

# Constraining $M_\nu$ with the Bispectrum I: Breaking Parameter Degeneracies

CHANGHOON HAHN,<sup>1,2,\*</sup> FRANCISCO VILLAESCUSA-NAVARRO,<sup>3</sup> EMANUELE CASTORINA,<sup>2,1</sup> AND  
ROMAN SCOCCIMARRO<sup>4</sup>

<sup>1</sup>*Lawrence Berkeley National Laboratory, 1 Cyclotron Rd, Berkeley CA 94720, USA*

<sup>2</sup>*Berkeley Center for Cosmological Physics, University of California, Berkeley, CA 94720, USA*

<sup>3</sup>*Center for Computational Astrophysics, Flatiron Institute, 162 5th Avenue, New York, NY 10010, USA*

<sup>4</sup>*Center for Cosmology and Particle Physics, Department of Physics, New York University, NY 10003, New York, USA*

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## ABSTRACT

Massive neutrinos suppress the growth of structure below their free-streaming scale and leave an imprint on large-scale structure. Measuring this imprint allows us to constrain the sum of neutrino masses,  $M_\nu$ , a key parameter in particle physics beyond the Standard Model. However, degeneracies among cosmological parameters, especially between  $M_\nu$  and  $\sigma_8$ , limit the constraining power of standard two-point clustering statistics. In this work, we investigate whether we can break these degeneracies and constrain  $M_\nu$  with the next higher-order correlation function — the bispectrum. We first examine the redshift-space halo bispectrum of 800  $N$ -body simulations from the HADES suite and demonstrate that the bispectrum helps break the  $M_\nu$ – $\sigma_8$  degeneracy. Then using 22,000  $N$ -body simulations of the Quijote suite, we quantify for the first time the full information content of the redshift-space halo bispectrum down to nonlinear scales using a Fisher matrix forecast of  $\{\Omega_m, \Omega_b, h, n_s, \sigma_8, M_\nu\}$ . For  $k_{\text{max}}=0.5 \ h/\text{Mpc}$ , the bispectrum provides  $\Omega_m$ ,  $\Omega_b$ ,  $h$ ,  $n_s$ , and  $\sigma_8$  constraints 1.9, 2.6, 3.1, 3.6, and 2.6 times tighter than the power spectrum. For  $M_\nu$ , the bispectrum improves the  $1\sigma$  constraint from 0.2968 to 0.0572 eV — over 5 times tighter than the power spectrum. Even with priors from *Planck*, the bispectrum improves  $M_\nu$  constraints by a factor of 2.7. Although we reserve marginalizing over a more complete set of bias parameters to the next paper of the series, these constraints are derived for a  $(1 \ h^{-1}\text{Gpc})^3$  box, a substantially smaller volume than upcoming surveys. Thus, our results demonstrate that the bispectrum offers significant improvements over the power spectrum, especially for constraining  $M_\nu$ .

*Keywords:* cosmology: cosmological parameters — cosmology: large-scale structure of Universe. — cosmology: theory

\* hahn.changhoon@gmail.com

## 1. INTRODUCTION

The lower bound on the sum of neutrino masses ( $M_\nu \gtrsim 0.06$  eV), discovered by neutrino oscillation experiments, provides conclusive evidence of physics beyond the Standard Model of particle physics (Forero et al. 2014; Gonzalez-Garcia et al. 2016). A more precise measurement of  $M_\nu$  has the potential to distinguish between the ‘normal’ and ‘inverted’ neutrino mass hierarchy scenarios and further reveal the physics of neutrinos. Neutrino oscillation experiments, however, are insensitive to the absolute neutrino mass scales. Other laboratory experiments sensitive to  $M_\nu$  (e.g. double beta decay and tritium beta decay experiments) have the potential to place upper bounds of  $M_\nu < 0.2$  eV in upcoming experiments (Bonn et al. 2011; Drexlin et al. 2013). However, these upper bound alone are not insufficient to distinguish between the mass hierarchies. Neutrinos, through the cosmic neutrino background, affect the expansion history and the growth of cosmic structure. Measuring these effects with cosmological observables provides complementary and potentially more precise measurements of  $M_\nu$ .

Neutrinos, in the early Universe, are relativistic and contribute to the energy density of radiation. Later as they become non-relativistic, they contribute to the energy density of matter. This transition affects the expansion history of the Universe and leaves imprints observable in the cosmic microwave background (CMB) anisotropy spectrum (Lesgourgues & Pastor 2012, 2014). Massive neutrinos also impact the growth of structure. On large scales, neutrino perturbations are indistinguishable from perturbations of cold dark matter (CDM). However, on scales smaller than their free-streaming scale, neutrinos do not contribute to the clustering and thereby reduce the amplitude of the total matter power spectrum. In addition, they also reduce the growth rate of CDM perturbations at late times. This combined suppression of the small-scale matter power spectrum leaves measurable imprints on the CMB as well as large-scale structure. For more on details the effect of neutrinos in cosmological observables, we refer readers to Lesgourgues & Pastor (2012, 2014); Gerbino (2018).

The tightest cosmological constraints on  $M_\nu$  currently come from combining CMB data with other cosmological probes. Temperature and large angle polarization data from the *Planck* satellite places an upper bound of  $M_\nu < 0.54$  eV with 95% confidence level (Planck Collaboration et al. 2018). Adding the Baryon Acoustic Oscillation (BAO) to the *Planck* likelihood breaks geometrical degeneracies (among  $M_\nu$ ,  $h$ ,  $\Omega_m$ ) and significantly tightens the upper bound to  $M_\nu < 0.16$  eV. CMB lensing further tightens the bound to  $M_\nu < 0.13$  eV, though not as significantly. Combining these probes, however, introduce a strong correlation between  $M_\nu$  and  $A_s$ , the scale amplitude. Since CMB experiments measure the combined quantity  $A_s e^{-2\tau}$ , where  $\tau$  is the optical depth of reionization, this leads to a strong degeneracy between  $M_\nu$  and  $\tau$  (Allison et al. 2015; Liu et al. 2016; Archidiacono et al. 2017).

The  $\tau - M_\nu$  degeneracy will continue to be a major bottleneck for  $M_\nu$  constraints from CMB data. The best constraints on  $\tau$  currently come from *Planck* —  $\tau \sim 0.06$ . However, upcoming ground-based CMB experiments (e.g. CMB-S4) will not observe scales larger than  $\ell < 30$ , and therefore will not directly constrain  $\tau$  (Abazajian et al. 2016). Meanwhile, proposed future space-based experiments

such as LiteBIRD<sup>1</sup> and CORe<sup>2</sup>, which have the potential to precisely measure  $\tau$ , have yet to be confirmed. CMB data, however, is not the only way to improve  $M_\nu$  constraints. The imprint of neutrinos on 3D clustering of galaxies can be measured to constrain  $M_\nu$  and with the sheer cosmic volumes mapped, upcoming surveys such as DESI<sup>3</sup>, PFS<sup>4</sup>, EUCLID<sup>5</sup>, and WFIRST<sup>6</sup> will be able tightly constrain  $M_\nu$  (Audren et al. 2013; Font-Ribera et al. 2014; Petracca et al. 2016; Sartoris et al. 2016; Boyle & Komatsu 2018).

A major limitation of using 3D clustering is obtaining accurate theoretical predictions beyond linear scales, for bias tracers, and in redshift space. Simulations have made huge strides in accurately and efficiently modeling nonlinear structure formation with massive neutrinos (*e.g.* Brandbyge et al. 2008; Villaescusa-Navarro et al. 2013; Castorina et al. 2015; Adamek et al. 2017; Emberson et al. 2017; Villaescusa-Navarro et al. 2018). In conjunction, new simulation based ‘emulation’ models that exploit the accuracy of  $N$ -body simulations while minimizing the computing budget have been applied to analyze small-scale galaxy clustering with remarkable success (*e.g.* Heitmann et al. 2009; Kwan et al. 2015; Euclid Collaboration et al. 2018; McClintock et al. 2018; Zhai et al. 2018; Wibking et al. 2019). Developments on these fronts have the potential to unlock the information content in nonlinear clustering to constrain  $M_\nu$ .

Various works have examined the impact of neutrino masses on nonlinear clustering of matter in real-space (*e.g.* Brandbyge et al. 2008; Saito et al. 2008; Wong 2008; Saito et al. 2009; Viel et al. 2010; Agarwal & Feldman 2011; Bird et al. 2012; Castorina et al. 2015; Banerjee & Dalal 2016) and in redshift-space (Marulli et al. 2011; Castorina et al. 2015; Upadhye et al. 2016). Most recently, using a suite of more than 1000 simulations, Villaescusa-Navarro et al. (2018) examined the impact of  $M_\nu$  on the redshift-space matter and halo power spectrum to find that the imprint of  $M_\nu$  and  $\sigma_8$  on the power spectrum are degenerate and differ by  $< 1\%$  (see also Figure 1). The strong  $M_\nu - \sigma_8$  degeneracy poses a serious limitation of constraining  $M_\nu$  with the power spectrum. However, information in the nonlinear regime cascade from the power spectrum to higher-order statistics — *e.g.* the bispectrum. In fact, the bispectrum has a comparable signal-to-noise ratio to the power spectrum on nonlinear scales (Sefusatti & Scoccimarro 2005; Chan & Blot 2017). Furthermore, although  $M_\nu$  is not included in their analyses, Sefusatti et al. (2006) and Yankelevich & Porciani (2019) have shown that including the bispectrum significantly improves constraints on cosmological parameters. Including  $M_\nu$ , Chudaykin & Ivanov (2019) find that the bispectrum significantly improves constraints for  $M_\nu$ . Their forecasts, however, do not include the constraining power on nonlinear scales (Section 4.2). No work to date has quantified the total information content and constraining power of the full bispectrum down to nonlinear scales — especially for  $M_\nu$ .

In this work, we examine the effect of massive neutrinos on the redshift-space halo bispectrum using  $\sim 23,000$   $N$ -body simulations with different neutrino masses. We first demonstrate that the bispectrum helps break the  $M_\nu - \sigma_8$  degeneracy found in the power spectrum. Then we present the

<sup>1</sup> <http://litebird.jp/eng/>

<sup>2</sup> <http://www.core-mission.org/>

<sup>3</sup> <https://www.desi.lbl.gov/>

<sup>4</sup> <https://pfs.ipmu.jp/>

<sup>5</sup> <http://sci.esa.int/euclid/>

<sup>6</sup> <https://wfirst.gsfc.nasa.gov/>

full information content of the bispectrum for all triangle configurations down to  $k_{\text{max}} = 0.5 \, h/\text{Mpc}$  using a Fisher forecast where we estimate the covariance matrix and derivatives with the large set of simulations from the Quijote suite (Villaescusa-Navarro et al. in preparation). This paper is the first of a series of papers that aim to demonstrate the potential of galaxy bispectrum analyses in constraining  $M_\nu$ . In particular, with this series, we aim to establish the potential for constraining cosmological parameters (especially  $M_\nu$ ) with galaxy bispectrum analyses that extend to nonlinear scales using simulation based emulation models for upcoming galaxy surveys. In the subsequent paper, we will include HOD parameters in our forecasts to quantify the full information content of the galaxy bispectrum. In the series, we will also present methods to tackle challenges that come with analyzing the full galaxy bispectrum, such as data compression for reducing the dimensionality of the bispectrum.

In Section 2, we describe the two  $N$ -body simulation suites, HADES and Quijote, and the halo catalogs constructed from them. We then describe in Section 3, how we measure the bispectrum of these simulations. Afterwards, we use the redshift-space halo bispectra to demonstrate the distinct imprint of  $M_\nu$  on the bispectrum, which allows it to break the degeneracy between  $M_\nu$  and  $\sigma_8$ , in Section 4.1. Finally, in Section 4.2 we present the full information content of the halo bispectrum with a Fisher forecast of cosmological parameters and demonstrate how the bispectrum significantly improves the constraints on the cosmological parameters:  $\Omega_m$ ,  $\Omega_b$ ,  $h$ ,  $n_s$ ,  $\sigma_8$ , and *especially*  $M_\nu$ .

## 2. HADES AND QUIJOTE SIMULATION SUITES

The HADES<sup>7</sup> and Quijote<sup>8</sup> suites are sets of, 43000 total,  $N$ -body simulations run on multiple cosmologies, including those with massive neutrinos ( $M_\nu > 0 \, \text{eV}$ ). In this work, we use a subset of the HADES and Quijote simulations. Below, we briefly describe these simulations; a summary of the simulations can be found in Table 1. The HADES simulations start from Zel’dovich approximated initial conditions generated at  $z = 99$  using the Zennaro et al. (2017) rescaling method and follow the gravitational evolution of  $N_{\text{cdm}} = 512^3$  CDM, plus  $N_\nu = 512^3$  neutrino particles for  $M_\nu > 0 \, \text{eV}$  cosmologies, to  $z = 0$ . They are run using the GADGET-III TreePM+SPH code (Springel 2005) in a periodic  $(1 \, h^{-1}\text{Gpc})^3$  box. All of the HADES simulations share the following cosmological parameter values, which are in good agreement with Planck constraints Ade et al. (2016):  $\Omega_m=0.3175$ ,  $\Omega_b=0.049$ ,  $\Omega_\Lambda=0.6825$ ,  $n_s=0.9624$ ,  $h=0.6711$ , and  $k_{\text{pivot}} = 0.05 \, h\text{Mpc}^{-1}$ .

The HADES suite includes  $N$ -body simulations with degenerate massive neutrinos of  $M_\nu = 0.06, 0.10$ , and  $0.15 \, \text{eV}$ . These simulations are run using the “particle method”, where neutrinos are described as a collisionless and pressureless fluid and therefore modeled as particles, same as CDM (Brandbyge et al. 2008; Viel et al. 2010). HADES also includes simulations with massless neutrino and different values of  $\sigma_8$  to examine the  $M_\nu - \sigma_8$  degeneracy. The  $\sigma_8$  values were chosen to match either  $\sigma_8^m$  or  $\sigma_8^c$  —  $\sigma_8$  computed with respect to total matter (CDM + baryons +  $\nu$ ) or CDM + baryons — of the massive neutrino simulations:  $\sigma_8 = 0.822, 0.818, 0.807$ , and  $0.798$ . Each model has 100 independent realizations and we focus on the snapshots saved at  $z = 0$ . Halos closely trace

<sup>7</sup> <https://franciscovillaescusa.github.io/hades.html>

<sup>8</sup> <https://github.com/franciscovillaescusa/Quijote-simulations>

the CDM+baryon field rather than the total matter field and neutrinos have negligible contribution to halo masses (*e.g.* Ichiki & Takada 2012; Castorina et al. 2014; LoVerde 2014; Villaescusa-Navarro et al. 2014). Hence, dark matter halos are identified in each realization using the Friends-of-Friends algorithm (FoF; Davis et al. 1985) with linking length  $b = 0.2$  on the CDM + baryon distribution. We limit the catalogs to halos with masses above  $M_{\text{lim}} = 3.2 \times 10^{13} h^{-1} M_{\odot}$ . We refer readers to Villaescusa-Navarro et al. (2018) for more details on the HADES simulations.

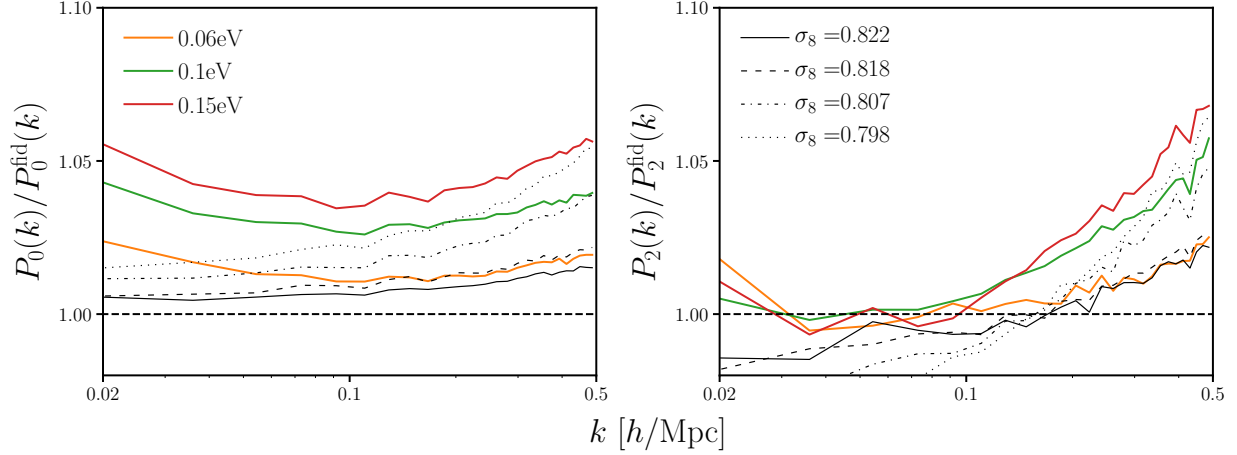
In addition to HADES, we use simulations from the Quijote suite, a set of 42,000  $N$ -body simulations that in total contain more than 8 trillion ( $8 \times 10^{12}$ ) particles over a volume of  $42000(h^{-1}\text{Gpc})^3$ . These simulations were designed to quantify the information content of different cosmological observables using Fisher matrix forecasting technique (Section 4.2). They are therefore constructed to accurately calculate the covariance matrices of observables and the derivatives of observables with respect to cosmological parameters:  $\Omega_{\text{m}}$ ,  $\Omega_{\text{b}}$ ,  $h$ ,  $n_s$ ,  $\sigma_8$ , and  $M_{\nu}$ .

