

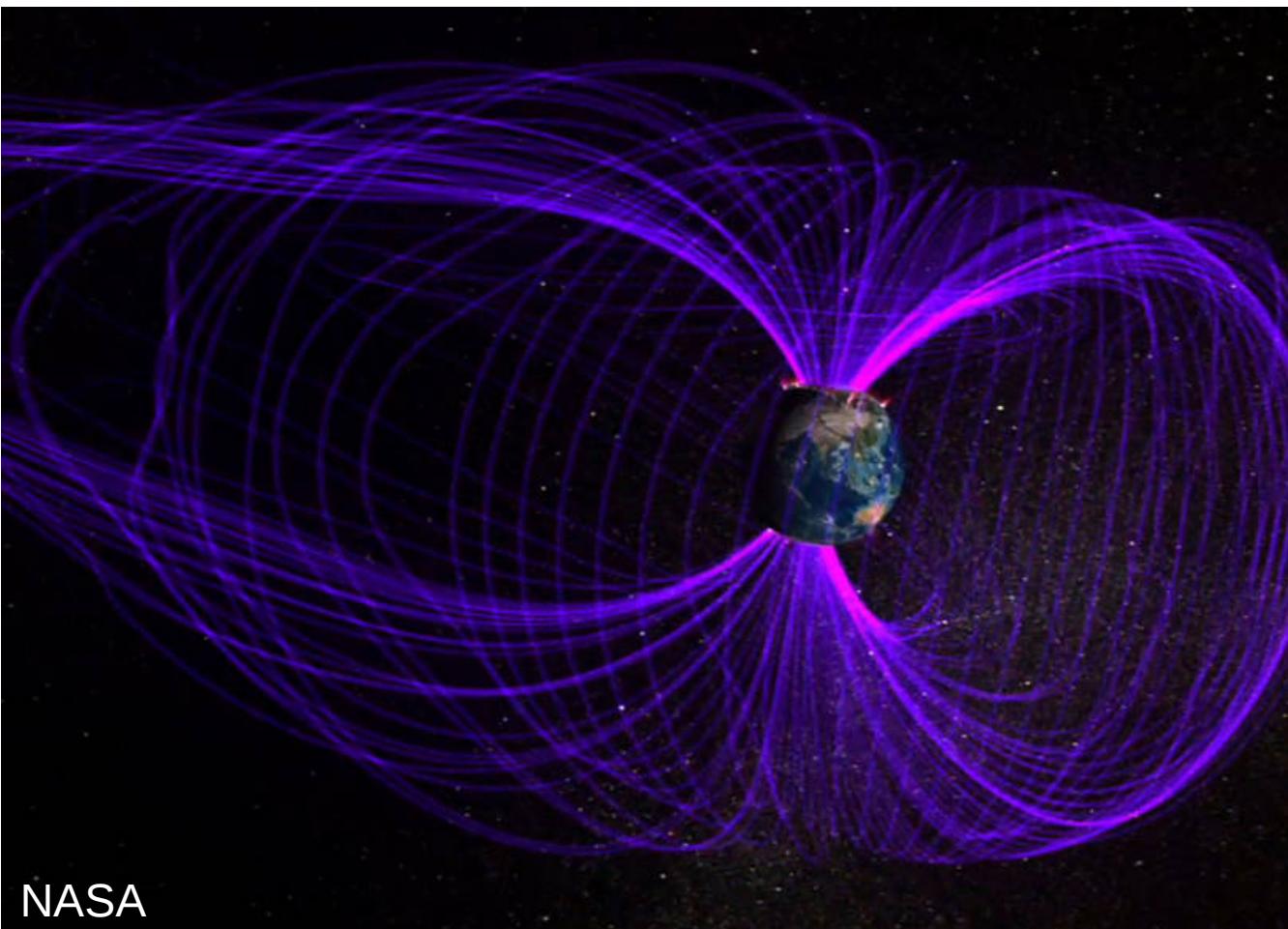
Field Line Resonance in Two and a Half Dimensions

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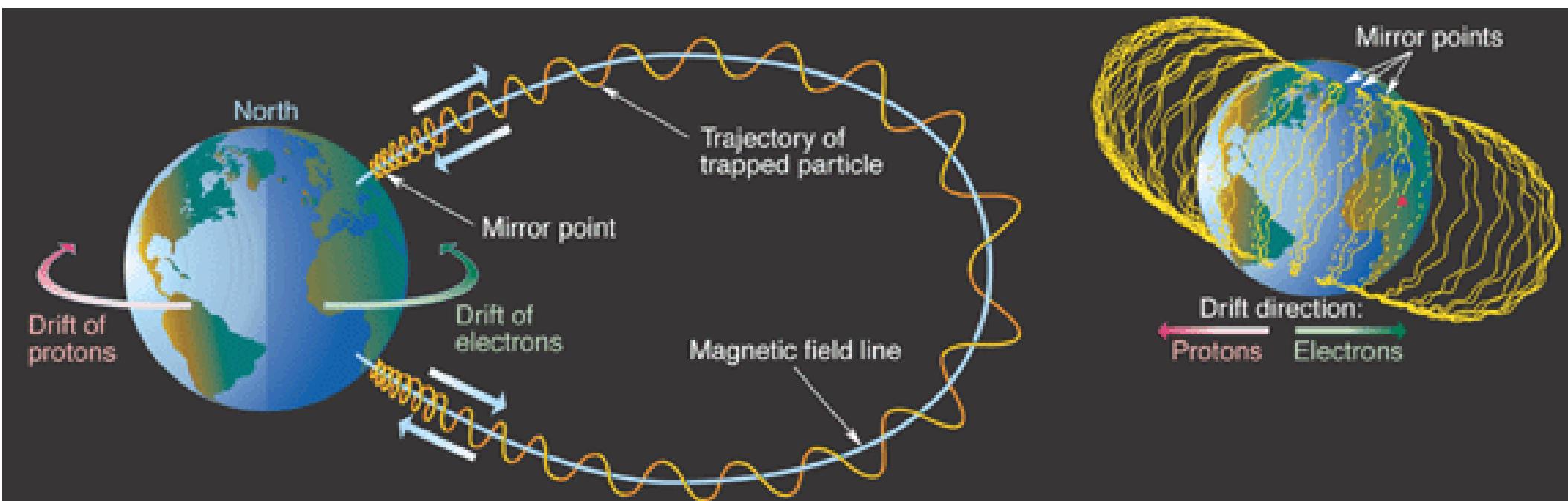
Earth's Magnetic Field

When shaken by the solar wind, etc, Earth's dipole field rattles.



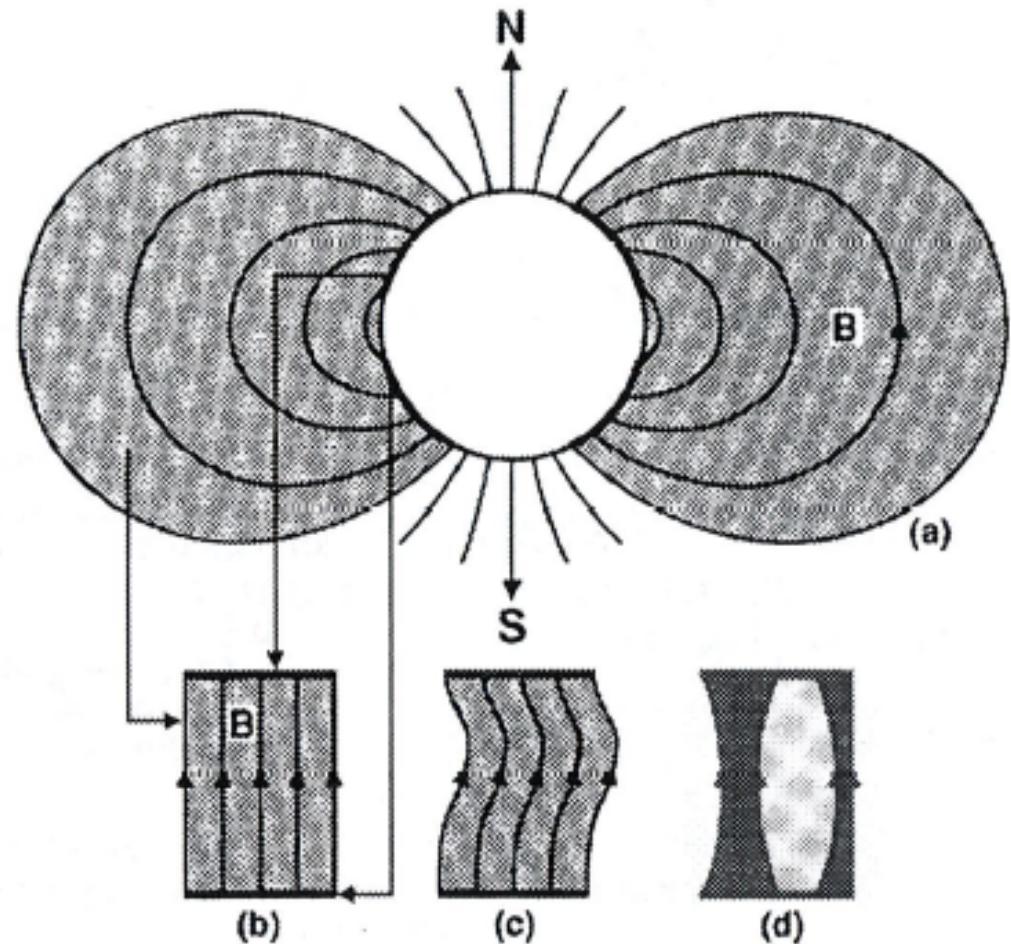
Particles in a Dipole Field

- Three fundamental motions: gyro, bounce, drift.
- Resonance happens when particle frequency matches wave frequency.



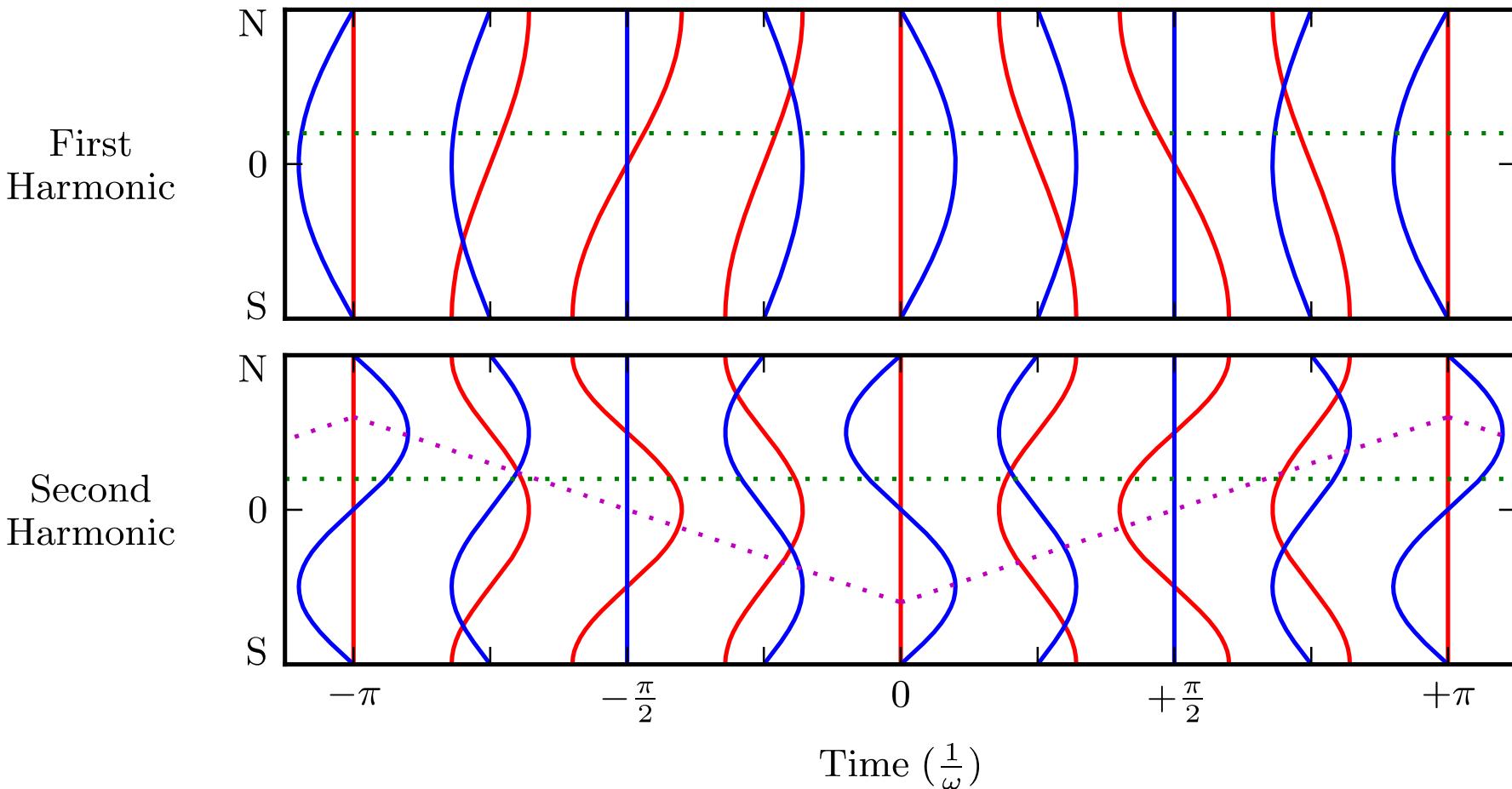
Alfvén Waves

- Shear Alfvén waves carry energy along the magnetic field.
- Compressional Alfvén waves can carry energy across field lines.



