The Quantum Linear Problem

And its Quantum Linear Solver (QLS) Algorithm

A PH4415 Final Year Project Presentation By

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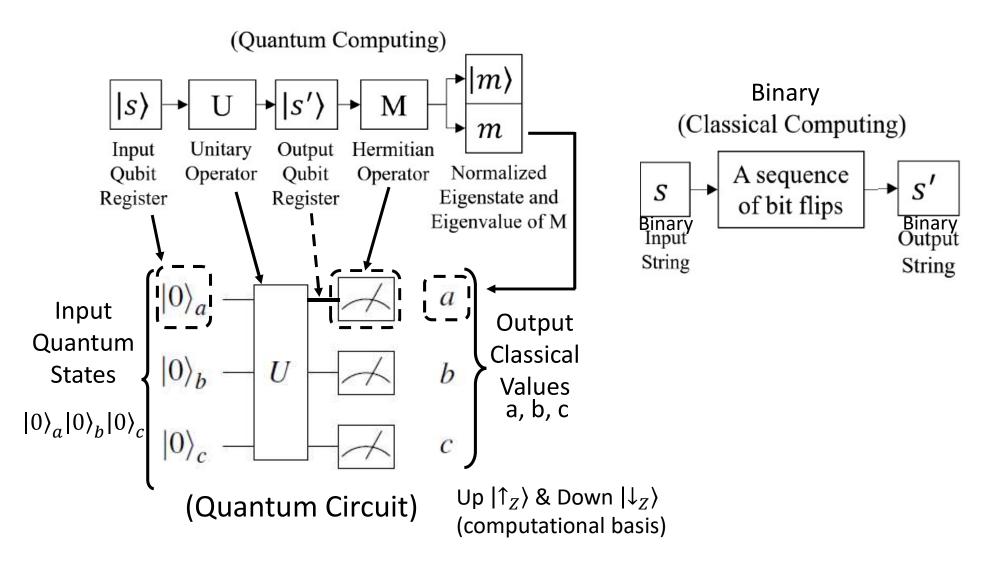
Contents

- 0. Quantum Computation
- 1. Quantum Linear Problem
- 2. Quantum Algorithmic Ideas
- 3. Original HHL QLS Algorithm
- 4. Full & Simplified QLS Algorithm
- 5. Worked Examples and Experiment Results

Abbreviations

- HHL: Aram Harrow, Avinatan Hassidim and Seth Lloyd
- QLS: Quantum Linear Solver

Quantum Computation



The Linear Problem

Resolving $|x\rangle = M^{-1}|b\rangle$

Problem: M^{-1} is Hermitian (still a measurement ops)

- Does not preserve norm → Mathematically incorrect
- Results in a measurement collapse → Physically incorrect

Solution:

• Renormalize : $|x\rangle = \frac{1}{d}M^{-1}|b\rangle = \frac{1}{d}\sum_{j=1}^{2^n} \frac{\beta_j}{\lambda_j}|m_j\rangle$

d: norm factor.

 $\left|m_{j}\right\rangle$: orthonormalized eigenstates of M.

 β_i : complex coefficients.

 λ_i : eigenvalues of M.

• Get $|x\rangle = \frac{1}{d} \sum_{j=1}^{2^n} \frac{\beta_j}{\lambda_j} |m_j\rangle$ by a measurement collapse of $U|b\rangle$.

> as implicitly developed by HHL (Algorithmic Ideas)

Goal: Get $|x\rangle = \frac{1}{d} \sum_{j=1}^{2^n} \frac{\beta_j}{\lambda_j} |m_j\rangle$ by a measurement collapse of $U|b\rangle$.

