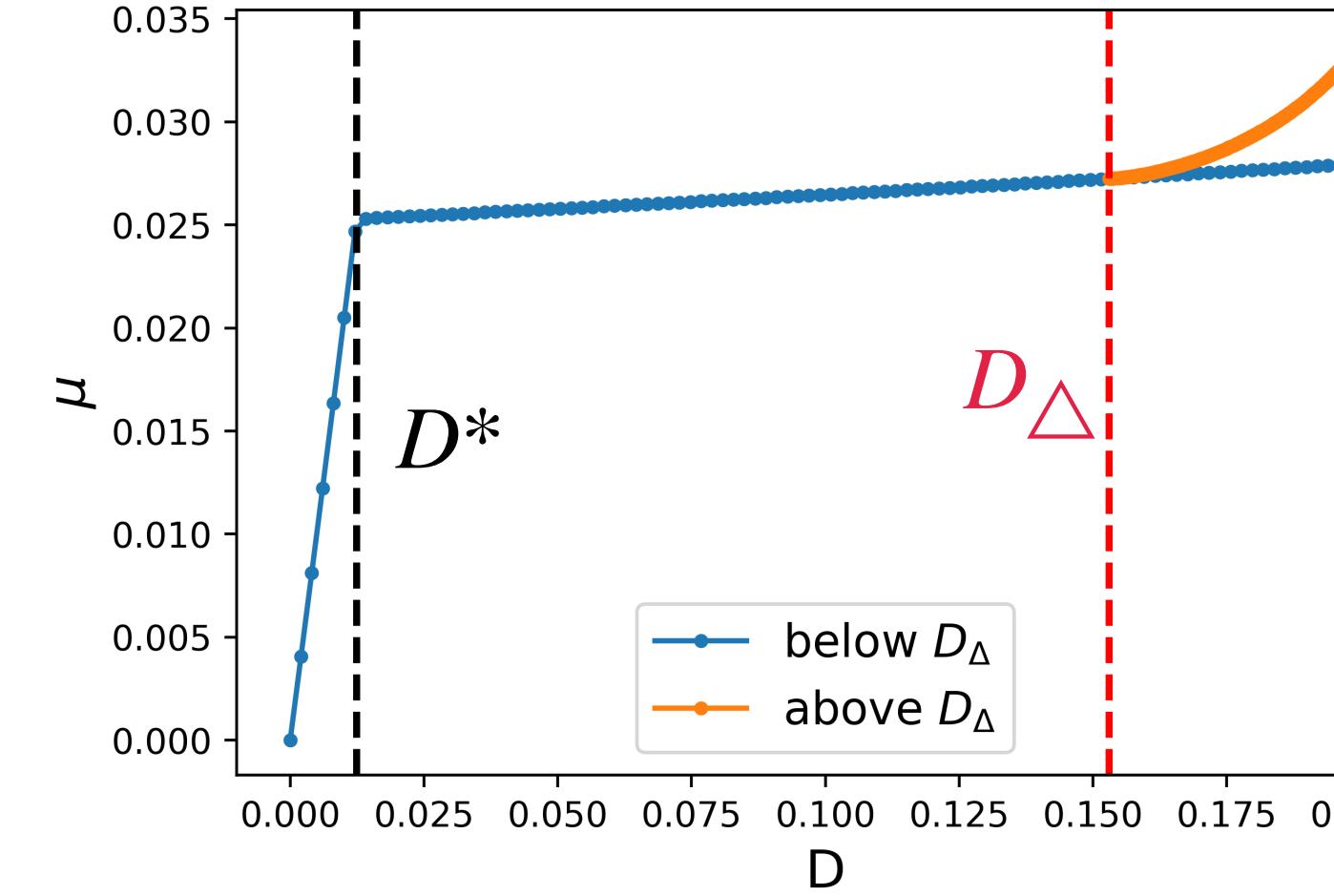
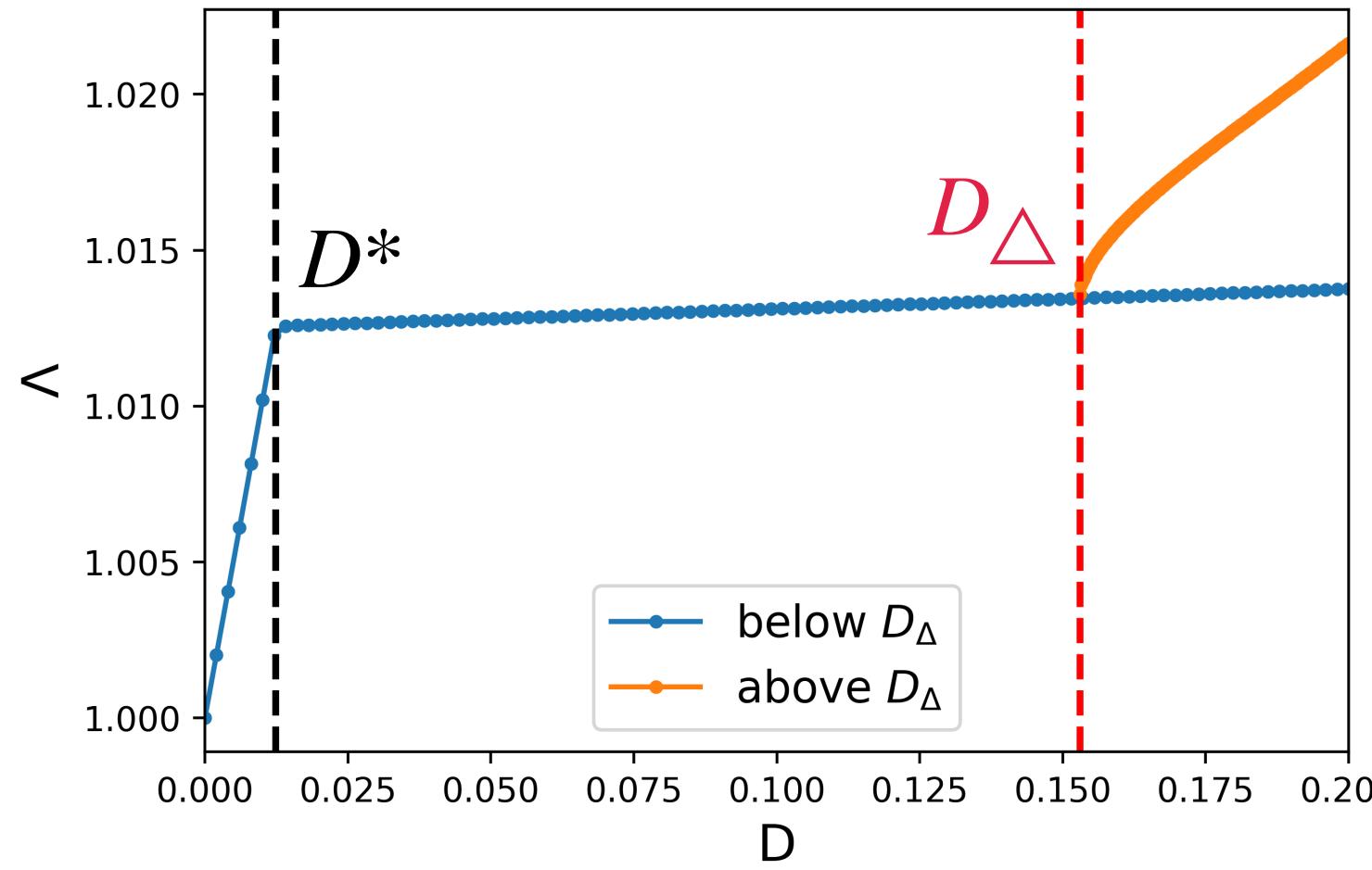
$$\frac{dI}{d\sigma'} = \cos\psi \left(1 + \frac{\dot{\psi}^2}{8\Xi^2(1 - 2s_*)^2} \right)$$