Odd and Even Harmonics

Electric (Blue) and Magnetic (Red) Harmonic Perturbations



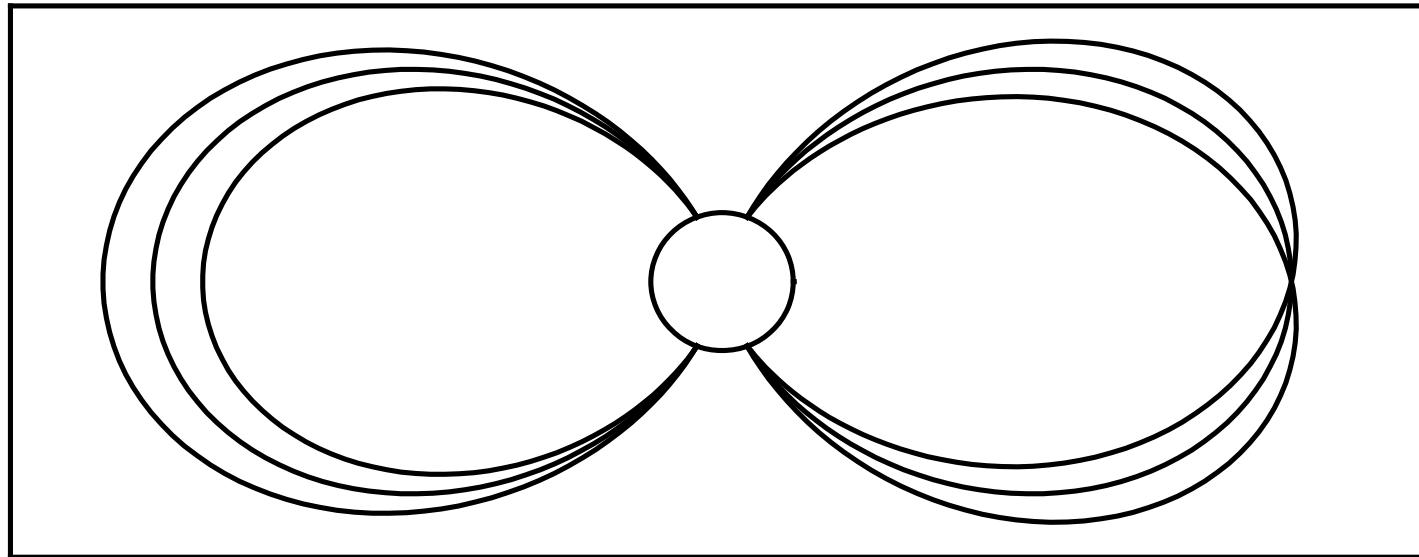
Poloidal and Toroidal Polarizations

First Harmonic

Second Harmonic

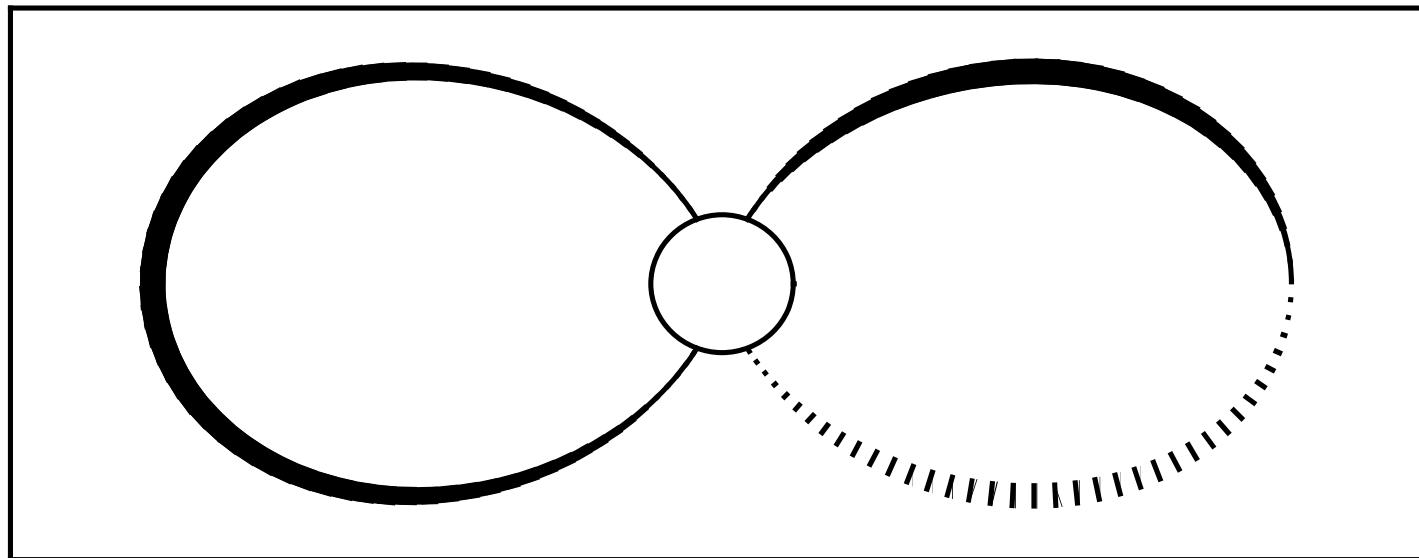
Poloidal

Z



Toroidal

Z

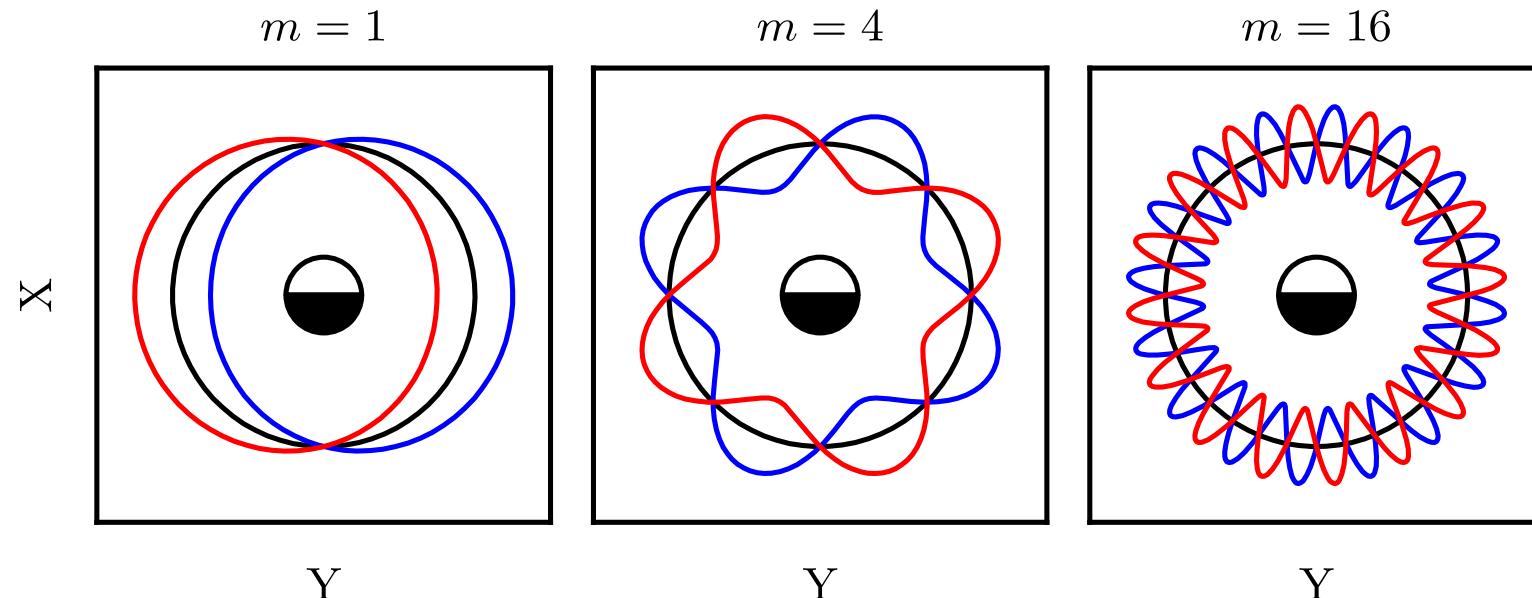


X

Azimuthal Modenumber

- Small m waves tend to be driven externally.
- Large m waves are driven internally.

Azimuthal Modenumbers Viewed from the Pole



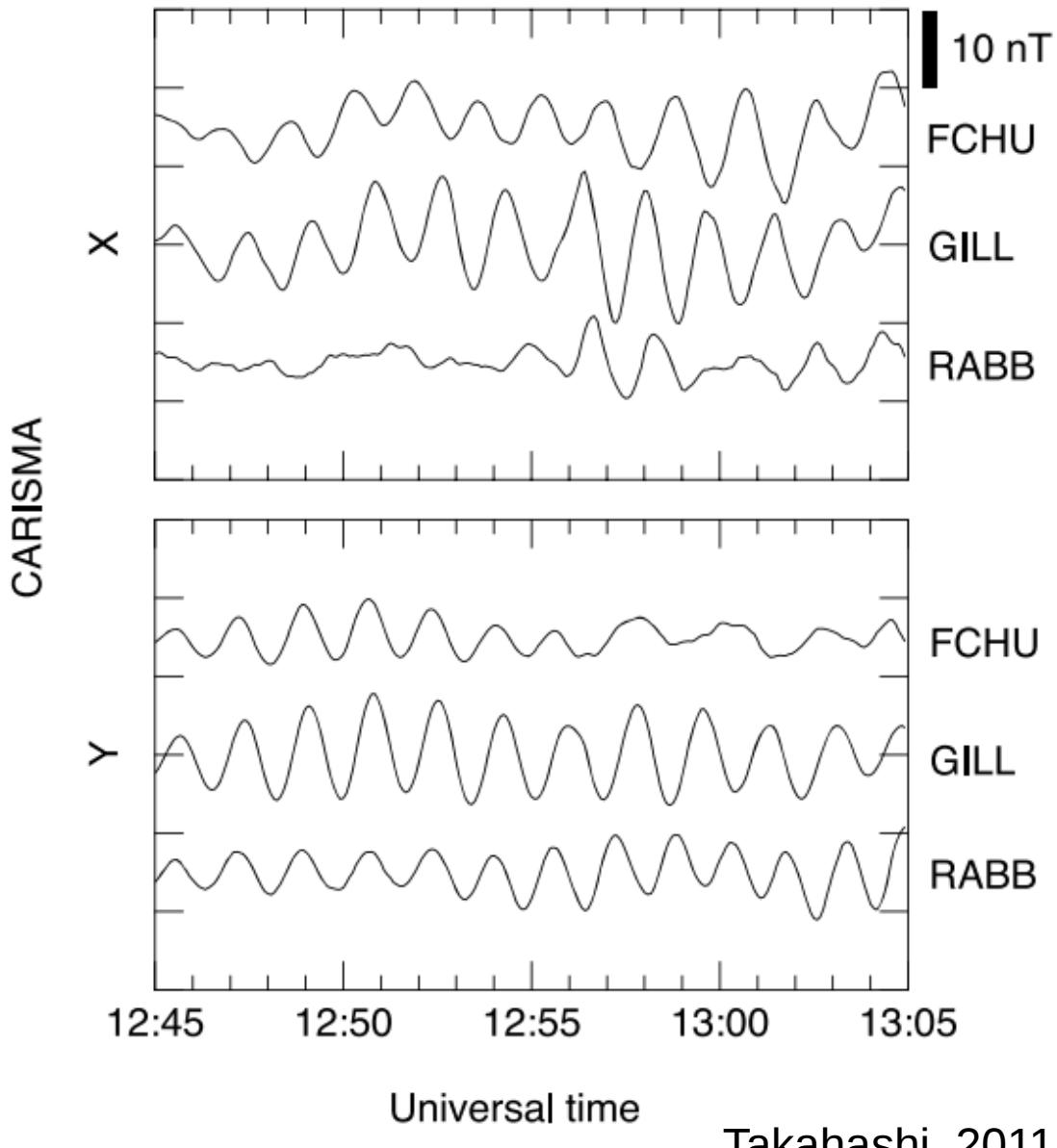
Pc4s and Giant Pulsations

Pc4s:

- FLR at 7mHz to 25mHz.
- Mostly toroidal.
- Peak near noon.

Pg:

- Odd poloidal Pc4.
- Peak near midnight.
- Striking ground signature.
- Typically $16 < m < 35$.