Algorithmic Ideas First Idea – a simple algorithm

- 1. Append an ancilla $|0\rangle_a$ to input state $|b\rangle_m \rightarrow |b\rangle_m |0\rangle_a$
- 2. Apply some unitary U (should contain M),

Change the amplitudes of
$$|m_j\rangle$$
 in $|b\rangle$ to $\frac{\beta_j c}{\lambda_j}$

$$U|b\rangle_m|0\rangle_a = \sum_{j=1}^{2^n} \frac{\beta_j C}{\lambda_j} |m_j\rangle_m|1\rangle_a + \sum_{j=1}^{2^n} \beta_j \sqrt{1 - \left(\frac{C}{\lambda_j}\right)^2} |m_j\rangle_m|0\rangle_a$$

$$d|x\rangle_m|1\rangle_a \qquad \text{Some other orthogonal state}$$

where C is a free real parameter (does not affect state)

- 3. Measure ancilla qubit in computational basis $\{|0\rangle, |1\rangle\}$.
- 4. Get $|x\rangle_m$ if ancilla = $|1\rangle_a$ with probability $\sum_{j=1}^{2^n} \left| \frac{\beta_j c}{\lambda_j} \right|^2$.

Algorithmic Ideas Second Idea

To implement U (Quantum If-Else):

Requires information about M

$$C[R_{y}(\theta_{j})]$$

• Apply controlled- $R_y(\theta_j)$ gate on $|0\rangle_a$, conditioned on eigenstates $|m_j\rangle_m$, where $\frac{\theta_j}{2} = \sin^{-1}\left(\frac{c}{\lambda_j}\right)$ (???).

$$|b\rangle_{m}|0\rangle_{a} \xrightarrow{C[R_{y}(\theta_{j})]} \sum_{j=1}^{2^{n}} \beta_{j} \sin\left(\frac{\theta_{j}}{2}\right) |m_{j}\rangle_{m} |1\rangle_{a} + \sum_{j=1}^{2^{n}} \beta_{j} \cos\left(\frac{\theta_{j}}{2}\right) |m_{j}\rangle_{m} |0\rangle_{a}$$

 $d|x\rangle_m|1\rangle_a$

Some other orthogonal state

Controlled- $R_y(\theta_i)$ on $|b\rangle|a\rangle$:

(Do in a quantum superposition)

- If $|b\rangle = |m_1\rangle$, then $R_{\nu}(\theta_1)|a\rangle$
- Else if $|b\rangle = |m_2\rangle$, then $R_y(\theta_2)|a\rangle$
- Else if ...

Note:
$$R_y(\theta_j) = \begin{pmatrix} \cos\left(\frac{\theta_j}{2}\right) & -\sin\left(\frac{\theta_j}{2}\right) \\ \sin\left(\frac{\theta_j}{2}\right) & \cos\left(\frac{\theta_j}{2}\right) \end{pmatrix}$$

(in the computational basis)

(Easier if using Boolean conditions and not eigenvalues $|m_i
angle$)

Algorithmic Ideas Third Idea

eigenvalue qubits
$$2 \rightarrow 10 \rightarrow |10\rangle$$
 binary

 $|\lambda_j\rangle_e$: qubit binary representation of eigenvalues of M. (e.g. Boolean qubits) $|1\rangle_e$: True, $|0\rangle_e$: False

Use $|\lambda_j\rangle_e$ (in $|1\rangle_e$, $|0\rangle_e$) as your Quantum If-Else conditions:

- Apply $C[R_y(\theta_j)]$ on $|0\rangle_a$ conditioned on $|\lambda_j\rangle_e$ in an appended eigenvalue register. \rightarrow (Ancillary Rotation step, AR)
- Get $|\lambda_j\rangle_e$ by Quantum Phase Estimation (QPE). $|m_j\rangle_m |\mathbf{0}\rangle_e \stackrel{QPE}{\longrightarrow} |m_j\rangle_m |\lambda_j\rangle_e$

$$|b\rangle_{m}|\lambda_{j}\rangle_{e}|0\rangle_{a} \xrightarrow{AR} \sum_{j=1}^{2^{n}} \beta_{j} \sin\left(\frac{\theta_{j}}{2}\right)|m_{j}\rangle_{m}|\lambda_{j}\rangle_{e}|1\rangle_{a} + \sum_{j=1}^{2^{n}} \beta_{j} \cos\left(\frac{\theta_{j}}{2}\right)|m_{j}\rangle_{m}|\lambda_{j}\rangle_{e}|0\rangle_{a}$$