$$\frac{ds_*}{d\sigma'} = 0$$

△ CELLS STIFFEN THE MONOLAYER

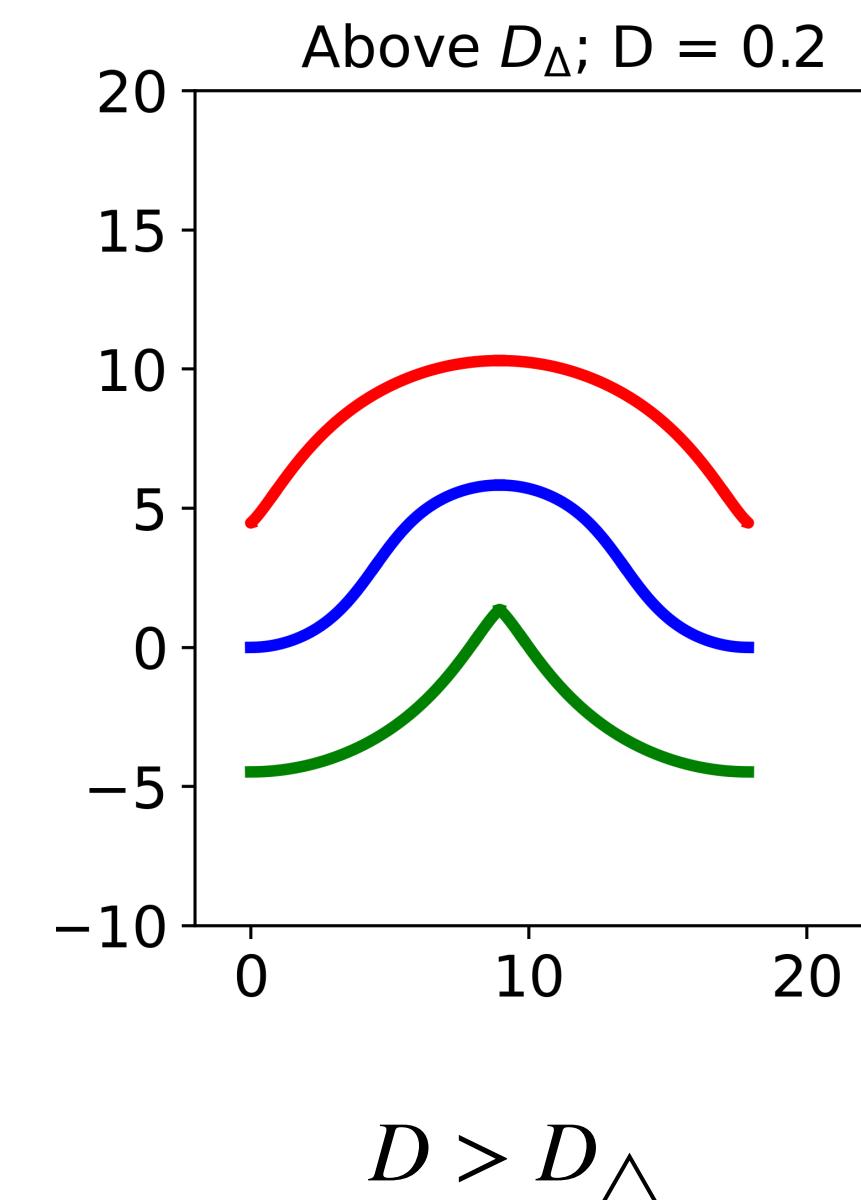
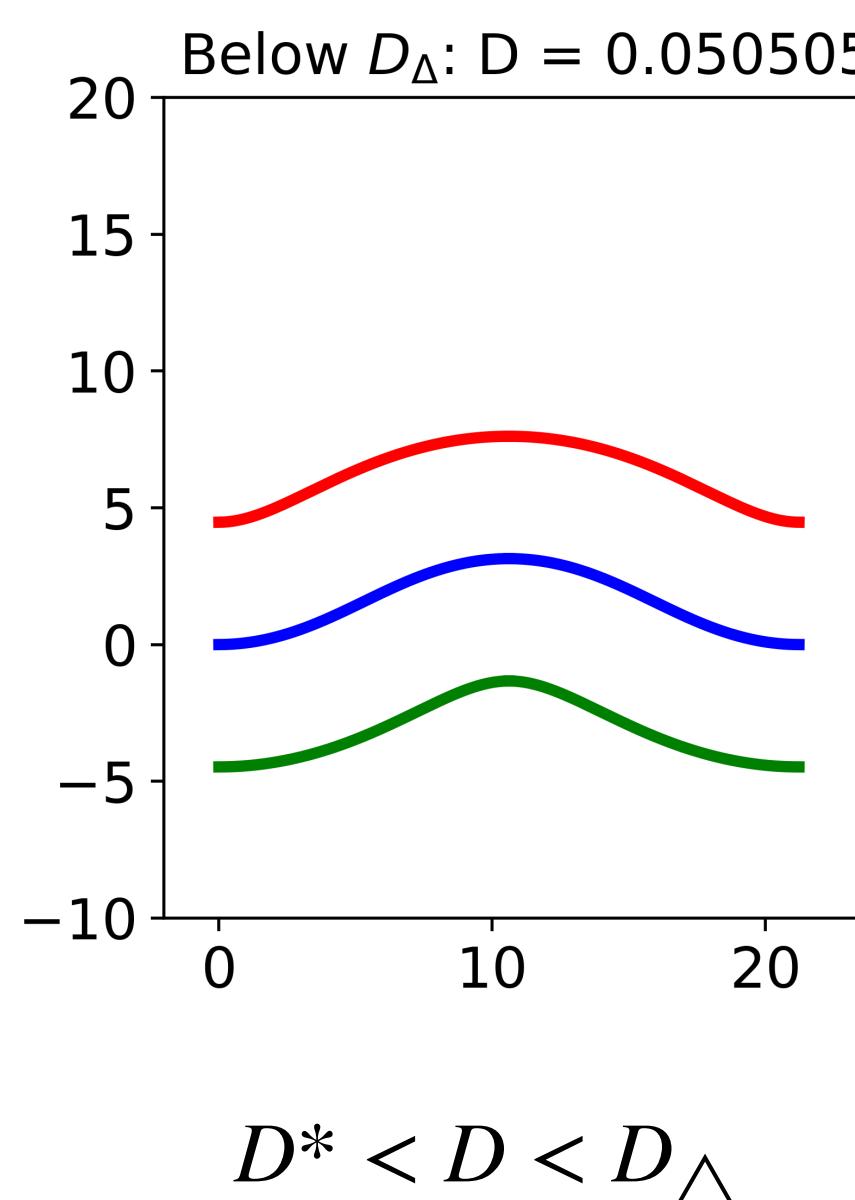
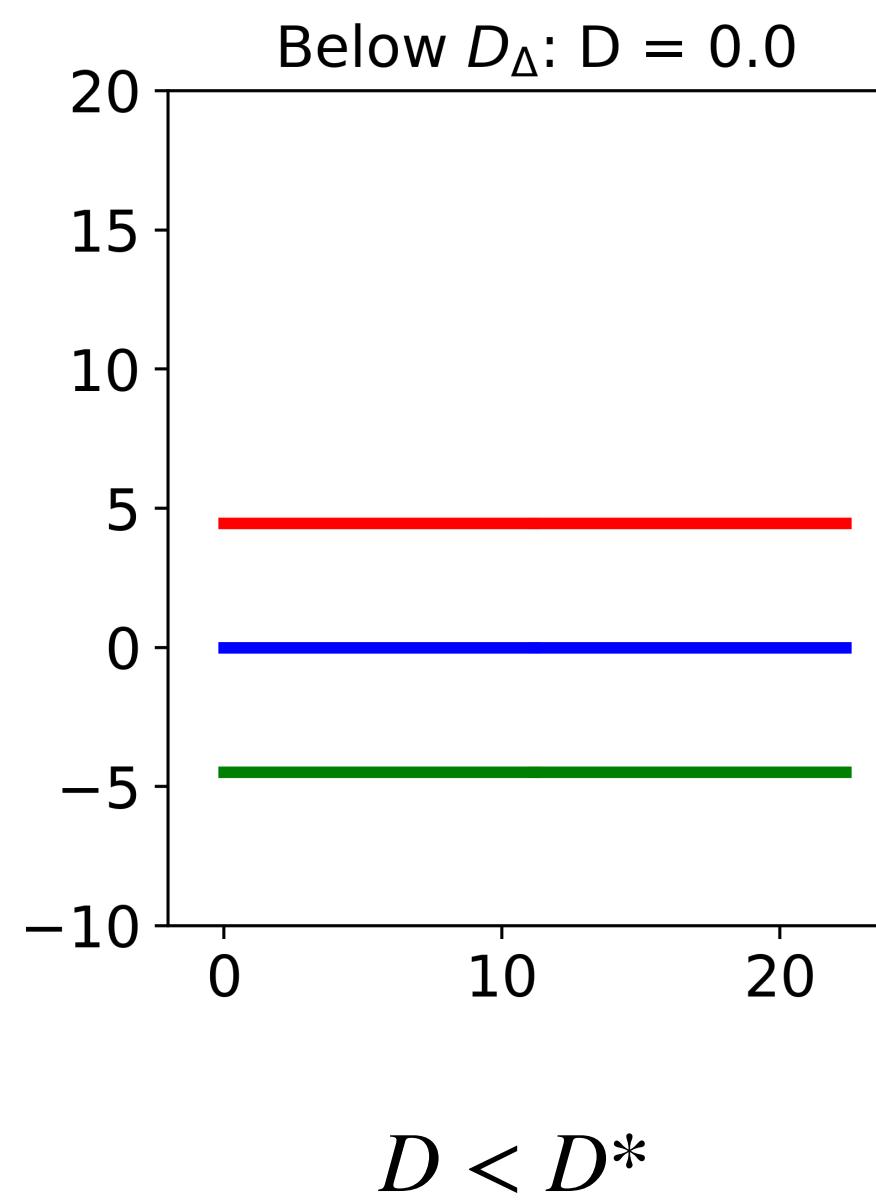