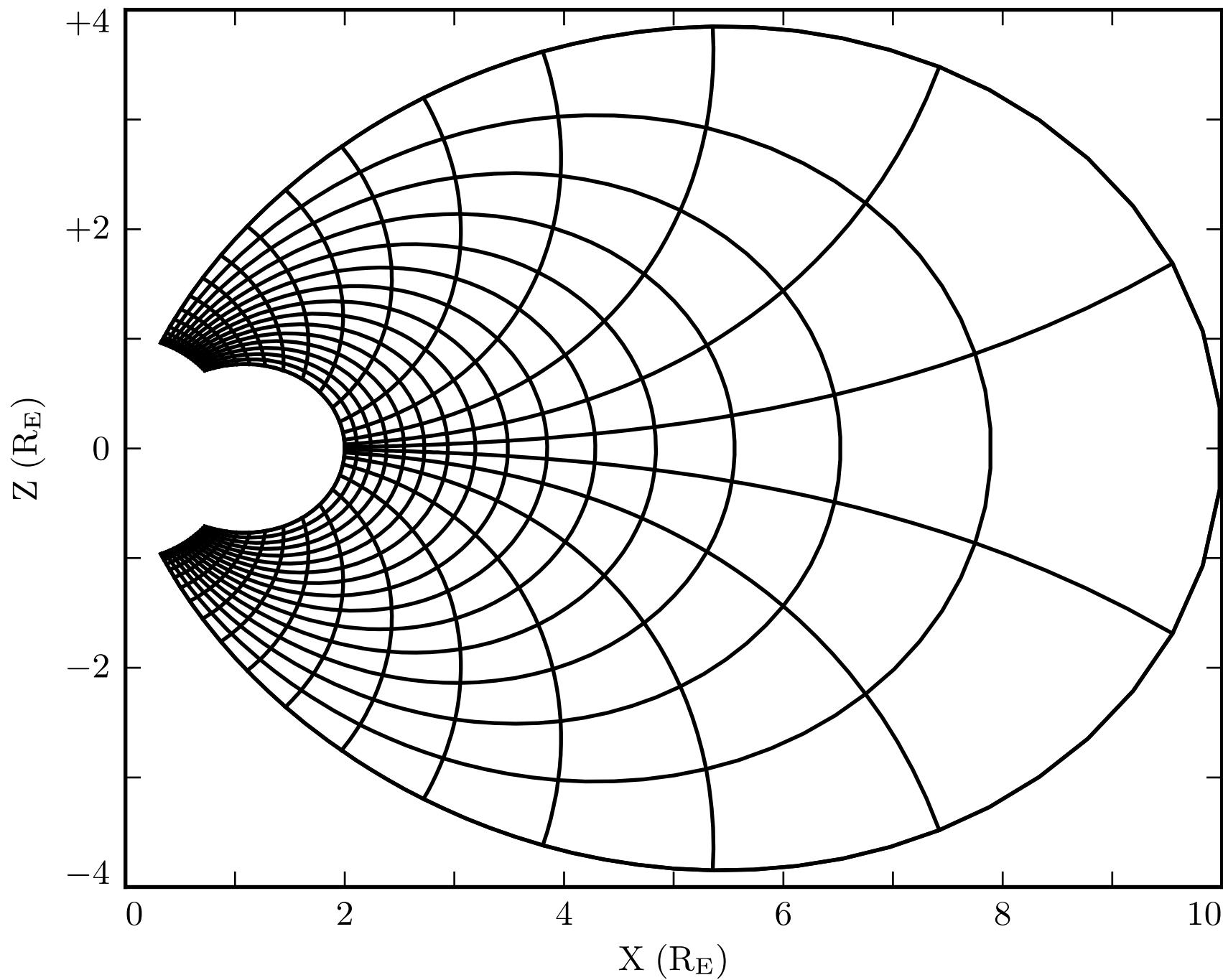
Outline

- Intro (done!)
- Numerical Model
- Numerical Results
- Van Allen Probe Observations

“Tuna Half” Dimensions

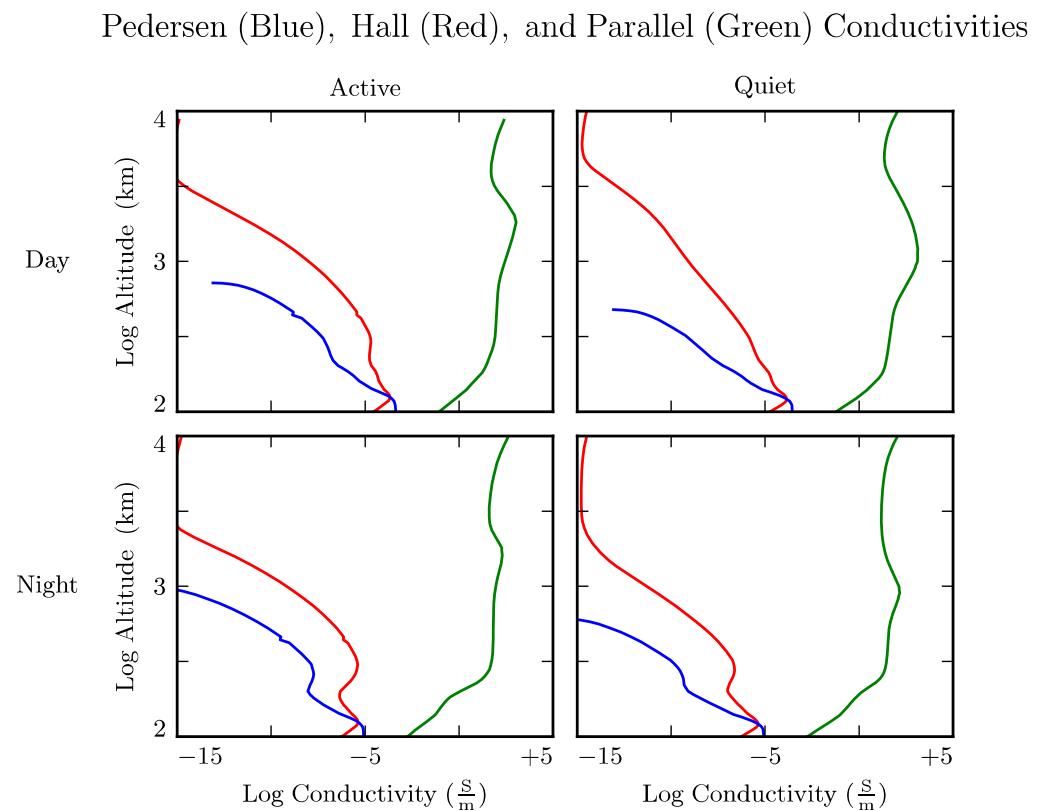
- The ionosphere is important to FLRs, but is also computationally demanding.
- High- m waves are of particular interest, but are hard to resolve in 3D.
- We don't actually need to look at the dayside and nightside at the same time.
- Say everything goes as $\exp(im\varphi)$.

Nonorthogonal Dipole Grid

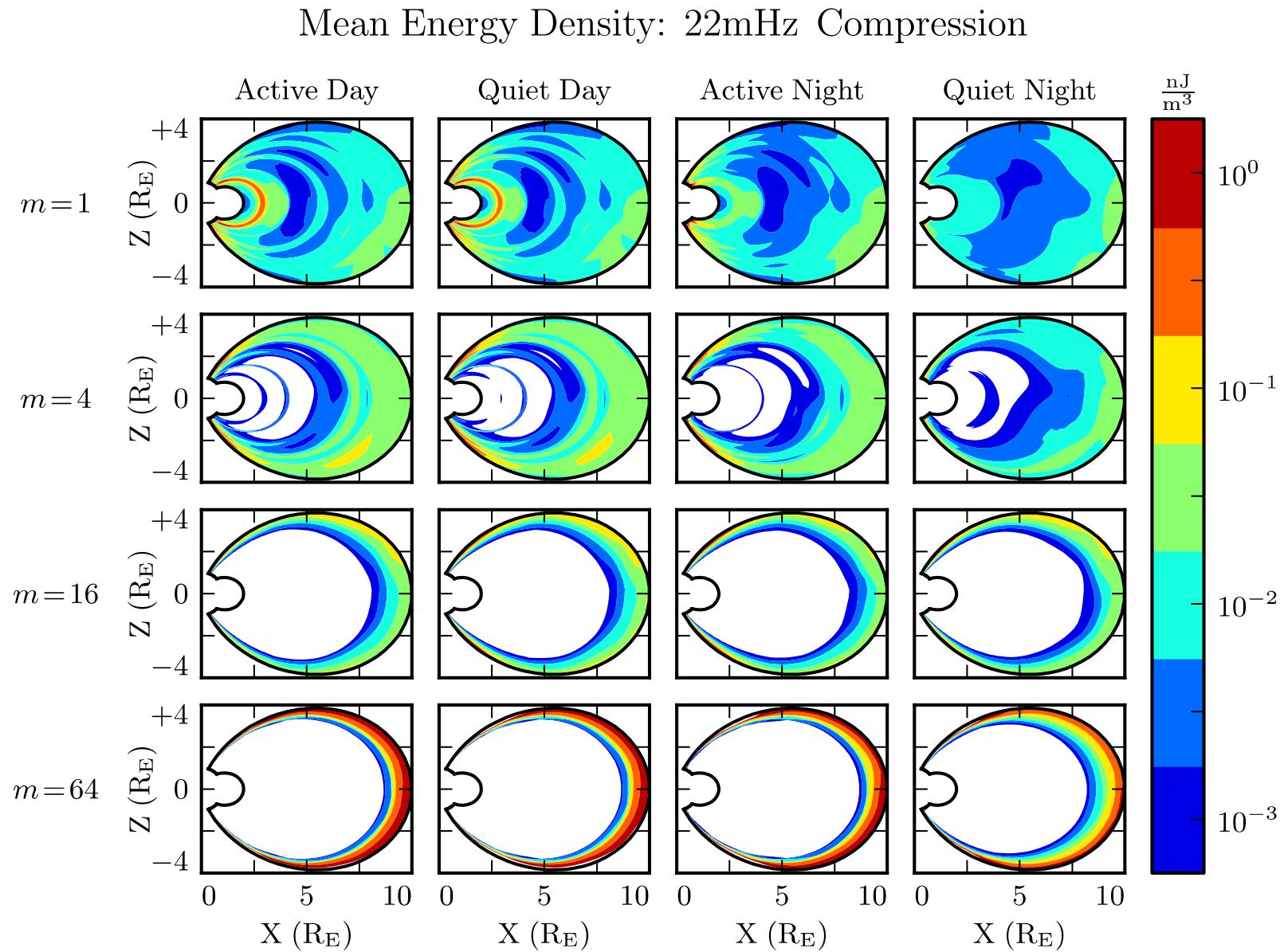


Physical Parameter Profiles

- Background dipole magnetic field.
- Height-resolved conductivity profiles.
- Density profile, including the plasmasphere.
- Alfvén speed.



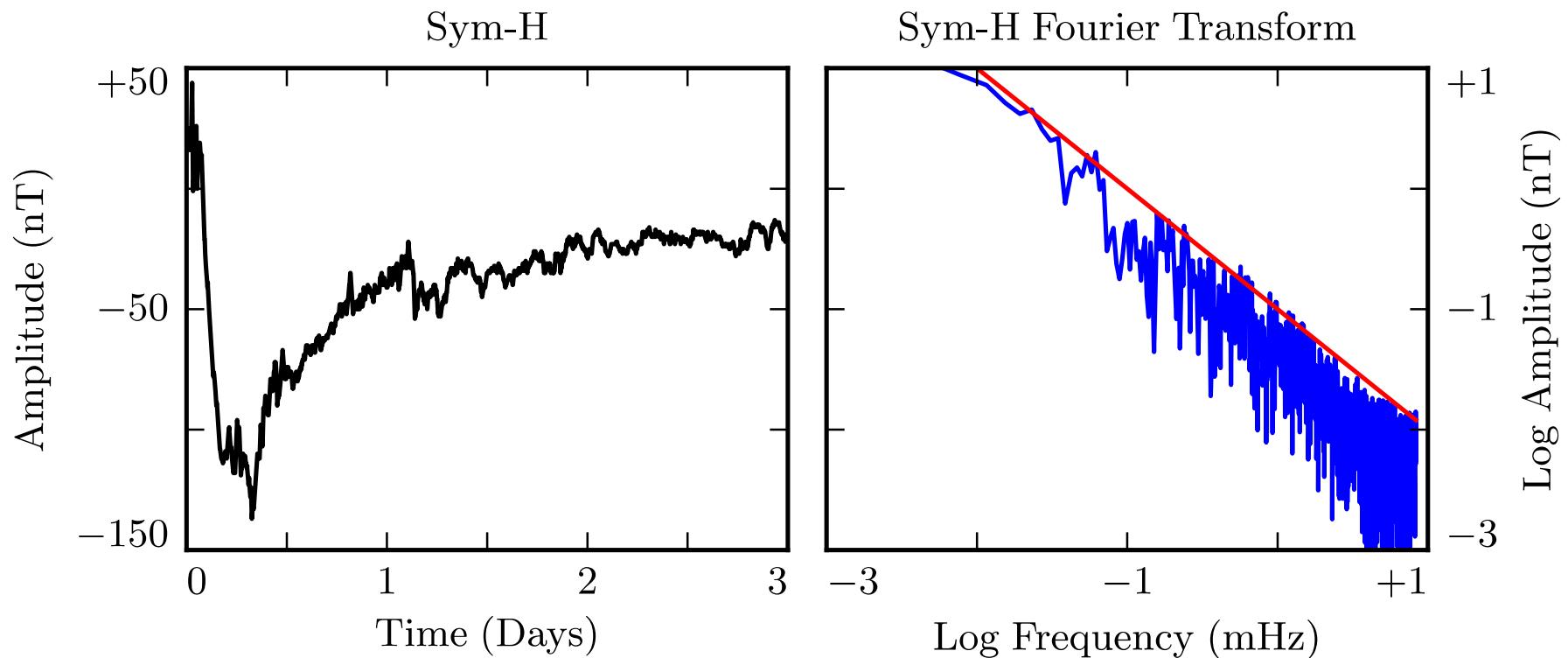
Driving with Compression?



