Networks and Random Processes Assignment 2

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1 Kingman's Coalescent

1.1 A

 N_t is the number of particles at time t with $N_0=L$. The process $(N_t:t\geq 0)$ has the state space $\{1,...,L\}$

1.1.1 Transition Rate of the process

$$r(n, n-1) = \binom{L}{2}, \ n \ge 2$$

QUESTION - WHAT ABOUT SAME STATE? r(n,n) = QUESTION - WHAT ABOUT OTHER STATES - HOW TO WRITE IT? $r(n,y) = , \ y \neq n,n-1$

1.1.2 Generator

This is a jump process, so the generator is

$$(\mathcal{L}f)(x) = \int_{\Re} r(x, y)[f(y) - f(x)]dy$$

For this process

$$(\mathcal{L}f)(n) = r(n, n-1)(f(n-1) - f(n))$$

$$(\mathcal{L}f)(n) = \binom{n}{2}(f(n-1) - f(n))$$

1.1.3 Master Equation

The master equation is

$$\frac{d}{dt}\pi_t(n) = \pi_t(n+1)r(n+1,n) - \pi_t(n)r(n,n-1)$$

$$\frac{d}{dt}\pi_t(n) = \pi_t(n+1)\binom{n+1}{2} - \pi_t(n)\binom{n}{2}$$

QUESTION - IS THIS RIGHT? QUESTION - IS THE NOTATION OKAY? QUESTION - WHAT ABOUT EDGES?

1.1.4 Ergodicity

The process is ergodic.

1.1.5 Absorbing States

The unique absorbing state is N=1.

1.1.6 Stationary Distributions

Let a distribution $\pi = [N = 1, N = 2, ..., N = L]$

The unique stationary distribution is

$$\pi_0 = [1, 0, ..., 0]$$

1.2 B - Mean Time to Asorption

The rate of coalescence, ie moving to the next state, for each state is

$$\lambda_n = r(n, n-1) = \binom{n}{2} = \frac{n(n-1)}{2}$$

The times in each state are expnentially distributed as

$$f_t(n) = \binom{n}{2} e^{-\binom{n}{2}t}$$

The expected time in each state, or the waiting time, is given by

$$\beta_n = \frac{1}{\lambda_n} = \frac{2}{n(n-1)}$$

The expected time to absorption is the sum of the expected waiting times in each of the states

$$E(T) = \sum_{n=2}^{L} \frac{2}{n(n-1)}$$

Bringing the 2 outside of the summand and splitting up into partial fractions

$$E(T) = 2\sum_{n=2}^{L} \frac{1}{n-1} - \frac{1}{n}$$

$$E(T) = 2(\frac{1}{1} - \frac{1}{2} + \frac{1}{2} - \frac{1}{3} + \frac{1}{3} + \dots - \frac{1}{L-1} + \frac{1}{L+1} - \frac{1}{L})$$

All but the first and last terms cancel giving

$$E(T) = 2(1 - \frac{1}{L})$$

1.3 C - Rescaled process

Rescale the process to N_t/L

$$(\mathcal{L}^L f)(n/L) = \frac{1}{L} \binom{n}{2} (f(\frac{n-1}{L}) - f(\frac{n}{L}))$$

Taylor expand and let $x = \frac{n}{L}$:

$$(\mathcal{L}^{L}f)(x) = \frac{1}{L}\frac{n(n-1)}{2}(f(x) - \frac{1}{L}f'(x) + \frac{1}{L^{2}}f''(x) + O(\frac{1}{L^{3}}) - f(x))$$

Cancel terms, substitute n = Lx and rearrange

$$(\mathcal{L}^{L}f)(x) = (\frac{x^{2}}{2} - \frac{x}{2L})(-f'(x) + \frac{1}{L}f''(x) + O(\frac{1}{L^{2}}))$$

$$\lim_{L \to \infty} (\mathcal{L}^{L}f)(x) = \frac{-x^{2}}{2}f'(x)$$

 $\lim_{L \to \infty} (L^{-}f)(x) = \frac{1}{2}f(x)$

The state space goes from $\frac{1}{L}$ to $\frac{L}{L}$. As $L \to \infty$, this is (0,1]. The initial condition is $X_0 = \frac{L}{L} = 1$,

1.3.1 Deterministic

The generator has no diffusion term, only drift. So there is no variance in the process and it must be entirely deterministic.