Kink due to \triangle cells

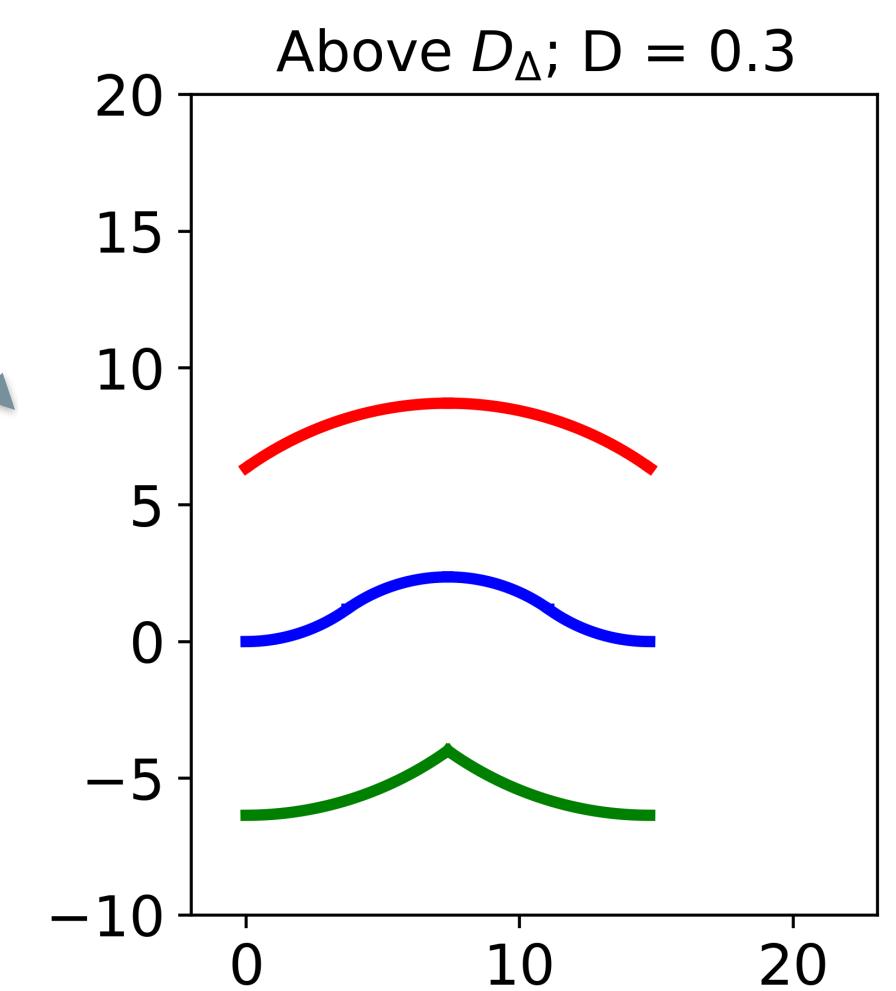
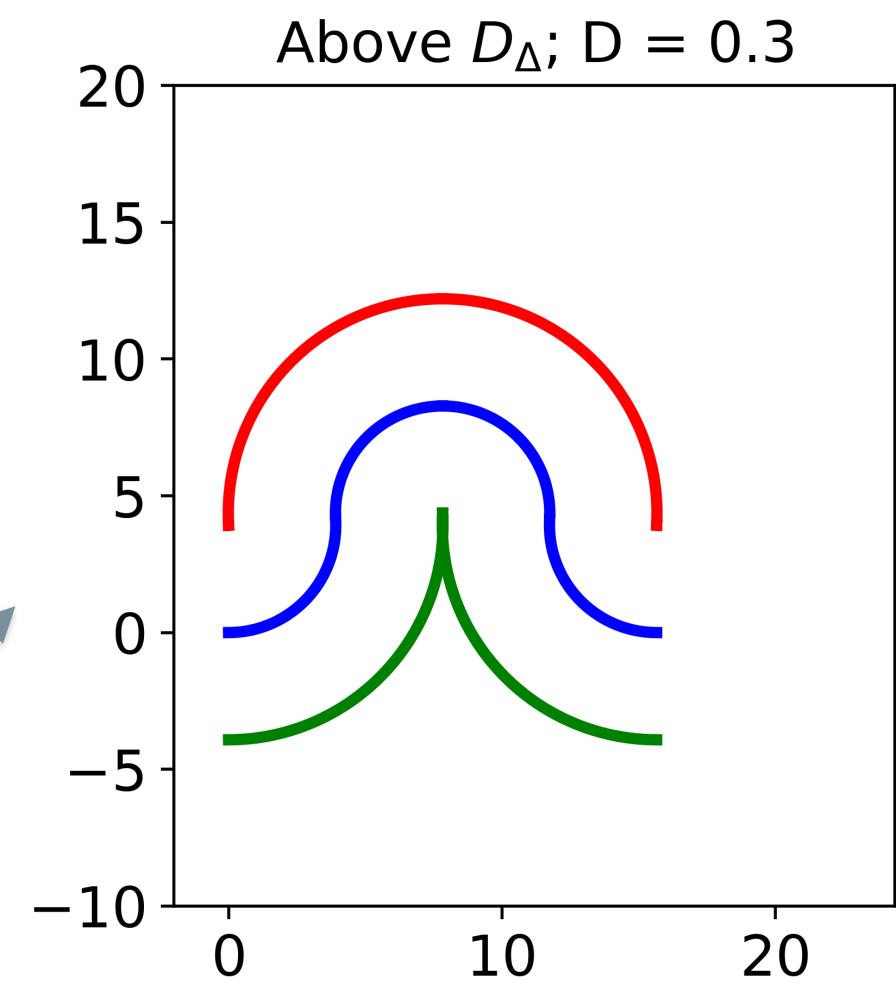


Monolayer stiffens once again!

COMPRESSING EVEN FURTHER ...