Driving with Current

- From Sym-H pink noise, estimate the strength of ring current perturbations at 10 mHz.

Sym-H for June 2013 Storm



Maxwell's Equations

$$\frac{\partial}{\partial t} \underline{B} = -\nabla \times \underline{E}$$

$$\underline{C} \equiv \nabla \times \underline{E}$$

$$B_1 \leftarrow B_1 - g_{11} \delta t C^1 - g_{13} \delta t C^3$$

$$B_2 \leftarrow B_2 - g_{22} \delta t C^2$$

$$\underline{\epsilon} \cdot \frac{\partial}{\partial t} \underline{E} = \frac{1}{\mu_0} \nabla \times \underline{B} - \underline{J} - \underline{\sigma} \cdot \underline{E}$$

$$\underline{F} \equiv \nabla \times \underline{B} - \mu_0 \underline{J}$$

$$B_3 \leftarrow B_3 - g_{31} \delta t C^1 - g_{33} \delta t C^3$$

$$\left(\underline{\Omega} + \underline{\underline{\epsilon}} \frac{\partial}{\partial t} \right) \cdot \underline{E} = \underline{\underline{V}}^2 \cdot \underline{F}$$

$$\underline{\underline{V}}^2 \equiv \frac{1}{\mu_0} \underline{\underline{\epsilon}}^{-1} \quad \text{and} \quad \underline{\Omega} \equiv \underline{\underline{\epsilon}}^{-1} \cdot \underline{\underline{\sigma}}$$

$$\begin{aligned} E_1 + \frac{g^{13}}{g^{11}} E_3 &\leftarrow E_1 \cos\left(\frac{-\sigma_H \delta t}{\epsilon_\perp}\right) \exp\left(\frac{-\sigma_P \delta t}{\epsilon_\perp}\right) \\ &+ E_2 \sin\left(\frac{-\sigma_H \delta t}{\epsilon_\perp}\right) \exp\left(\frac{-\sigma_P \delta t}{\epsilon_\perp}\right) \sqrt{\frac{g^{22}}{g^{11}}} \\ &+ E_3 \cos\left(\frac{-\sigma_H \delta t}{\epsilon_\perp}\right) \exp\left(\frac{-\sigma_P \delta t}{\epsilon_\perp}\right) \frac{g^{13}}{g^{11}} \\ &+ F^1 \cos\left(\frac{-\sigma_H \delta t}{2\epsilon_\perp}\right) \exp\left(\frac{-\sigma_P \delta t}{2\epsilon_\perp}\right) \frac{v_A^2 \delta t}{g^{11}} \\ &+ F^2 \sin\left(\frac{-\sigma_H \delta t}{2\epsilon_\perp}\right) \exp\left(\frac{-\sigma_P \delta t}{2\epsilon_\perp}\right) \frac{v_A^2 \delta t}{\sqrt{g^{11} g^{22}}} \end{aligned}$$

$$E_3 \leftarrow E_3 \exp\left(\frac{-\sigma_0 \delta t}{\epsilon_0}\right) + c^2 \delta t (g_{31} F^1 + g_{33} F^3) \exp\left(\frac{-\sigma_0 \delta t}{2\epsilon_0}\right)$$

$$\begin{aligned} E_2 &\leftarrow -E_1 \sin\left(\frac{-\sigma_H \delta t}{\epsilon_\perp}\right) \exp\left(\frac{-\sigma_P \delta t}{\epsilon_\perp}\right) \sqrt{\frac{g^{11}}{g^{22}}} \\ &+ E_2 \cos\left(\frac{-\sigma_H \delta t}{\epsilon_\perp}\right) \exp\left(\frac{-\sigma_P \delta t}{\epsilon_\perp}\right) \\ &- E_3 \sin\left(\frac{-\sigma_H \delta t}{\epsilon_\perp}\right) \exp\left(\frac{-\sigma_P \delta t}{\epsilon_\perp}\right) \frac{g^{13}}{\sqrt{g^{11} g^{22}}} \\ &- F^1 \sin\left(\frac{-\sigma_H \delta t}{2\epsilon_\perp}\right) \exp\left(\frac{-\sigma_P \delta t}{2\epsilon_\perp}\right) \frac{v_A^2 \delta t}{\sqrt{g^{11} g^{22}}} \\ &+ F^2 \cos\left(\frac{-\sigma_H \delta t}{2\epsilon_\perp}\right) \exp\left(\frac{-\sigma_P \delta t}{2\epsilon_\perp}\right) \frac{v_A^2 \delta t}{g^{22}} \end{aligned}$$

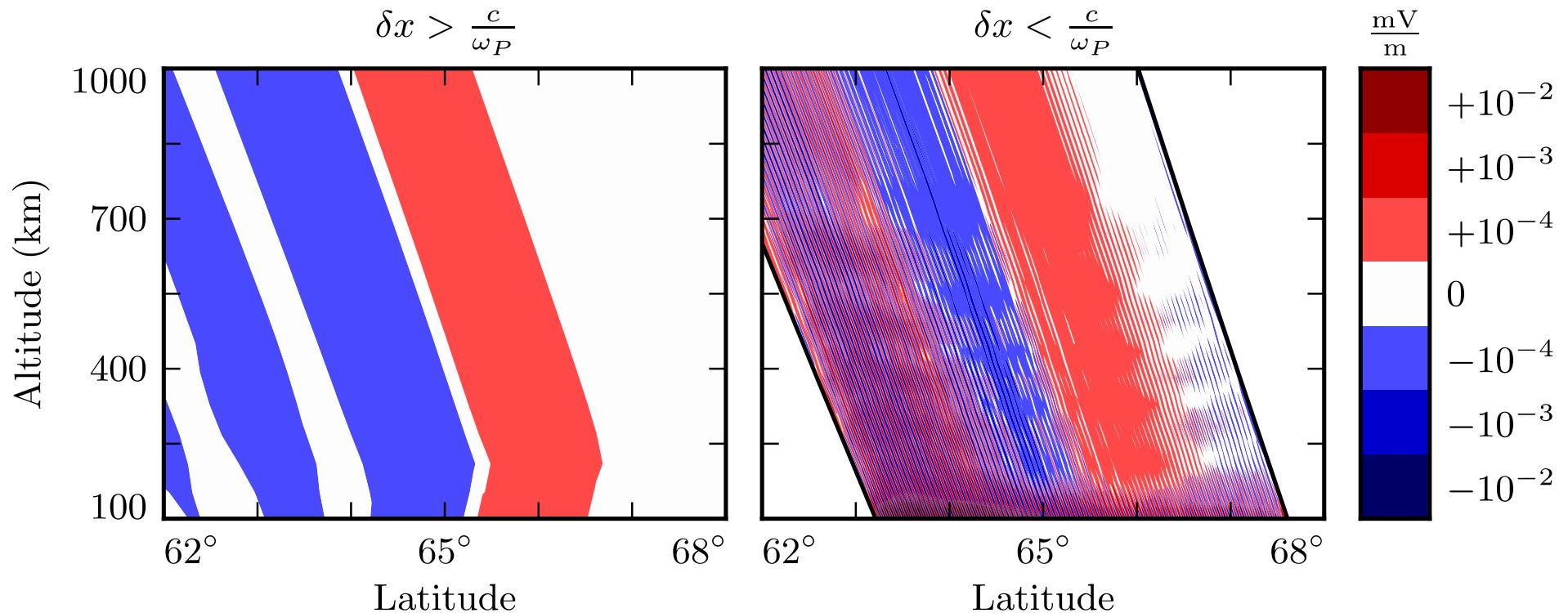
Electron Inertial Effects

$$0 = \sigma_0 E_z - J_z$$

becomes

$$\frac{1}{\nu} \frac{\partial}{\partial t} J_z = \sigma_0 E_z - J_z$$

Parallel Electric Fields: Quiet Day, 100s of 16mHz Current, $m = 16$



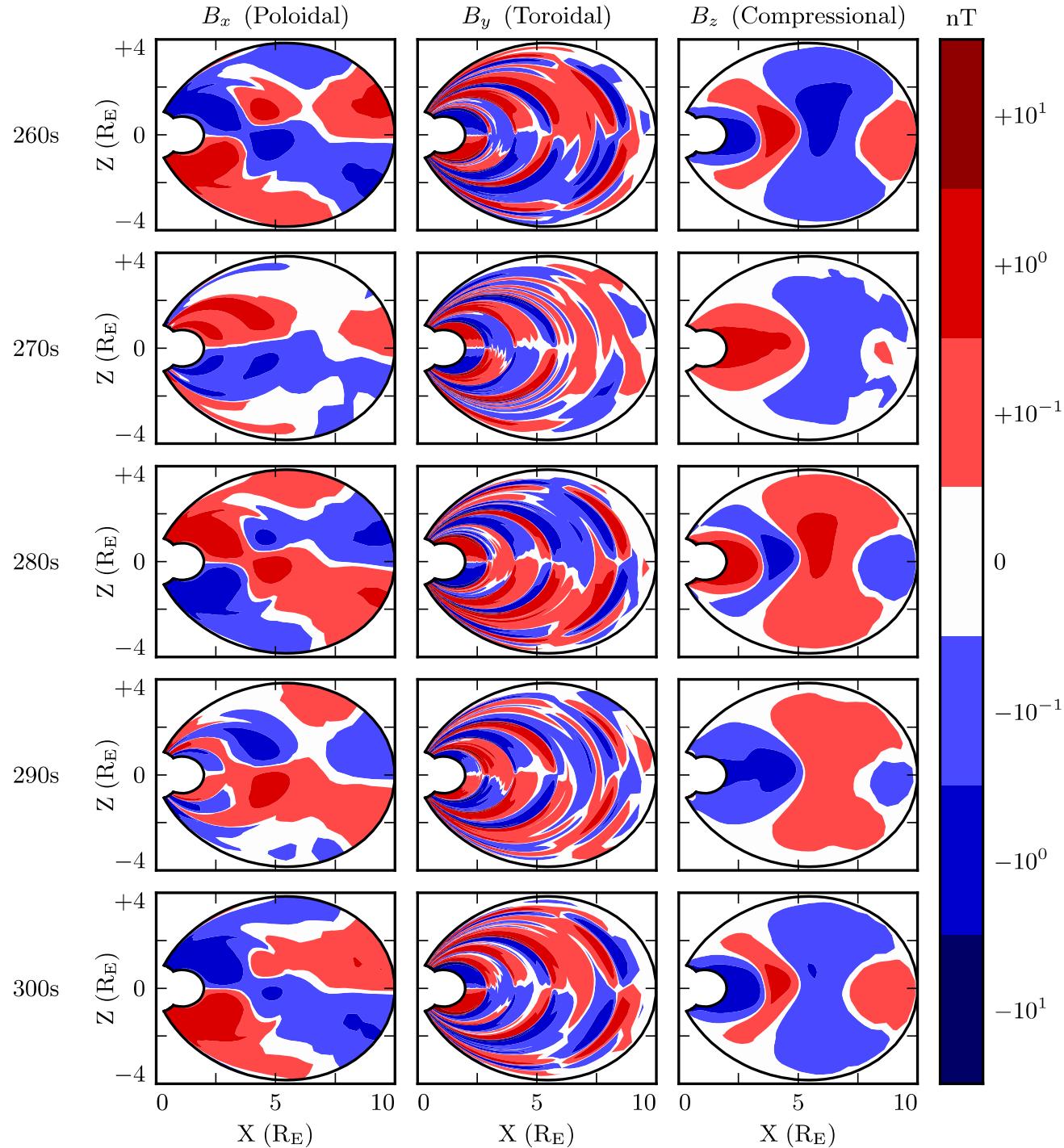
Coupling to the Atmosphere

$$\nabla \times \underline{B} = 0 \quad \text{and} \quad \nabla \cdot \underline{B} = 0$$