To calculate covariance matrices, Quijote includes  $N_{\text{cov}}=15,000$   $N$ -body simulations at a fiducial cosmology ( $\Omega_{\text{m}}=0.3175$ ,  $\Omega_{\text{b}}=0.049$ ,  $h=0.6711$ ,  $n_s=0.9624$ ,  $\sigma_8=0.834$ , and  $M_{\nu}=0.0$  eV). It also includes sets of 500  $N$ -body simulations run at different cosmologies where only one parameter is varied from the fiducial cosmology at a time for the derivatives. Along  $\Omega_{\text{m}}$ ,  $\Omega_{\text{b}}$ ,  $h$ ,  $n_s$ , and  $\sigma_8$ , the fiducial cosmology is adjusted by either a small step above or below the fiducial value:  $\{\Omega_{\text{m}}^+, \Omega_{\text{m}}^-, \Omega_{\text{b}}^+, \Omega_{\text{b}}^-, h^+, h^-, n_s^+, n_s^-, \sigma_8^+, \sigma_8^-\}$ . Along  $M_{\nu}$ , because  $M_{\nu} \geq 0.0$  eV and the derivative of certain observable with respect to  $M_{\nu}$  is noisy, Quijote includes sets of 500 simulations for  $M_{\nu} = 0.1$ , 0.2, and 0.4 eV. Table 1 lists the cosmologies included in the Quijote suite.

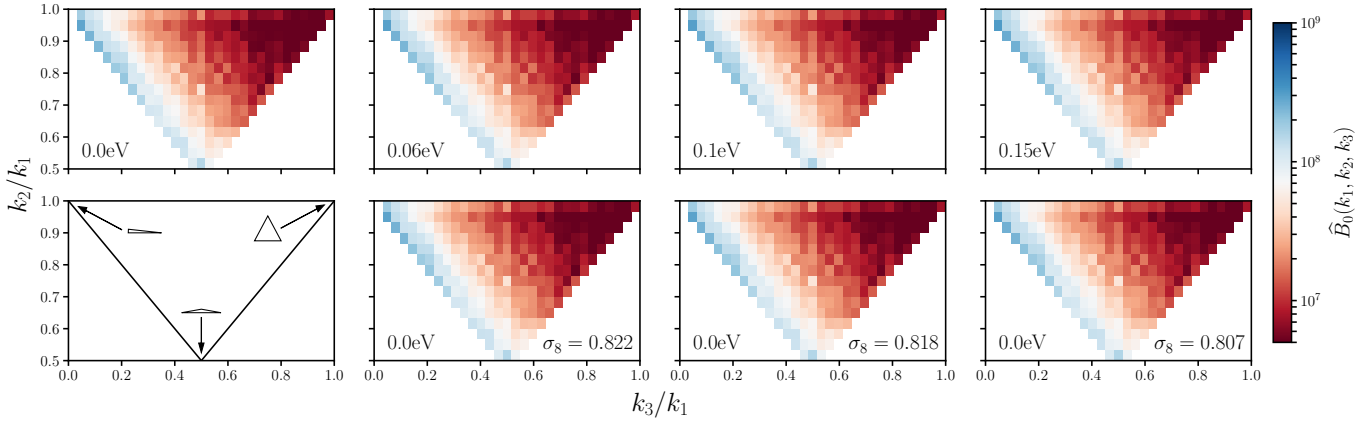
The initial conditions for all Quijote simulations were generated at  $z = 127$  using 2LPT for simulations with massless neutrinos and the Zel'dovich approximation for massive neutrinos. The suite also includes a set of simulations at the fiducial  $M_{\nu} = 0$  eV cosmology with Zel'dovich approximation initial conditions, which we use later in Section 4.2. Like HADES, the initial conditions of simulations with massive neutrinos take their scale-dependent growth factors/rates into account using the Zennaro et al. (2017) method. From the initial conditions, all of the simulations follow the gravitational evolution of  $512^3$  dark matter particles, and  $512^3$  neutrino particles for massive neutrino models, to  $z = 0$  using GADGET-III TreePM+SPH code (same as HADES). Halos are then identified using the same FoF scheme and mass limit as HADES. For the fiducial cosmology, the halo catalogs have  $\sim 156,000$  halos ( $\bar{n} \sim 1.56 \times 10^{-4} h^3 \text{Gpc}^{-3}$ ) with  $\bar{n}P_0(k = 0.1) \sim 3.23$ . We refer readers to Villaescusa-Navarro et al. (in preparation) for further details on the Quijote simulations.

### 3. BISPECTRUM

We are interested in breaking parameter degeneracies that limit the constraining power on  $M_{\nu}$  of two-point clustering analyses using three-point clustering statistics — *i.e.* the bispectrum. In this section, we describe the bispectrum estimator used throughout the paper. We focus on the bispectrum monopole ( $\ell = 0$ ) and use an estimator that exploits the speed of Fast Fourier Transforms (FFTs). Our estimator is similar to the estimators described in Sefusatti (2005), Scoccimarro (2015), and Sefusatti et al. (2016); we also follow their formalism in our description below. Although Sefusatti et al. (2016) and Scoccimarro (2015) respectively describe estimators in real- and redshift-space, since we focus on the bispectrum monopole, we note that there is no difference.



**Figure 1.** Impact of  $M_\nu$  and  $\sigma_8$  on the redshift-space halo power spectrum monopole,  $P_0$ , and quadrupole,  $P_2$ , of the HADES simulation suite. The  $P_\ell$  measurements at each cosmology is averaged over 100  $N$ -body realizations.  $M_\nu$  and  $\sigma_8$  produce almost identical effects on halo clustering on small scales ( $k > 0.1 h/\text{Mpc}$ ). This degeneracy can be partially broken through the quadrupole; however,  $M_\nu$  and  $\sigma_8$  produce almost the same effect on two-point clustering — within a few percent.



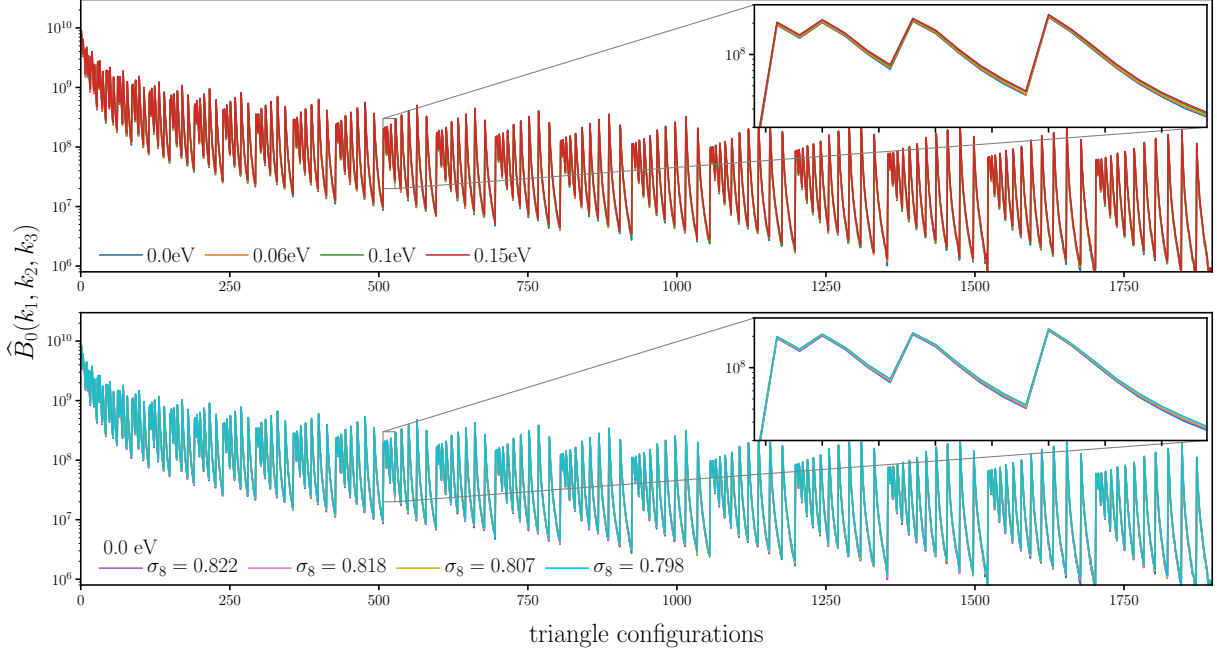
**Figure 2.** The redshift-space halo bispectrum,  $\hat{B}_0(k_1, k_2, k_3)$ , as a function of triangle configuration shape for  $M_\nu = 0.0, 0.06, 0.10$ , and  $0.15 \text{ eV}$  (upper panels) and  $M_\nu = 0.0 \text{ eV}$ ,  $\sigma_8 = 0.822, 0.818$ , and  $0.807$  (lower panels). The  $\hat{B}_0$  for each cosmology (each panel) is averaged over 100  $N$ -body realizations. The HADES simulations of the top and bottom panels in the three right-most columns, have matching  $\sigma_8$  values (Section 2). We describe the triangle configuration shape by the ratio of the triangle sides:  $k_3/k_1$  and  $k_2/k_1$ . As we describe schematically in the lower leftmost panel, in each panel, the upper left bin contains squeezed triangles ( $k_1 = k_2 \gg k_3$ ), the upper right bin contains equilateral triangles ( $k_1 = k_2 = k_3$ ), and the bottom center bin contains folded triangles ( $k_1 = 2k_2 = 2k_3$ ). We include all 1898 triangle configurations with  $k_1, k_2, k_3 \leq k_{\text{max}} = 0.5 h/\text{Mpc}$  and use the  $\hat{B}_0$  estimator in Section 3.

**Table 1.** Specifications of the HADES and Quijote simulation suites.

Name	$M_\nu$ (eV)	$\Omega_m$	$\Omega_b$	$h$	$n_s$	$\sigma_8^m$ ( $10^{10}h^{-1}M_\odot$ )	$\sigma_8^c$ ( $10^{10}h^{-1}M_\odot$ )	ICs	realizations
HADES suite									
Fiducial	0.0	0.3175	0.049	0.6711	0.9624	0.833	0.833	Zel'dovich	100
	0.06	0.3175	0.049	0.6711	0.9624	0.819	0.822	Zel'dovich	100
	0.10	0.3175	0.049	0.6711	0.9624	0.809	0.815	Zel'dovich	100
	0.15	0.3175	0.049	0.6711	0.9624	0.798	0.806	Zel'dovich	100
	0.0	0.3175	0.049	0.6711	0.9624	0.822	0.822	Zel'dovich	100
	0.0	0.3175	0.049	0.6711	0.9624	0.818	0.818	Zel'dovich	100
	0.0	0.3175	0.049	0.6711	0.9624	0.807	0.807	Zel'dovich	100
	0.0	0.3175	0.049	0.6711	0.9624	0.798	0.798	Zel'dovich	100
Quijote suite									
Fiducial	0.0	0.3175	0.049	0.6711	0.9624	0.834	0.834	2LPT	15,000
Fiducial ZA	0.0	0.3175	0.049	0.6711	0.9624	0.834	0.834	Zel'dovich	500
$M_\nu^+$	<u>0.1</u>	0.3175	0.049	0.6711	0.9624	0.834	0.834	Zel'dovich	500
$M_\nu^{++}$	<u>0.2</u>	0.3175	0.049	0.6711	0.9624	0.834	0.834	Zel'dovich	500
$M_\nu^{+++}$	<u>0.4</u>	0.3175	0.049	0.6711	0.9624	0.834	0.834	Zel'dovich	500
$\Omega_m^+$	0.0	<u>0.3275</u>	0.049	0.6711	0.9624	0.834	0.834	2LPT	500
$\Omega_m^-$	0.0	<u>0.3075</u>	0.049	0.6711	0.9624	0.834	0.834	2LPT	500
$\Omega_b^+$	0.0	0.3175	<u>0.051</u>	0.6711	0.9624	0.834	0.834	2LPT	500
$\Omega_b^-$	0.0	0.3175	<u>0.047</u>	0.6711	0.9624	0.834	0.834	2LPT	500
$h^+$	0.0	0.3175	0.049	<u>0.6911</u>	0.9624	0.834	0.834	2LPT	500
$h^-$	0.0	0.3175	0.049	<u>0.6511</u>	0.9624	0.834	0.834	2LPT	500
$n_s^+$	0.0	0.3175	0.049	0.6711	<u>0.9824</u>	0.834	0.834	2LPT	500
$n_s^-$	0.0	0.3175	0.049	0.6711	<u>0.9424</u>	0.834	0.834	2LPT	500
$\sigma_8^+$	0.0	0.3175	0.049	0.6711	0.9624	<u>0.849</u>	<u>0.849</u>	2LPT	500
$\sigma_8^-$	0.0	0.3175	0.049	0.6711	0.9624	<u>0.819</u>	<u>0.819</u>	2LPT	500

**Top:** The HADES suite includes sets of 100  $N$ -body simulations with degenerate massive neutrinos of  $M_\nu = 0.06, 0.10$ , and  $0.15$  eV as well as sets of simulations with  $M_\nu = 0.0$  eV and  $\sigma_8 = 0.822, 0.818, 0.807$ , and  $0.798$  to examine the  $M_\nu - \sigma_8$  degeneracy. **Bottom:** The Quijote suite includes 15,000  $N$ -body simulations at the fiducial cosmology to accurately estimate the covariance matrices. It also includes sets of 500 simulations at 13 different cosmologies, where only one parameter is varied from the fiducial value (underlined), to estimate derivatives of observables along the cosmological parameters.





**Figure 3.** The redshift-space halo bispectrum,  $\hat{B}_0(k_1, k_2, k_3)$ , as a function of triangle configurations for  $M_\nu = 0.0, 0.06, 0.10$ , and  $0.15$  eV (top panel) and  $M_\nu = 0.0$  eV,  $\sigma_8 = 0.822, 0.818, 0.807$ , and  $0.798$  (lower panel). Each  $\hat{B}_0$  is averaged over 100  $N$ -body realizations. We include all possible triangle configurations with  $k_1, k_2, k_3 \leq k_{\max} = 0.5$   $h/\text{Mpc}$  where we order the configurations by looping through  $k_3$  in the inner most loop and  $k_1$  in the outer most loop satisfying  $k_1 \geq k_2 \geq k_3$  (see Appendix A for more details). In the insets of the panels we zoom into triangle configurations with  $k_1 = 0.320$ ,  $0.170 \leq k_2 \leq 0.302$ , and  $0.094 \leq k_3 \leq 0.302$   $h/\text{Mpc}$ .

To measure the bispectrum of our halo catalogs, we begin by interpolating the halo positions to a grid,  $\delta(\mathbf{x})$ , and Fourier transforming the grid to get  $\delta(\mathbf{k})$ . We use a fourth-order interpolation to get interlaced grids, which has advantageous anti-aliasing properties that allow unbiased measurements up to the Nyquist frequency (Hockney & Eastwood 1981; Sefusatti et al. 2016). Then using  $\delta(\mathbf{k})$  we measure the bispectrum monopole as

$$\hat{B}_{\ell=0}(k_1, k_2, k_3) = \frac{1}{V_B} \int_{k_1} d^3 q_1 \int_{k_2} d^3 q_2 \int_{k_3} d^3 q_3 \delta_D(\mathbf{q}_{123}) \delta(\mathbf{q}_1) \delta(\mathbf{q}_2) \delta(\mathbf{q}_3) - B_{\ell=0}^{\text{SN}}. \quad (1)$$

$\delta_D$  is the Dirac delta function and hence  $\delta_D(\mathbf{q}_{123}) = \delta_D(\mathbf{q}_1 + \mathbf{q}_2 + \mathbf{q}_3)$  ensures that the  $\mathbf{q}_i$  triplet actually forms a closed triangle. Each of the integrals above represent an integral over a spherical shell in  $k$ -space with radius  $\delta k$  centered at  $\mathbf{k}_i$ :

$$\int_{k_i} d^3 q \equiv \int_{k_i - \delta k/2}^{k_i + \delta k/2} dq q^2 \int d\Omega. \quad (2)$$



$V_B$  is a normalization factor proportional to the number of triplets  $\mathbf{q}_1$ ,  $\mathbf{q}_2$ , and  $\mathbf{q}_3$  that can be found in the triangle bin defined by  $k_1$ ,  $k_2$ , and  $k_3$  with width  $\delta k$ :

$$V_B = \int_{k_1} d^3 q_1 \int_{k_2} d^3 q_2 \int_{k_3} d^3 q_3 \delta_D(\mathbf{q}_{123}). \quad (3)$$

Lastly,  $B_{\ell=0}^{\text{SN}}$  is the correction for the Poisson shot noise, which contributes due to the self-correlation of individual objects:

$$B_{\ell=0}^{\text{SN}}(k_1, k_2, k_3) = \frac{1}{\bar{n}} (P_0(k_1) + P_0(k_2) + P_0(k_3)) + \frac{1}{\bar{n}^2} \quad (4)$$

where  $\bar{n}$  is the number density of objects (halos) and  $P_0$  is the power spectrum monopole.

