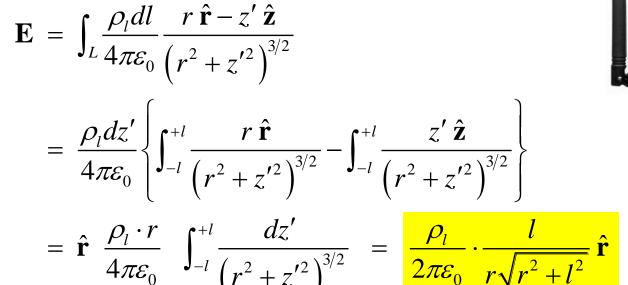
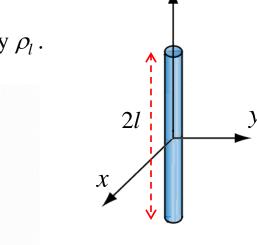
Example: Line Charge, Finite



Calculate the electric field **E** at any point *P* in the r- ϕ plane due to a line charge of length 2l, centered on the origin and extending along the z axis, with a constant charge density ρ_l .

$$\mathbf{R} - \mathbf{R}' = r \,\hat{\mathbf{r}} - z' \,\hat{\mathbf{z}}$$
, $dl = dz'$





$$\mathbf{E} = \frac{1}{4\pi\varepsilon_0} \int_L \rho_l dl \frac{\mathbf{R} - \mathbf{R}'}{|\mathbf{R} - \mathbf{R}'|^3}$$

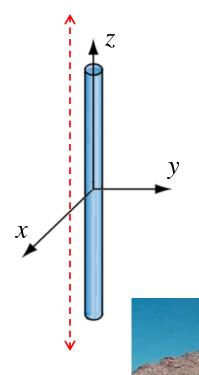
$$= \frac{1}{4\pi\varepsilon_0} \int_S \rho_s dS \frac{\mathbf{R} - \mathbf{R}'}{|\mathbf{R} - \mathbf{R}'|^3}$$

$$= \frac{1}{4\pi\varepsilon_0} \int_V \rho_v dV \frac{\mathbf{R} - \mathbf{R}'}{|\mathbf{R} - \mathbf{R}'|^3}$$

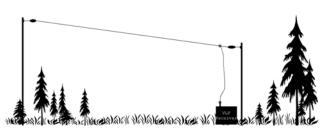
Example: Line Charge, Infinite



Calculate the electric field **E** at any point $P(r, \phi, z)$ due to a line charge of <u>infinite</u> length along the z axis, with a constant charge density ρ_l .



Very-Low-Frequency (VLF) antennas





Dr. Gregory J. Mazzaro Spring 2015

ELEC 318 – Electromagnetic Fields

Lecture 4(b)

Electrostatic Fields:

Gauss' Law & Electric Potential

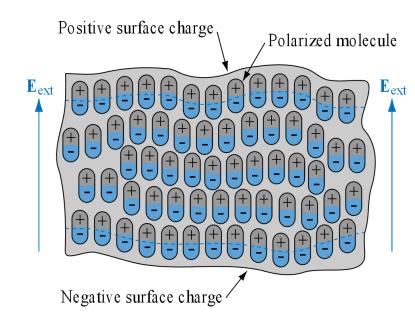
Electric Flux Density

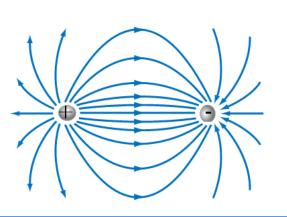


electric flux density, D (C/m²)

$$\mathbf{D} = \varepsilon \mathbf{E}$$

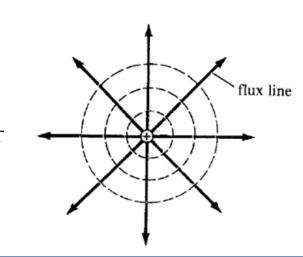
- -- similar to electric *field intensity*, but it encompasses permittivity, ε
- -- for a given **E** field, a measure of how *polarized* a particular material is
- -- useful for particular electrostatic calculations (e.g. for analyzing fields at boundaries between material media)
- -- 1 line of flux begins on a 1-C positive charge & terminates on a 1-C negative charge