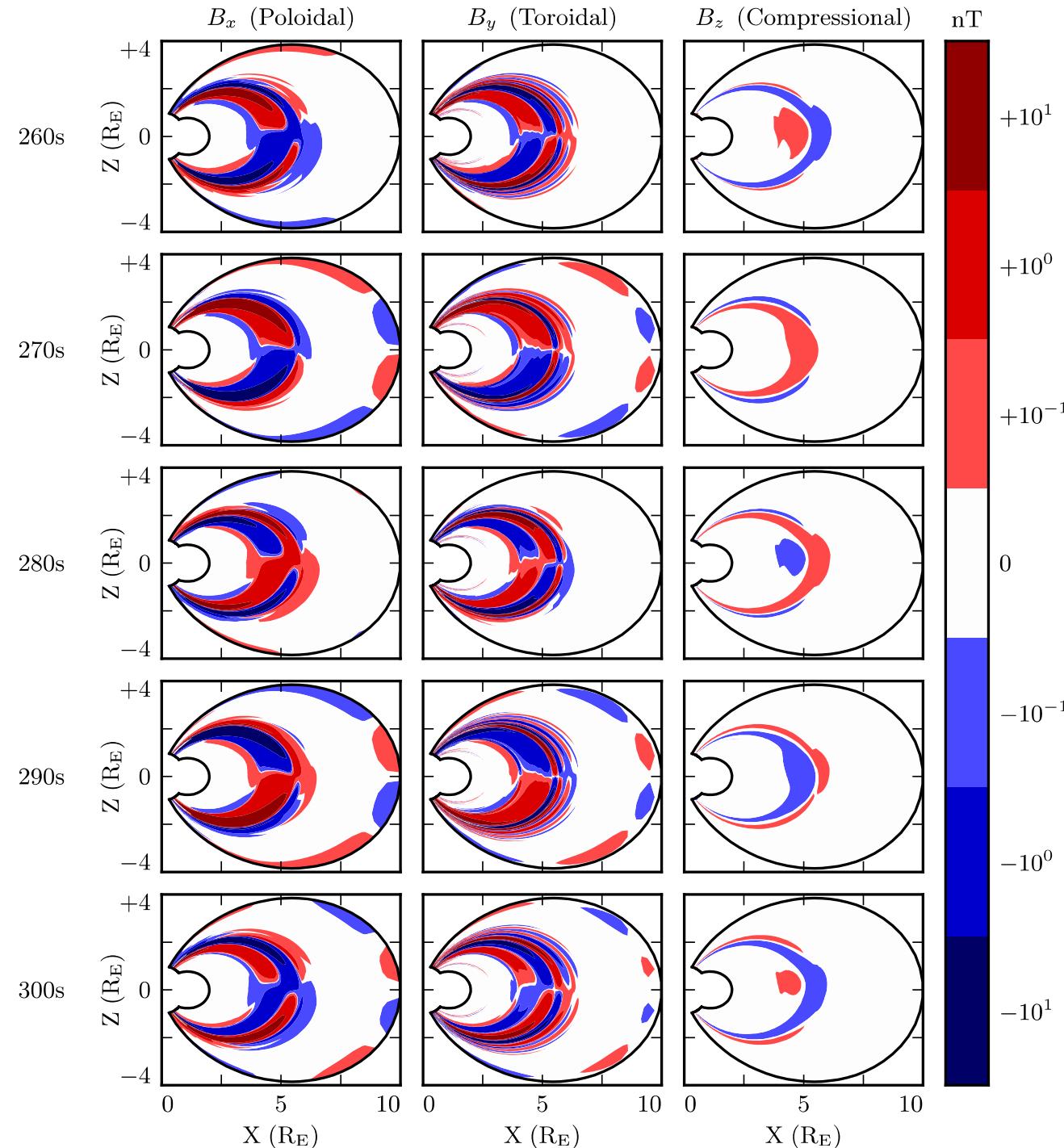
so $\underline{B} = \nabla \Psi$ where $\nabla^2 \Psi = 0$

$$\underline{\Sigma} \cdot \underline{E} = \frac{1}{\mu_0} \lim_{\delta r \rightarrow 0} \left[\hat{r} \times \underline{B} \right]_{R_I - \delta r}^{R_I + \delta r}$$

Magnetic Field Snapshots: Quiet Day , 22mHz Current, $m = 2$



Magnetic Field Snapshots: Quiet Day , 22mHz Current, $m = 32$



Poloidal-to-Toroidal Rotation

- Poloidal waves rotate to the toroidal mode over time.
- Low- m waves rotate faster.

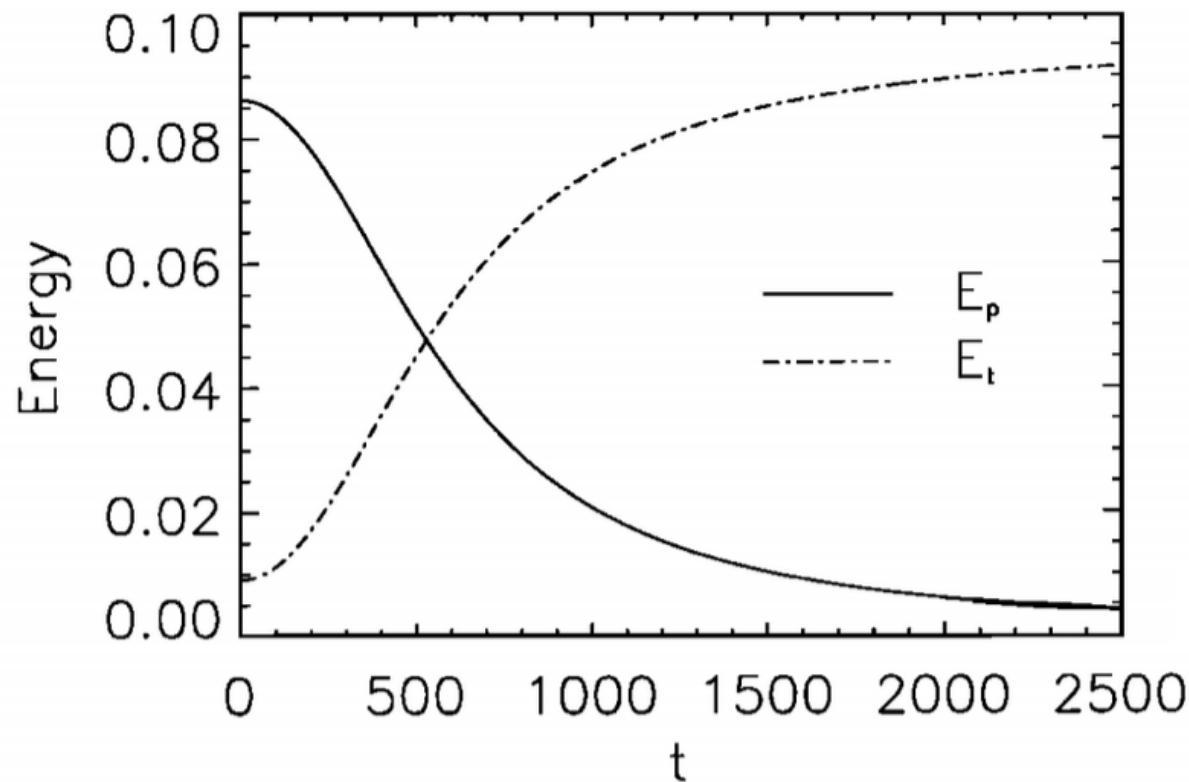
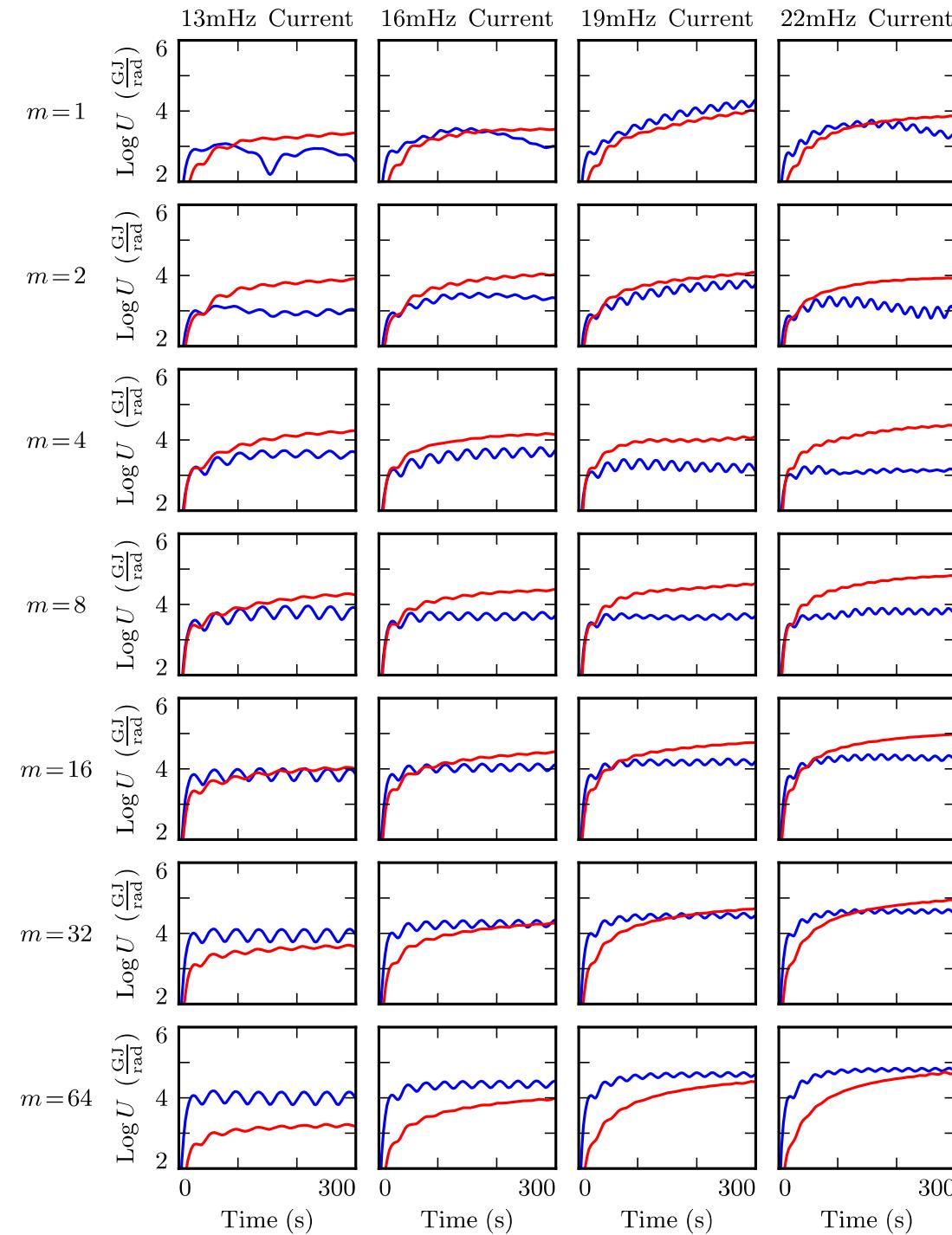


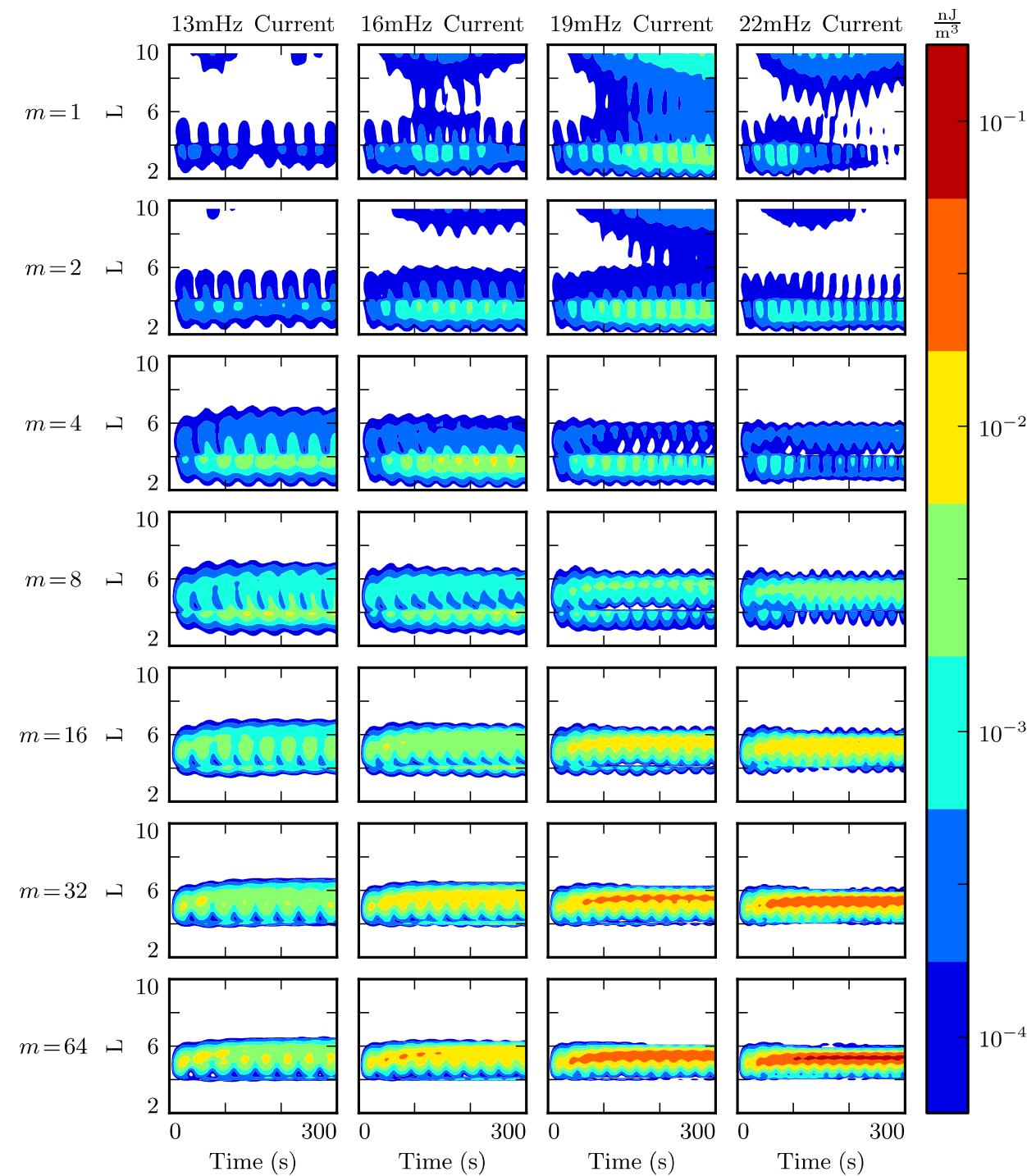
Figure 3. Time evolution of the \hat{x} integrated energy densities $E_p = \int_0^1 e_p \, dx$ and $E_t = \int_0^1 e_t \, dx$. The polarization rotation from poloidal to toroidal is clearly shown.