In order to evaluate the integrals in Eq. 1, we take advantage of the plane-wave representation of the Dirac delta function and rewrite the equation as

$$\widehat{B}_{\ell=0}(k_1, k_2, k_3) = \frac{1}{V_B} \int \frac{d^3 x}{(2\pi)^3} \int_{k_1} d^3 q_1 \int_{k_2} d^3 q_2 \int_{k_3} d^3 q_3 \delta(\mathbf{q}_1) \delta(\mathbf{q}_2) \delta(\mathbf{q}_3) e^{i\mathbf{q}_{123} \cdot \mathbf{x}} - B_{\ell=0}^{\text{SN}} \quad (5)$$

$$= \frac{1}{V_B} \int \frac{d^3 x}{(2\pi)^3} \prod_{i=1}^3 I_{k_i}(\mathbf{x}) - B_{\ell=0}^{\text{SN}} \quad (6)$$

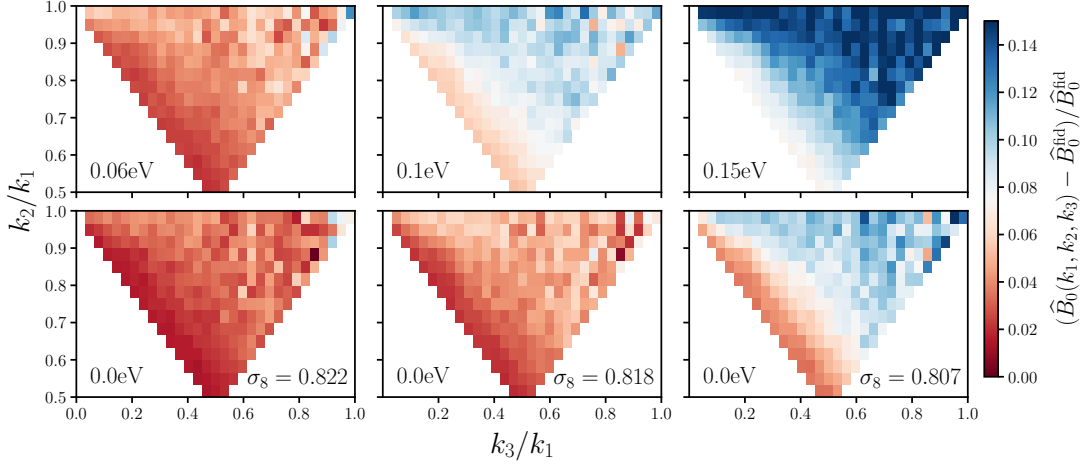
where

$$I_{k_i}(\mathbf{x}) = \int_k d^3 q \delta(\mathbf{q}) e^{i\mathbf{q} \cdot \mathbf{x}}. \quad (7)$$

At this point, we measure  $\widehat{B}_{\ell=0}(k_1, k_2, k_3)$  by calculating the  $I_{k_i}$ s with inverse FFTs and summing over in real space. For  $\widehat{B}_{\ell=0}$  measurements throughout the paper, we use  $\delta(\mathbf{x})$  grids with  $N_{\text{grid}} = 360$  and triangle configurations defined by  $k_1, k_2, k_3$  bins of width  $\Delta k = 3k_f = 0.01885 \text{ } h/\text{Mpc}$ , three times the fundamental mode  $k_f = 2\pi/(1000 \text{ } h^{-1}\text{Mpc})$  given the box size<sup>9</sup>.

We present the redshift-space halo bispectrum of the HADES simulations measured using the estimator above in two ways: one that emphasizes the triangle shape dependence (Figure 2) and the other that emphasizes the amplitude (Figure 3). In Figure 2, we plot  $\widehat{B}_0(k_1, k_2, k_3)$  as a function of  $k_2/k_1$  and  $k_3/k_1$ , which describe the triangle configuration shapes. In each panel, the colormap of the  $(k_2/k_1, k_3/k_1)$  bins represent the weighted average  $\widehat{B}_0$  amplitude of all triangle configurations in the bins. The upper left bins contain squeezed triangles ( $k_1 = k_2 \gg k_3$ ), the upper right bins contain equilateral triangles ( $k_1 = k_2 = k_3$ ), and the bottom center bins contain folded triangles ( $k_1 = 2k_2 = 2k_3$ ) as schematically highlighted in the lower leftmost panel. We include all possible 1898 triangle configurations with  $k_1, k_2, k_3 < k_{\text{max}} = 0.5 \text{ } h/\text{Mpc}$ .  $\widehat{B}_0$  in the upper panels are HADES models with  $(M_\nu, \sigma_8) = (0.0 \text{ eV}, 0.833)$  (fiducial),  $(0.06 \text{ eV}, 0.822)$ ,  $(0.10 \text{ eV}, 0.815)$ , and  $(0.15 \text{ eV}, 0.806)$ .  $\widehat{B}_0$  in the lower panels are HADES models with  $M_\nu = 0.0 \text{ eV}$  and  $\sigma_8 = 0.822, 0.818$ , and  $0.807$ . The top and bottom panels of the three right-most columns have matching  $\sigma_8$  values (Section 2).

<sup>9</sup> The code that we use to evaluate  $\widehat{B}_{\ell=0}$  is publicly available at <https://github.com/changhoonhahn/pySpectrum>



**Figure 4.** The shape dependence of the  $M_\nu$  and  $\sigma_8$  imprint on the redshift-space halo bispectrum,  $\Delta\hat{B}_0/\hat{B}_0^{\text{fid}}$ . We align the  $M_\nu = 0.06, 0.10$ , and  $0.15$  eV HADES simulations in the upper panels with  $M_\nu = 0.0$  eV  $\sigma_8 = 0.822, 0.818$ , and  $0.807$  simulations on the bottom such that the top and bottom panels in each column have matching  $\sigma_8^c$ , which produce mostly degenerate imprints on the redshift-space power spectrum.  $\hat{B}_0^{\text{fid}}$  is measured from the fiducial HADES simulation:  $M_\nu = 0.0$  eV and  $\sigma_8 = 0.833$ . The difference between the top and bottom panels highlight that  $M_\nu$  leaves a imprint distinct from  $\sigma_8$  on elongated and isosceles triangles, bins along the bottom left and bottom right edges, respectively. *The imprint of  $M_\nu$  has a distinct shape dependence on the bispectrum that cannot be replicated by varying  $\sigma_8$ .*

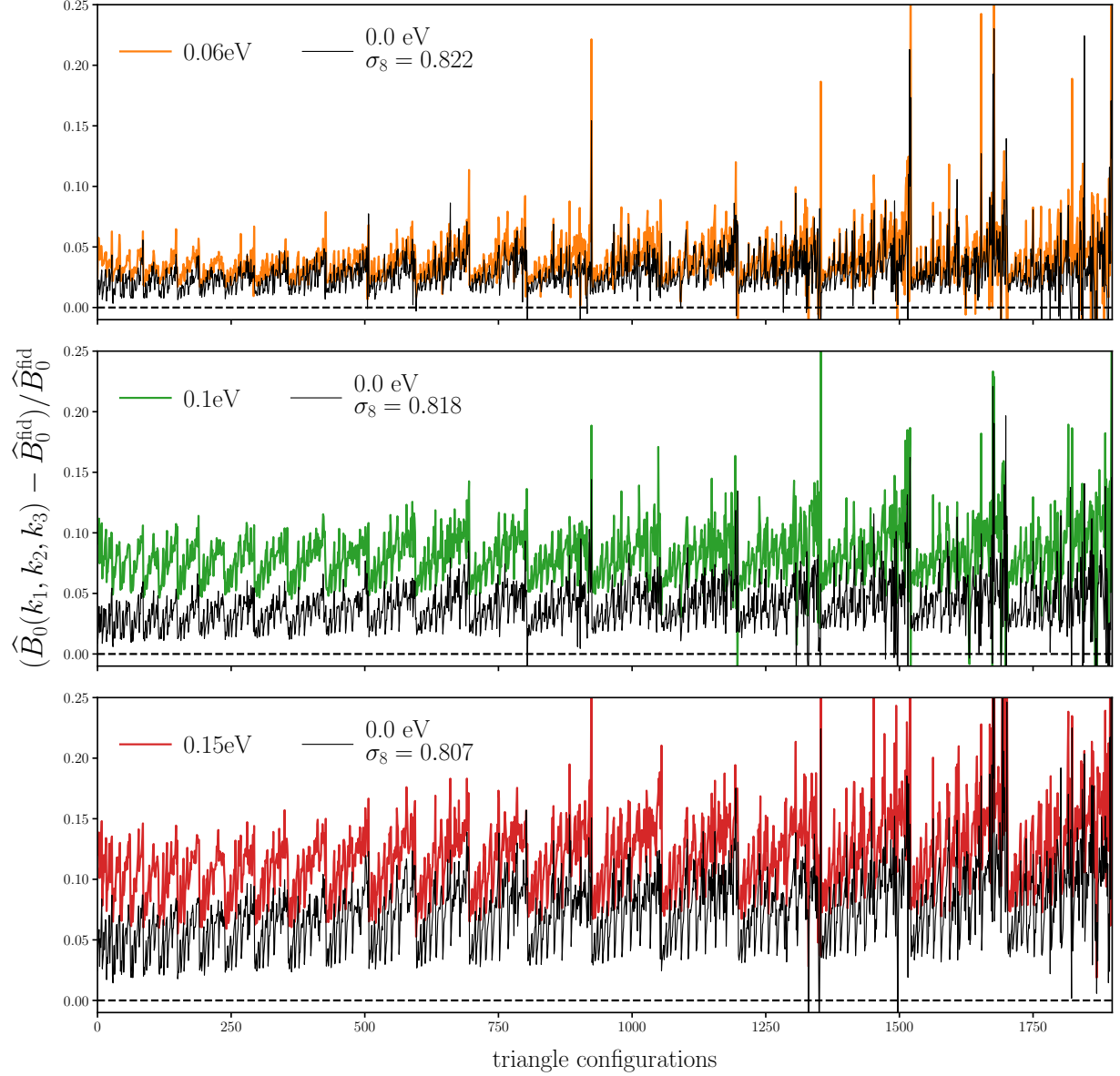
Next, in Figure 3 we plot  $\hat{B}_0(k_1, k_2, k_3)$  for all possible triangle configurations with  $k_1, k_2, k_3 < k_{\text{max}} = 0.5$  h/Mpc where we order the configurations by looping through  $k_3$  in the inner most loop and  $k_1$  in the outer most loop with  $k_1 \geq k_2 \geq k_3$ . In the top panel, we present  $\hat{B}_0$  of HADES models with  $M_\nu = 0.0, 0.06, 0.10$ , and  $0.15$  eV; in the lower panel, we present  $\hat{B}_0$  of HADES models with  $M_\nu = 0.0$  eV and  $\sigma_8 = 0.822, 0.818$ , and  $0.807$ . We zoom into triangle configurations with  $k_1 = 0.320$ ,  $0.170 \leq k_2 \leq 0.302$ , and  $0.094 \leq k_3 \leq 0.302$  h/Mpc in the insets of the panels. For further details on the redshift-space bispectrum, we refer to Appendix A.

## 4. RESULTS

### 4.1. Breaking the $M_\nu - \sigma_8$ degeneracy

One major bottleneck of constraining  $M_\nu$  with the power spectrum alone is the strong  $M_\nu - \sigma_8$  degeneracy. The imprint of  $M_\nu$  and  $\sigma_8$  on the power spectrum are degenerate and for models with the same  $\sigma_8^c$ , the power spectrum only differ by  $< 1\%$  (see Figure 1 and Villaescusa-Navarro et al. 2018). The HADES suite, which has simulations with  $M_\nu = 0.0, 0.06, 0.10$ , and  $0.15$  eV as well as  $M_\nu = 0.0$  eV simulations with matching  $\sigma_8^c$ , provides an ideal set of simulations to disentangle the impact of  $M_\nu$  and examine the degeneracy between  $M_\nu$  and  $\sigma_8$  (Section 2 and Table 1). We measure the bispectrum of the HADES simulations (Figure 2 and 3) and present how the bispectrum can significantly improve  $M_\nu$  constraints by breaking the  $M_\nu - \sigma_8$  degeneracy.

We begin by examining the triangle shape dependent imprint of  $M_\nu$  on the redshift-space halo bispectrum versus  $\sigma_8$  alone. In Figure 4, we present the fractional residual,  $(\Delta\hat{B}_0 = \hat{B}_0 - \hat{B}_0^{\text{fid}})/\hat{B}_0^{\text{fid}}$ ,



**Figure 5.** The impact of  $M_\nu$  and  $\sigma_8$  on the redshift-space halo bispectrum,  $\Delta\hat{B}_0/\hat{B}_0^{\text{fid}}$ , for all 1898 triangle configurations with  $k_1, k_2, k_3 \leq 0.5 h/\text{Mpc}$ . We compare  $\Delta\hat{B}_0/\hat{B}_0^{\text{fid}}$  of the  $M_\nu = 0.06$  (top),  $0.10$  (middle), and  $0.15$  eV (bottom) HADES simulations to  $\Delta\hat{B}_0/\hat{B}_0^{\text{fid}}$  of  $M_\nu = 0.0$  eV,  $\sigma_8 = 0.822$ ,  $0.818$ , and  $0.807$  simulations. The imprint of  $M_\nu$  on the bispectrum has a significantly different amplitude than the imprint of  $\sigma_8$ . For instance,  $M_\nu = 0.15$  eV (red) has a  $\sim 5\%$  stronger impact on the bispectrum than  $M_\nu = 0.0$  eV  $\sigma_8 = 0.798$  (black) even though their power spectra only differ by  $< 1\%$  (Figure 1). *The distinct imprint of  $M_\nu$  on the bispectrum illustrate that the bispectrum can break the degeneracy between  $M_\nu$  and  $\sigma_8$  that degrade constraints from two-point analyses.*

as a function of  $k_2/k_1$  and  $k_3/k_1$  for  $M_\nu = 0.06, 0.10$ , and  $0.15$  eV in the upper panels and  $0.0$  eV  $\sigma_8 = 0.822, 0.818$ , and  $0.807$  in the bottom panels. The simulations in the top and bottom panels of each column have matching  $\sigma_8^c$ . Overall, as  $M_\nu$  increases, the amplitude of the bispectrum increases

for all triangle shapes (top panels). This increase is due to halo bias (Villaescusa-Navarro et al. 2018, see also Figure 1). Since we impose a fixed  $M_{\text{lim}}$  on our halos so lower values of  $\sigma_8$  translate to a larger halo bias, which boosts the amplitude of the bispectrum. Within the overall increase in amplitude, however, there is a significant triangle dependence. Equilateral triangles (upper left) have the largest increase. For  $M_\nu = 0.15$  eV, the bispectrum is  $\sim 15\%$  higher than  $\hat{B}_0^{\text{fid}}$  for equilateral triangles while only  $\sim 8\%$  higher for folded triangles (lower center). The noticeable difference in  $\Delta\hat{B}_0/\hat{B}_0^{\text{fid}}$  between equilateral and squeezed triangles (upper left) is roughly consistent with the comparison in Figure 7 of Ruggeri et al. (2018). They, however, fix  $A_s$  in their simulations and measure the real-space halo bispectrum so we refrain from any detailed comparisons.

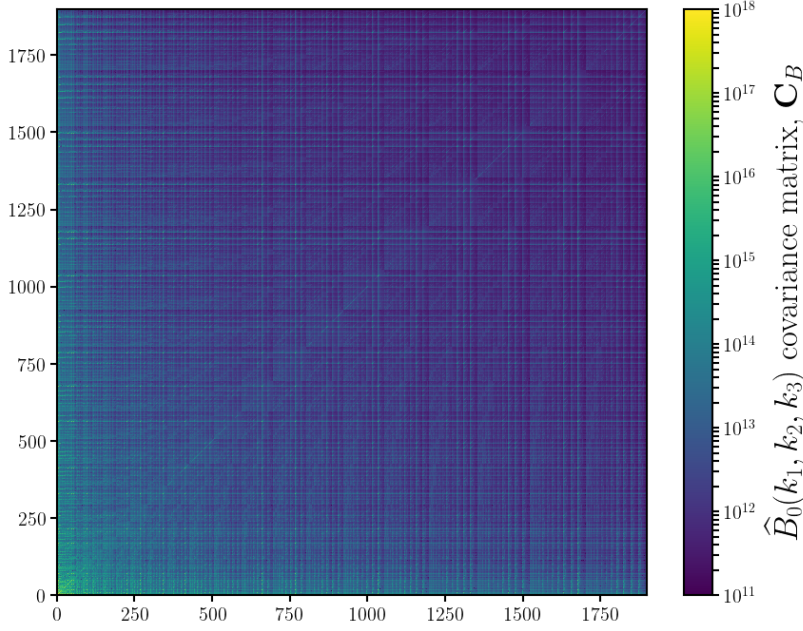
As  $\sigma_8$  increases with  $M_\nu = 0.0$  eV, the bispectrum also increases overall for all triangle shapes (bottom panels). However, the comparison of the top and bottom panels in each column reveals significant differences in  $\Delta\hat{B}_0/\hat{B}_0^{\text{fid}}$  for  $M_\nu$  versus  $\sigma_8$  alone. Between  $M_\nu = 0.15$  eV and  $\{0.0 \text{ eV}, \sigma_8 = 0.807\}$  cosmologies, there is an overall  $\gtrsim 5\%$  difference in the bispectrum. In addition, the shape dependence of the  $\Delta\hat{B}_0/\hat{B}_0^{\text{fid}}$  increase is different for  $M_\nu$  than  $\sigma_8$ . This is particularly clear in the differences between 0.1 eV (top center panel) and  $\{0.0 \text{ eV}, \sigma_8 = 0.807\}$  (bottom right panel): near equilateral triangles in the two panels have similar  $\Delta\hat{B}_0/\hat{B}_0^{\text{fid}}$  while triangle shapes near the lower left edge from the squeezed to folded triangles have significantly different  $\Delta\hat{B}_0/\hat{B}_0^{\text{fid}}$ . Hence,  $M_\nu$  leaves an imprint on the bispectrum with a distinct triangle shape dependence than  $\sigma_8$  alone. In other words, the triangle shape dependent imprint of  $M_\nu$  on the bispectrum cannot be replicated by varying  $\sigma_8$  — unlike the power spectrum.