near a single point charge q in free space $\varepsilon_0 \dots \mathbf{E} = \frac{q}{4\pi\varepsilon_0} \frac{\mathbf{R} - \mathbf{R'}}{|\mathbf{R} - \mathbf{R'}|^3}$

$$\mathbf{D} = \frac{q}{4\pi} \frac{\mathbf{R} - \mathbf{R}'}{\left|\mathbf{R} - \mathbf{R}'\right|^3}$$



Gauss' Law



Gauss's Law: -- 1 of 4 of Maxwell's Equations

-- 1 of 2 that govern the behavior of *electric fields*

-- can be written in 2 ways:

 $\nabla \cdot \mathbf{D} = \rho_{v}$

differential form

$$\int \rho_{v} dV = 0$$

$$\int_{V} \rho_{V} dV = Q$$

$$\oint_{S} \mathbf{D} \cdot d\mathbf{S} = \int_{V} \nabla \cdot \mathbf{D} \, dV$$

(Divergence Theorem)



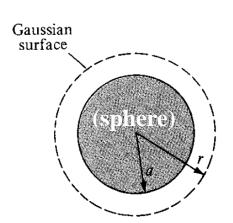
integral form

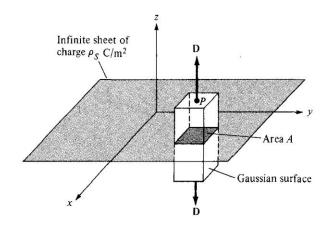
Q = charge contained in a "Gaussian" surface

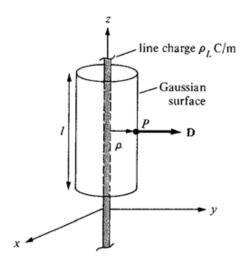
 \mathbf{D} = electric flux density

dS is normal to the surface and directed outward

 ρ_V = volume charge density (C/m³)







an alternative to Coulomb's Law for determining electric field, under symmetry

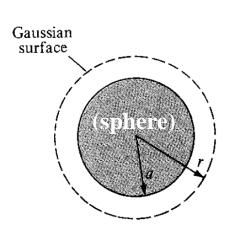
Gauss' Law, Integral Form: "How-To"

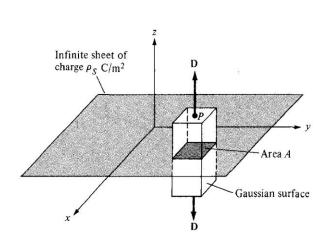


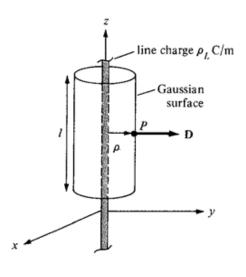
(1) Identify the symmetry of the problem (e.g. planar, cylindrical, spherical) .

$$Q = \oint_{S} \mathbf{D} \cdot d\mathbf{S}$$

- (2) Express **D** as a vector according to the symmetry (e.g. $\mathbf{D} = \mathbf{D}_z \mathbf{z}$, $\mathbf{D} = \mathbf{D}_R \mathbf{R}$).
- (3) Draw a Gaussian surface in accordance with the symmetry, which (a) passes through the observation point and (b) is perfectly perpendicular or parallel to **D**.
- (4) Solve for charge enclosed by the Gaussian surface as the (line/surface/volume) integral of charge.
- (5) Evaluate the flux integral, set the result equal to the charge enclosed, and solve for the scalar component(s) of \mathbf{D} .
- (6) Write **D** as a vector using (2). Solve for **E** by dividing the result by ε .