1.3.2 Computing X_t

$$\frac{d}{dt}E(X_t) = E(-\frac{X_t^2}{2})$$

 X_t is deterministic, so

$$\frac{d}{dt}X_t = -\frac{X_t^2}{2}$$

$$\frac{dX_t}{X_t^2} = -\frac{1}{2}dt$$

$$\frac{1}{X_t} = \frac{1}{2}t + c$$

$$X_t = \frac{1}{\frac{t}{2} + c}$$

From the initial conditions, $X_0 = 1$, c = 1.

$$X_t = \frac{1}{\frac{t}{2} + 1}$$

1.3.3 Comparing to result from b

The result from part b was

$$E(T) = 2(1 - 1T)$$

For the rescaled process:

$$X_t = \frac{1}{\frac{t}{2} + 1}$$

Consider when it reaches the absorbing state, i.e. $X_t = \frac{1}{L}$:

$$\frac{1}{L} = \frac{1}{\frac{t}{2} + 1}$$

$$L = \frac{t}{2} + 1$$

$$t = 2L(1 - \frac{1}{L})$$

This is the same as the result from part b with a factor of L. This arises because we slowed down time by a factor of L during the rescaling process.

1.4 D - Simulation

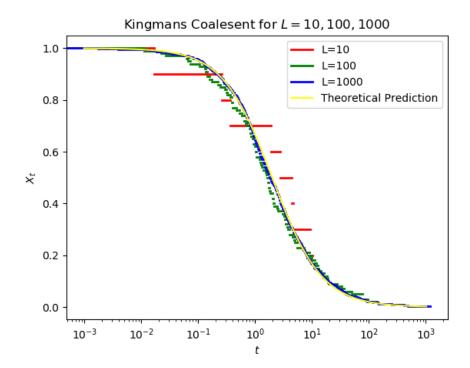


Figure 1: The Kingsman's coalescent rescaled process simulated for L=10, L=100, L=1000. The theoretical prediction of $X_t=\frac{1}{\frac{t}{2}+1}$ is also shown.

Figure 1 shows sample paths of the rescaled Kingsman's coalescent process for L=10,100 and 1000. The theoretical prediction for $L \to \infty$ is also shown. The simulations line up very well with the theoretical prediction.

2 Ornstein-Uhlenbeck process

2.1 A

The mean:

$$\frac{dm}{dt} = \frac{d}{dt}E(X_t) = E(-\alpha X_t) = -\alpha E(X_t)$$

$$\frac{dm}{dt} = -\alpha m(t)$$

 X_t^2 :

$$\frac{d}{dt}E(X_t^2) = E(-\alpha 2X_t^2 + \sigma^2)$$

$$\frac{d}{dt}E(X_t^2) = -2\alpha E(X_t^2) + \sigma^2$$

The variance:

$$v(t) = E(X_t^2) - m(t)^2$$

$$\frac{dv}{dt} = \frac{d}{dt}E(X_t^2) - 2m\frac{dm}{dt}$$

$$\frac{dv}{dt} = -2\alpha E(X_t^2) + \sigma^2 + 2\alpha m^2$$

2.2 B

2.2.1 Solution of m(t)

Solving for m(t), and considering the initial conditions, $c=x_0$

$$m(t) = x_0 e^{-\alpha t}$$

2.2.2 Solution of v(t)

Solving for v(t)

$$\frac{dv}{dt} = -2\alpha(v+m^2) + \sigma^2 + 2\alpha m^2$$

$$\frac{dv}{dt} + 2\alpha v = \sigma^2$$

Homogenous solution:

$$v_h o m = c_2 e^{-2\alpha t}$$

Particular solution:

$$v_p art = \frac{\sigma^2}{2\alpha}$$

Full solution:

$$v(t) = c_2 e^{-2\alpha t} + \frac{\sigma^2}{2\alpha}$$

From initial conditions, v(0) = 0:

$$v(t) = \frac{\sigma^2}{2\alpha} (1 - e^{-2\alpha t})$$

2.2.3 Distribution

As a Gaussian process, the distribution is fully described by the mean and variance.

$$f(X_t) = \frac{1}{\sqrt{\frac{2\pi\sigma^2(1 - e^{-2\alpha t})}{2\alpha}}} exp(\frac{-\alpha(x - x_0 e^{-\alpha t})^2}{\sigma^2(1 - e^{-2\alpha t})})$$

IS THAT RIGHT?

