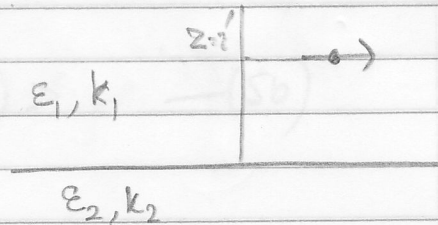


The Equivalent of a horizontal electric dipole above a half-space



$$\vec{J} = \hat{x} J_x \delta(\vec{r}-\vec{r}') = \hat{x} J_x \delta(x-x') \delta(z-z') \delta(y-y') \quad (1)$$

The corresponding spectral representation of the source is,

$$\vec{J} = \hat{x} J_x \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \delta(x-x') \delta(y-y') \delta(z-z') e^{jk_x x} e^{jk_y y} dk_x dk_y \quad (2)$$

$$= \hat{x} J_x e^{-jk_x(x-x')} e^{-jk_y(y-y')} \delta(z-z') \quad (3)$$

See
Dr.
Neel's
Note

Here $\vec{r} = \hat{x}x + \hat{y}y$

$$\vec{r}' = \hat{x}x' + \hat{y}y'$$

$$\vec{p} = (x-x')\hat{x} + (y-y')\hat{y}$$

$$p = \sqrt{(x-x')^2 + (y-y')^2} \quad \vec{p} = p \sin \phi \hat{x} + p \cos \phi \hat{y}$$

Also $\vec{k}_p = k_x \hat{x} + k_y \hat{y}$
 $k_x = k_p \cos \theta$
 $k_y = k_p \sin \theta$

$$\text{So } k_x(x-x') + k_y(y-y') = k_p p \cos(\theta - \phi) \quad (4)$$

We'll use this result when inverse F.T ing in the polar coordinates.

$$\vec{E}(\vec{r}) = < \vec{G}^{EJ} ; \vec{J}(\vec{r}') > \quad \text{--- (5a)}$$

$$\vec{H}(\vec{r}) = < \vec{G}^{HJ} ; \vec{J}(\vec{r}') > \quad \text{--- (5b)}$$

Transverse components:

$$\vec{E}_t = \hat{x} V^{TM} ; \quad \vec{H}_t = \hat{y} I^{TM} \quad \text{--- (6)}$$

The subsequent T.L. equations:

$$\frac{dV^{TM}}{dz} = -jk_2 Z^{TM} I^{TM} \quad \text{--- (7a)}$$

$$\frac{dI^{TM}}{dz} = -jk_2 Y^{TM} V^{TM} + i^{TM} \quad \text{--- (7b)}$$

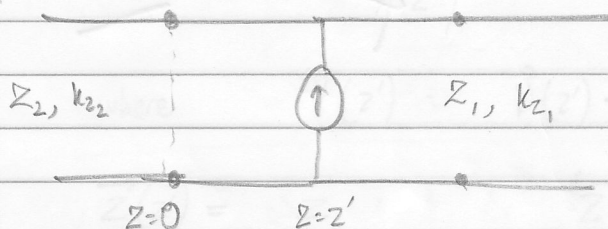
$$Z^{TM} = \frac{1}{Y^{TM}} = \frac{k_2}{\omega \epsilon_0 \epsilon_r}$$

where

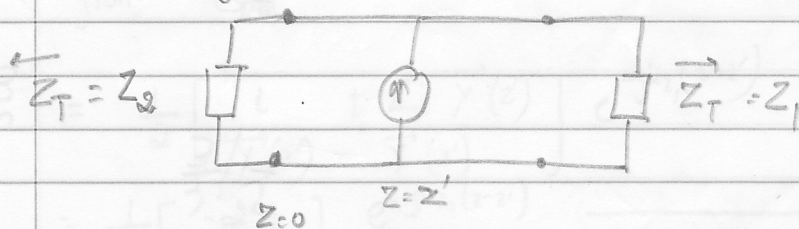
$$i^{TM} = -\frac{\hat{z} \times \vec{J}_x}{J_x} \quad \text{--- (8)}$$

