Perturbative Gadgets

Charles Hadfield UC Berkeley

June 2018

summary

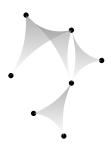
perturbative gadgets following Jordan and Farhi, arxiv 0802.1874

background and idea of problem simplified setting perturbation theory: general perturbation theory: simplified setting numerical simulation commentary

some other maths

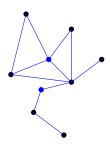
geometry of AdS/CFT spectral theory geodesics error correcting codes





$$H^{\text{comp}} = \sum_{s=1}^{r} c_s H_s$$

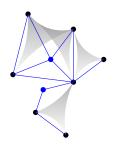
$$H_s = \sigma_{s,1} \sigma_{s,2} \cdots \sigma_{s,k}$$



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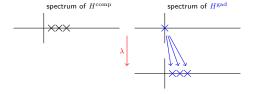
$$H^{\mathrm{gad}} = \sum \mathrm{two\ body}$$

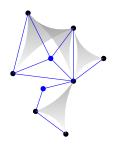


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two body

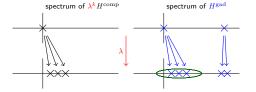


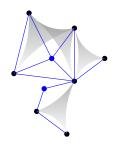


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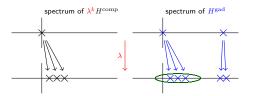




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$$\text{low.energy.sec}(H^{\text{gad}}) \sim \textcolor{red}{\lambda^k} H^{\text{comp}} + \mathcal{O}(\textcolor{blue}{\lambda^{k+1}})$$

▶ *k*-local Hamiltonian is QMA complete.

Kempe**-Kitaev*-Regev**
$$k = 2$$
, 2005 * $k = 5$ in 2002, ** $k = 3$ in 2003,

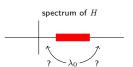
adiabatic computation = circuit model

Aharonov et al
$$k = 3$$
, 2004
KKR $k = 2$

Oliveira-Terhal 2005, Jordan-Farhi 2008, Cao-Kais 2016.

.ao-Kais 2016,

Cubitt-Montanaro-Piddock 2017





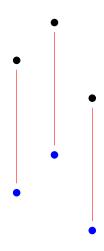
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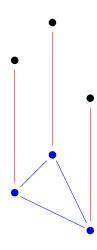
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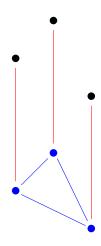
$$V := \sum_{i=1}^{3} \sigma_i X_i$$



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$$H^{\mathrm{anc}} := \sum_{i < j} \frac{1}{2} (1 - Z_i Z_j)$$



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$$H^{\mathrm{gad}} := H^{\mathrm{anc}} + \lambda V$$

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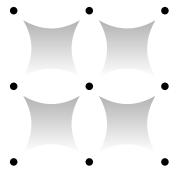
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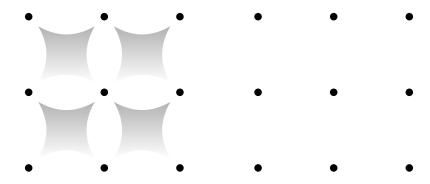
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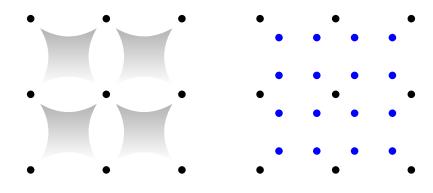
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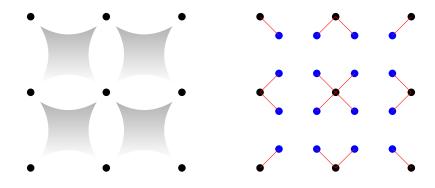
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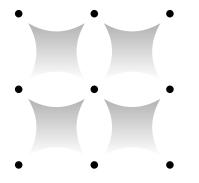
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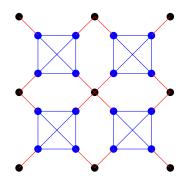