2.2.4 Stationary Distribution

Given enough time, the process will converge to a Gaussian stationary distribution.

$$\lim_{t \to \infty} m(t) = 0$$

$$\lim_{t\to\infty}v(t)=\frac{\sigma^2}{2\alpha}$$

The stationary distribution is $\sim N(0,\frac{\sigma^2}{2\alpha})$

$$f_0(X_t) = \frac{1}{\sqrt{\frac{\pi\sigma^2}{\alpha}}} exp(\frac{-\alpha x^2}{\sigma^2})$$

2.3 Simulation

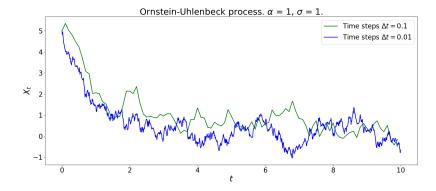


Figure 2: The Ornstein-Uhlenbeck process with $\alpha=1,\ \sigma^2=1$ and $X_0=5$. Simulated for 10 seconds with timsteps $\Delta t=0.1$ and 0.01

Figure 2 shows the Ornstein-Uhlenbeck process simulated. The process begins at $X_0=5$, and expereinces a "force" pulling it towards zero. As time progresses,

the process moves towards $X_t = 0$ and the noise term begins to dominate the behaviour. Both choices of timestep give a similar result.

3 Moran Model and Wright-Fisher diffusion

3.1 A

3.1.1 State space

Let the total possible L types be

$$T = \{1, 2, ..., L\}$$

Each of the L individuals can have any of those types, so the state space is

$$S = \{1, 2, 3, ..., L\}^{L}$$

3.1.2 Irreducibility

It is not irreducible because there are absorbing states.

3.1.3 Stationary distributions

The absorbing states are where all individuals have the same type

$$x_k = [k, k, ..., k] \quad \forall k \in \{1, 2, ..., L\}$$

The stationary distributions are any linear combination of the absorbing states that sum to 1.

$$\pi(y) = \sum_{k=1}^{L} \alpha_k \pi_k(y)$$

$$\sum_{k=1}^{L} \alpha_k = 1$$

The coefficients, α_k , are determined by the initial conditions and can be thought of as the probability of each type "winning" and taking over all individuals.

3.2 B

3.2.1 Markov process

 $N_t: t \geq 0$ is a Markov process. It's future distribution is determined only by it's current state, not the specific history.

3.2.2 State space

Each type can have any integer between 0 and L types, and there are L types, so the state space is

$$S = \{0, 1, 2, ..., L\}^{L}$$

3.2.3 Generator

For each type, we assume that the number of individuals of that type, n, can only increase and decrease by 1 at one moment in time, i.e only one event happens at a time. The rates of gain and loss can be described as:

$$r(n, n+1) = \frac{n(L-n)}{L-1}$$

$$r(n, n-1) = \frac{n(L-n)}{L-1}$$

These are symmetrical.

This is a jump process, so the generator is

$$(\mathcal{L}f)(x) = \int_{\Re} r(x, y)[f(y) - f(x)]dy$$

For this process

$$(\mathcal{L}f)(n) = r(n, n-1)(f(n-1) - f(n)) + r(n, n+1)(f(n+1) - f(n))$$

$$(\mathcal{L}f)(n) = \frac{(L-n)n}{L-1}(f(n-1) + f(n+1) - 2f(n))$$

3.2.4 Irreducibility

The process is not irreducible because there are absorbing states.

3.2.5 Stationary dsitributions

The absorbing states are where one type has $N_k = L$ and the rest $N_k = 0$. i.e

$$\pi_k(y) = \delta_{y,k} = \begin{cases} L & , & y = k \\ 0 & , & \text{otherwise} \end{cases}$$

?? IS THAT RIGHT?

3.2.6 Limiting distribution

As $t \to \infty$ with initial condition $N_0 = 1$. By symmetry each type has the same chance of "winning", therefore the limiting distribution is

$$\pi_{\infty} = [\frac{L-1}{L}, 0, 0, ..., 0, \frac{1}{L}]$$

3.3 C

3.3.1 $m_1(t)$

$$\frac{d}{dt}(E(N_t)) = E(\frac{n(L+n)}{L-1}(n+1+n-1-2n))$$

$$\frac{d}{dt}(E(N_t)) = 0$$

 $N_0 = n$, therefore $E(N_t) = n$.