We next examine the amplitude of the  $M_\nu$  imprint on the redshift-space halo bispectrum versus  $\sigma_8$  alone as a function of all triangle configurations. We present  $\Delta\hat{B}_0/\hat{B}_0^{\text{fid}}$  for all 1898 possible triangle configurations with  $k_1, k_2, k_3 < k_{\text{max}} = 0.5 \text{ h/Mpc}$  in Figure 5. We compare  $\Delta\hat{B}_0/\hat{B}_0^{\text{fid}}$  of the  $M_\nu = 0.06, 0.10$ , and  $0.15$  eV HADES models to the  $\Delta\hat{B}_0/\hat{B}_0^{\text{fid}}$  of  $M_\nu = 0.0$  eV  $\sigma_8 = 0.822, 0.818$ , and  $0.807$  models in the top, middle, and bottom panels, respectively. The comparison confirms the difference in overall amplitude of varying  $M_\nu$  and  $\sigma_8$  (Figure 4). For instance,  $M_\nu = 0.15$  eV (red) has a  $\sim 5\%$  stronger impact on the bispectrum than  $M_\nu = 0.0$  eV  $\sigma_8 = 0.798$  (black) even though their power spectra differ by  $< 1\%$  (Figure 1).

The comparison in the panels of Figure 5 confirms the difference in the configuration dependence in  $\Delta\hat{B}_0/\hat{B}_0^{\text{fid}}$  between  $M_\nu$  versus  $\sigma_8$ . The triangle configurations are ordered by looping through  $k_3$  in the inner most loop and  $k_1$  in the outer most loop such that  $k_1 \geq k_2 \geq k_3$ . In this ordering,  $k_1$  increases from left to right (see Figure 10 and Appendix A).  $\Delta\hat{B}_0/\hat{B}_0^{\text{fid}}$  of  $M_\nu$  expectedly increases with  $k_1$ : for small  $k_1$  (on large scales), neutrinos behave like CDM and therefore the impact is reduced. However,  $\Delta\hat{B}_0/\hat{B}_0^{\text{fid}}$  of  $M_\nu$  has a smaller  $k_1$  dependence than  $\Delta\hat{B}_0/\hat{B}_0^{\text{fid}}$  of  $\sigma_8$ . The distinct imprint of  $M_\nu$  on the redshift-space halo bispectrum illustrates that the bispectrum can break the degeneracy between  $M_\nu$  and  $\sigma_8$ . Therefore, by including the bispectrum, we can more precisely constrain  $M_\nu$  than with the power spectrum.

#### 4.2. $M_\nu$ and other Cosmological Parameter Forecasts

We demonstrate in the previous section with the HADES simulations, that the bispectrum helps break the  $M_\nu$ - $\sigma_8$  degeneracy, a major challenge in precisely constraining  $M_\nu$  with the power spectrum.



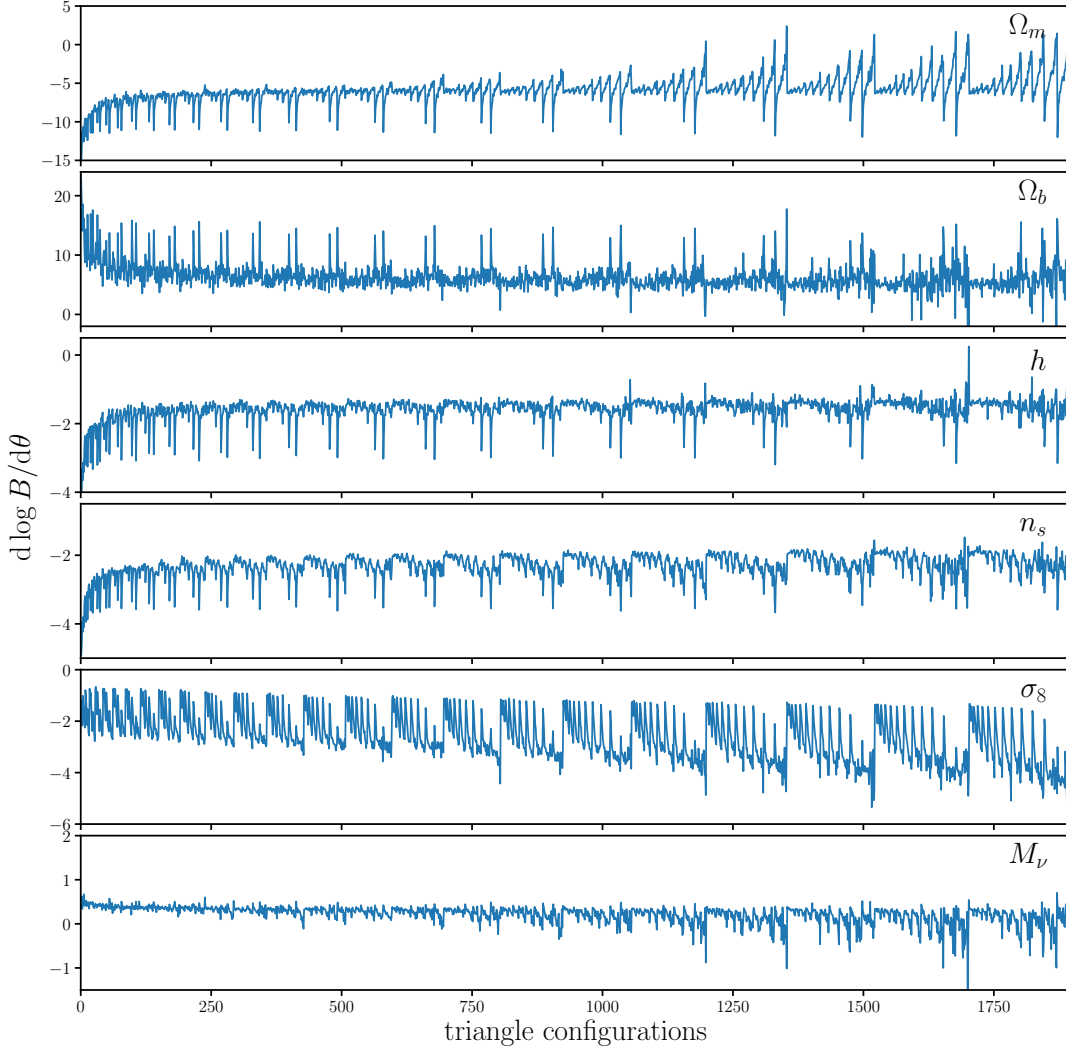
**Figure 6.** Covariance matrix of the redshift-space halo bispectrum estimated using  $N_{\text{cov}} = 15,000$  realizations of the Quijote simulation suite at the fiducial cosmology:  $\Omega_{\text{m}}=0.3175$ ,  $\Omega_{\text{b}}=0.049$ ,  $h=0.6711$ ,  $n_{\text{s}}=0.9624$ ,  $\sigma_8=0.834$ , and  $M_{\nu}=0.0$  eV. We include all possible triangle configurations with  $k_1, k_2, k_3 \leq k_{\text{max}} = 0.5 h/\text{Mpc}$  and order the configurations (bins) in the same way as Figures 3 and 5. We use the covariance matrix above for the Fisher matrix forecasts presented in Section 4.2.

While this establishes the bispectrum as a promising probe for  $M_{\nu}$ , we are ultimately interested in determining the constraining power of the bispectrum for an analysis that include cosmological parameters beyond  $M_{\nu}$  and  $\sigma_8$ — *i.e.*  $\Omega_{\text{m}}$ ,  $\Omega_{\text{b}}$ ,  $h$ , and  $n_{\text{s}}$ . The Quijote simulation suite is *specifically* designed to answer this question through Fisher matrix forecast.

First, the Quijote suite includes  $N_{\text{cov}} = 15,000$   $N$ -body realizations run at a fiducial cosmology:  $M_{\nu}=0.0\text{eV}$ ,  $\Omega_{\text{m}}=0.3175$ ,  $\Omega_{\text{b}}=0.049$ ,  $n_{\text{s}}=0.9624$ ,  $h=0.6711$ , and  $\sigma_8=0.834$  (see Table 1). This allows us to robustly estimate the  $1898 \times 1898$  covariance matrix of the *full* bispectrum,  $\mathbf{C}$  (Figure 6). Second, the Quijote suite includes 500  $N$ -body realizations evaluated at 13 different cosmologies, each a small step away from the fiducial cosmology parameter values along one parameter (Section 2 and Table 1). We apply redshift-space distortions along 3 different directions for these 500 realizations, which then effectively gives us  $N_{\text{deriv.}}=1,500$  realizations. These simulations allow us to precisely estimate the derivatives of the bispectrum with respect to each of the cosmological parameters.

Since their introduction to cosmology over two decades ago, Fisher information matrices have been ubiquitously used to forecast the constraining power of future experiments (*e.g.* Jungman et al. 1996; Tegmark et al. 1997; Dodelson 2003; Heavens 2009; Verde 2010). Defined as

$$F_{ij} = -\left\langle \frac{\partial^2 \ln \mathcal{L}}{\partial \theta_i \partial \theta_j} \right\rangle, \quad (8)$$



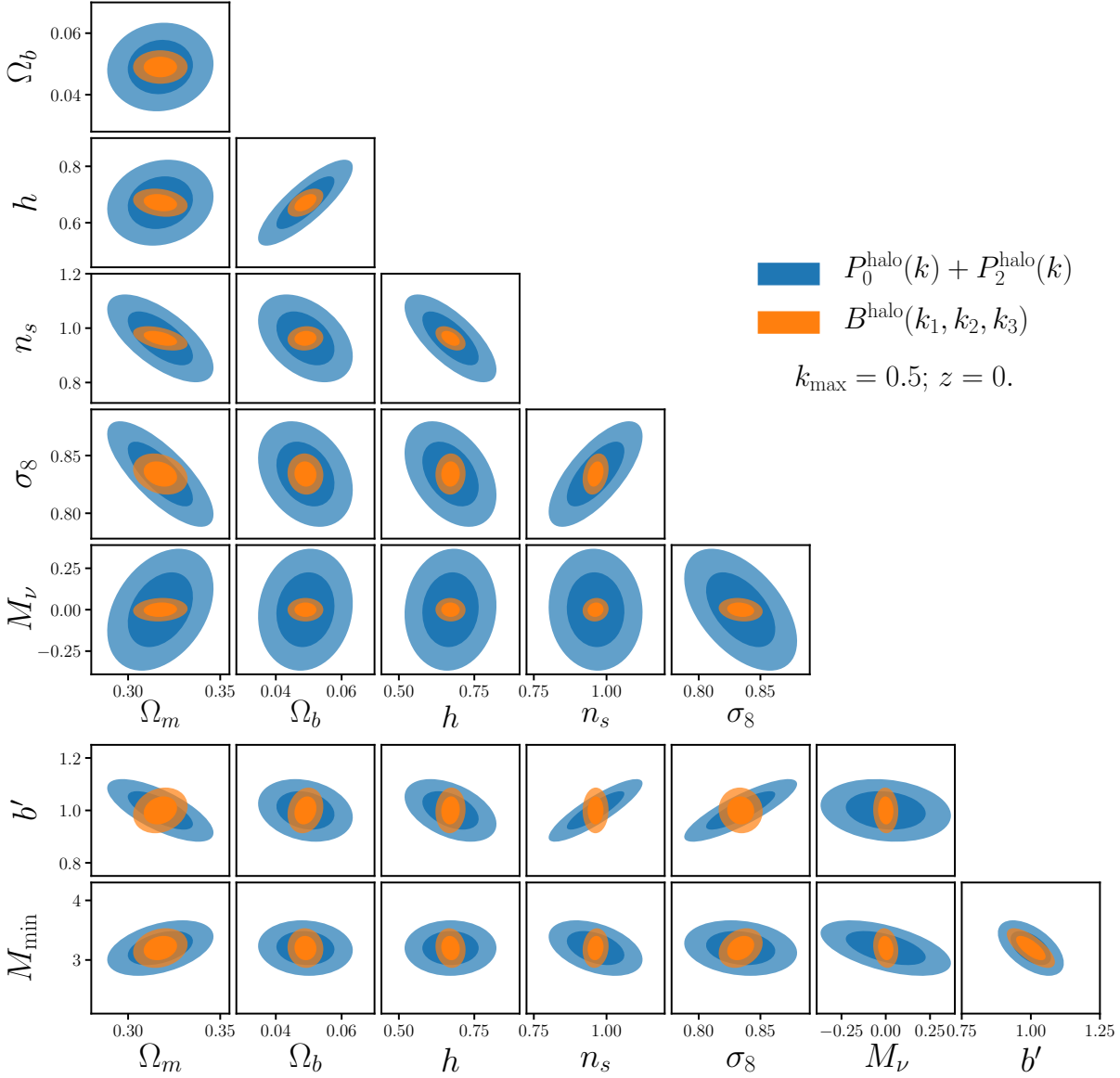
**Figure 7.** Derivatives of the redshift-space halo bispectrum,  $d \log B_0 / d\theta$  with respect to  $\Omega_m$ ,  $\Omega_b$ ,  $h$ ,  $n_s$ ,  $\sigma_8$ , and  $M_\nu$  as a function of all 1898 triangle configurations with  $k_1, k_2, k_3 \leq 0.5 \, h/\text{Mpc}$  (top to bottom panels). We estimate the derivatives at the fiducial parameters using  $N_{\text{deriv}} = 1,500$   $N$ -body realizations from the Quijote suite. The configurations above are ordered in the same way as Figures 3 and 5 (Appendix A). By using  $N$ -body simulations for the derivatives, we rely on fewer assumptions and approximations than analytic methods (*i.e.* perturbation theory). Furthermore, with their accuracy on small scales, these derivatives enable us to quantify, for the first time, the information content of the bispectrum in the nonlinear regime.

where  $\mathcal{L}$  is the likelihood, the Fisher matrix for the bispectrum can be written as

$$F_{ij} = \frac{1}{2} \text{Tr} \left[ \mathbf{C}^{-1} \frac{\partial \mathbf{C}}{\partial \theta_i} \mathbf{C}^{-1} \frac{\partial \mathbf{C}}{\partial \theta_j} + \mathbf{C}^{-1} \left( \frac{\partial B_0}{\partial \theta_i} \frac{\partial B_0^T}{\partial \theta_j} + \frac{\partial B_0^T}{\partial \theta_i} \frac{\partial B_0}{\partial \theta_j} \right) \right]. \quad (9)$$

Since we assume that the  $B_0$  likelihood is Gaussian, including the first term in Eq. 9 runs the risk of incorrectly including information from the covariance already included in the mean (Carron 2013).





**Figure 8.** Fisher matrix constraints for  $M_\nu$  and other cosmological parameters for the redshift-space halo bispectrum monopole (orange). We include Fisher parameter constraints for the redshift-space halo power spectrum monopole and quadrupole in blue for comparison. The contours mark the 68% and 95% confidence intervals. We set  $k_{\text{max}} = 0.5 \, h/\text{Mpc}$  for both power spectrum and bispectrum. We include in our forecasts  $b'$  and  $M_{\text{min}}$ , a free amplitude scaling factor and halo mass limit, respectively. They serve as a simplistic bias model and we marginalize over them so that our constraints do not include extra constraining power from the difference in bias/number density in the different Quijote cosmologies. The bispectrum *substantially* improves constraints on all of the cosmological parameters over the power spectrum. Constraints on  $\Omega_m$ ,  $\Omega_b$ ,  $h$ ,  $n_s$ , and  $\sigma_8$  improve by factors of 1.9, 2.6, 3.1, 3.6, and 2.6, respectively. For  $M_\nu$ , the bispectrum improves  $\sigma_{M_\nu}$  from 0.2968 to 0.0572 eV — over a factor of  $\sim 5$  improvement over the power spectrum.