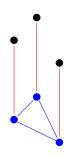






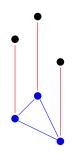






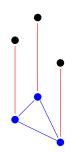
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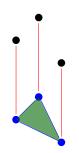
$$H^{\mathrm{gad}} := H^{\mathrm{anc}} + \frac{\lambda V}{V} : \mathbb{C}^{64} \to \mathbb{C}^{64}$$



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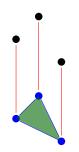


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$$X_1X_2X_3$$

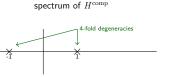


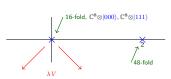


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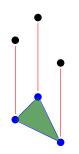
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$$[H^{\mathrm{gad}}, X_1 X_2 X_3] = 0$$





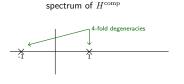
spectrum of $H^{\rm anc}$

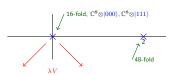


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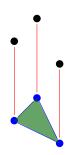
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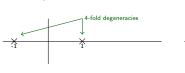
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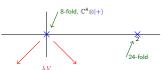
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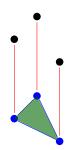
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spectrum of H^{comp}



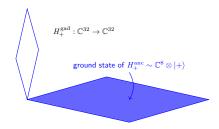
spectrum of $H_{\perp}^{\rm anc}$

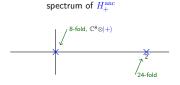


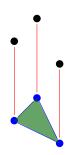
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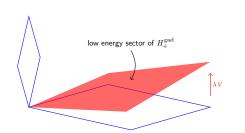


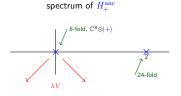


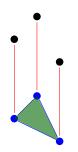
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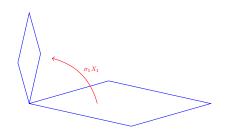


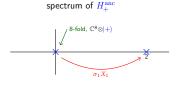


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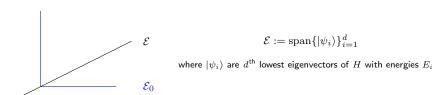




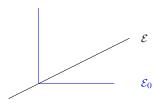
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where $|\psi_i\rangle$ are d^{th} lowest eigenvectors of H with energies E_i define effective hamiltonian which captures low energy sector \mathcal{E}

 $\mathcal{E} := \operatorname{span}\{|\psi_i\rangle\}_{i=1}^d$

$$H_{\mathrm{eff}}(H,d) := \sum_{i=1}^{d} E_i |\psi_i\rangle\langle\psi_i|$$

$$H:\mathbb{C}^n\to\mathbb{C}^n$$

$$I := H_0 +$$

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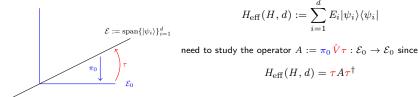
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 \mathcal{E}_0

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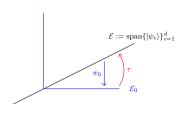
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$$H_{\text{eff}}(H,d) = {}^{} A {}^{}$$

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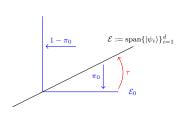
need to study the operator $A:=\pi_0\,\hat{V}{\pmb{ au}}:{\cal E}_0 o{\cal E}_0$ since

$$H_{\text{eff}}(H,d) = \tau A \tau^{\dagger}$$

now start calculating to obtain perturbative formulae for

$$\tau := \sum_{k=0}^{\infty} \lambda^k \tau^{(k)} \qquad A := \sum_{k=1}^{\infty} \lambda^k A^{(k)}$$

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$$S^{0} := -\pi_{0}$$

$$S^{k} := (-1)^{k} H_{0}^{-k} (1 - \pi_{0})$$

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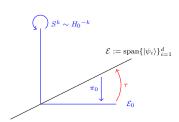
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$$H_{\mathrm{eff}}(H,d) = \tau A \tau^\dagger$$