$$-i^{TM} + i^{TM}$$

$$\frac{dH_t}{dz} = -\hat{z} \times \vec{J}_x$$



Representing the above TL with terminal impedances,

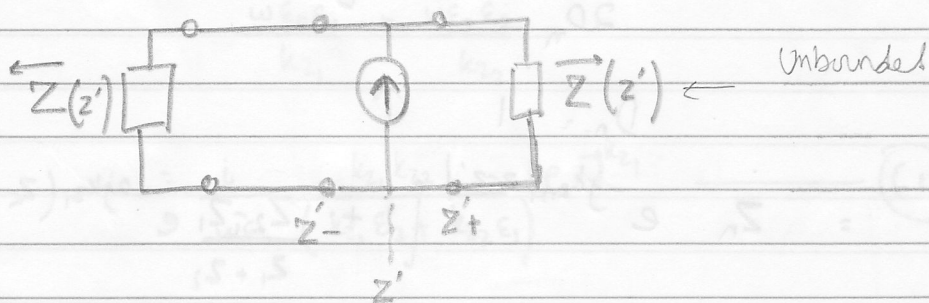


The voltage and current can be expressed as.
for $z > z'$

$$V(z, z') = \frac{1}{2} [V(z', z') + Z I(z', z')] e^{-jk_2(z-z')} \quad \text{--- (9a)}$$

$$I(z, z') = \frac{1}{2} [V(z', z') + Z I(z', z')] e^{-jk_2(z-z')} \quad \text{--- (9b)}$$

The voltage and current $V(z', z')$ and $I(z', z')$ are found by ~~trans~~ modifying the TL.



$$V(z', z') = \frac{i}{\overleftarrow{Y}(z')} \quad ; \quad I(z', z') = \frac{i \overrightarrow{Y}(z')}{\overleftarrow{Y}(z')}$$

$$\text{where } \overrightarrow{Y}(z') = \overrightarrow{Y}(z') + \overleftarrow{Y}(z') = \frac{1}{\overrightarrow{Z}(z')} + \frac{1}{\overleftarrow{Z}(z')}$$

$$\overrightarrow{Z}(z') = Z_1, \quad \overleftarrow{Z}(z') = Z_1 \frac{Z_2 + jZ_1 \tan k_{21} z'}{Z_1 + jZ_2 \tan k_{21} z'}$$

from (9a)

$$\begin{aligned} V(z, z') &= \frac{1}{2} \left[\frac{i}{\overleftarrow{Y}(z')} + \frac{Z_1 i \overrightarrow{Y}(z')}{\overleftarrow{Y}(z')} \right] e^{-jk_{21}(z-z')} \\ &= \frac{1}{2} \left[\frac{2i}{\overleftarrow{Y}(z')} \right] e^{-jk_{21}(z-z')} \quad \text{--- (10)} \end{aligned}$$

$$I(z, z') = \frac{Y_1 i}{\overleftarrow{Y}(z')} e^{-jk_{21}(z-z')}$$

when the source lies on the interface

$$z = z' = 0$$

$$\bar{Z}(0) = Z_2$$

$$z > 0 \quad V(z, 0) = \frac{i}{\frac{1}{Z_1} + \frac{1}{Z_2}} e^{-jk_2 z}$$

$$= \frac{i}{\frac{\omega \epsilon_0 \epsilon_1}{k_{z1}} + \frac{\omega \epsilon_0 \epsilon_2}{k_{z2}}} e^{-jk_2 z}$$

$$= \frac{i}{\omega \epsilon_0} \frac{k_{z1} k_{z2}}{(k_{z1} \epsilon_2 + k_{z2} \epsilon_1)} e^{-jk_2 z} \quad (11)$$

The current is expressed as:

$$I(z, 0) = Y_1 V(z, 0) = \frac{i}{\omega \epsilon_0} \frac{\omega \epsilon_0 \epsilon_1}{k_{z1}} \frac{k_{z1} k_{z2}}{(k_{z1} \epsilon_2 + k_{z2} \epsilon_1)} e^{-jk_2 z}$$

$$I(z, 0) = \frac{i}{\omega \epsilon_0} \frac{k_{z2} \epsilon_1}{(k_{z1} \epsilon_2 + k_{z2} \epsilon_1)} e^{-jk_2 z} \quad (12)$$

$$= \frac{i}{2} \left(1 + \frac{k_{z2} \epsilon_1 - k_{z1} \epsilon_2}{k_{z2} \epsilon_1 + k_{z2} \epsilon_1} \right) e^{-jk_2 z}$$

This form is equivalent to Michalski's forms.