We, therefore, conservatively neglect the first term and calculate the Fisher matrix as,

$$F_{ij} = \frac{1}{2} \text{Tr} \left[ \mathbf{C}^{-1} \left( \frac{\partial B_0}{\partial \theta_i} \frac{\partial B_0}{\partial \theta_j}^T + \frac{\partial B_0}{\partial \theta_i}^T \frac{\partial B_0}{\partial \theta_j} \right) \right], \quad (10)$$

directly with  $\mathbf{C}$  and  $\partial B_0/\partial \theta_i$  along each cosmological parameter from the Quijote simulations.

For  $\Omega_m$ ,  $\Omega_b$ ,  $h$ ,  $n_s$ , and  $\sigma_8$ , we estimate

$$\frac{\partial B_0}{\partial \theta_i} \approx \frac{B_0(\theta_i^+) - B_0(\theta_i^-)}{\theta_i^+ - \theta_i^-}, \quad (11)$$

where  $B_0(\theta_i^+)$  and  $B_0(\theta_i^-)$  are the average bispectrum of the 1,500 realizations at  $\theta_i^+$  and  $\theta_i^-$ , respectively. Meanwhile, for  $M_\nu$ , where the fiducial value is 0.0 eV and we cannot have negative  $M_\nu$ , we use the Quijote simulations at  $M_\nu^+$ ,  $M_\nu^{++}$ ,  $M_\nu^{+++} = 0.1, 0.2, 0.4$  eV (Table 1) to estimate

$$\frac{\partial B_0}{\partial M_\nu} \approx \frac{-21B_0(\theta_{\text{fid}}^{\text{ZA}}) + 32B_0(M_\nu^+) - 12B_0(M_\nu^{++}) + B_0(M_\nu^{+++})}{1.2}, \quad (12)$$

which provides a  $\mathcal{O}(\delta M_\nu^2)$  order approximation. Since the simulations at  $M_\nu^+$ ,  $M_\nu^{++}$ , and  $M_\nu^{+++}$  are generated from Zel'dovich initial conditions, we use simulations at the fiducial cosmology also generated from Zel'dovich initial conditions ( $\theta_{\text{fid}}^{\text{ZA}}$ ). By using these  $N$ -body simulations, instead of analytic methods (*e.g.* perturbation theory), we exploit the accuracy of the simulations in the nonlinear regime and rely on fewer assumptions and approximations. In fact, these  $N$ -body simulation estimated derivatives are key in enabling us to quantify, for the first time, the total information content of the redshift-space bispectrum in the non-linear regime. We present the bispectrum derivatives for all triangle configurations in Figure 7 and discuss subtleties of the bispectrum derivatives and tests of convergence and stability in Appendix B.

We present the constraints on  $M_\nu$  and other cosmological parameters  $\{\Omega_m, \Omega_b, h, n_s, \sigma_8\}$  derived from the redshift-space halo bispectrum Fisher matrix for  $k_{\text{max}} = 0.5 h/\text{Mpc}$  in Figure 8. We include Fisher constraints for the redshift-space halo power spectrum monopole and quadrupole with the same  $k_{\text{max}}$  for comparison (blue). The shaded contours mark the 68% and 95% confidence intervals. We include in our Fisher constraints the following nuisance parameters:  $b'$  and  $M_{\text{min}}$ .  $b'$  is a scaling factor on the bispectrum amplitude, analogous to linear bias. And  $M_{\text{min}}$  is the halo mass limit, which we choose as a nuisance parameter to address the difference in the number densities among the Quijote cosmologies, which impacts the derivatives  $\partial B_0/\partial \theta_i$ . For instance, the  $\sigma_8^+$  and  $\sigma_8^-$  cosmologies have halo  $\bar{n} = 1.586 \times 10^{-4}$  and  $1.528 \times 10^{-4} h^3 \text{Mpc}^{-3}$ . These parameters serve as a simplistic bias model and by marginalizing over them we aim to ensure that our Fisher constraints do not include extra constraining power from the difference in bias or number density.  $b'$  is a multiplicative factor so  $\partial B_0/\partial b' = B_0$ .  $\partial B_0/\partial M_{\text{min}}$ , we estimate numerically using  $B_0$  evaluated at  $M_{\text{min}}^+ = 3.3 \times 10^{13} h^{-1} M_\odot$  and  $M_{\text{min}}^- = 3.1 \times 10^{13} h^{-1} M_\odot$  with all other parameters set to the fiducial value.

*The bispectrum substantially improves constraints on all parameters over the power spectrum.* For  $k_{\text{max}} = 0.5 h/\text{Mpc}$ , the bispectrum tightens the marginalized  $1\sigma$  constraints,  $\sigma_\theta$ , of  $\Omega_m$ ,  $\Omega_b$ ,  $h$ ,  $n_s$ , and  $\sigma_8$  by factors of  $\sim 1.9, 2.6, 3.1, 3.6, 2.6$  over the power spectrum. *For  $M_\nu$ , the bispectrum improves*

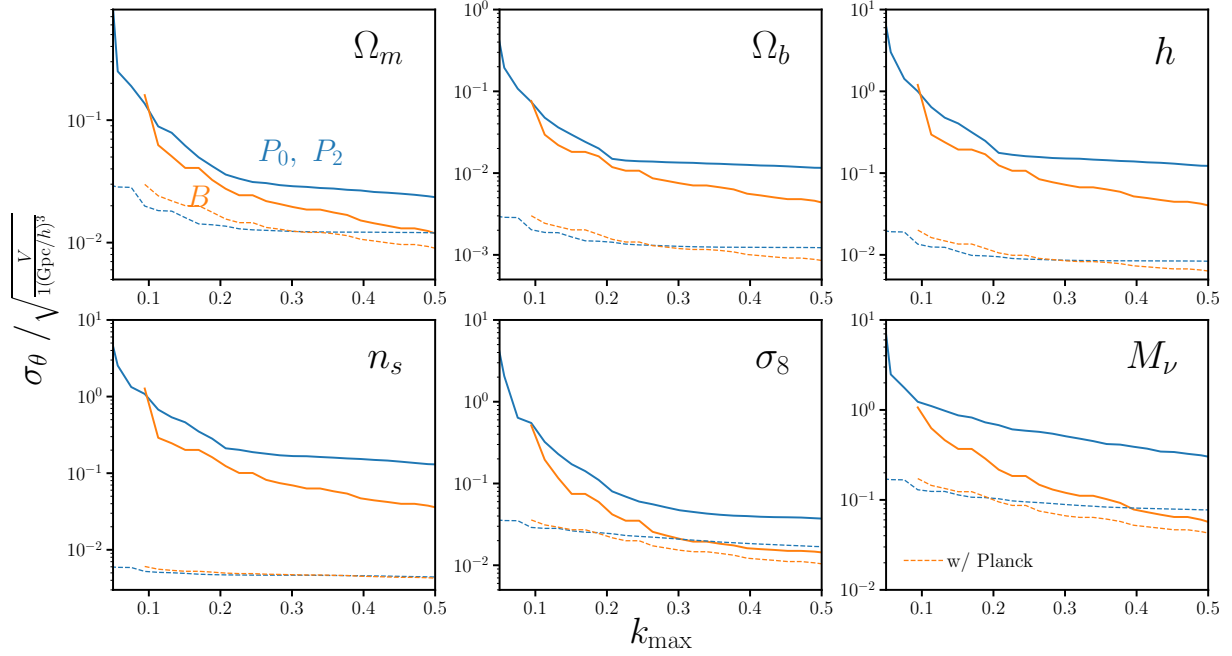
**Table 2.** Marginalized Fisher parameter constraints from the redshift-space halo power spectrum (top) and bispectrum (bottom) for different  $k_{\text{max}}$ . We list constraints for cosmological parameters  $M_\nu$ ,  $\Omega_m$ ,  $\Omega_b$ ,  $h$ ,  $n_s$ , and  $\sigma_8$  as well as nuisance parameters  $b'$  and  $M_{\text{min}}$ .

	$k_{\text{max}}$ ( $h/\text{Mpc}$ )	$M_\nu$ (eV)	$\Omega_m$	$\Omega_b$	$h$	$n_s$	$\sigma_8$	$b'$	$M_{\text{min}}$ ( $10^{13} h^{-1} M_\odot$ )
		0.0	0.3175	0.049	0.6711	0.9624	0.834	1.	3.2
$P_{\ell=0,2}$	0.2	$\pm 0.711$	$\pm 0.037$	$\pm 0.015$	$\pm 0.179$	$\pm 0.213$	$\pm 0.085$	$\pm 0.223$	$\pm 1.827$
	0.3	$\pm 0.509$	$\pm 0.029$	$\pm 0.013$	$\pm 0.151$	$\pm 0.167$	$\pm 0.047$	$\pm 0.117$	$\pm 0.725$
	0.4	$\pm 0.383$	$\pm 0.026$	$\pm 0.012$	$\pm 0.137$	$\pm 0.152$	$\pm 0.039$	$\pm 0.102$	$\pm 0.470$
	0.5	$\pm \mathbf{0.297}$	$\pm 0.023$	$\pm 0.012$	$\pm 0.123$	$\pm 0.130$	$\pm 0.037$	$\pm 0.096$	$\pm 0.376$
$B_0$	0.2	$\pm 0.256$	$\pm 0.029$	$\pm 0.014$	$\pm 0.142$	$\pm 0.138$	$\pm 0.049$	$\pm 0.264$	$\pm 1.299$
	0.3	$\pm 0.130$	$\pm 0.021$	$\pm 0.008$	$\pm 0.077$	$\pm 0.074$	$\pm 0.023$	$\pm 0.143$	$\pm 0.654$
	0.4	$\pm 0.078$	$\pm 0.015$	$\pm 0.006$	$\pm 0.052$	$\pm 0.047$	$\pm 0.016$	$\pm 0.088$	$\pm 0.367$
	0.5	$\pm \mathbf{0.057}$	$\pm 0.012$	$\pm 0.004$	$\pm 0.040$	$\pm 0.036$	$\pm 0.014$	$\pm 0.070$	$\pm 0.268$

the constraint from  $\sigma_{M_\nu} = 0.2968$  to  $0.0572$  eV — over a factor of 5 improvement over the power spectrum. This  $\sigma_{M_\nu} = 0.0572$  eV constraint is from the *bispectrum alone* and only for a  $1h^{-1}\text{Gpc}$  box. For a larger volume,  $V$ ,  $\sigma_\theta$  scales as  $\propto 1/\sqrt{V}$ . We list the precise marginalized Fisher parameter constraints of both cosmological and nuisance parameters for  $P_\ell$  and  $B_0$  in Table 2.

Even below  $k_{\text{max}} < 0.5$   $h/\text{Mpc}$ , the bispectrum significantly improves cosmological parameter constraints. We compare  $\sigma_\theta$  of  $\Omega_m$ ,  $\Omega_b$ ,  $h$ ,  $n_s$ ,  $\sigma_8$ , and  $M_\nu$  as a function of  $k_{\text{max}}$  for  $B_0$  (orange) and  $P_{\ell=0,2}$  (blue) in Figure 9. We only include the  $k_{\text{max}}$  range where the Fisher forecast is well defined — *i.e.* more data bins than the number of parameters:  $k_{\text{max}} > 5 k_f \approx 0.03$   $h/\text{Mpc}$  for  $P_\ell$  and  $k_{\text{max}} > 12 k_f \approx 0.075$   $h/\text{Mpc}$  for  $B_0$ . Figure 9 reveals that the improvement of the bispectrum  $\sigma_\theta$  over the power spectrum  $\sigma_\theta$  is larger at higher  $k_{\text{max}}$ . Although limited by the  $k_{\text{max}}$  range, the figure suggests that on large scales ( $k_{\text{max}} \lesssim 0.1$   $h/\text{Mpc}$ )  $\sigma_\theta$  of  $P_\ell$  crosses over  $\sigma_\theta$  of  $B_0$  so  $P_\ell$  has more constraining power than  $B_0$ , as expected on linear scales. On slightly smaller scales,  $k_{\text{max}} = 0.2$   $h/\text{Mpc}$ , we find that the bispectrum improves  $\sigma_\theta$  by factors of  $\sim 1.3, 1.1, 1.3, 1.5, 1.7$ , and  $2.8$  for  $\Omega_m$ ,  $\Omega_b$ ,  $h$ ,  $n_s$ ,  $\sigma_8$ , and  $M_\nu$  respectively.

Our forecasts demonstrate that the bispectrum has significant constraining power beyond the power spectrum in the weakly nonlinear regime ( $k > 0.1$   $h/\text{Mpc}$ ). This constraining power comes from the bispectrum breaking degeneracies among the cosmological and nuisance parameters. This is evident when we compare the unmarginalized constraints from  $P_\ell$  and  $B_0$ :  $1/\sqrt{F_{ii}}$  where  $F_{ii}$  is a diagonal element of the Fisher matrix. For  $k < 0.4$   $h/\text{Mpc}$ , the unmarginalized constraints from  $P_\ell$  are tighter than those from  $B_0$ . Yet, once we marginalize the constraints over the other parameters, the  $P_\ell$  constraints are degraded and the  $B_0$  constraints are tighter for  $k > 0.1$   $h/\text{Mpc}$ . The derivatives,  $\partial B_0/\partial \theta_i$ , also shed light on how  $B_0$  breaks parameter degeneracies. The parameter degeneracies in the  $P_\ell$  forecasts of Figure 8 are consistent with similarities in the shape and scale dependence of  $P_\ell$



**Figure 9.** Marginalized  $1\sigma$  constraints,  $\sigma_\theta$ , of the cosmological parameters  $\Omega_m$ ,  $\Omega_b$ ,  $h$ ,  $n_s$ ,  $\sigma_8$ , and  $M_\nu$  as a function of  $k_{\max}$  for the redshift-space halo bispectrum (orange) and power spectrum ( $\ell = 0, 2$ ; blue). The constraints are marginalized over the nuisance parameters  $b'$  and  $M_{\min}$  in our forecast (Section 4.2). We impose  $k_{\max} > 5 k_f$  for  $P_\ell$  and  $k_{\max} > 12 k_f$  for  $B_0$ , the  $k$  ranges where we have more data bins than number of parameters. We also include  $\sigma_\theta$  constraints with *Planck* priors (dotted). Even below  $k_{\max} < 0.5 h/\text{Mpc}$ , the bispectrum significantly improves cosmological parameter constraints. The improvement, however, is larger for higher  $k_{\max}$ . At  $k_{\max} = 0.2 h/\text{Mpc}$ , the bispectrum improves constraints on  $\Omega_m$ ,  $\Omega_b$ ,  $h$ ,  $n_s$ ,  $\sigma_8$ , and  $M_\nu$  by factors of  $\sim 1.3$ ,  $1.1$ ,  $1.3$ ,  $1.5$ ,  $1.7$ , and  $2.8$  over the power spectrum. Even with *Planck* priors,  $B_0$  significantly improves  $\sigma_\theta$  for  $k_{\max} \gtrsim 0.2 h/\text{Mpc}$ . While the constraining power of  $P_\ell$  saturates at  $k_{\max} = 0.2 h/\text{Mpc}$ , the constraining power of  $B_0$  continues to increase out to  $k_{\max} = 0.5 h/\text{Mpc}$ .

derivatives  $\partial P_0/\partial\theta$  and  $\partial P_2/\partial\theta$ . On the other hand, the  $B_0$  derivatives with respect to the parameters have significant different scale and triangle shape dependencies, as seen in Figure 7.

By exploiting the massive number of  $N$ -body simulations of the Quijote suite, we present for the first time the total information content of the full redshift-space bispectrum beyond the linear regime. The information content of the bispectrum has previously been examined using perturbation theory. Previous works, for instance, measure the signal-to-noise ratio (SNR) of the bispectrum derived from covariance matrices estimated using perturbation theory (Sefusatti & Scoccimarro 2005; Sefusatti et al. 2006; Chan & Blot 2017). More recently, Chan & Blot (2017), used covariance matrices that include non-Gaussian contributions calibrated with  $N$ -body simulations to find that the cumulative SNR of the halo bispectrum is  $\sim 30\%$  of the SNR of the halo power spectrum at  $k_{\max} \sim 0.1 h/\text{Mpc}$  and increases to  $\sim 40\%$  at  $k_{\max} \sim 0.35 h/\text{Mpc}$ . While these simple SNR measurements cannot be easily compared to Fisher analysis (Repp et al. 2015; Blot et al. 2016), we note that they are roughly consistent with the unmarginalized constraints, which loosely represent the SNRs of the derivatives. Also, when we measure the halo power spectrum and bispectrum SNRs using our

covariance matrices (Figure 6), we find a relation between the SNRs consistent with Chan & Blot (2017). Beyond the  $k$  range explored by Chan & Blot (2017),  $k_{\text{max}} > 0.35 \text{ h/Mpc}$ , we find that the SNR of  $B_0$  continues to increase at higher  $k_{\text{max}}$  in contrast to the  $P_0$  SNR, which saturates at  $k_{\text{max}} \sim 0.1 \text{ h/Mpc}$ . At  $k_{\text{max}} = 0.75 \text{ h/Mpc}$ , the largest  $k$  we measure the  $B_0$ , the SNR of  $B_0$  is  $\sim 75\%$  of the SNR of  $P_0$ .