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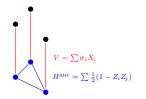
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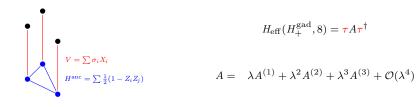
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$$p_i \ge 0$$
 for all p_i
 $p_1 + \dots + p_{k'} \ge k'$ for all $k' \le k$
 $p_1 + \dots + p_k = k$

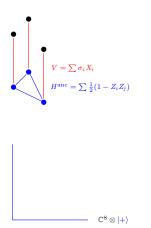


$$H_{\mathrm{eff}}(H_{+}^{\mathrm{gad}}, 8) = \tau A \tau^{\dagger}$$

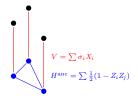


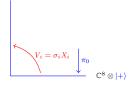




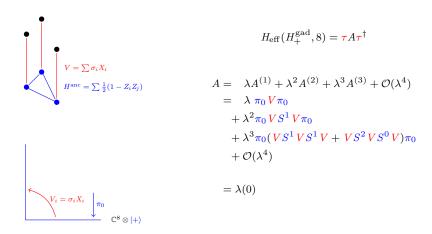


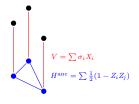
$$\begin{split} H_{\mathrm{eff}}(H_{+}^{\mathrm{gad}},8) &= \tau A \tau^{\dagger} \\ A &= \lambda A^{(1)} + \lambda^{2} A^{(2)} + \lambda^{3} A^{(3)} + \mathcal{O}(\lambda^{4}) \\ &= \lambda \, \pi_{0} \, V \pi_{0} \\ &+ \lambda^{2} \pi_{0} \, V S^{1} \, V \pi_{0} \\ &+ \lambda^{3} \pi_{0} (V S^{1} \, V S^{1} \, V + V S^{2} \, V S^{0} \, V) \pi_{0} \\ &+ \mathcal{O}(\lambda^{4}) \end{split}$$

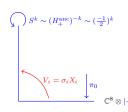




$$\begin{split} H_{\mathrm{eff}}(H_{+}^{\mathrm{gad}},8) &= \tau A \tau^{\dagger} \\ A &= \lambda A^{(1)} + \lambda^{2} A^{(2)} + \lambda^{3} A^{(3)} + \mathcal{O}(\lambda^{4}) \\ &= \lambda \pi_{0} V \pi_{0} \\ &+ \lambda^{2} \pi_{0} V S^{1} V \pi_{0} \\ &+ \lambda^{3} \pi_{0} (V S^{1} V S^{1} V + V S^{2} V S^{0} V) \pi_{0} \\ &+ \mathcal{O}(\lambda^{4}) \end{split}$$







$$H_{\text{eff}}(H_{+}^{\text{gad}}, 8) = \tau A \tau^{\dagger}$$

$$A = \lambda A^{(1)} + \lambda^{2} A^{(2)} + \lambda^{3} A^{(3)} + \mathcal{O}(\lambda^{4})$$

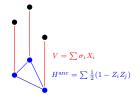
$$= \lambda \pi_{0} V \pi_{0}$$

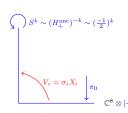
$$+ \lambda^{2} \pi_{0} V S^{1} V \pi_{0}$$

