

Hello!
 Here is some math.

$$\alpha\beta\gamma\delta\Gamma\Delta\Lambda\Theta abcdABCD$$

Now here are some symbols.

$$\int_{-\infty}^{\infty} \sin \theta = \sqrt{\frac{e^{i\pi}}{\sum_{i=0} \epsilon \Gamma \Lambda \cdot i}}$$

Here is a compatibility test. Lowercase letters:

$$\alpha\alpha\beta b\gamma\gamma\delta d\epsilon\epsilon\epsilon\eta\eta\theta o i\iota\kappa\kappa\lambda\ell\mu\mu\nu\nu\rho\rho\sigma\sigma\tau\tau\nu\nu\phi\phi\chi x\omega\omega$$

Uppercase letters:

$$\Gamma\text{F}\Delta\text{A}\Theta\text{O}\Lambda\text{T}\Xi\text{E}\Sigma\text{X}\Upsilon\Upsilon\text{O}\Phi\text{I}\Psi\text{U}\Omega\text{O}$$

Large symbols, text-size:

$$[(\langle\{\int S \Sigma E\}\rangle)]$$

Display style:

$$\left[\left(\left\langle\left\{\int S \Sigma E\right\}\right\rangle\right)\right]$$

Common math symbols:

$$a + \frac{2}{\pi} \neq 15 \implies A \in \Pi, \forall A \approx \nabla.$$

Unbound Orbits: Deflection of Light by the Sun

Consider a particle or photon approaching the sun from very great distances. At infinity the metric is Minkowskian, that is, $A(\infty) = B(\infty) = 1$, and we expect motion on a straight line at constant velocity V

$$\begin{aligned} b &\simeq r \sin(\phi - \phi_\infty) \simeq r(\phi - \phi_\infty) \\ -V &\simeq \frac{d}{dt}(r \cos(\phi - \phi_\infty)) \simeq \frac{dr}{dt} \end{aligned}$$

where b is the “impact parameter” and ϕ_∞ is the incident direction. We see that they do satisfy the equations of motion at infinity, where $A = B = 1$, and that the constants of motion are:

$$J = bV^2 \tag{1}$$

$$E = 1 - V^2 \tag{2}$$

(Of course a photon has $V = 1$, and as we have already seen, this gives $E = 0$.) It is often more convenient to express J in terms of the distance r_0 of closest approach to the sun, rather than the impact parameter b . At r_0 , $dr/d\phi$ vanishes, so our earlier equations give:

$$J = r_0 \left(\frac{1}{B(r_0)} - 1 + V^2 \right)^{1/2}$$

The orbit is then described by

$$\phi(r) = \phi_\infty + \int_r^\infty \left\{ \frac{A^{\frac{1}{2}}(r)dr}{r^2 \left(\frac{1}{r_0^2} \left[\frac{1}{B(r)-1+V^2} \right] \left[\frac{1}{B(r)-1+V^2} \right]^{-1} - \frac{1}{r^2} \right)^{\frac{1}{2}}} \right\}.$$

The total change in ϕ as r decreases from infinity to its minimum value r_0 and then increases again to infinity is just twice its change from ∞ to r_0 , that is, $2|\phi(r_0) - \phi'_\infty|$. If the trajectory were a straight line, this would equal just π ;

$$\Delta\phi = 2|\phi(r_0) - \phi_\infty| - \pi.$$

If this is positive, then the angle ϕ changes by more than 180° , that is, the trajectory is bent *toward* the sun; if $\Delta\phi$ is negative then the trajectory is bent away from the sun.

The end.