Beyond these signal-to-noise calculations, a number of previous works have quantified the information content of the bispectrum (Scoccimarro et al. 2004; Sefusatti et al. 2006; Sefusatti & Komatsu 2007; Song et al. 2015; Tellarini et al. 2016; Yamauchi et al. 2017; Karagiannis et al. 2018; Yankelevich & Porciani 2019; Chudaykin & Ivanov 2019). While most of these works fix most cosmological parameters and focus solely on forecasting constraints of primordial non-Gaussianity and bias parameters, Sefusatti et al. (2006), Yankelevich & Porciani (2019), and Chudaykin & Ivanov (2019) provide bispectrum forecasts for full sets of cosmological parameters. In Sefusatti et al. (2006), they present likelihood analysis forecasts for  $\omega_d$ ,  $\omega_b$ ,  $\Omega_\Lambda$ ,  $n_s$ ,  $A_s$ ,  $w$ ,  $\tau$ . For  $\Lambda$ CDM, with fixed bias parameters, and  $k_{\text{max}} = 0.3 \text{ h/Mpc}$ , they find constraints on  $\Omega_m$ ,  $\Omega_b$ ,  $h$ ,  $n_s$ , and  $\sigma_8$  from WMAP,  $P_0$ , and  $B_0$  is  $\sim 1.5$  times tighter than constraints from WMAP and  $P_0$ . In comparison, for  $k_{\text{max}} = 0.3 \text{ h/Mpc}$  our  $B_0$  constraints are tighter than  $P_0$  constraints by factors of 1.4, 1.7, 2.0, 2.3, 2.0, and 3.9. Both Sefusatti et al. (2006) and our analysis find significantly tighter constraints with the bispectrum. They however include the WMAP likelihood in their forecast and use perturbation theory models, which is limited to larger scales than used here (Scoccimarro et al. 1998, 1999; Sefusatti et al. 2010; Pollack et al. 2012; Gil-Marín et al. 2014; Lazanu et al. 2016; Eggemeier et al. 2019).

Yankelevich & Porciani (2019) present Fisher forecasts for  $\Omega_{\text{cdm}}$ ,  $\Omega_b$ ,  $h$ ,  $n_s$ ,  $A_s$ ,  $w_0$ , and  $w_0$  for a Euclid-like survey (Laureijs et al. 2011) in 14 non-overlapping redshift bins over  $0.65 < z < 2.05$ . They use the full redshift-space bispectrum, rather than just the monopole, and a more sophisticated bias expansion than Sefusatti et al. (2006) but use a perturbation theory bispectrum model, which consequently limit their forecast to  $k_{\text{max}} = 0.15 \text{ h/Mpc}$ . They find similar constraining power on cosmological parameters from  $B$  alone as  $P$ . They also find that combining the bispectrum with the power spectrum only moderately improves parameter constraints because posterior correlations are similar for  $P$  and  $B$ . While this seemingly conflicts with the results we present, there are significant differences between our forecasts. For instance, they forecast the Euclid survey (*i.e.*  $z > 0.7$ ), while our forecasts are for  $z = 0$ . They also forecast the *galaxy*  $P$  and  $B$  and marginalizes over 56 nuisance parameters (14  $z$  bins each with 3 bias parameters and 1 RSD parameter). They also neglect non-Gaussian contributions to the  $B$  covariance matrix, which play a significant role on small scales (Chan & Blot 2017) and may impact the constraints. Despite differences, Yankelevich & Porciani (2019) find that the constraining power of  $B$  relative to  $P_\ell$  increases for higher  $k_{\text{max}}$ , consistent with our forecasts as a function of  $k_{\text{max}}$  (Figure 9). Also, consistent with their results, for  $k_{\text{max}} = 0.15 \text{ h/Mpc}$ , we find similar posterior correlations between the  $P_\ell$  and  $B_0$  constraints. At  $k_{\text{max}} = 0.5 \text{ h/Mpc}$ , however, we find the posterior correlations are no longer similar, which contribute to the constraining power of  $B_0$  (Figure 8).

Finally, Chudaykin & Ivanov (2019) present power spectrum and bispectrum forecasts for  $\omega_{\text{cdm}}$ ,  $\omega_b$ ,  $h$ ,  $n_s$ ,  $A_s$ ,  $n_s$ , and  $M_\nu$  of a Euclid-like survey in 8 non-overlapping redshift bins over  $0.5 < z < 2.1$ .

They use a one-loop perturbation theory model for the redshift-space power spectrum multipoles ( $\ell = 0, 2, 4$ ) and a tree-level bispectrum monopole model. Also, rather than imposing a  $k_{\text{max}}$  cutoff to restrict their forecasts to scales where their perturbation theory model can be trusted, they use a theoretical error covariance model approach from [Baldauf et al. \(2016\)](#). They find  $\sim 1.4, 1.5, 1.2, 1.5$ , and  $1.3$  times tighter constraints from  $P_\ell$  and  $B_0$  than from  $P_\ell$  alone. For  $M_\nu$ , they find a factor of  $1.4$  improvement, from  $0.038$  eV to  $0.028$  eV. Overall, [Chudaykin & Ivanov \(2019\)](#) find more modest improvements from including  $B_0$  than the improvements we find in our  $k_{\text{max}} = 0.5$   $h/\text{Mpc}$   $B_0$  constraints over the  $P_\ell$  constraints.

Among the differences from our forecast, most notably, [Chudaykin & Ivanov \(2019\)](#) marginalize over 64 parameters of their bias model (4 bias and 4 counterterm normalization parameters at each redshift bin). They also include the Alcock-Paczynski (AP) effect for  $P_\ell$ , which significantly improve the  $P_\ell$  parameter constraints (*e.g.* tightens  $M_\nu$  constraints by  $\sim 30\%$ ). They however, do not include AP effects in  $B_0$ . Furthermore, [Chudaykin & Ivanov \(2019\)](#) use theoretical error covariance to quantify the uncertainty of their perturbation theory models. Because they use the tree-level perturbation theory model for  $B_0$ , theoretical errors for  $B_0$  quickly dominate at  $k \gtrsim 0.1$   $h/\text{Mpc}$ , where the one- and two-loop contribute significantly (*e.g.* [Lazanu & Liguori 2018](#)). Hence, their forecasts do not include the constraining power on nonlinear scales. Also different from our analysis, [Chudaykin & Ivanov \(2019\)](#) use an Markov-Chain Monte-Carlo (MCMC) approach, which derives more accurate parameter constraints than our Fisher approach. They, however, neglect non-Gaussian contributions to both the  $P_\ell$  and  $B_0$  covariance matrices and do not include the covariance between  $P_\ell$  and  $B_0$  for their joint constraints. We find that neglecting the covariance between  $P_\ell$  and  $B_0$  overestimates constraints by  $\sim 20\%$  for  $k_{\text{max}} = 0.2$   $h/\text{Mpc}$ . Nonetheless, [Chudaykin & Ivanov \(2019\)](#) and our results both find significant improves in cosmological parameter constraints from including  $B_0$ . In fact, taking the theoretical errors into account, the improvements from  $B_0$  they find are loosely consistent with our results at  $k_{\text{max}} \sim 0.2$   $h/\text{Mpc}$ .

Various differences between our forecast and previous work prevent more thorough comparisons. However, crucial aspects of our simulation based approach distinguish our forecasts from other works. We present the first bispectrum forecasts for a full set of cosmological parameters using bispectrum measured entirely from  $N$ -body simulations. By using the simulations, we go beyond perturbation theory models and accurately model the redshift-space bispectrum to the nonlinear regime. Furthermore, by exploiting the immense number of simulations, we accurately estimate the full high-dimensional covariance matrix of the bispectrum. With these advantages, we present the first forecast of cosmological parameters from the bispectrum down to nonlinear scales and demonstrate the constraining power of the bispectrum for  $M_\nu$ . Below, we underline a few caveats of our forecasts.

Our forecasts are derived from Fisher matrices. Such forecasts make the assumption that the posterior is approximately Gaussian and, as a result, they underestimate the constraints for posteriors that are highly non-elliptical or asymmetric ([Wolz et al. 2012](#)). Fisher matrices also rely on the stability, and in our case also convergence, of numerical derivatives. We examine the stability of the  $B_0$  derivatives with respect to  $M_\nu$  by comparing the derivatives computed using  $N$ -body simulations at three different sets of cosmologies: (1)  $\{\theta_{\text{fid}}^{\text{ZA}}, M_\nu^+, M_\nu^{++}, M_\nu^{+++}\}$  (Eq. 12), (2)  $\{\theta_{\text{fid}}^{\text{ZA}}, M_\nu^+, M_\nu^{++}\}$ ,



and (3)  $\{\theta_{\text{fid}}^{\text{ZA}}, M_\nu^+\}$  (see Appendix B; Figure 13). The derivatives computed using the different set of cosmologies, do not impact the  $\Omega_m$ ,  $\Omega_b$ ,  $h$ ,  $n_s$ , and  $\sigma_8$  constraints. They do however affect the  $M_\nu$  constraints; but because  $P_\ell$  and  $B_0$  derivatives are affected by the same factor, the relative improvement of the  $B_0$   $M_\nu$  constraint over the  $P_\ell$  constraints is *not* impacted. In addition to the stability, because we use  $N$ -body simulations, we test whether the convergence of our covariance matrix and derivatives impact our forecasts by varying the number of simulations used to estimate them:  $N_{\text{cov}}$  and  $N_{\text{deriv}}$ , respectively. For  $N_{\text{cov}}$ , we find  $< 5\%$  variation in the Fisher matrix elements,  $F_{ij}$ , for  $N_{\text{cov}} > 5000$  and  $< 1\%$  variation in  $\sigma_\theta$  for  $N_{\text{cov}} > 12000$ . For  $N_{\text{deriv}}$ , we find  $< 5\%$  variation in the  $F_{ij}$  elements and  $< 5\%$  variation in  $\sigma_\theta$  for  $N_{\text{deriv}} > 1200$ . Since our constraints vary by  $< 10\%$  for sufficient  $N_{\text{cov}}$  and  $N_{\text{deriv}}$ , the convergence of the covariance matrix and derivatives do not impact our forecasts to the accuracy level of Fisher forecasting. We refer readers to Appendix B for a more details on the robustness of our results to the stability of the derivatives and convergence.

We argue that the constraining power of the bispectrum and its improvement over the power spectrum come from breaking degeneracies among the cosmological parameters. However, numerical noise can impact our forecasts when we invert the Fisher matrix. Since *Planck* constrain  $\{\Omega_m, \Omega_b, h, n_s, \sigma_8\}$  tighter than either  $P_\ell$  or  $B_0$  alone, the elements of the *Planck* prior matrix are larger than the elements of  $P_\ell$  and  $B_0$  Fisher matrices. Including *Planck* priors (*i.e.* adding the prior matrix to the Fisher matrix) increases the numerical stability of the matrix inversion. It also reveals whether the bispectrum still improves parameter constraints once we include CMB constraints. With *Planck* priors and  $P_\ell$  to  $k_{\text{max}} = 0.5$   $h/\text{Mpc}$ , we derive the following constraints:  $\sigma_{\Omega_m} = 0.0120$ ,  $\sigma_{\Omega_b} = 0.0012$ ,  $\sigma_h = 0.0084$ ,  $\sigma_{n_s} = 0.0044$ , and  $\sigma_{\sigma_8} = 0.0169$ . Including *Planck* priors expectedly tighten the constraints from  $P_\ell$ . Meanwhile, with *Planck* priors and  $B_0$  to  $k_{\text{max}} = 0.5$   $h/\text{Mpc}$ , we get  $\sigma_{\Omega_m} = 0.0090$ ,  $\sigma_{\Omega_b} = 0.0009$ ,  $\sigma_h = 0.0064$ ,  $\sigma_{n_s} = 0.0043$ , and  $\sigma_{\sigma_8} = 0.0104$ , 1.3, 1.4, 1.3, 1.0, and 1.6 times tighter constraints. For  $M_\nu$ ,  $\sigma_{M_\nu} = 0.0773$  eV for  $P_\ell$  and  $\sigma_{M_\nu} = 0.0430$  eV for  $B_0$ , a factor of 1.8 improvement. Since we find substantial improvements in parameter constraints with the *Planck* prior, the improvement from  $B_0$  are numerically robust. Furthermore,  $\sigma_\theta$  as a function of  $k_{\text{max}}$  with *Planck* priors reveal that while the constraining power of  $P_\ell$  saturates at  $k_{\text{max}} = 0.2$   $h/\text{Mpc}$ , the constraining power of  $B_0$  continues to increase out to  $k_{\text{max}} = 0.5$   $h/\text{Mpc}$  (dotted; Figure 9).

Our forecasts are derived using the power spectrum and bispectrum in *periodic boxes*. We do not consider a realistic geometry or radial selection function of actual observations from galaxy surveys. A realistic selection function will smooth out the triangle configuration dependence and consequently degrade the constraining power of the bispectrum. In Sefusatti & Scoccimarro (2005), for instance, they find that the signal-to-noise of the bispectrum is significantly reduced once survey geometry is included in their forecast. We also do not account for super-sample covariance (*e.g.* Hamilton et al. 2006; Sefusatti et al. 2006; Takada & Hu 2013; Li et al. 2018), which may also degrade constraints. Survey geometry and super-sample covariance, however, also degrades the signal-to-noise of their power spectrum forecasts. Hence, with the substantial improvement in the  $M_\nu$  constraints of the bispectrum, even with survey geometry we expect the bispectrum will significantly improve  $M_\nu$  constraints over the power spectrum.

We include the nuisance parameter  $M_{\min}$  in our forecasts to address the difference in halo bias and number densities among the Quijote cosmologies. Although we marginalize over  $M_{\min}$ , this may not fully account for the extra information from  $\bar{n}$  and nonlinear bias leaking into the derivatives. To test this, we include extra nuisance parameters,  $\{A_{\text{SN}}, B_{\text{SN}}, b_2, \gamma_2\}$ , and examine their impact on our forecasts.  $A_{\text{SN}}$  and  $B_{\text{SN}}$  are multiplicative factors of the first and second terms of Eq. 4, which we include to account for any  $\bar{n}$  dependence that may be introduced from the shotnoise correction (Eq. 4).  $b_2$  and  $\gamma_2$  are the quadratic bias and nonlocal bias parameters (Chan et al. 2012; Sheth et al. 2013) to account for information from nonlinear bias. Marginalizing over  $b'$ ,  $M_{\min}$ ,  $A_{\text{SN}}$ ,  $B_{\text{SN}}$ ,  $b_2$  and  $\gamma_2$ , we obtain the following constraints for  $B_0$  with  $k_{\max} = 0.5 \text{ h/Mpc}$ :  $\sigma_{\Omega_m} = 0.0129$ ,  $\sigma_{\Omega_b} = 0.0044$ ,  $\sigma_h = 0.0404$ ,  $\sigma_{n_s} = 0.0456$ ,  $\sigma_{\sigma_8} = 0.0228$ , and  $\sigma_{M_\nu} = 0.0575$ . While constraints on  $n_s$  and  $\sigma_8$  are broadened from our fiducial forecasts, by 27% and 60%, the other parameters, especially  $M_\nu$ , are not significantly impacted by marginalizing over the extra nuisance parameters. As another test, we calculate derivatives using halo catalogs from Quijote  $\theta^-$  and  $\theta^+$  cosmologies with fixed  $\bar{n}$ . We similarly find no significant impact on the  $B_0$  parameter constraints. Forecasts using additional nuisance parameters and with fixed  $\bar{n}$  derivatives, both support the robustness of our forecast. Yet these tests do not ensure that our forecast entirely marginalizes over halo bias.

In this paper, we focus on the halo bispectrum and power spectrum. However, constraints on  $M_\nu$  will ultimately be derived from the distribution of galaxies. Besides the cosmological parameters, bias and nuisance parameters that allow us to marginalize over galaxy bias need to be incorporated to forecast  $M_\nu$  and other cosmological parameter constraints for the galaxy bispectrum. Although we include a *naïve* bias model through  $b'$  and  $M_{\min}$ , and even  $b_2$  and  $\gamma_2$  in our tests, this is insufficient to describe how galaxies trace matter. A more realistic bias model such as a halo occupation distribution (HOD) model involve extra parameters that describe the distribution of central and satellite galaxies in halos (*e.g.* Zheng et al. 2005; Leauthaud et al. 2012; Tinker et al. 2013; Zentner et al. 2016; Vakili & Hahn 2019). We, therefore, refrain from a more exhaustive investigation of the impact of halo bias on our results and focus quantifying the constraining power of the galaxy bispectrum in the next paper of this series.

Marginalizing over galaxy bias parameters, will likely reduce the constraining power at high  $k$ . Hand et al. (2017), for instance, with their 13 parameter model only find a 15-30% improvement in  $f\sigma_8$  when they extend their power spectrum multipole analysis from  $0.2 \text{ h/Mpc}$  to  $0.4 \text{ h/Mpc}$ . However, jointly analyzing power spectrum and bispectrum will help break parameter degeneracies and improve constraints on cosmological parameters (Sefusatti et al. 2006; Yankelevich & Porciani 2019; Chudaykin & Ivanov 2019). We again emphasize that the constraints we present is for a  $1h^{-1}\text{Gpc}$  box, a substantially smaller volume than upcoming surveys. Thus, even if the constraining power at high  $k$  is reduced, our forecasts suggest that the bispectrum offers significant improvements over the power spectrum, especially for constraining  $M_\nu$ .