3.3.2
$$m_2(t)$$

$$\frac{d}{dt}(E(N^2)) = E(\frac{n(L+n)}{L-1}((n+1)^2 + (n-1)^2 - 2n^2)$$

$$\frac{d}{dt}(E(N^2)) = E(\frac{n(L+n)}{L-1}(n^2 + 2n + 1 + n^2 - 2n + 1 - 2n^2)$$

$$\frac{d}{dt}(E(N^2)) = E(\frac{2n(L+n)}{L-1})$$

$$\frac{d}{dt}(E(N^2)) = \frac{2L}{L-1}n + \frac{2}{L-1}E(N^2)$$

Solution in the form:

$$E(N^2) = Ae^{\frac{-2t}{L-1}} + B$$

From initial condition, $A = N_0^2$

Substitute everything into the original differential equation:

$$\frac{-2n}{L-1}e^{\frac{-2t}{L-1}} = \frac{2Ln}{L-1} + \frac{2n}{L-1}e^{\frac{-2t}{L-1}} + \frac{2B}{L-1}$$
$$B = Ln(1 - ne^{\frac{-2t}{L-1}})$$

$$E(N^2) = n^2 e^{\frac{-2t}{L-1}} + Ln(1 - e^{\frac{-2t}{L-1}})$$

QUESTION - I HAVE A FACTOR OF n in front of exponential in the bracket for B - I don't think that I should have

3.3.3 Absorption probabilities

By symmetry, each individual has an equal chance of winning?? NOT SURE ABOUT THIS

3.3.4 Absorption time scales with system size L

The expected value of $E(N_t) = n$.

The variance is given by

$$Var(N_t) = n^2 e^{\frac{-2t}{L-1}} + Ln(1 - ne^{\frac{-2t}{L-1}}) - n^2$$

$$Var(N_t) = (Ln - n^2)(1 - e^{\frac{-2t}{L-1}})$$

If we assume that we start in the middle, so $n = \frac{L}{2}$

$$Var(N_t) = \frac{L^2}{4}(1 - e^{\frac{-2t}{L-1}})$$

As the standard deviation grows, the probability of being absorbed at either N=0 or N=L also grows. The rate at which the standard deviation grows with time scales approximately with the length of the state space.

3.4 D

3.4.1 Rescaling limit

$$(\mathcal{L}f)(n) = \frac{(L-n)n}{L-1}(f(n+1) + f(n-1) - 2f(n))$$

Rescale by dividing state space by L, multiplying time by L^{α} and changing the variable to $x = \frac{n}{L}$:

$$(\mathcal{L}f)(x) = \frac{(L-xL)xL}{L-1}L^{\alpha}(f(x+\frac{1}{L})+f(x+\frac{1}{L})-2f(x))$$

Taylor expand and simplify:

$$(\mathcal{L}f)(x) = \frac{(L-xL)xL}{L-1}L^{\alpha}(\frac{1}{L^{2}}f''(x) - O(\frac{1}{L^{3}})$$

$$(\mathcal{L}f)(x) = \frac{L^{\alpha}}{L} \frac{xL^2 - x^2L^2}{1 - \frac{1}{L}} (\frac{1}{L^2}f''(x) - O(\frac{1}{L^3})$$

Set $\alpha = 1$ to get a non-trivial scalining limit:

$$(\mathcal{L}f)(x) = (\frac{x - x^2}{1 - \frac{1}{L}})f''(x) - O(\frac{1}{L})$$

Take limit $L \to \infty$:

$$(\mathcal{L}^L f)(x) = \lim_{L \to \infty} (\mathcal{L}f)(x) = (x - x^2) f''(x)$$
$$(\mathcal{L}^L f)(x) = x(1 - x) f''(x)$$

3.4.2 Fokker-Planck Equation

$$\frac{\partial}{\partial t}p_t(x,y) = \frac{\partial^2}{\partial u^2}(y(1-y)p_t(x,y))$$

WHAT DOES THIS MEAN?

3.5 E

3.5.1 m(t)

$$\frac{d}{dt}E(M) = 0$$

No change in $E(M_t)$ with time. It's a diffusion process with no fridt, so $E(M_t) = 0$

3.5.2 v(t)

$$v(t) = E(M_t^2)$$

Let $f(m) = m^2$:

$$\frac{d}{dt}E(M^2) = E(2M - 2M^2) = -2E(M^2)$$

$$v(t) = E(M^2) = c_1 e^{-2t}$$

Set $v(0) = v_0$

$$v(t) = E(M^2) = v_0 e^{-2t}$$

 $v(t) \to 0$ as $t \to \infty$. Process is absorbed on the edges, so variance drops to zero

QUESTION - IS THIS RIGHT??

It is not Gaussian because it has a limited state space, Gaussians go to infinity.

3.6 F - Simulation

N_t Number of individuals of each species out of total population size L=40 t=40 t=40

Figure 3: The Moran model process simulated for a population of L=40. N_t is the number of individuals of each type, plotted against continuous time.

Figure 3 shows a simulation of the Moran model process. The number of each individuals of each type is shown a a function of continuous time. Gains in one type must conincide with losses of another type, as they must always sum to 40. This can be seen in the symmetry in the patterns between the two surviving types after around t=18.