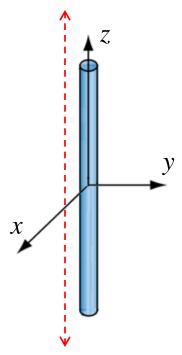
→ an alternative to Coulomb's Law for determining *electric field*, under <u>symmetry</u>

Example: Line Charge, Infinite



Calculate the electric field **E** at any point $P(r, \phi, z)$ due to a line charge of <u>infinite</u> length along the z axis, with a constant charge density ρ_l , in free space.

$$Q = \oint_{S} \mathbf{D} \cdot d\mathbf{S}$$



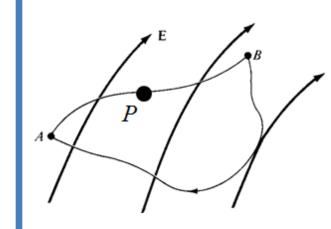
Electric (Scalar) Potential



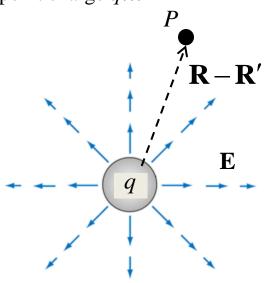
electric potential (V)

$$W = -\int_A^B \mathbf{F} \cdot d\mathbf{l} = -q \int_A^B \mathbf{E} \cdot d\mathbf{l}$$

$$V = \frac{W}{q} \quad \Rightarrow \quad V_{AB} = -\int_{A}^{B} \mathbf{E} \cdot d\mathbf{l}$$



for a point charge q...



$$V_{\text{point}} = -\int_{\infty}^{|\mathbf{R} - \mathbf{R}'|} \mathbf{E}_{\text{point}} \cdot d\mathbf{I}$$

$$= -\int_{\infty}^{|\mathbf{R} - \mathbf{R}'|} \left(\frac{q}{4\pi\varepsilon_0} \frac{\mathbf{R} - \mathbf{R}'}{|\mathbf{R} - \mathbf{R}'|^3} \right) \cdot \left(\frac{\mathbf{R} - \mathbf{R}'}{|\mathbf{R} - \mathbf{R}'|} \right) dr$$

$$= \frac{q}{4\pi\varepsilon_0} \cdot \left\{ \frac{1}{|\mathbf{R} - \mathbf{R}'|} - \frac{1}{\infty} \right\}$$

$$V_{\text{formulation of the point of$$

$$V_{\text{point}} = \frac{q}{4\pi\varepsilon_0 |\mathbf{R} - \mathbf{R}'|}$$

V for a charge distribution:

$$dV = \frac{dq}{4\pi\varepsilon_0 \left| \mathbf{R} - \mathbf{R'} \right|}$$

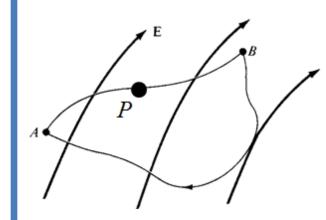
Electric (Scalar) Potential



electric potential (V)

$$W = -\int_A^B \mathbf{F} \cdot d\mathbf{l} = -q \int_A^B \mathbf{E} \cdot d\mathbf{l}$$

$$V = \frac{W}{q} \quad \Rightarrow \quad V_{AB} = -\int_{A}^{B} \mathbf{E} \cdot d\mathbf{l}$$



for a **static** electric field, V_{AB} is independent of the path chosen from A to B

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$
 no time variation (i.e. *static* fields)

$$V_{AP} - V_{PB} = V_{AB} \quad \forall P(x, y, z)$$

$$\nabla \times \mathbf{E} = 0$$

$$\oint_L \mathbf{E} \cdot d\mathbf{l} = 0$$

(Stokes' Theorem)

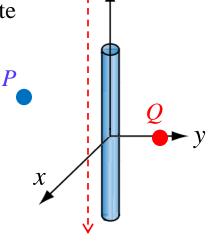
$$\int_{S} (\nabla \times \mathbf{E}) \cdot d\mathbf{S} = \oint_{L} \mathbf{E} \cdot d\mathbf{I}$$

Example: V from E & Work



Calculate the difference in electric potential from point P(x = 4, y = 0, z = 4) to point Q(x = 0, y = 2, z = 0) in the presence of a line charge of infinite length along the z axis with a constant charge density of 16 pC/m.