$$+ \lambda^{3} \pi_{0} (V S^{1} V S^{1} V + V S^{2} V S^{0} V) \pi_{0}$$

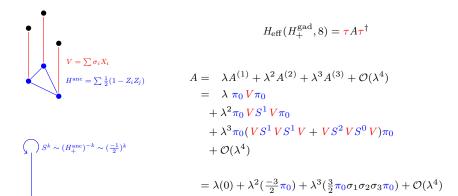
$$+ \mathcal{O}(\lambda^{4})$$

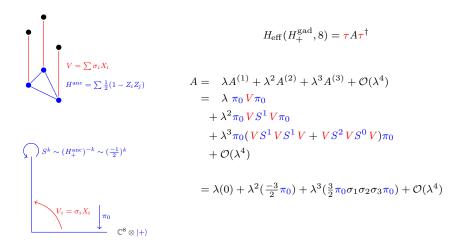
$$= \lambda(0)$$



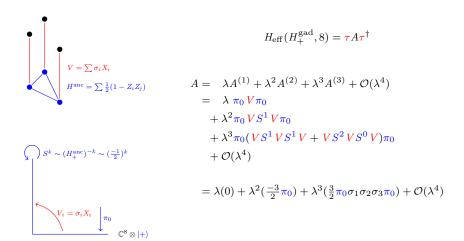


$$\begin{split} H_{\mathrm{eff}}(H_{+}^{\mathrm{gad}},8) &= \tau A \tau^{\dagger} \\ A &= \lambda A^{(1)} + \lambda^{2} A^{(2)} + \lambda^{3} A^{(3)} + \mathcal{O}(\lambda^{4}) \\ &= \lambda \pi_{0} V \pi_{0} \\ &+ \lambda^{2} \pi_{0} V S^{1} V \pi_{0} \\ &+ \lambda^{3} \pi_{0} (V S^{1} V S^{1} V + V S^{2} V S^{0} V) \pi_{0} \\ &+ \mathcal{O}(\lambda^{4}) \\ &= \lambda(0) + \lambda^{2} (\frac{-3}{2} \pi_{0}) \end{split}$$



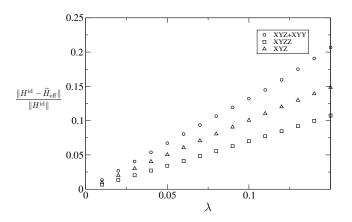


$$H_{\rm eff}(H_+^{\rm gad}, 8) = \frac{-3}{2} \lambda^2 \pi_{\rm low\ energy} + \frac{3}{2} \lambda^3 H^{\rm comp} \otimes \pi_0 + \mathcal{O}(\lambda^4)$$



$$\tilde{H}_{\mathrm{eff}}(H_{+}^{\mathrm{gad}}, 8, \Delta) = \frac{3}{2} \lambda^{3} H^{\mathrm{comp}} \otimes \pi_{0} + \mathcal{O}(\lambda^{4})$$

numerical simulation



$$\tilde{H}_{\mathrm{eff}}(H^{\mathrm{gad}}_+, 2^k, \Delta) = c \lambda^k H^{\mathrm{comp}} \otimes \pi_0 + \mathcal{O}(\lambda^{k+1})$$

spectral shift

$$H_{\mathrm{eff}}(H_{+}^{\mathrm{gad}},8) = \frac{-3}{2} \lambda^2 \pi_{\mathrm{low\ energy}} + \frac{3}{2} \lambda^3 H^{\mathrm{comp}} \otimes \pi_0 + \mathcal{O}(\lambda^4)$$

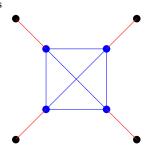
spectral shift

$$\mathit{H}_{\mathrm{eff}}(\mathit{H}_{+}^{\mathrm{gad}},8) = \frac{-3}{2} \lambda^2 \pi_{\mathrm{low\ energy}} + \frac{3}{2} \lambda^3 \mathit{H}^{\mathrm{comp}} \otimes \pi_0 + \mathcal{O}(\lambda^4)$$



spectral shift

$$\mathit{H}_{\mathrm{eff}}(\mathit{H}_{+}^{\mathrm{gad}},8) = \frac{-3}{2} \lambda^2 \pi_{\mathrm{low\ energy}} + \frac{3}{2} \lambda^3 \mathit{H}^{\mathrm{comp}} \otimes \pi_0 + \mathcal{O}(\lambda^4)$$