## 5. SUMMARY

A precise measurement of  $M_\nu$  can distinguish between the ‘normal’ and ‘inverted’ neutrino mass hierarchy scenarios and reveal physics beyond the Standard Model of particle physics. The total neutrino mass, through impact the expansion history and the growth of cosmic structure in the

Universe, can be measured using cosmological observables (*e.g.* CMB and large-scale structure). In fact, cosmological probes have the potential to more precisely measure  $M_\nu$  than current and upcoming laboratory experiments. The current tightest cosmological  $M_\nu$  constraints come from combining CMB data with other cosmological probes. The degeneracy between  $M_\nu$  and  $\tau$ , the optical depth of reionization, however, is a major bottleneck for  $M_\nu$  constraints from CMB data that will *not* be addressed by upcoming ground-based CMB experiments. Measuring the imprint of neutrinos on the 3D clustering of galaxies provides a promising alternative to tightly constrain  $M_\nu$ , especially with the unprecedented cosmic volumes that will be mapped by upcoming surveys.

Recent developments in simulation based emulation methods have improved the accuracy of theoretical predictions beyond linear scales and address the challenges of unlocking the information content in nonlinear clustering to constrain  $M_\nu$ . Yet, for power spectrum analyses, the strong degeneracy between the imprints of  $M_\nu$  and  $\sigma_8$  poses a serious limitation for constraining  $M_\nu$ . Information in the nonlinear regime, however, cascade from the power spectrum to higher-order statistics such as the bispectrum. Previous works have demonstrated that the bispectrum has comparable signal-to-noise as the power spectrum on nonlinear scales (Sefusatti & Scoccimarro 2005; Chan & Blot 2017) and that it improves constraints on cosmological parameters (Sefusatti et al. 2006; Yankelevich & Porciani 2019; Chudaykin & Ivanov 2019). No work to date has quantified the total information content and constraining power of the full bispectrum down to nonlinear scales.

In this work, we examined the effect of massive neutrinos on the redshift-space halo bispectrum using  $\sim 23,000$   $N$ -body simulations with massive neutrinos from the HADES and Quijote suite. Using  $N$ -body simulations and with such a massive number of them, we directly address key challenges of accurately modeling the bispectrum in the nonlinear regime and estimating its high dimensional covariance matrix. More specifically,

- We examine the imprint of  $M_\nu$  and  $\sigma_8$  on the redshift-space halo bispectrum,  $B_0$ , using two sets of HADES simulations. One with massive neutrinos,  $M_\nu = 0.0, 0.06, 0.1$ , and  $0.15$  eV; the other with  $M_\nu = 0.0$  eV and matching  $\sigma_8^c$ .  $M_\nu$  and  $\sigma_8$  leave distinct imprints on  $B_0$  with significantly different scale and triangle shape dependencies. Thus, we demonstrate that  $B_0$  helps break the  $M_\nu$ - $\sigma_8$  degeneracy found in the power spectrum.
- We quantify the total information content of  $B_0$  using a Fisher matrix forecast of  $\{\Omega_m, \Omega_b, h, n_s, \sigma_8, M_\nu\}$ . With bispectrum measured from 22,000  $N$ -body simulations from the Quijote suite, we are able to derive stable and converged covariance matrix and derivatives for our forecast. Furthermore, our simulation based approach allows us to extend beyond perturbation theory and use all 1898 triangle configurations of the bispectrum down to  $k_{\max} = 0.5$   $h/\text{Mpc}$ .
- For  $k_{\max}=0.5$   $h/\text{Mpc}$ , the bispectrum produces  $\Omega_m$ ,  $\Omega_b$ ,  $h$ ,  $n_s$ , and  $\sigma_8$  constraints 1.9, 2.6, 3.1, 3.6, and 2.6 times tighter than the power spectrum. For  $M_\nu$ , we derive  $1\sigma$  constraint of 0.0572 eV — over 5 times tighter than the power spectrum. Even with priors from *Planck*, the bispectrum improves  $M_\nu$  constraints by a factor of 1.8. These constraints are derived for a  $(1\ h^{-1}\text{Gpc})^3$  box, a substantially smaller volume than upcoming surveys.

While our results clearly showcase the advantages of the bispectrum for more precisely constraining  $M_\nu$ , as well as the cosmological parameters, a number of assumptions go into our forecast. Fisher matrix forecasts assume that the posterior is approximately Gaussian and, as a result, they overestimate the constraints for posteriors that are highly non-elliptical or asymmetric. Furthermore, our forecasts are derived using the power spectrum and bispectrum in periodic boxes, instead of a realistic geometry or radial selection function of galaxy surveys, which degrades the constraining power. Since we use periodic boundary conditions, we also do not account for super-sample covariance, which in practice results from a non-trivial survey geometry. Lastly, our forecast focuses on the halo bispectrum, while the *galaxy* bispectrum is what is actually measured from observations. We marginalize over a simplistic bias model through  $b'$  and  $M_{\text{min}}$ ; however, such a model is insufficient to flexibly describe how galaxies trace matter. A more realistic bias model involve extra parameters that describe the distribution of central and satellite galaxies. In the next paper of the series, we include HOD parameters in our forecasts and quantify the full information content of the redshift-space galaxy bispectrum.

The constraints from our forecasts are derived for a  $(1\ h^{-1}\text{Gpc})^3$  volume. These constraints roughly scale as  $\propto 1/\sqrt{V}$  with volume. Upcoming spectroscopic galaxy surveys will map out vastly larger cosmic volumes: PFS  $\sim 9\ h^{-3}\text{Gpc}^3$  and DESI  $\sim 50\ h^{-3}\text{Gpc}^3$  (Takada et al. 2014; Collaboration et al. 2016). Euclid and WFIRST, space-based surveys, will expand these volumes to higher redshifts. A key science goal in these surveys will be constraining  $M_\nu$ . Scaling our 0.0572 eV  $M_\nu$  constraint from the bispectrum to the survey volumes gives 0.0191 eV and 0.0081 eV  $1\sigma$  constraints for PFS and DESI. With such precision, we would detect  $M_\nu$  by  $> 3\sigma$  and distinguish between the normal and inverted neutrino mass hierarchy scenarios by  $> 2\sigma$  from the bispectrum alone. Such naive projections, however, should be taken with more than a grain of salt. Nevertheless, the substantial improvement we find in constraints with the bispectrum over the power spectrum, strongly advocate for analyzing future surveys with more than the power spectrum. Our results demonstrate the potential of the bispectrum to tightly constrain  $M_\nu$  with unprecedented precision.

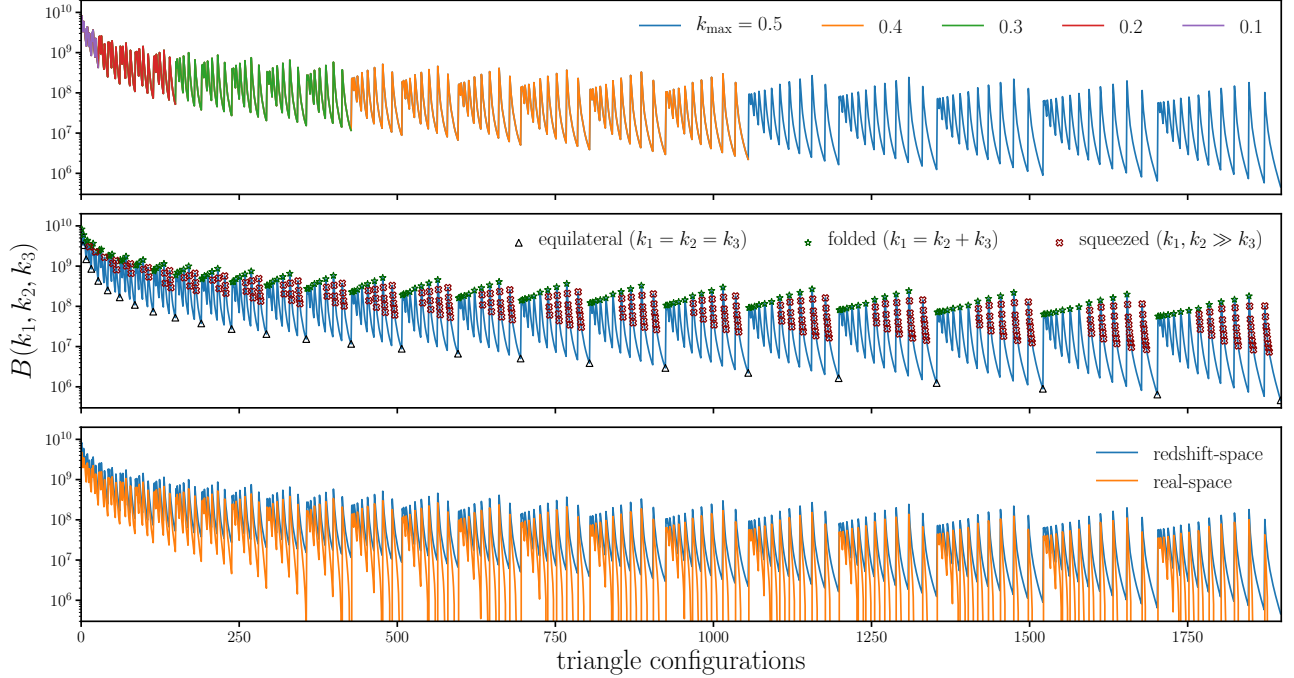
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## APPENDIX

### A. REDSHIFT-SPACE BISPECTRUM

In this work, we present how the redshift-space bispectrum helps break degeneracies among cosmological parameters, especially between  $M_\nu$  and  $\sigma_8$ , and improve constraints on  $M_\nu$  as well as the other cosmological parameters. We measure the redshift-space bispectrum monopole,  $\hat{B}_0$ , of the HADES



**Figure 10.** *Top:* We highlight the triangle configurations of the bispectrum that fall within:  $k_{\max} < 0.1$  (purple), 0.2 (red), 0.3 (green), 0.4 (orange), and 0.5  $h/\text{Mpc}$  (blue). The number of triangle configurations increase significantly at higher  $k_{\max}$ , which contributes to the significant constraining power on nonlinear scales. *Center:* We mark triangle configurations with three different shapes: equilateral (black triangle), folded (green star), and squeezed (red cross) triangles. Equilateral triangles have  $k_1 = k_2 = k_3$ ; folded triangles have  $k_1 = k_2 + k_3$ ; squeezed triangles have  $k_1, k_2 \gg k_3$ , above we mark configurations with  $k_1, k_2 > 2.4k_3$  as squeezed. *Bottom:* Comparison of the real (orange) and redshift-space (blue) halo bispectrum for fiducial Quijote simulations. The redshift-space  $\hat{B}_0$  has a significantly higher overall amplitude than the real-space  $\hat{B}_0$ , with a shape and  $k_{\max}$  dependence. As a result, the redshift-space  $\hat{B}_0$  has a higher signal-to-noise and more constraining power than the real-space  $\hat{B}_0$ .

and Quijote simulations (Section 2) using an FFT based estimator similar to the ones in Sefusatti (2005), Scoccimarro (2015), and Sefusatti et al. (2016) (see Section 3 for details). We use triangle configurations defined by  $k_1$ ,  $k_2$ , and  $k_3$  bins of width  $\Delta k = 3k_f = 0.01885 h/\text{Mpc}$ . For  $k_{\max} = 0.5 h/\text{Mpc}$  we have  $\hat{B}_0$  for 1898 triangle configurations. For  $k_{\max} = 0.4, 0.3, 0.2$ , and  $0.1 h/\text{Mpc}$ , we have 1056, 428, 150, and 28 triangle configurations respectively. The number of triangle configurations sharply increases for higher  $k_{\max}$ . This contributes to the significant increase in constraining power as we go to increasingly nonlinear scales (Figure 9). We highlight the triangle configurations at different  $k_{\max}$  for the  $\hat{B}_0$  of the Quijote simulation at the fiducial cosmology in the top panel of Figure 10. We mark the configurations within  $k_{\max} = 0.1, 0.2, 0.3, 0.4$ , and  $0.5 h/\text{Mpc}$  in purple, red, green, orange, and blue respectively.

In Figure 10, as well as in Figures 3, 5, 7, we order the triangle configurations by looping through  $k_3$  in the inner most loop and  $k_1$  in the outer most loop such that  $k_1 \geq k_2 \geq k_3$ . This ordering is different from the ordering in Gil-Marín et al. (2017), which loops through  $k_3$  in the inner most loop



and  $k_1$  in the outer most increasing loop but with  $k_1 \leq k_2 \leq k_3$ . As the top panel demonstrates, our ordering clearly reflects the  $k_{\max}$  range of the configurations. Furthermore, the repeated configuration sequences in our ordering cycles through all available triangle configurations for a given  $k_1$ . In the center panel of Figure 10, we mark equilateral (black triangle), folded (green star), and squeezed (red cross) triangle configurations. Equilateral triangles have  $k_1 = k_2 = k_3$ , folded triangles have  $k_1 = k_2 + k_3$ , and squeezed triangles have  $k_1, k_2 \gg k_3$  —  $k_1, k_2 > 2.4k_3$  in Figure 10. These shapes correspond to the three vertices of the bispectrum shape plots of Figures 2 and 4.

Lastly, in the bottom panel of Figure 10 we compare the redshift-space bispectrum (blue) to the real-space bispectrum of the same Quijote simulations at the fiducial cosmology (orange). Overall, the redshift-space  $\hat{B}_0$  has a higher amplitude than the real-space  $\hat{B}_0$  with a significant triangle shape dependence: equilateral triangles have the largest difference in amplitude while folded triangles have the smallest. The relative amplitude difference also depends significantly on  $k_{\max}$  where the difference is larger for configurations with higher  $k_1$ . With its higher amplitude, the redshift-space  $\hat{B}_0$  has a higher signal-to-noise and more constraining power than the real-space  $\hat{B}_0$ .

## B. FISHER FORECASTS USING $N$ -BODY SIMULATIONS

The two key elements in calculating the Fisher matrices in our bispectrum forecasts are the bispectrum covariance matrix ( $\mathbf{C}$ ; Figure 6) and the derivatives of the bispectrum along the cosmological and nuisance parameters ( $\partial B_0/\partial\theta$ ; Figure 7). We compute both these elements directly using the  $N$ -body simulations of the Quijote suite (Section 2). This exploits the accuracy of  $N$ -body simulations in the nonlinear regime and allows us to accurately quantify the constraining power of the bispectrum beyond perturbation theory models. However, we must ensure that both  $\mathbf{C}$  and  $\partial B_0/\partial\theta$  have converged and that numerical effects do not introduce any biases that impact our results. Below, we test the convergence of  $\mathbf{C}$  and  $\partial B_0/\partial\theta$  and discuss some of the subtleties and caveats of our  $\partial B_0/\partial\theta$  calculations.

We use 15,000 Quijote  $N$ -body simulations at the fiducial cosmology to estimate  $\mathbf{C}$ . This is a *significantly* larger number of simulations than any previous bispectrum analyses; however, we also consider 1898 triangle configurations out to  $k_{\max} = 0.5 h/\text{Mpc}$ . For reference, Gil-Marín et al. (2017) recently used 2048 simulations to estimate the covariance matrix of the bispectrum with 825 configurations. To check the convergence of covariance matrix, we vary the number of simulations used to estimate  $\mathbf{C}$ ,  $N_{\text{cov}}$ , and determine whether this significantly impacts the elements of the Fisher matrix,  $F_{ij}$ , or the final marginalized Fisher parameter constraints,  $\sigma_\theta$ . In the left panel of Figure 11, we present the ratio between  $F_{ij}(N_{\text{cov}})$ ,  $F_{ij}$  derived from  $\mathbf{C}$  calculated with  $N_{\text{cov}}$  simulations, and  $F_{ij}(N_{\text{cov}} = 15,000)$  for all 36 elements of the Fisher matrix. We shade  $\pm 5\%$  deviations in the ratios for reference. The  $F_{ij}$  elements vary by  $\lesssim 5\%$  for  $N_{\text{cov}} > 5000$  and  $\lesssim 1\%$  for  $N_{\text{cov}} > 10,000$ . Next, we present ratio between  $\sigma_\theta(N_{\text{cov}})$ , the marginalized  $1\sigma$  constraints for  $\{\Omega_m, \Omega_b, h, n_s, \sigma_8, M_\nu\}$  derived from  $\mathbf{C}$  calculated with  $N_{\text{cov}}$  simulations, and  $\sigma_\theta(N_{\text{cov}} = 15,000)$  in the left panel of Figure 12. The constraints vary by  $\lesssim 5\%$  for  $N_{\text{cov}} > 5000$  and  $\lesssim 1\%$  for  $N_{\text{cov}} > 12000$ . Hence,  $N_{\text{cov}} = 15,000$  is sufficient to accurately estimate  $\mathbf{C}$  and its convergence does not impact our forecasts.