$$\sum_{j=1}^{2^{n}} \frac{\beta_{j}}{\lambda_{j}}|m_{j}\rangle_{m}|\lambda_{j}\rangle_{e}|1\rangle_{a} \xrightarrow{\text{uncompute}} |\lambda_{j}\rangle_{e}$$

$$|\lambda_{j}\rangle_{e}$$

$$|\lambda_{j}\rangle_{e}$$

Algorithmic Ideas

Summary

1. Get $|x\rangle = \frac{1}{d} \sum_{j=1}^{2^n} \frac{\beta_j}{\lambda_j} |m_j\rangle$ by a measurement collapse of $U|b\rangle$.

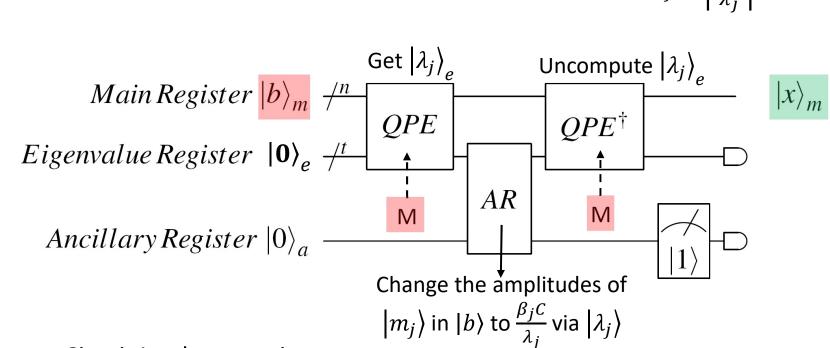
Recall: (Main Purpose of U)

Change the amplitudes of $|m_j\rangle$ in $|b\rangle$ to $\frac{\beta_j c}{\lambda_j}$.

- 2. Construct U operator Using Quantum If-Else.
- 3. Use qubit-binary $\left|\lambda_{j}\right\rangle_{e}$, eigenvalues of M as your lf-Else conditions.

Original HHL QLS Algorithm

For 2^n quantum linear problem $d|x\rangle = M^{-1}|b\rangle$, given M and $|b\rangle$, run the following quantum algorithm to obtain $|x\rangle$ with probability $\sum_{j=1}^{2^n} \left|\frac{\beta_j c}{\lambda_j}\right|^2$.



Quantum Circuit Implementation:

- QPE: (Quantum Phase Estimation) By Alexey Kitaev and Richard Cleve et al.
- AR: (Ancillary Rotation) Not given by HHL.

Original QLS Problems

- HHL's main purpose was to prove its exponential computation speedup by
 - Assuming M is sparse.
 - Assuming input $|b\rangle$ can be efficiently prepared.
 - Treating QPE and AR as black boxes oracle subroutines.
- However, the original paper did not elaborate on implementing AR subroutine.
- Naïve and Direct AR requires one to know the eigenvalues λ_j to calculate rotation angles $\frac{\theta_j}{2} = \sin^{-1}\left(\frac{C}{\lambda_j}\right)$
- QLS modifications (Rule-breaking Implementations):
 - Eigenvalue Marking Modification. → Simplified QLS
 - 1st order Taylor Approximation of AR. \rightarrow Full QLS

Motivation:

- λ_i Eigenvalues are not given by the problem.
- Need to find an AR method that implements controlled- $R_y(\theta_j)$ on $|0\rangle_a$ conditioned on $|\lambda_j\rangle_e$ without need of any information of M or $|b\rangle$ except C.