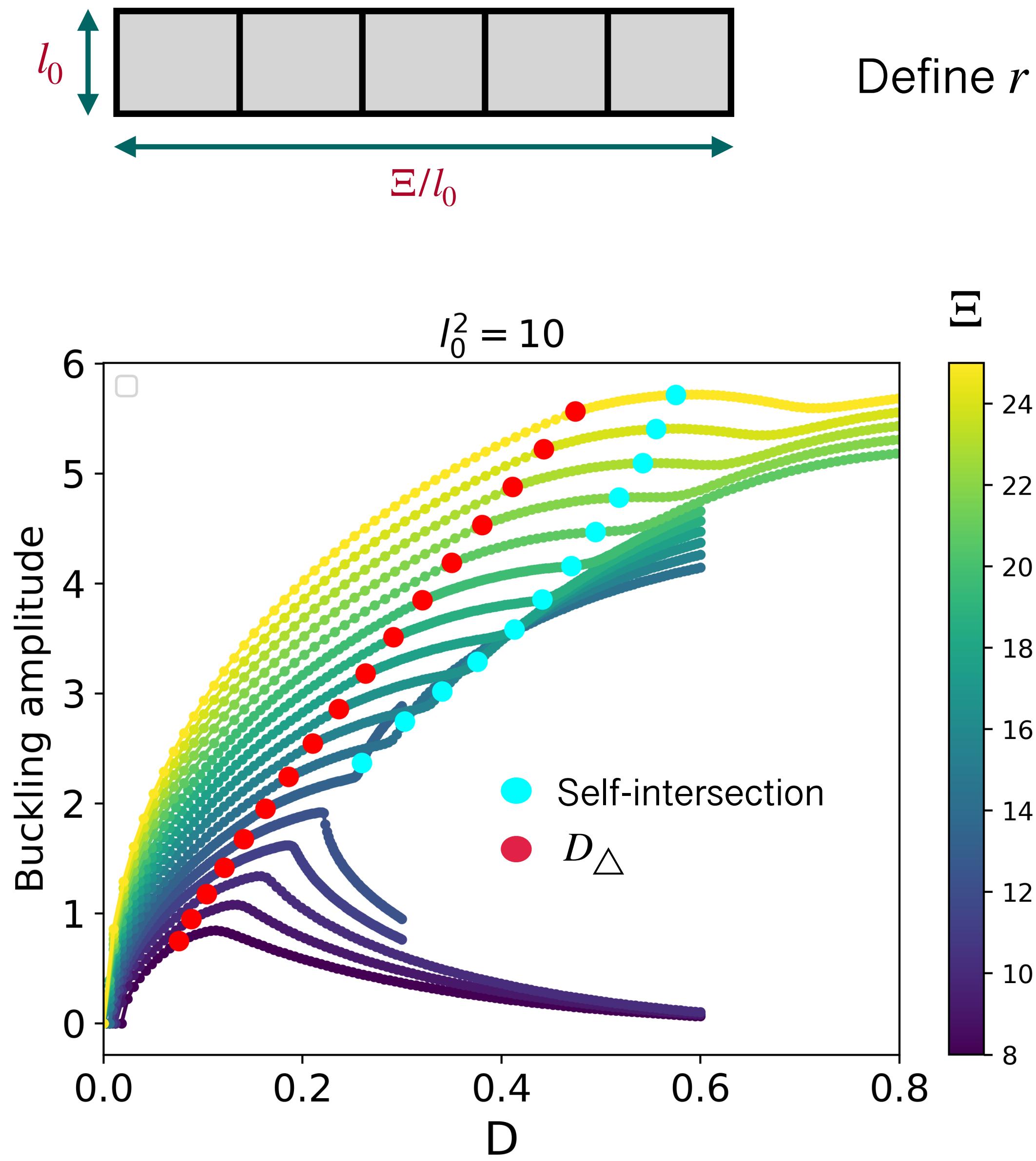


Thin monolayer
Thick monolayer

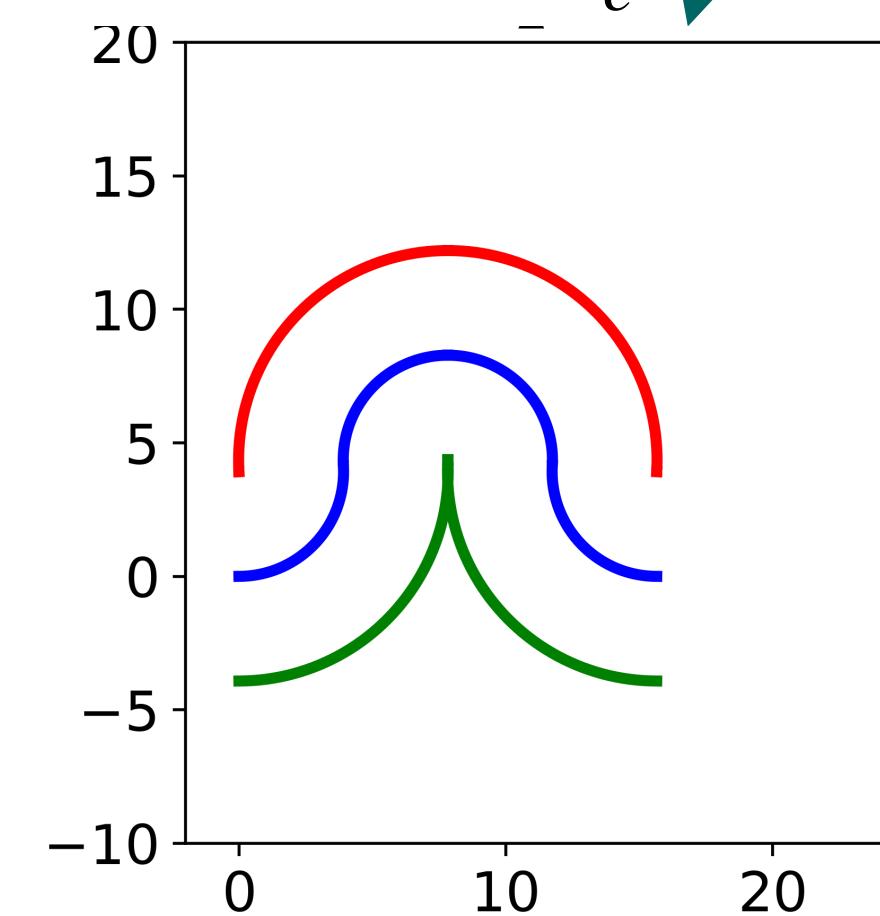
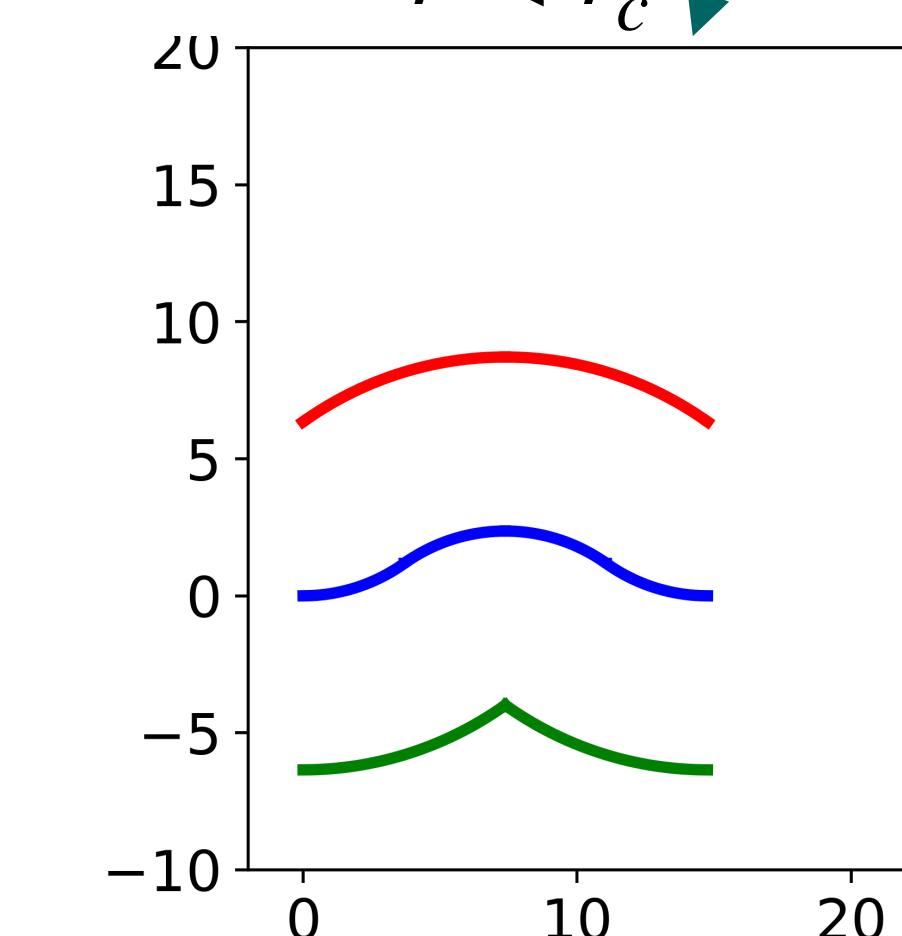
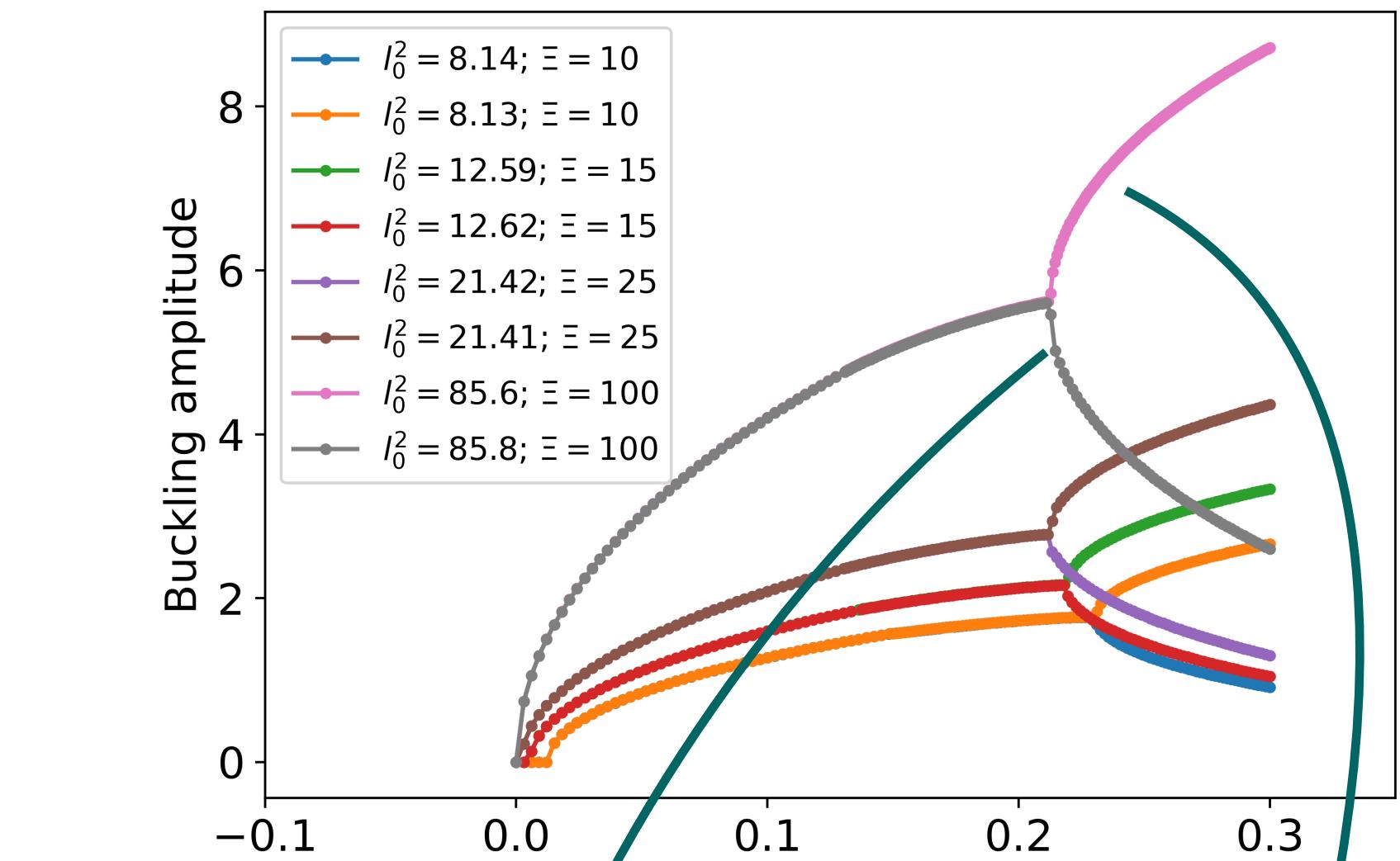


Two different behaviours emerge!