Poloidal (Blue) and Toroidal (Red) Energy: Quiet Day



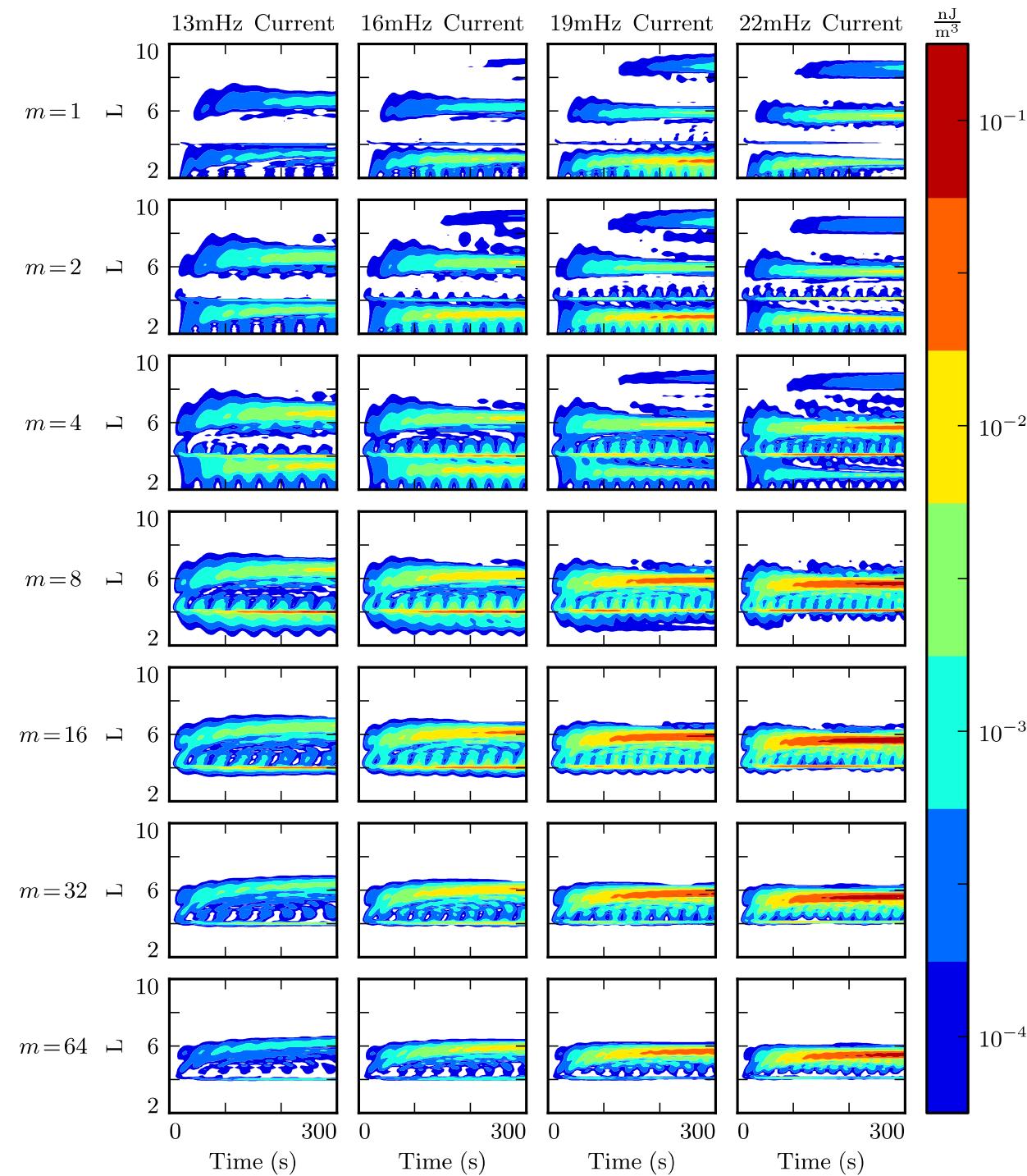
- Slowest rotation at high m .
- Faster rotation at moderate m .
- Low m : ???

Poloidal Energy Density by L-Shell: Quiet Day



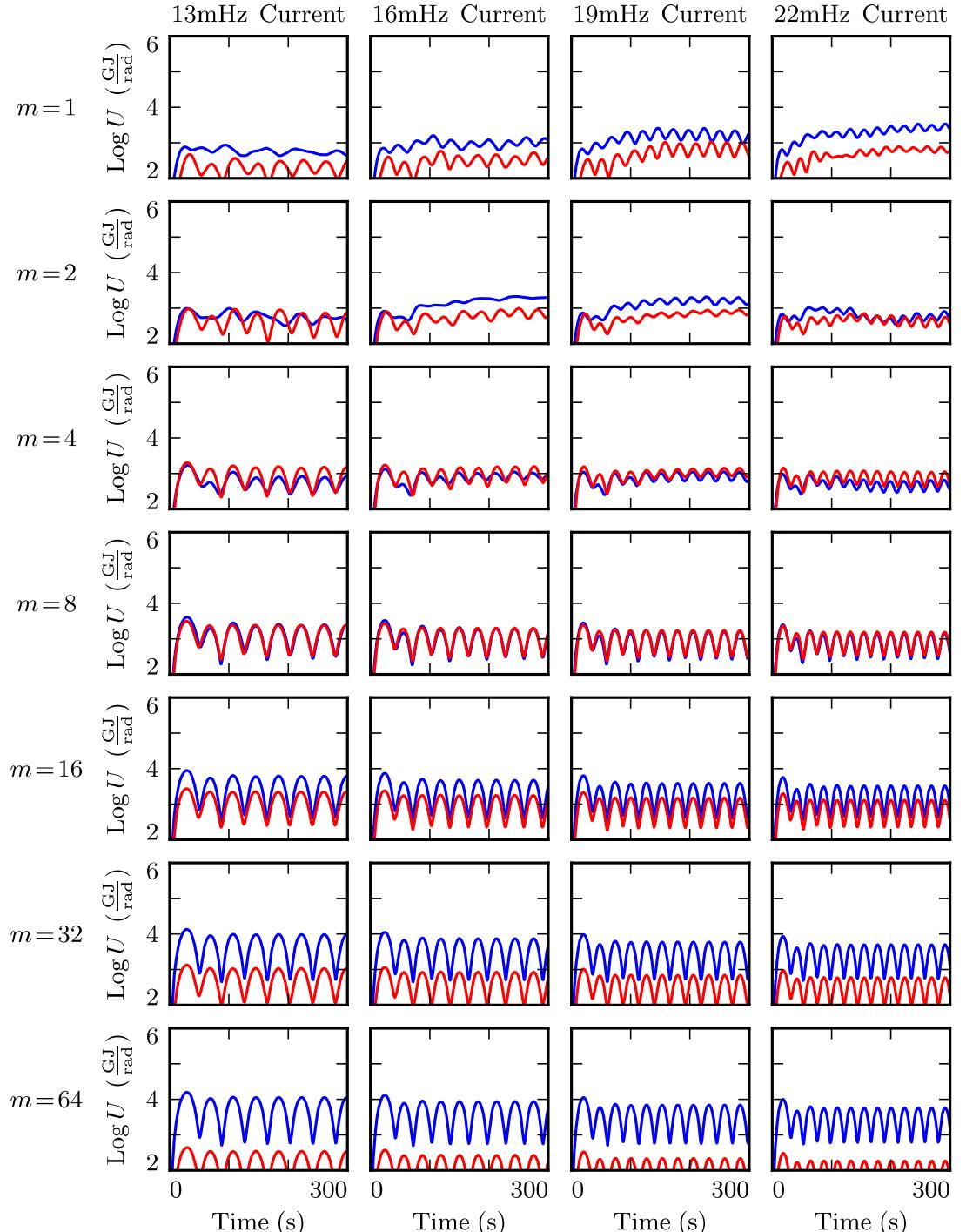
- Guided at high m .
- Escape at low m .
- Buildup near boundary likely nonphysical.

Toroidal Energy Density by L-Shell: Quiet Day



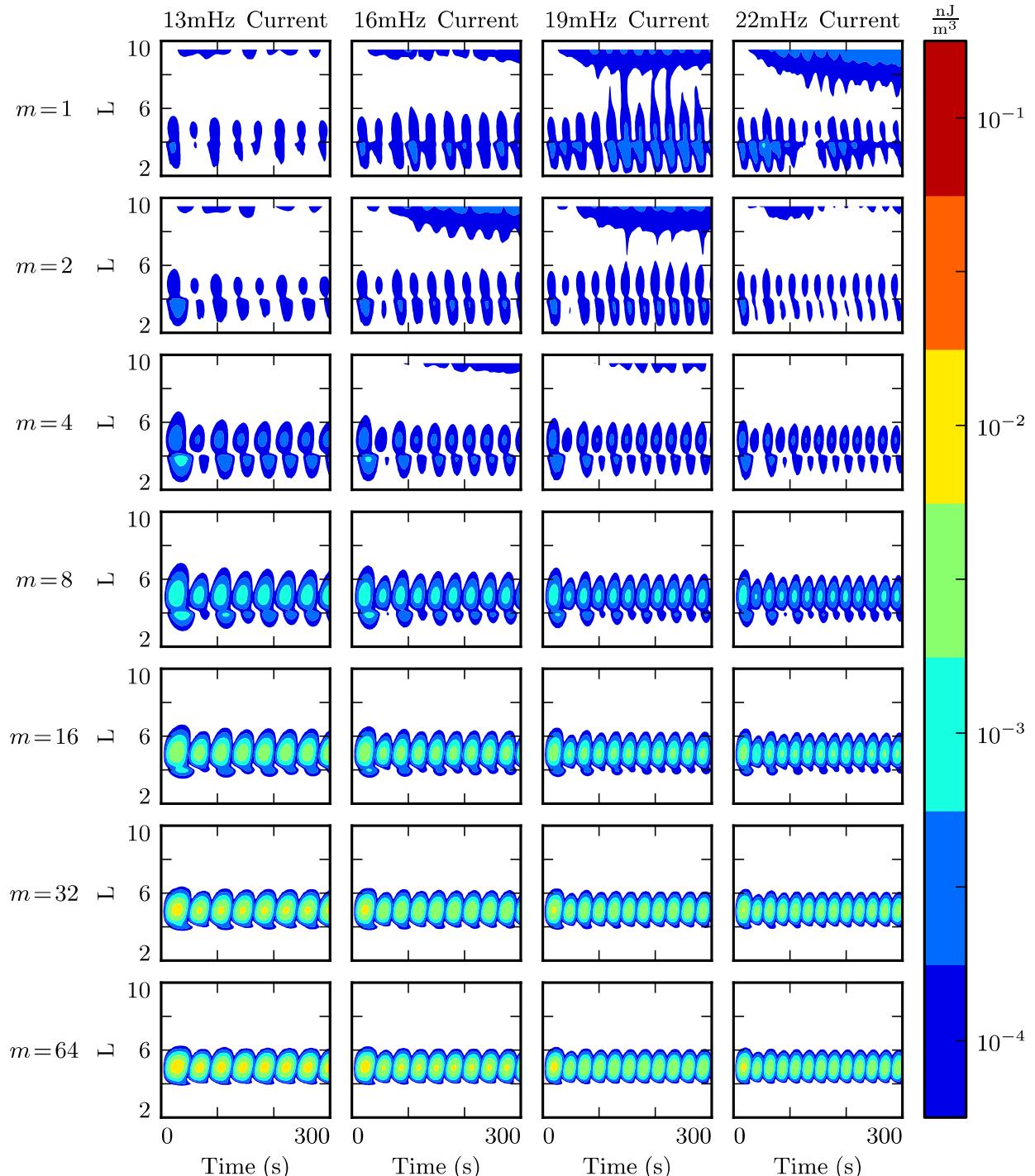
- Sharp in L, regardless of m .
- Asymptotically exceeds poloidal energy in most runs.

