

# Math 207B Partial Differential Equations: Ch 4

## Techniques for Solving important examples:

### Other ways to represent solutions

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#### Turing instability

#### Chapter 4.2 Similarity Solutions

When investigating partial differential equations, it is often profitable to look for specific solutions  $u$ , the form of which reflects various symmetries in the structure of the PDE. We have already seen this idea in our derivation of the fundamental solutions for Laplace's equation and the heat equation in §2.2.1 and §2.3.1 and our discovery of rarefaction waves for conservation laws in §3.4.4. Following are some further applications of this important method.

#### Chapter 4.3 Transformation Methods

In this section we develop some of the theory for the Fourier transform  $\mathcal{F}$ , the Radon transform  $\mathcal{R}$  and the Laplace transform  $\mathcal{L}$ . These provide extremely powerful tools for converting certain linear partial differential equations into either algebraic equations or else differential equations involving fewer variables.

**Chapter 4.3.1 Fourier Transform** In this section all functions are complex-valued, and  $\bar{\cdot}$  denotes the complex conjugate

**Definition** If  $u \in L^1(\mathbb{R}^n)$ , we define its Fourier transform  $\mathcal{F}u = \hat{u}$  by

$$\hat{u}(y) := \frac{1}{(2\pi)^{n/2}} \int_{\mathbb{R}^n} e^{-ixy} u(x) dx \quad (y \in \mathbb{R}^n)$$

and its inverse Fourier transform  $\mathcal{F}^{-1}u = \hat{u}$  by

$$\hat{u}(y) := \frac{1}{(2\pi)^{n/2}} \int_{\mathbb{R}^n} e^{ixy} u(x) dx \quad (y \in \mathbb{R}^n)$$

**Theorem 1 (Plancherel's Theorem).** Assume  $u \in L^1(\mathbb{R}^n) \cap L^2(\mathbb{R}^n)$ . Then  $\hat{u}, \bar{\hat{u}} \in L^2(\mathbb{R}^n)$  and

$$\|\hat{u}\|_{L^2(\mathbb{R}^n)} = \|\bar{\hat{u}}\|_{L^2(\mathbb{R}^n)} = \|u\|_{L^2(\mathbb{R}^n)}$$

**Theorem 2 (Properties of Fourier Transform)** Assume  $u, v \in L^2(\mathbb{R}^n)$ . Then

- (i)  $\int_{\mathbb{R}^n} u, \hat{v} dx = \int_{\mathbb{R}^n} \hat{u}, \bar{\hat{v}} dy$
- (ii)  $(D^\alpha u)^\wedge = (iy)^\alpha \hat{u}$  for each multiindex  $\alpha$  such that  $D^\alpha u \in L^2(\mathbb{R}^n)$ .
- (iii) If  $u, v \in L^1(\mathbb{R}^n) \cap L^2(\mathbb{R}^n)$ , then  $(u * v)^\wedge = (2\pi)^{n/2} \hat{u} \hat{v}$

(iv) Furthermore,  $u = (\hat{u})^\sim$ . Assertion (iv) is the Fourier inversion formula which represents a function  $u$  in terms of the exponential plane waves  $e^{ixy}$ , provided  $\hat{u} \in L^1(R^n)$

$$u(x) = \frac{1}{(2\pi)^{n/2}} \int_{R^n} e^{ixy} \hat{u}(y) dy$$

**Definition** The Radon transform  $\mathcal{R}u = \tilde{u}$  of a function  $u \in C_c^\infty(R^n)$  is

$$\tilde{u}(x, \omega) := \int \prod(s, \omega) u dS \quad (s \in R, \omega \in S^{n-1})$$

The term on the right is the integral over the plane  $\prod(s\omega)$  with respect to (n-1)-dimensional surface measure