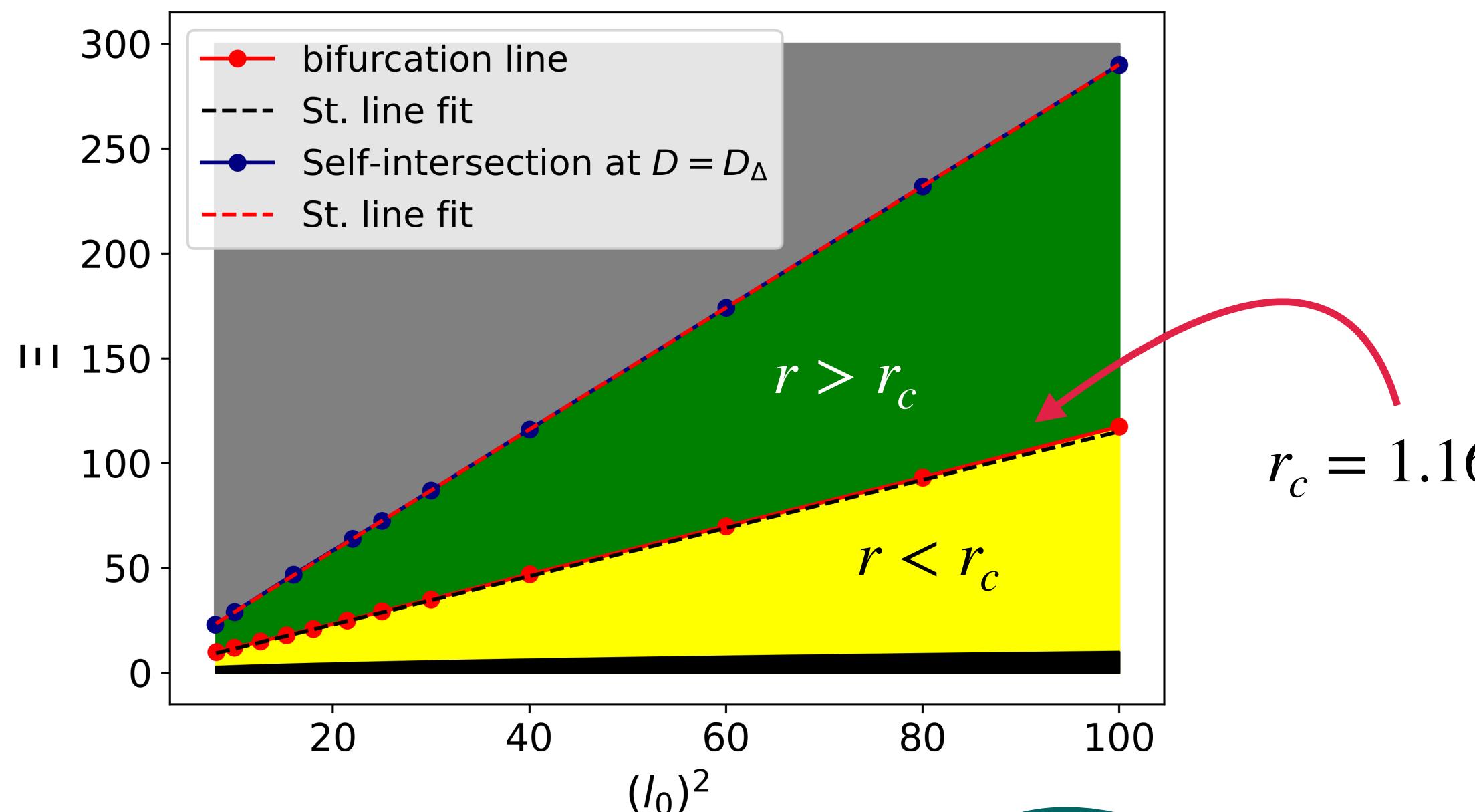
PITCHFORK BIFURCATION



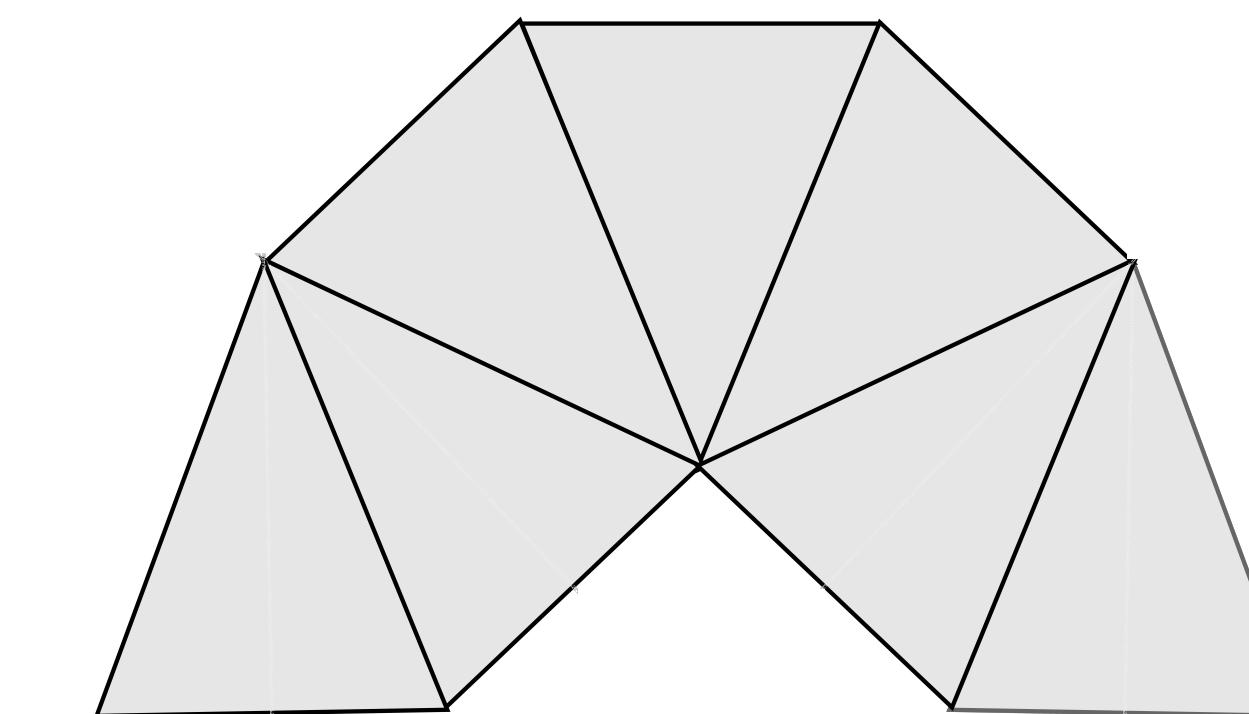
Define $r = \Xi/l_0^2$



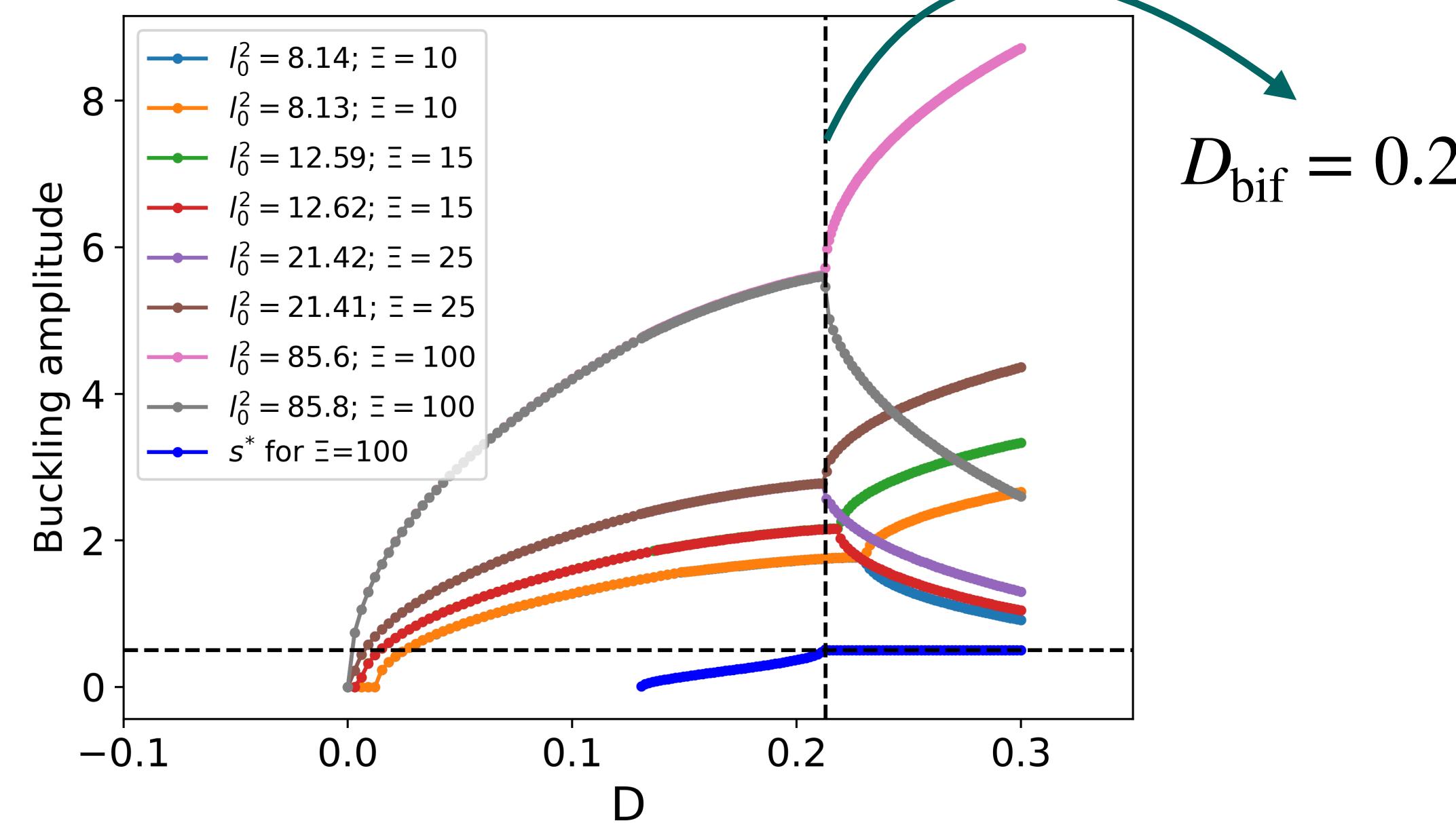
NUMERICS MATCH PRELIMINARY ANALYTICS (?)