Poloidal (Blue) and Toroidal (Red) Energy: Quiet Night



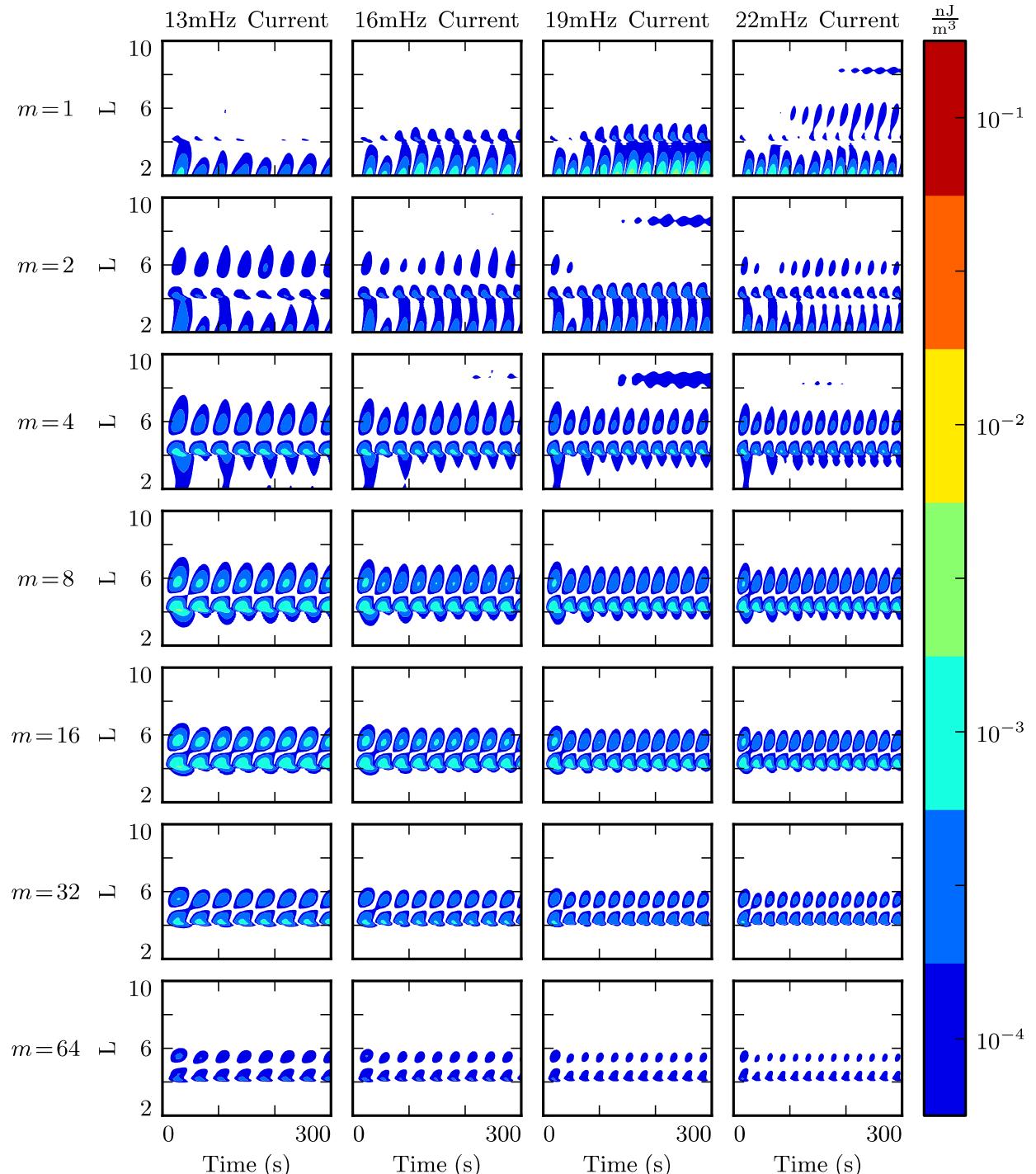
- Ionosphere is the only change.
- Low m : still weird.
- Moderate m : still best for toroidal.
- High m : toroidal very weak.

Poloidal Energy Density by L-Shell: Quiet Night



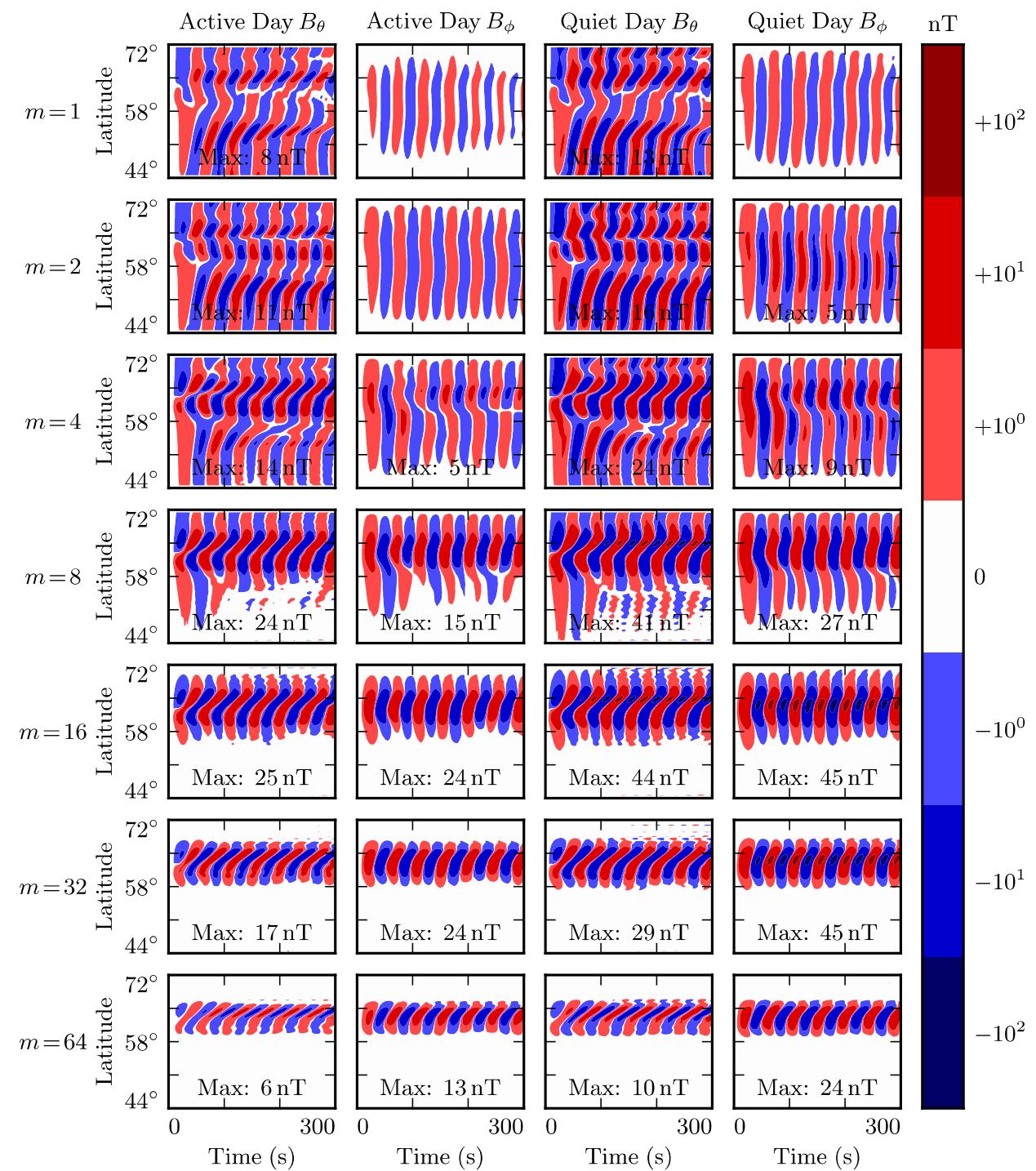
- Low m : still escaping.
- Higher $m \rightarrow$ more guided, stronger.
- Down 10x from the dayside.

Toroidal Energy Density by L-Shell: Quiet Night



- Low to moderate m : weaker than dayside, but qualitatively similar.
- High m : barely exists. Rotation slower than Joule dissipation.

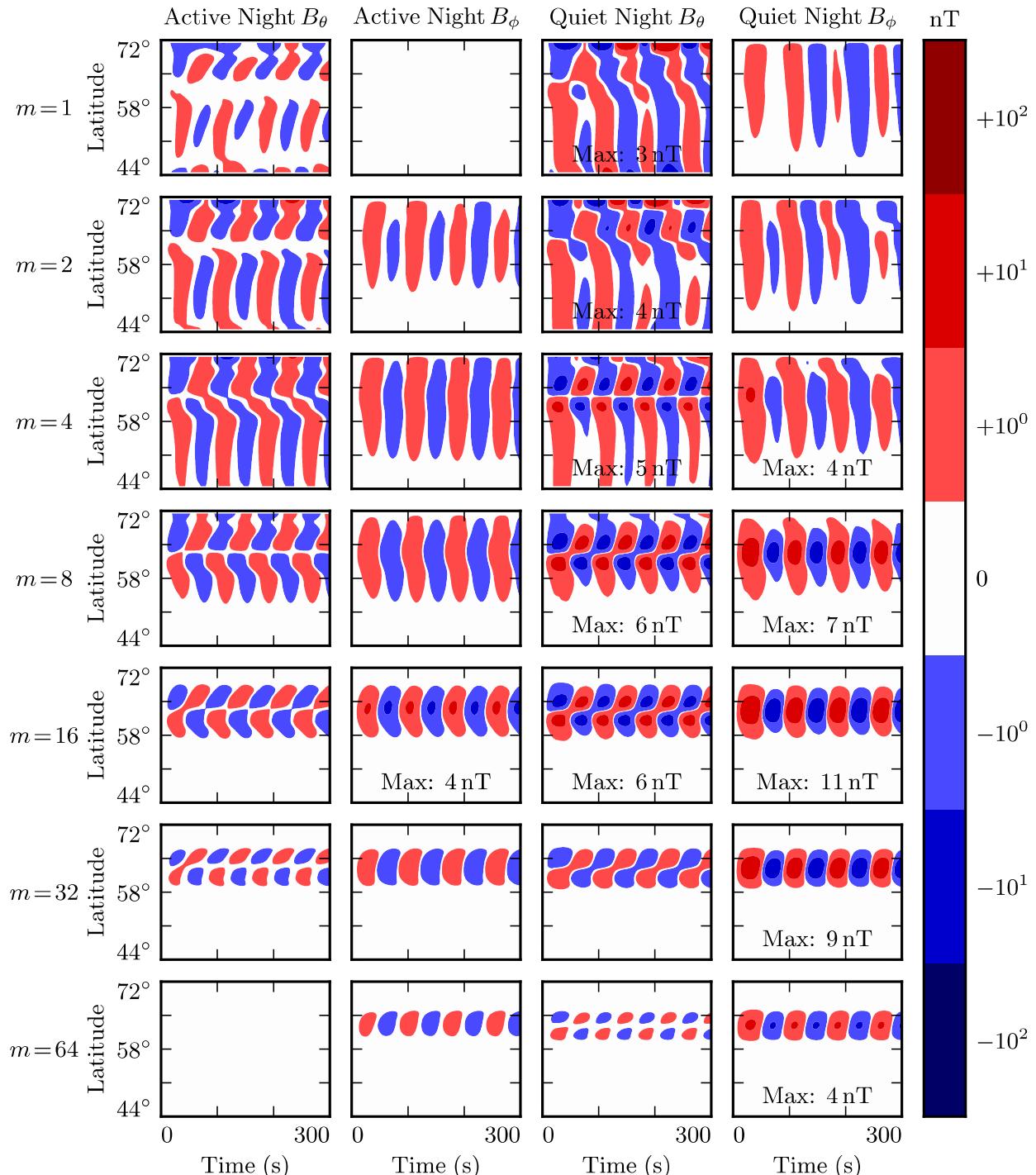
Dayside Magnetic Ground Signatures: 22mHz Current



- Peak at $m=16, 32$.
- Mostly east-west polarized.
- Strongest for quiet profile.
- Low latitude CCW, high latitude CW.

- Weaker by a factor of 4.
- Pg properties still apparent.

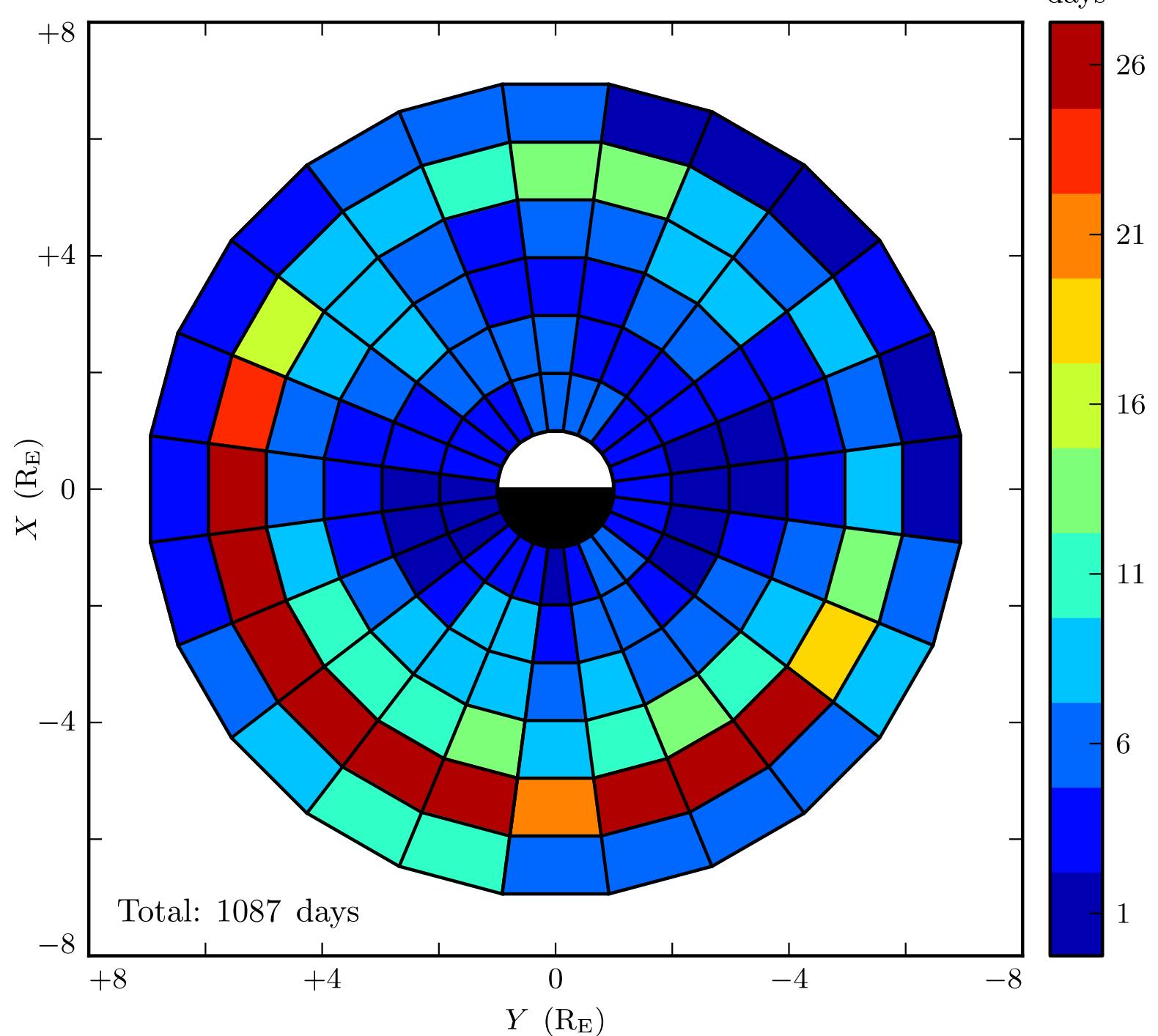
Nightside Magnetic Ground Signatures: 13mHz Current



Comparison to Observations?