Determine the work required to move 3 mC from P to Q.



$$V_{PQ} = -\int_{P}^{Q} \mathbf{E} \cdot d\mathbf{l}$$

$$V = \frac{W}{q}$$

$$\mathbf{E}_{\text{line charge}}^{\text{infinite}} = \frac{\rho_l}{2\pi\varepsilon_0 r} \,\hat{\mathbf{r}}$$

Electric (Scalar) Potential

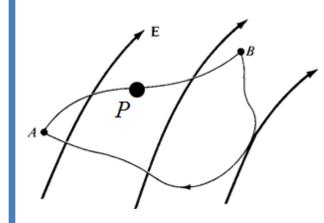


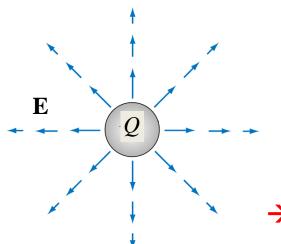
electric potential (V)

-- a measure of *work* required to move a unit charge through an **E** field: \mathbf{c}^B

$$W = -\int_A^B \mathbf{F} \cdot d\mathbf{l} = -q \int_A^B \mathbf{E} \cdot d\mathbf{l}$$

$$V = \frac{W}{q} \implies V_{AB} = -\int_{A}^{B} \mathbf{E} \cdot d\mathbf{l}$$





$$dV = -\mathbf{E} \cdot d\mathbf{l} = -E_{x}dx - E_{y}dy - E_{z}dz$$

$$dV = \frac{\partial V}{\partial x}dx + \frac{\partial V}{\partial y}dy + \frac{\partial V}{\partial z}dz$$

$$\mathbf{E} = -\frac{\partial E_{x}}{\partial x} - \frac{\partial E_{z}}{\partial y} - \frac{\partial E_{z}}{\partial z}$$

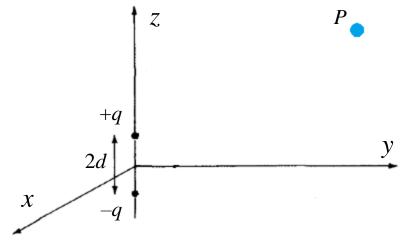
$$\mathbf{E} = -\nabla V \quad \text{``E from } V\text{''}$$

 \rightarrow It's usually easier to calculate V, then take its gradient rather than to compute E directly from Coulomb's Law.

Example: E from V, Dipole



For the dipole depicted, determine **E** at all points $P(R, \theta, \phi)$. You may assume that R >> 2d.



$$\mathbf{E} = -\nabla V$$

$$\nabla V = \frac{\partial V}{\partial x} \hat{\mathbf{x}} + \frac{\partial V}{\partial y} \hat{\mathbf{y}} + \frac{\partial V}{\partial z} \hat{\mathbf{z}}$$

$$= \frac{\partial V}{\partial r} \hat{\mathbf{r}} + \frac{1}{r} \frac{\partial V}{\partial \phi} \hat{\boldsymbol{\phi}} + \frac{\partial V}{\partial z} \hat{\mathbf{z}}$$

$$= \frac{\partial V}{\partial R} \hat{\mathbf{R}} + \frac{1}{R} \frac{\partial V}{\partial \theta} \hat{\boldsymbol{\theta}} + \frac{1}{R \sin \theta} \frac{\partial V}{\partial \phi} \hat{\boldsymbol{\phi}}$$

To be studied outside of class



- dipole moment
- Poisson's equation & Laplace's equation