Bifurcation due to geometric singularity!



All cells are \triangle ($s_* = 1/2$)



Preliminary Analytics

$$\tan r_c = 2r_c \quad ? \quad D_{\text{bif}} = 1 - \frac{\sin r_c}{r_c}$$

$$r_c \approx 1.165$$

$$D_{\text{bif}} \approx 0.211$$

FUTURE DIRECTIONS

Project 1

- Bifurcation analysis
- Consequences on tissue mechanics
(eg., monolayer stiffening)

Project 2 : More realistic model

Relaxation of assumptions like
“isosceles trapezoids”

—> δ in lowest order —> “wrinkles”

Project 3 : Introduce fluctuations

Fluctuations \Leftrightarrow Biology

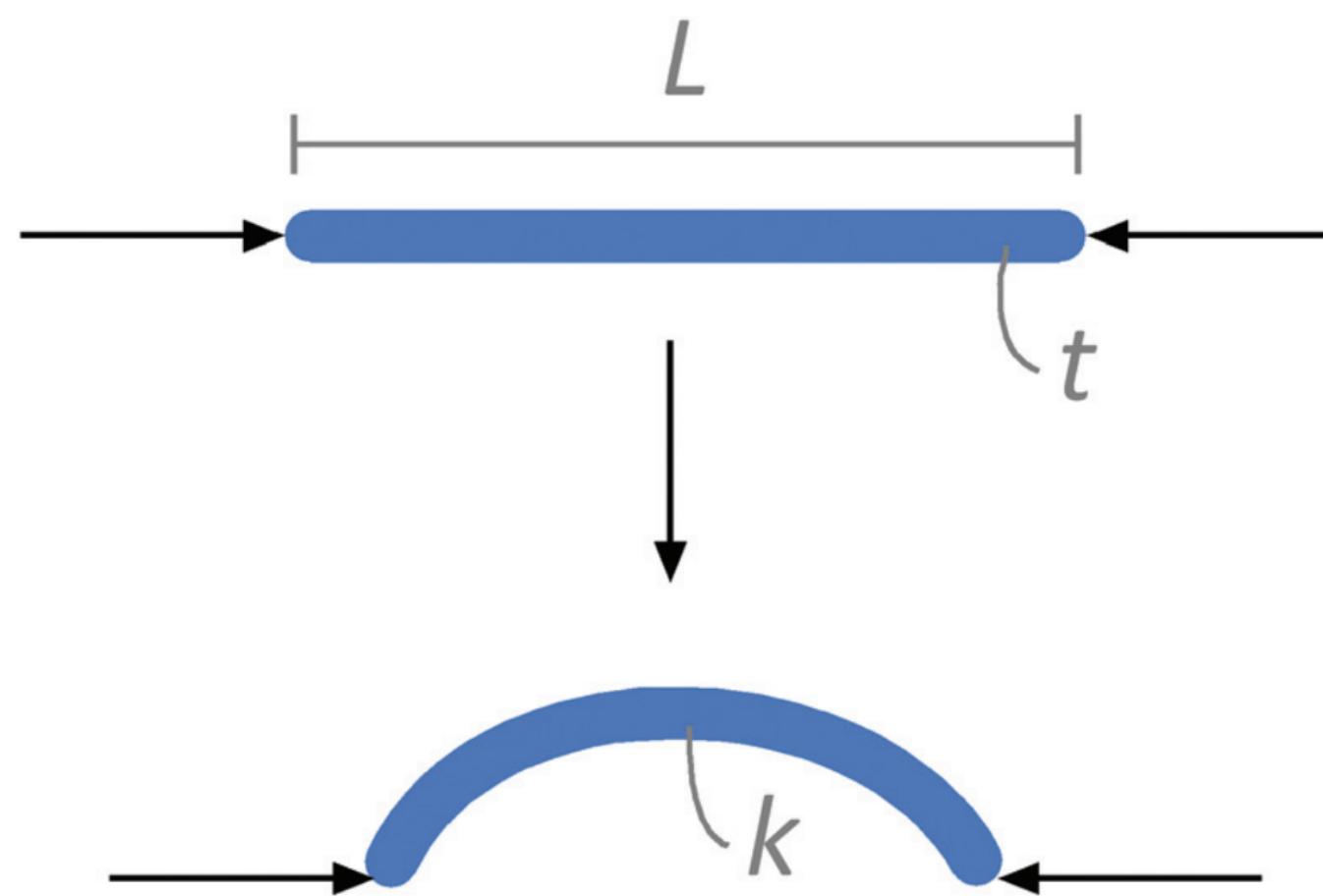
Static noise, i.e., variability

THANK YOU!