- How does the distribution of Pgs line up with that of odd poloidal Pc4s overall?
- Do the distributions of odd poloidal Pc4s and odd toroidal Pc4s look similar – with the toroidal ones skewed dayward?
- How about even Pc4s?

Distribution of Usable Data: 2012-10-01 to 2015-08-24

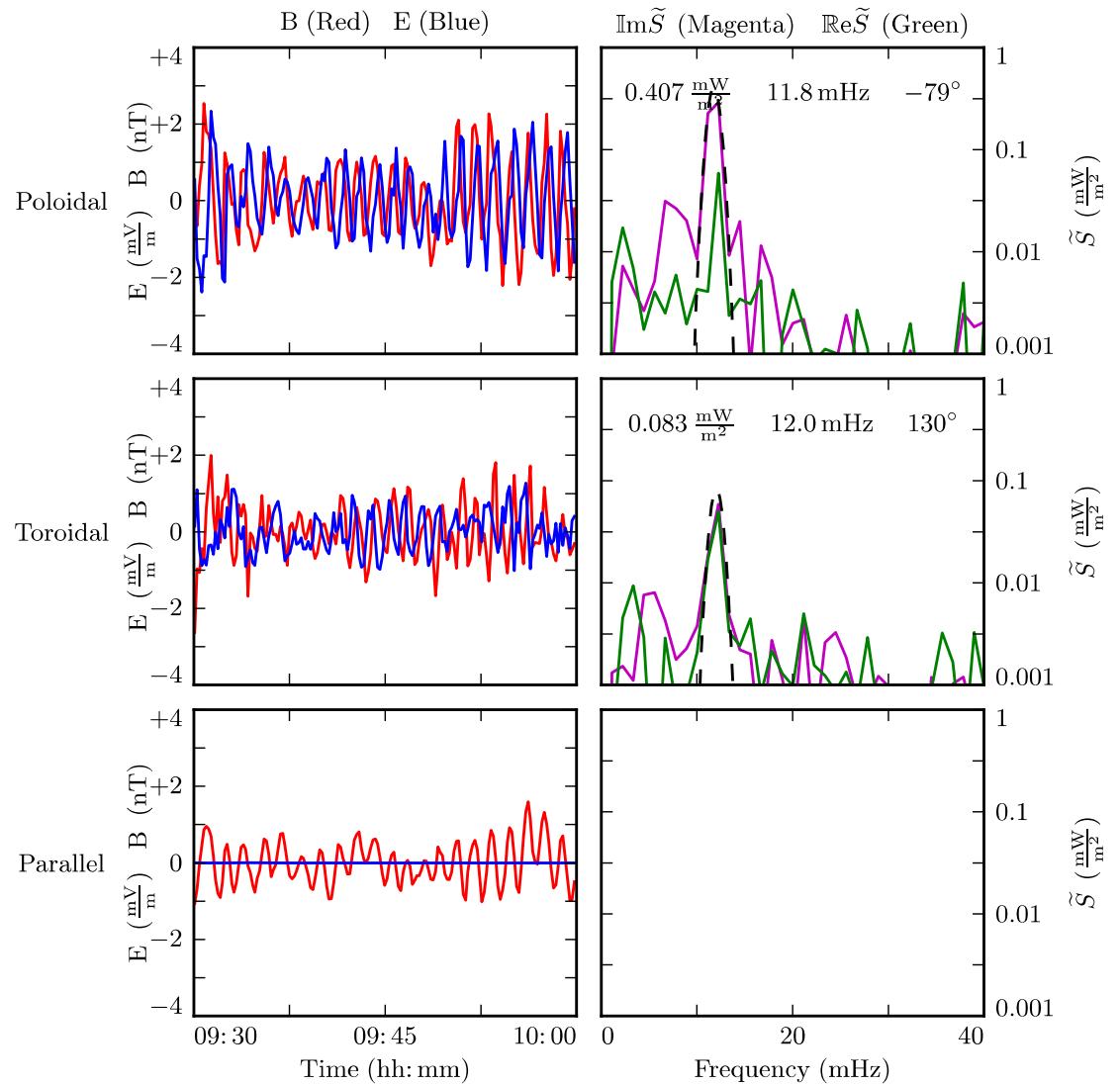


Event Selection

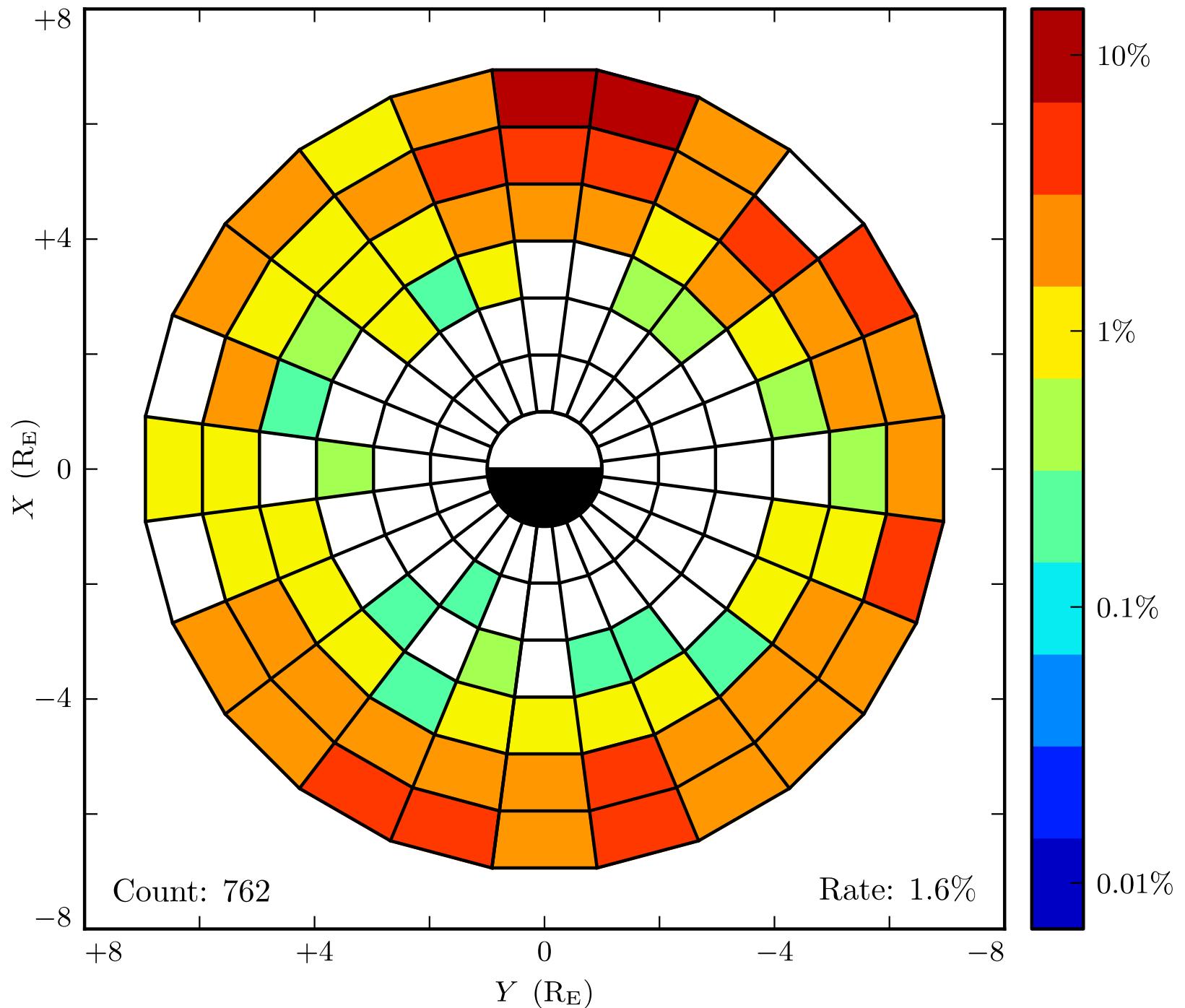
For each channel:

- Fit the spectrum.
- Check frequency.
- Check magnitude.
- Check coherence.
- Check parity.

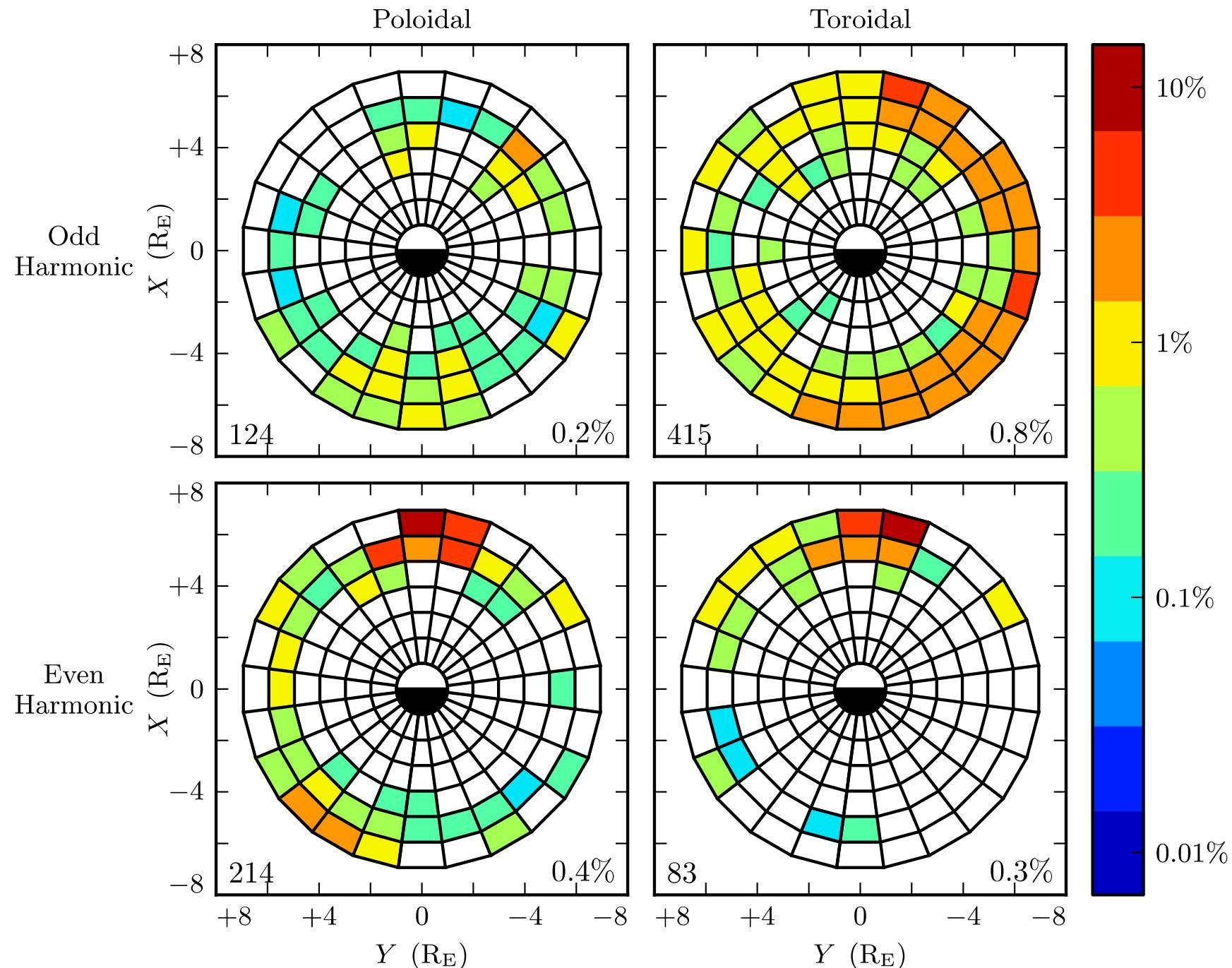
Waveforms and Spectra: Odd Poloidal Wave and Odd Toroidal Wave



Pc4 Observation Rate



Pc4 Observation Rate by Mode



Summary and Conclusions

- Past work has described Pc4s in terms of a mishmash of different properties.
- Tuna is a code designed to simulate Pc4s effectively, even at high modenumbers.
- Numerical results suggest novel connections between known Pc4 properties.
- In situ survey and numerical results compliment one another.