Solution:

 Use 1st order Taylor Approximation of AR inspired by Yudong Cao et al

$$\theta_j = 2 \sin^{-1} \left(\frac{C}{\lambda_j} \right) \approx 2 \left(\frac{C}{\lambda_j} \right), \qquad \left| \frac{C}{\lambda_j} \right| \ll 1$$

- Assume eigenvalues λ_i are powers of two.
 - (Practical and Easier Implementation only, otherwise not needed).

1st order approximation of Ancillary Rotation (AR)

- Idea: Break up Quantum If-Else, $C[R_y(\theta_j)]$, so that angles need not be calculated
- Apply 1st order approximation: $\theta_j = 2 \sin^{-1} \left(\frac{c}{\lambda_j} \right) \approx 2 \left(\frac{c}{\lambda_j} \right)$.
- Then, Rotation-Y gate is approximated: $R_y(\theta_j) \approx R_y(\frac{2C}{\lambda})$.
- Let $C=\frac{\pi}{2^r}$ and $y=\frac{1}{\lambda}$, where in binary $y=y_t2^0+\cdots+y_j2^{j-t}+\cdots+y_12^{1-t}$. *Introduce a new parameter r $y_j=\{0,1\}$
- Substitute into $R_{\mathcal{Y}}\left(\frac{2C}{\lambda_i}\right)$ and <u>break up</u> using binary,

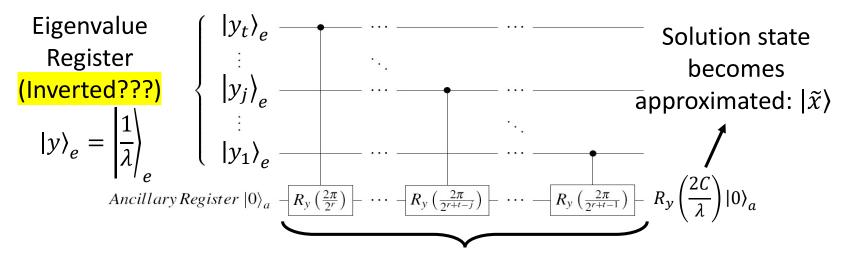
$$R_{y}\left(\frac{2C}{\lambda}\right) = R_{y}\left(\frac{2\pi}{2^{r}}y_{t} \cdot 2^{0}\right) \cdots R_{y}\left(\frac{2\pi}{2^{r}}y_{j} \cdot 2^{j-t}\right) \cdots R_{y}\left(\frac{2\pi}{2^{r}}y_{1} \cdot 2^{1-t}\right)$$

1st order approximation of AR

• Create a broken up and approximated Quantum If-Else,

$$C\left[R_{y}\left(\frac{2C}{\lambda}\right)\right] = \cdots C\left[R_{y}\left(\frac{2\pi}{2^{r}}y_{j}\cdot 2^{j-t}\right)\right]\cdots$$

• Then, 1st order approximation of AR can be implemented as follows.



Quantum Gates are independent of eigenvalues

Eigenvalue Inversion (EI)

- The 1st order AR demands inverted eigenvalues $|y\rangle_e = \left|\frac{1}{\lambda}\right|_e$.
- Therefore, there is a need to use general <u>quantum inversion</u>
 - -> A collection of quantum arithmetic subroutines. (Very Difficult)

Simplification:

- Use the simplest quantum inversion -
 - -> Powers of two inversion.

Idea:

 Inverting an eigenvalue that is a power of two is the same as swapping the positions of the binary sequence about the middle.