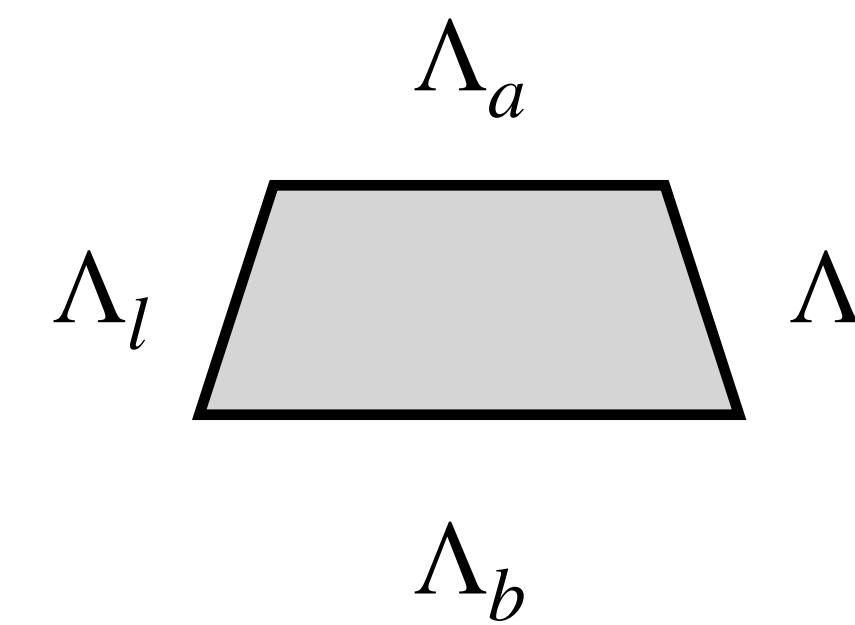
MECHANICAL PROPAGATIO OF GEOMETRIC

Tend to be **thin**

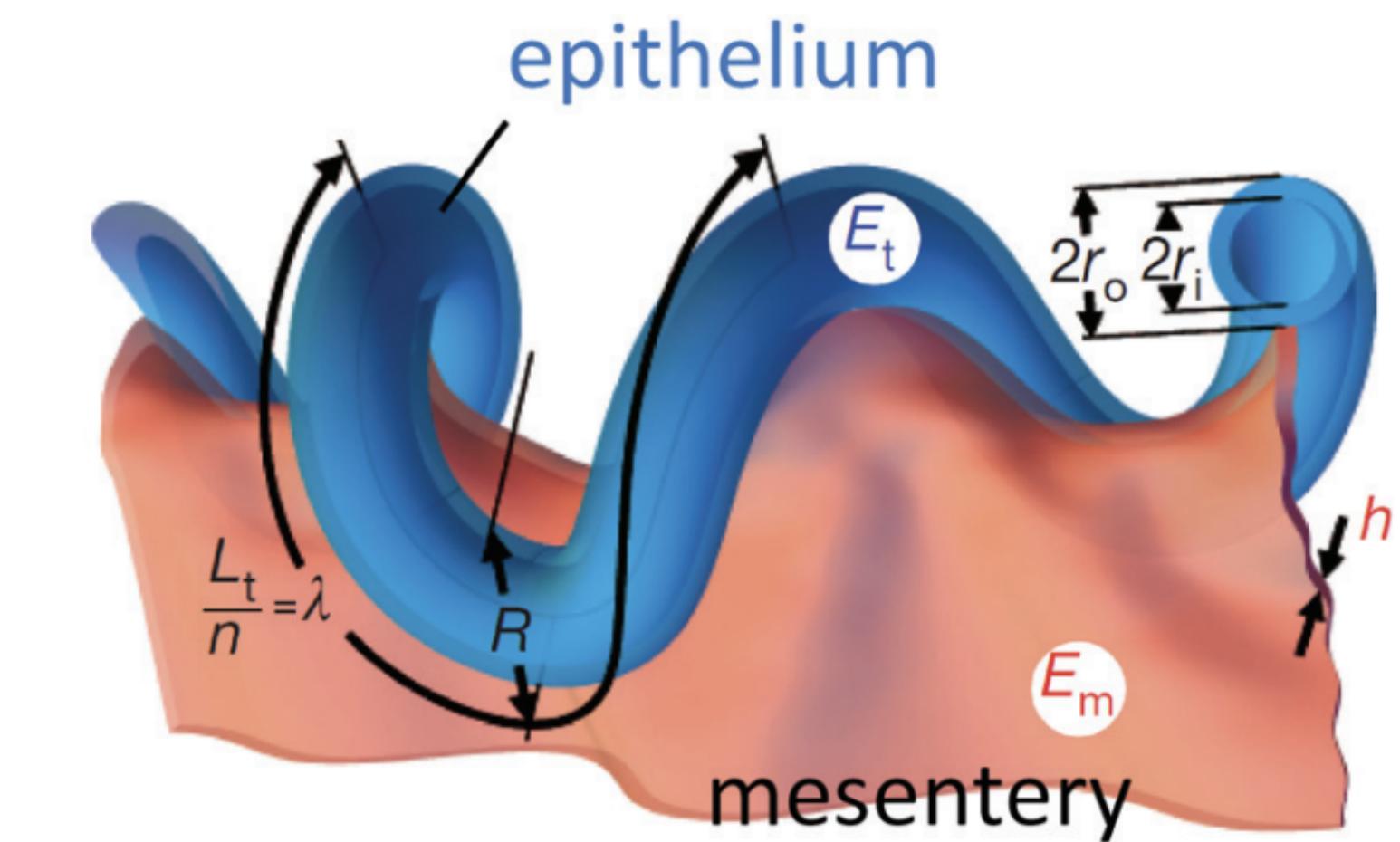
$L \gg t$: more susceptible to instabilities



Cell polarity: Spatial difference leads to differential material properties



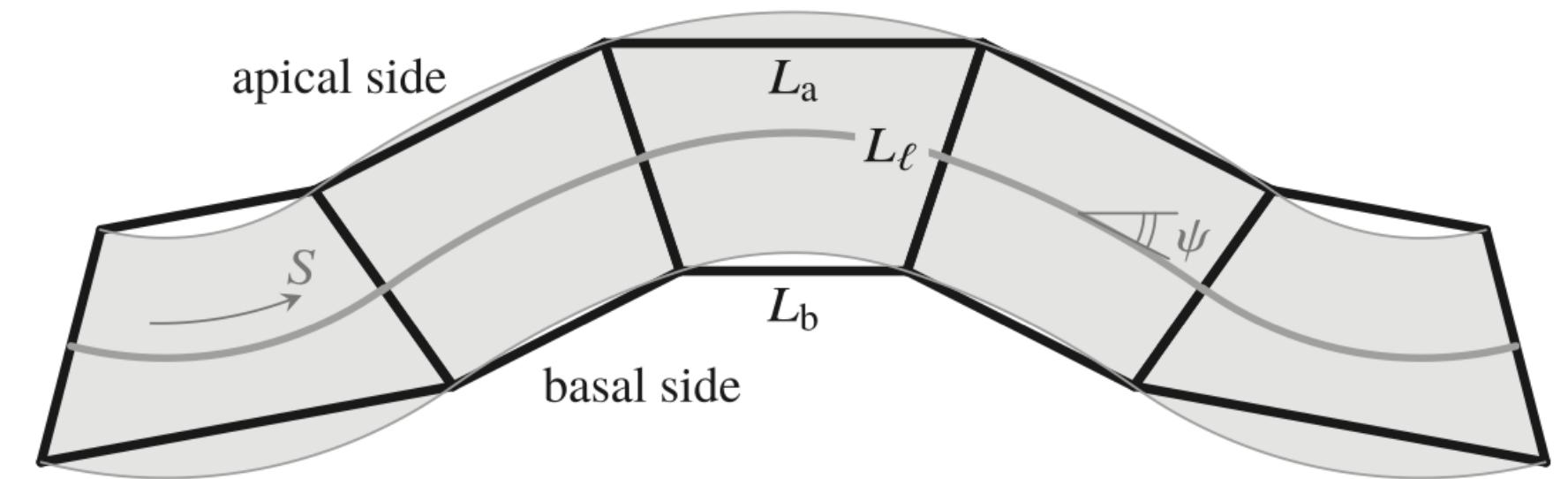
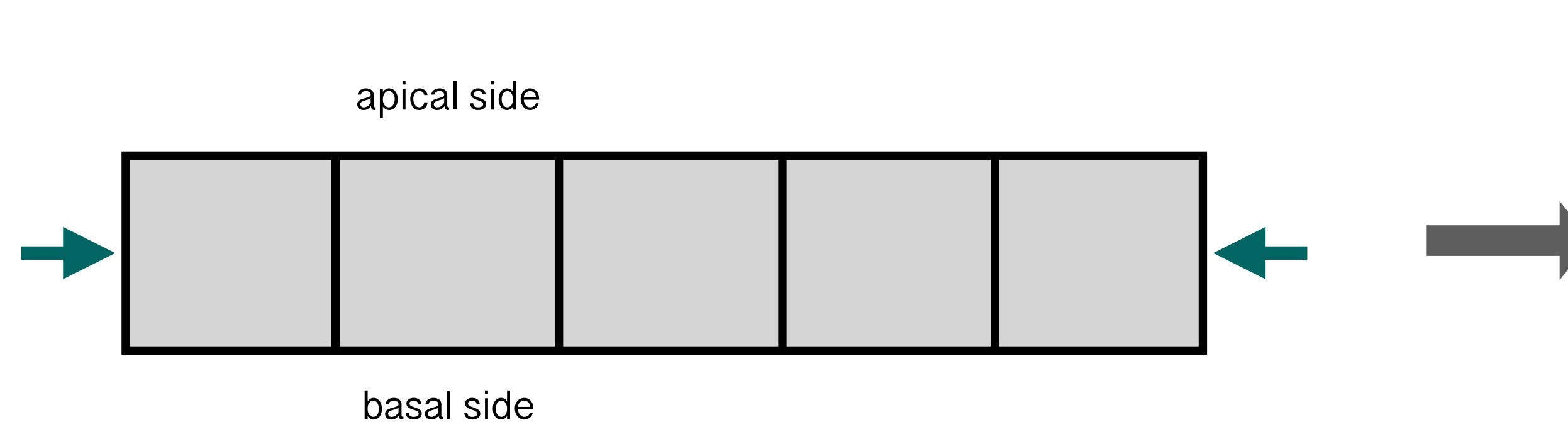
Effective **compression** due to differential growth rates of epithelium and foundation



Nelson, On Buckling Morphogenesis (2016)

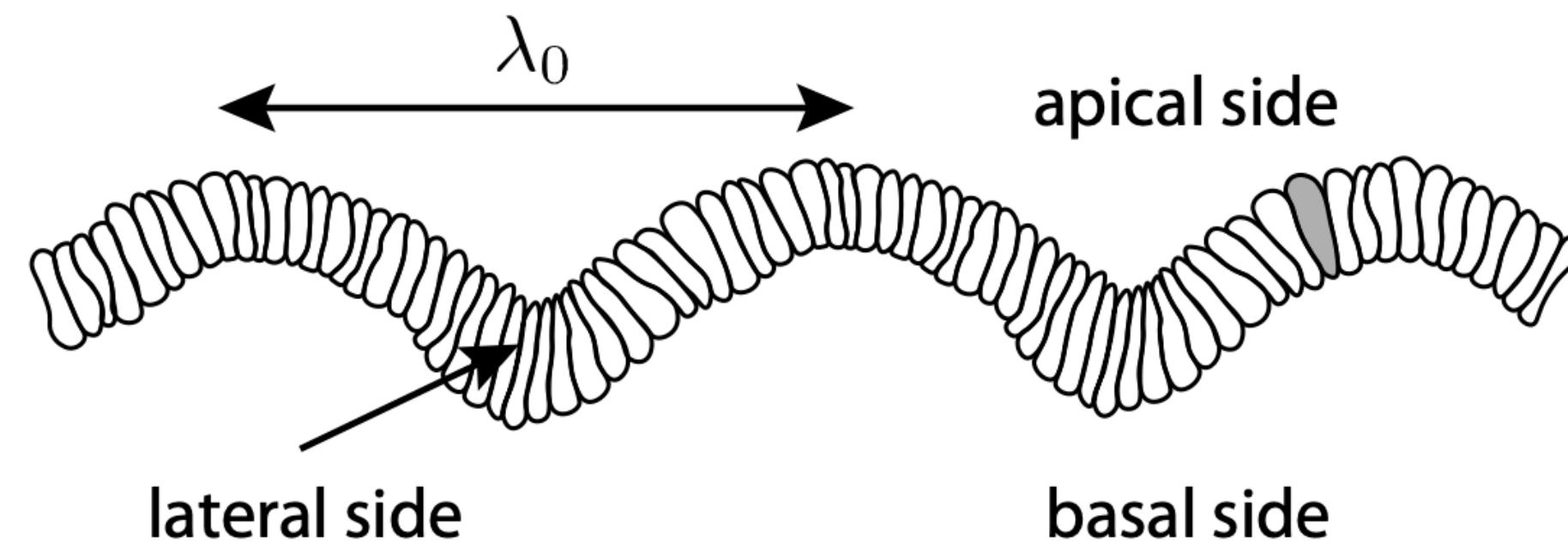
MECHANICAL INSTABILITIES

Buckling:



Haas et al., Nonlinear and nonlocal elasticity in coarse-grained differential-tension models of epithelia, PRE (2019)

Wrinkling:

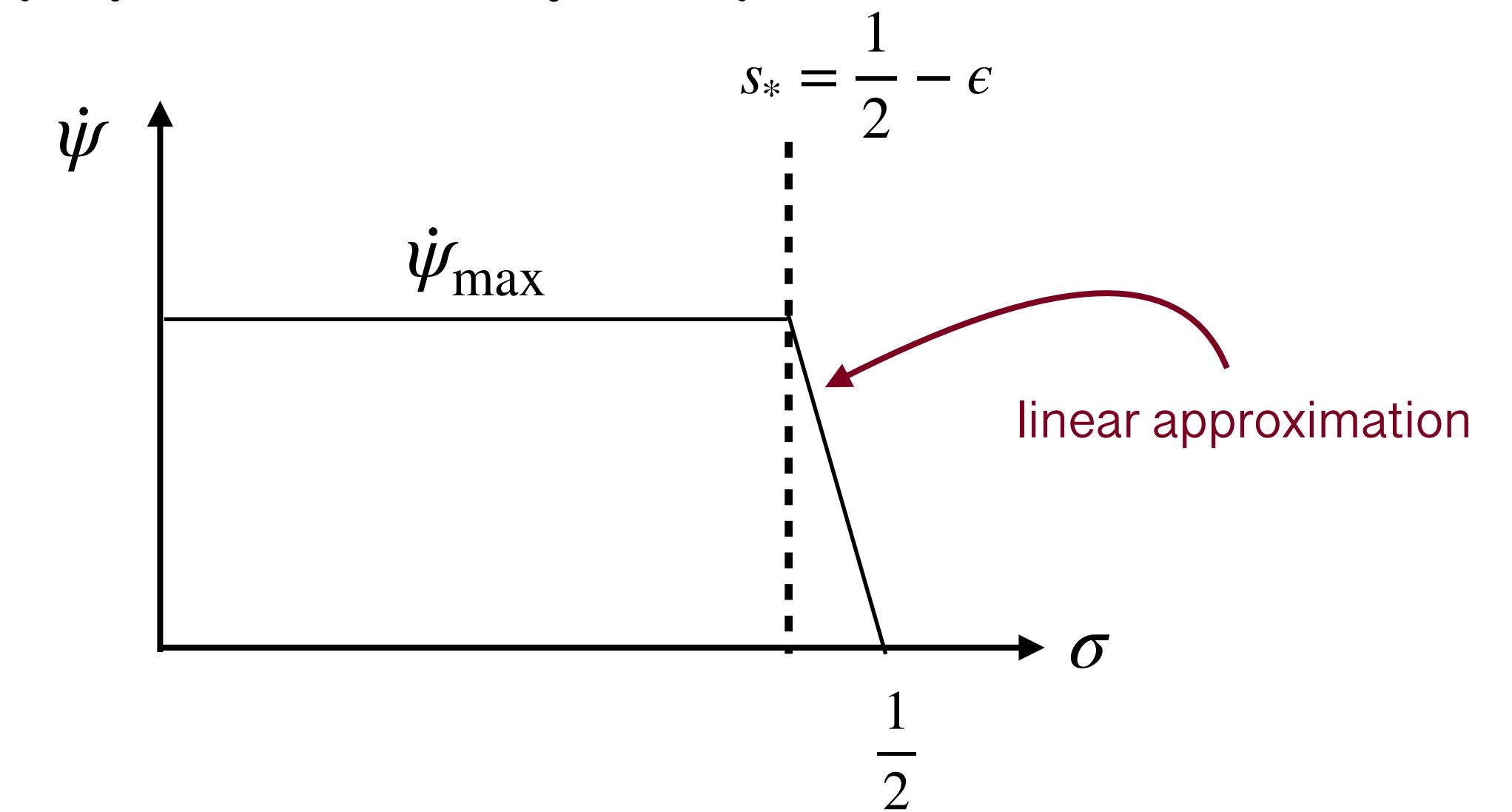
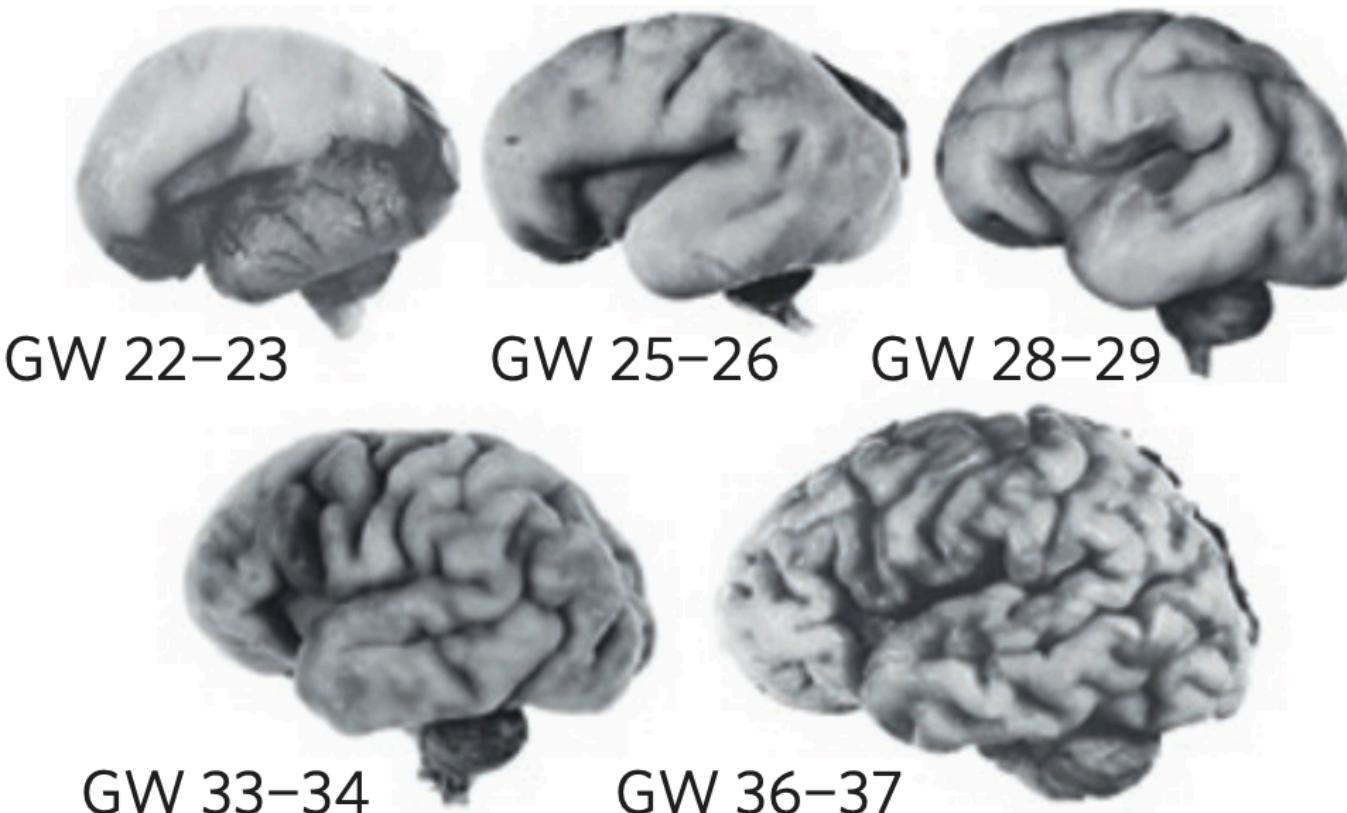


Andrenšek et al., Wrinkling Instability in Unsupported Epithelial Sheets, PRL (2023)

GOAL: INVESTIGATE LARGE COMPRESSIONS

Why?

- Large compressions in biology
- Interesting for continuum mechanics



How?

- write governing ems
- Solve for shapes
- Modification required for large compressions