**Theorem 3 (Properties of Radon transform)** Assume  $u \in C_c^\infty(R^n)$ . Then

- (i)  $\tilde{u}(-s, -\omega) = \tilde{u}(s, \omega)$
- (ii)  $(D^\alpha u)^\sim = \omega^\alpha \frac{\partial^{|\alpha|}}{\partial s^{|\alpha|}} \tilde{u}$  for each multiindex  $\alpha$
- (iii)  $(\Delta u)^\sim = \omega^\alpha \frac{\partial^2}{\partial s^2} \tilde{u}$

If  $u = 0$  in  $R^n - B(0, \mathcal{R})$ , then  $\tilde{u}(s, \omega) = 0$  for  $|s| \geq \mathcal{R}$

**Theorem 4 (Radon and Fourier transforms)** Assume that  $u \in C_c^\infty(R^n)$ . Then

$$\bar{u}(r, \omega) := \int_R \tilde{u}(s, \omega) e^{-irs} ds = (2\pi)^{n/2} \hat{u}(r\omega) \quad (r \in R, \omega \in S^{n-1})$$

where  $\hat{u} = \mathcal{F}u$  is the Fourier transform.

**Theorem 5 (Inverting the Radon transform)**

- (i) We have

$$u(x) = \frac{1}{2(2\pi)^n} \int_R \int_{S^{n-1}} \bar{u}(r, \omega) r^{n-1} e^{ir\omega \cdot x} dS dr$$

the function  $\bar{u}$  defined by (31)

- (ii) If  $n = 2k + 1$  is odd, then

$$u(x) = \int_{S^{n-1}} r(x \cdot \omega, \omega) dS$$

for

$$r(s, \omega) := \frac{(-1)^k}{2(2\pi)^{2k}} \frac{\partial^{2k}}{\partial s^{2k}} \tilde{u}(s, \omega).$$

Formulas (33) and (34) provide an elegant and useful decompositions of  $u$  into plane waves.

**Definition** If  $u \in L^1(R_+)$ , we define its Laplace transform  $\mathcal{L}u = u\#$  to be

$$u\#(s) := \int_0^\infty e^{-st} u(t) dt \quad (s \geq 0)$$

**Chapter 4.4 Converting Nonlinear into Linear PDE**

**Chapter 4.5 Asymptotics**

**Chapter 4.6 Power Series**

**Definition** We say the surface  $\Gamma$  is noncharacteristic for the partial differential equation (1) provided

$$\sum_{|\alpha|=k} a_\alpha \nu^\alpha \neq 0 \quad \text{on } \Gamma$$

for all values of the arguments of the coefficients  $a_\alpha (|\alpha| = k)$

**Theorem 1 (Cauchy data and Noncharacteristic surfaces)** Assume that  $\Gamma$  is noncharacteristic for the PDE (1). Then if  $u$  is a smooth solution of (1) and  $u$  satisfies the Cauchy conditions (2), we can uniquely compute all the partial differential equations of  $u$  along  $\Gamma$  in terms of  $\Gamma$ , the function  $g_0, \dots, g_{k-1}$  and the coefficients  $a_\alpha (|\alpha| = k), a_0$

**Definition** A function  $f : R^n \rightarrow R$  is called (real) analytic near  $x_0$  if there exist  $r > 0$  and constants  $\{f_\alpha\}$  such that

$$f(x) = \sum_\alpha f_\alpha (x - x_0)^\alpha \quad (|x - x_0| < r),$$

the sum taken over all multiindex  $\alpha$

**Definition** Let

$$f = \sum_\alpha f_\alpha x^\alpha, g = \sum_\alpha g_\alpha x^\alpha$$

be two power series. We say  $g$  majorizes  $f$ , written

$$g \gg f$$

provided

$$g_\alpha \leq |f_\alpha| \quad \text{for all multiindices } \alpha$$

**Theorem 2 (Cauchy-Kovalevskaya Theorem)** Assume  $\{B_j\}_{j=1}^{n-1}$  and  $c$  are real analytics functions. Then there exist  $r > 0$  and a real analytic function

$$u = \sum_\alpha u_\alpha x^\alpha$$

solving the boundary-value problem (15)