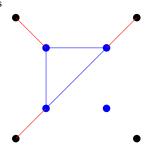
spectral shift

$$\mathit{H}_{\mathrm{eff}}(\mathit{H}_{+}^{\mathrm{gad}},8) = \frac{-3}{2} \lambda^2 \pi_{\mathrm{low\ energy}} + \frac{3}{2} \lambda^3 \mathit{H}^{\mathrm{comp}} \otimes \pi_0 + \mathcal{O}(\lambda^4)$$



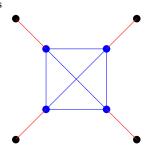
spectral shift

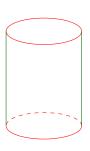
$$\mathit{H}_{\mathrm{eff}}(\mathit{H}_{+}^{\mathrm{gad}},8) = \frac{-3}{2} \lambda^2 \pi_{\mathrm{low\ energy}} + \frac{3}{2} \lambda^3 \mathit{H}^{\mathrm{comp}} \otimes \pi_0 + \mathcal{O}(\lambda^4)$$



spectral shift

$$\mathit{H}_{\mathrm{eff}}(\mathit{H}_{+}^{\mathrm{gad}},8) = \frac{-3}{2} \lambda^2 \pi_{\mathrm{low\ energy}} + \frac{3}{2} \lambda^3 \mathit{H}^{\mathrm{comp}} \otimes \pi_0 + \mathcal{O}(\lambda^4)$$



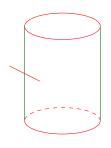


boundary $\sim \mathbb{R} \times \mathbb{S}^{d-1}$



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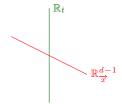
bulk $\sim \mathrm{AdS}_{d+1}$

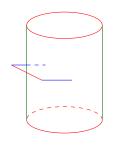


boundary $\sim \mathbb{R} \times \mathbb{S}^{d-1}$

boundary metric $\sim -dt^2 + dx^2$

bulk $\sim \mathrm{AdS}_{d+1}$



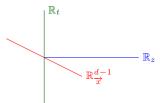


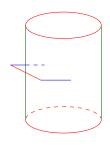
boundary $\sim \mathbb{R} \times \mathbb{S}^{d-1}$

boundary metric $\sim -dt^2 + dx^2$

bulk $\sim \mathrm{AdS}_{d+1}$

bulk metric $\sim \frac{dz^2 - dt^2 + dx^2}{z^2}$





boundary
$$\sim \mathbb{R} \times \mathbb{S}^{d-1}$$

boundary metric
$$\sim -dt^2 + dx^2$$

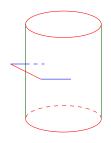
bulk
$$\sim AdS_{d+1}$$

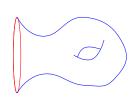
bulk metric
$$\sim \frac{dz^2 - dt^2 + dx^2}{z^2}$$

constant time slice (or wick rotation) gives hyperbolic geometry

Poincaré metric
$$\sim \frac{dz^2 + dx^2}{z^2}$$







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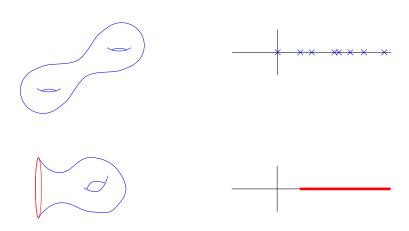
bulk metric
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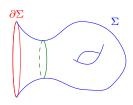
Poincaré metric
$$\sim \frac{dz^2 + dx^2}{z^2}$$

spectral theory

spectrum of laplacian $\boldsymbol{\Delta}$



geodesics



$$\zeta_{\text{Ruelle}}(s) = \prod_{\text{geodesics}} (1 - e^{-s(\text{lengths})})$$

 $\zeta_R(s) \sim s^{\dim H^1(\Sigma, \partial \Sigma)}$.torsion + h.o.t.