We estimate  $\partial B_0/\partial\theta$  using  $N_{\text{deriv}} = 1,500$   $N$ -body simulations at 13 different cosmologies listed in Table 1. To check the convergence of  $\partial B_0/\partial\theta$  and its impact on our results, we examine the



ratio between  $F_{ij}(N_{\text{deriv}})$ , the Fisher matrix element derived from  $\partial B_0/\partial\theta$  calculated with  $N_{\text{deriv}}$  simulations, and  $F_{ij}(N_{\text{deriv}} = 1,500)$  for all 36 elements of the Fisher matrix in the right panel of Figure 11. For  $N_{\text{deriv}} > 1000$ ,  $F_{ij}$  elements vary by  $\lesssim 5\%$ . Next, we present the ratio between  $\sigma_\theta(N_{\text{deriv}})$ , the marginalized  $1\sigma$  constraints for  $\{\Omega_m, \Omega_b, h, n_s, \sigma_8, M_\nu\}$  derived from  $\partial B_0/\partial\theta$  calculated with  $N_{\text{deriv}}$  simulations, and  $\sigma_\theta(N_{\text{deriv}} = 1,500)$  in the right panel of Figure 12. Unlike  $\sigma_\theta(N_{\text{cov}})$ ,  $\sigma_\theta(N_{\text{deriv}})$  depend significantly on  $\theta$ . For instance,  $\sigma_{\sigma_8}$  and  $\sigma_{\Omega_m}$  vary by  $\lesssim 10\%$  for  $N_{\text{deriv}} > 600$  and  $\lesssim 1\%$  for  $N_{\text{deriv}} > 1200$ .  $\sigma_\theta$  for the other parameter vary significantly more. Nonetheless, for  $N_{\text{deriv}} > 800$  and  $1200$  they vary by  $\lesssim 10$  and  $5\%$ , respectively.

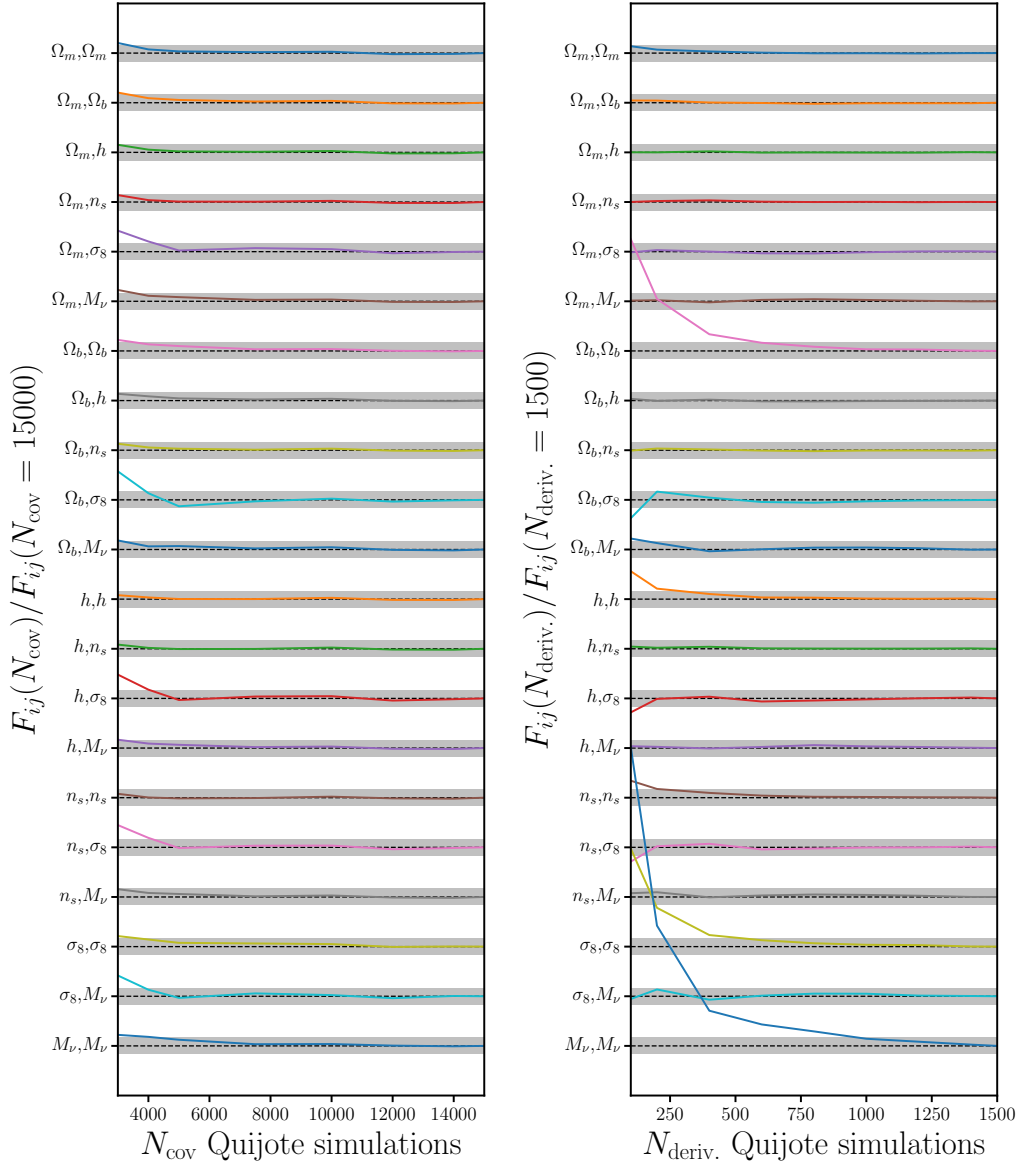
For  $\Omega_m, \Omega_b, h, n_s, \sigma_8$ , and also the nuisance parameter  $M_{\text{lim}}$  we estimate  $\partial B_0/\partial\theta$  using a centered difference approximation (Eq. 11). However, for  $M_\nu$  we cannot have values below  $0.0$  eV and, thus, cannot estimate the derivative with the same method. If we use the forward difference approximation,

$$\frac{\partial \bar{B}_0}{\partial M_\nu} \approx \frac{\bar{B}_0(\theta_{\text{fid}}^{\text{ZA}} + \delta M_\nu) - \bar{B}_0(M_\nu^{\text{fid}})}{\delta M_\nu}, \quad (\text{B1})$$

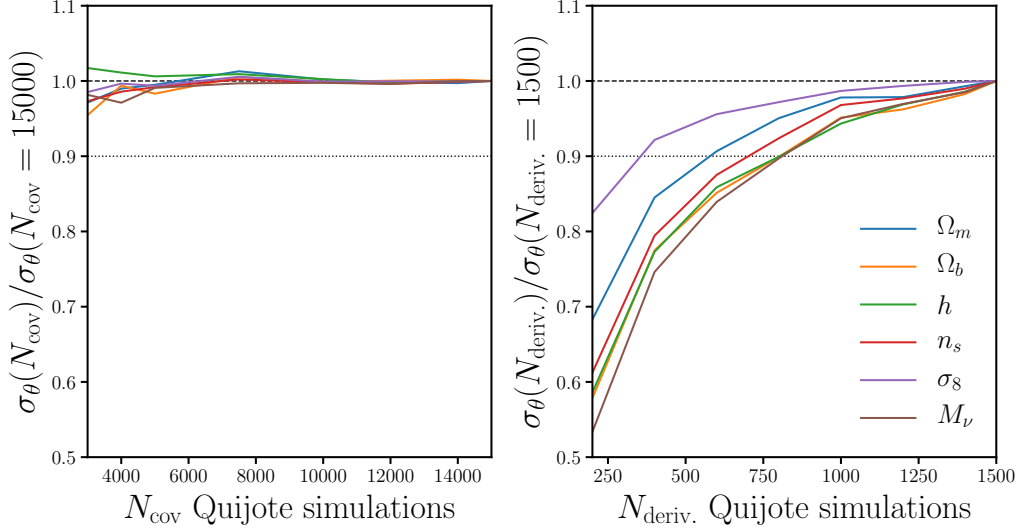
the error goes as  $\mathcal{O}(\delta M_\nu)$ . Instead, we use a finite difference approximation with the Quijote simulations at  $M_\nu^+, M_\nu^{++}, M_\nu^{+++}$ , and the fiducial cosmology, a  $\mathcal{O}(\delta M_\nu^2)$  order approximation (Eq. 12). We can also use a finite difference approximation with simulations at  $M_\nu^+, M_\nu^{++}$ , and the fiducial cosmology. We examine the stability of  $\partial \log B_0/\partial M_\nu$  (right) and  $\partial \log P_\ell/\partial M_\nu$  (left) by comparing the derivatives computing using simulations at  $\{\theta_{\text{fid}}^{\text{ZA}}, M_\nu^+, M_\nu^{++}, M_\nu^{+++}\}$  (blue),  $\{\theta_{\text{fid}}^{\text{ZA}}, M_\nu^+, M_\nu^{++}\}$  (orange), and  $\{\theta_{\text{fid}}^{\text{ZA}}, M_\nu^+\}$  (green) in Figure 13. The three  $\partial \log B_0/\partial M_\nu$  approximations differ from one another by  $\sim 10\%$  with Eq. 12 producing the largest estimate for both  $P_\ell$  and  $B_0$ . If we use the  $\{\theta_{\text{fid}}^{\text{ZA}}, M_\nu^+, M_\nu^{++}\}$  derivative and  $\{\theta_{\text{fid}}^{\text{ZA}}, M_\nu^+\}$  derivative instead of Eq. 12 for our Fisher forecasts, the marginalized constraint on  $M_\nu$  for  $k_{\text{max}} = 0.5$  h/Mpc increases to  $0.390$  and  $0.682$  eV for  $P_\ell$  and  $0.0754$  and  $0.1354$  eV for  $B_0$ . Compared to our  $\sigma_{M_\nu} = 0.0572$  eV  $B_0$  forecast, these correspond to a  $\sim 30$  and  $130\%$  relative increase. While the derivative estimated from  $\{\theta_{\text{fid}}^{\text{ZA}}, M_\nu^+\}$  significantly impact the forecasts, we emphasize that this is a  $\mathcal{O}(\delta M_\nu)$  approximation, unlike the other  $\mathcal{O}(\delta M_\nu^2)$  approximations. In fact, the  $\{\theta_{\text{fid}}^{\text{ZA}}, M_\nu^+\}$  derivatives are better  $\mathcal{O}(\delta M_\nu^2)$  estimates for  $\partial P_\ell/\partial M_\nu$  and  $\partial B_0/\partial M_\nu$  at  $0.05$  eV. When we compare the derivative of the linear theory power spectrum,  $P^{(\text{LT})}$ , we find that  $\partial P^{(\text{LT})}/\partial M_\nu$  at  $0.0$  is larger than at  $0.05$  eV. Hence, the differences between the  $\{\theta_{\text{fid}}^{\text{ZA}}, M_\nu^+\}$  derivatives and the other derivatives are not solely due to numerical stability. Moreover, because the discrepancies in the derivative propagate similarly to both  $P_\ell$  and  $B_0$  forecasts, the relative improvement of  $B_0$  over  $P_\ell$  remains roughly the same. Hence, we conclude that the derivatives with respect to  $M_\nu$  are sufficiently stable and robust for our Fisher forecasts.

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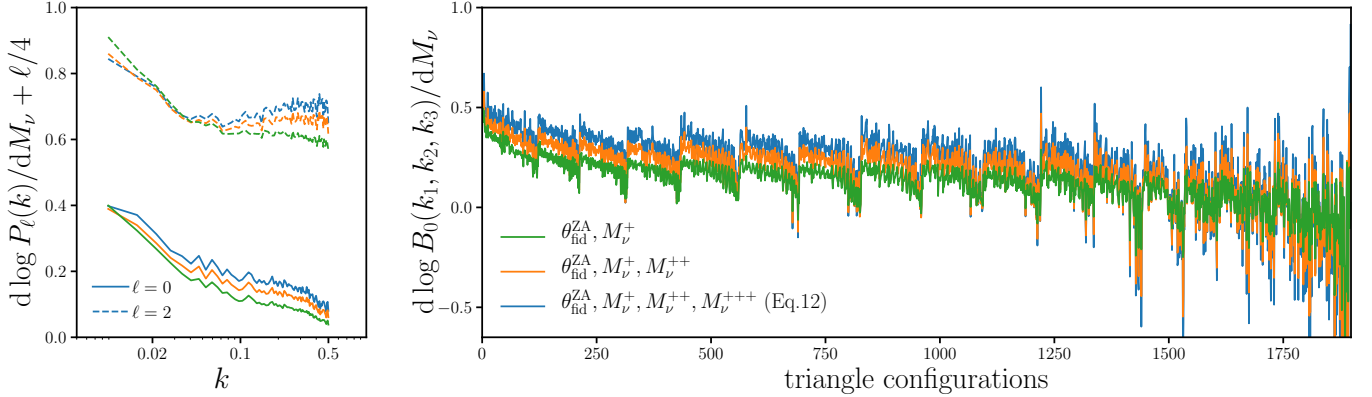
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**Figure 11.** *Left:* The convergence of all 36 Fisher matrix elements,  $F_{ij}$ , as a function of  $N_{\text{cov}}$ , the number of  $N$ -body simulations used to estimate the covariance matrix,  $\mathbf{C}$ . We present the ratio between  $F_{ij}(N_{\text{cov}})$ ,  $F_{ij}$  derived from  $\mathbf{C}$  calculated with  $N_{\text{cov}}$  simulations, and  $F_{ij}(N_{\text{cov}} = 15,000)$ . We mark  $\pm 5\%$  deviations in the ratios with the shaded regions for reference. All  $F_{ij}$  elements vary by  $\lesssim 5\%$  for  $N_{\text{cov}} > 5000$  and  $\lesssim 1\%$  for  $N_{\text{cov}} > 10,000$ . *Right:* The convergence of  $F_{ij}$  as a function of  $N_{\text{deriv}}$ , the number of  $N$ -body simulations used to estimate the derivatives,  $\partial B_0/\partial\theta$ . We plot the ratios between  $F_{ij}(N_{\text{deriv}})$  and  $F_{ij}(N_{\text{deriv}} = 1,500)$ . All  $F_{ij}$  elements vary by  $\lesssim 5\%$  for  $N_{\text{deriv}} > 1000$ . Hence,  $N_{\text{cov}} = 15,000$  and  $N_{\text{deriv}} = 1,500$  are sufficient and the convergence of  $\mathbf{C}$  and  $\partial B_0/\partial\theta$  does not impact  $F_{ij}$ .



**Figure 12.** *Left:* The convergence of the marginalized  $1\sigma$  constraints for  $\{\Omega_m, \Omega_b, h, n_s, \sigma_8, M_\nu\}, \sigma_\theta$ , as a function of  $N_{\text{cov}}$ , the number of Quijote simulations used to estimate  $\mathbf{C}$ . We plot the ratio between  $\sigma_\theta(N_{\text{cov}})$ , derived from  $\mathbf{C}$  with  $N_{\text{cov}}$  simulations, and  $\sigma_\theta(N_{\text{cov}} = 15,000)$ .  $\sigma_\theta$  vary by  $\lesssim 5$  and  $1\%$  for  $N_{\text{cov}} > 5000$  and  $12,000$ , respectively. *Right:* The convergence of  $\sigma_\theta$  as a function of  $N_{\text{deriv.}}$ , the number of simulated used to estimate  $\partial B_0/\partial M_\nu$ . We plot the ratio between  $\sigma_\theta(N_{\text{deriv.}})$ , derived from  $\partial B/\partial M_\nu$  with  $N_{\text{deriv.}}$  simulations, and  $\sigma_\theta(N_{\text{deriv.}} = 1,500)$ . Although  $\sigma_\theta(N_{\text{deriv.}})/\sigma_\theta(N_{\text{deriv.}} = 1,500)$  vary among the parameters, the ratio vary by  $\lesssim 10$  and  $5\%$  for  $N_{\text{deriv.}} > 800$  and  $1200$ , respectively. Hence, *we have a sufficient number of simulations to estimate  $\mathbf{C}$  and the derivatives of the bipsectrum and our forecasts are robust to their convergence.*



**Figure 13.** Comparison of  $\partial \log B_0(k_1, k_2, k_3)/\partial M_\nu$  (right) and  $\partial \log P_\ell(k)/\partial M_\nu$  (left), computed using Quijote simulations at  $\{\theta_{\text{fid}}^{\text{ZA}}, M_\nu^+, M_\nu^{++}, M_\nu^{+++}\}$  (blue),  $\{\theta_{\text{fid}}^{\text{ZA}}, M_\nu^+, M_\nu^{++}\}$  (orange), and  $\{\theta_{\text{fid}}^{\text{ZA}}, M_\nu^+\}$  (green). The derivative approximations differ from one another by  $\sim 10\%$  with Eq. 12 producing the largest estimate for both  $P_\ell$  and  $B_0$ . Using the  $\{\theta_{\text{fid}}^{\text{ZA}}, M_\nu^+, M_\nu^{++}\}$  derivatives instead of the Eq. 12 derivatives, increases the marginalized constraint on  $M_\nu$  by  $\sim 30\%$ . However, the differences in the derivatives propagate similarly to the  $P_\ell$  and  $B_0$  forecasts so the relative improvement of  $B_0$  over  $P_\ell$  remains the same. Hence, we conclude that *the derivatives with respect to  $M_\nu$  are sufficiently stable and robust for our Fisher forecasts.*

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