For Example,

$$4 \to \frac{1}{4}$$
, $100 \to 0.01$ $8 \to \frac{1}{8}$, $1000 \to 0.001$

Eigenvalue Register of
$$t$$
 qubits $\begin{vmatrix} \lambda_t \rangle_e \\ |\lambda_{t-1} \rangle_e \end{vmatrix}$ $\begin{vmatrix} \lambda_t \rangle_e \\ |\lambda_2 \rangle_e \end{vmatrix}$ $\begin{vmatrix} \lambda_1 \rangle \end{vmatrix}$ $\begin{vmatrix} \lambda_1 \rangle_e \end{vmatrix}$

4D Case – Worked Example

Obtain $|x\rangle = \frac{1}{d}M^{-1}|b\rangle$ approximately,

Given that

- $|b\rangle = |+\rangle \otimes |+\rangle$
- $M = \Gamma \Lambda \Gamma^{\dagger}$, where $\Gamma = (U_2 \otimes U_2)$, $\Lambda = diag(\lambda_i)$
- Eigenvalues $\lambda_j = 2^{j+1}$, j = 0,1,2,3.

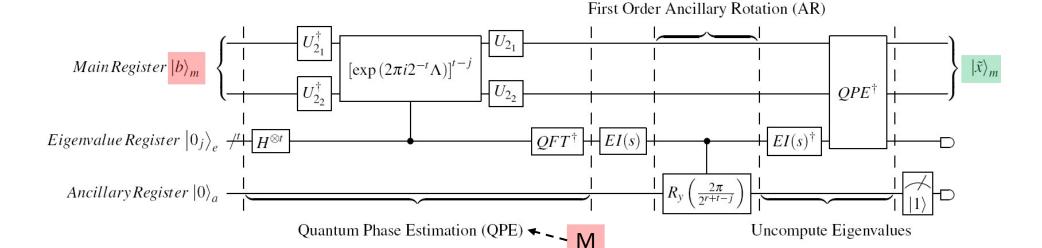
• Let
$$C = \frac{\pi}{2^r}$$

3.
$$|\widetilde{x}\rangle = \sqrt{\frac{1}{\sum_{j=1}^{4} \sin^{2}\left(\frac{\pi}{2^{r+j}}\right)}} \begin{pmatrix} -\frac{i}{2} \sin\left(\frac{\pi}{2^{r+4}}\right) + \frac{1}{2} \sin\left(\frac{\pi}{2^{r+3}}\right) + \frac{1}{2} \sin\left(\frac{\pi}{2^{r+2}}\right) + \frac{i}{2} \sin\left(\frac{\pi}{2^{r+1}}\right) \\ \frac{1}{2} \sin\left(\frac{\pi}{2^{r+4}}\right) - \frac{i}{2} \sin\left(\frac{\pi}{2^{r+3}}\right) + \frac{i}{2} \sin\left(\frac{\pi}{2^{r+2}}\right) + \frac{1}{2} \sin\left(\frac{\pi}{2^{r+1}}\right) \\ \frac{1}{2} \sin\left(\frac{\pi}{2^{r+4}}\right) + \frac{i}{2} \sin\left(\frac{\pi}{2^{r+3}}\right) - \frac{i}{2} \sin\left(\frac{\pi}{2^{r+2}}\right) + \frac{1}{2} \sin\left(\frac{\pi}{2^{r+1}}\right) \\ \frac{i}{2} \sin\left(\frac{\pi}{2^{r+4}}\right) + \frac{1}{2} \sin\left(\frac{\pi}{2^{r+3}}\right) + \frac{1}{2} \sin\left(\frac{\pi}{2^{r+2}}\right) - \frac{i}{2} \sin\left(\frac{\pi}{2^{r+1}}\right) \end{pmatrix}$$

$$M = \frac{1}{2} \begin{pmatrix} 15 & -5i & -9i & -3 \\ 5i & 15 & 3 & -9i \\ 9i & 3 & 15 & -5i \\ -3 & 9i & 5i & 15 \end{pmatrix}$$

$$U_2 = \begin{pmatrix} \frac{1}{\sqrt{2}} & \frac{i}{\sqrt{2}} \\ \frac{i}{\sqrt{2}} & \frac{1}{\sqrt{2}} \end{pmatrix}$$

$$|x\rangle = \frac{1}{2\sqrt{85}} \begin{pmatrix} 6+7i\\ 9+2i\\ 9-2i\\ 6-7i \end{pmatrix}$$



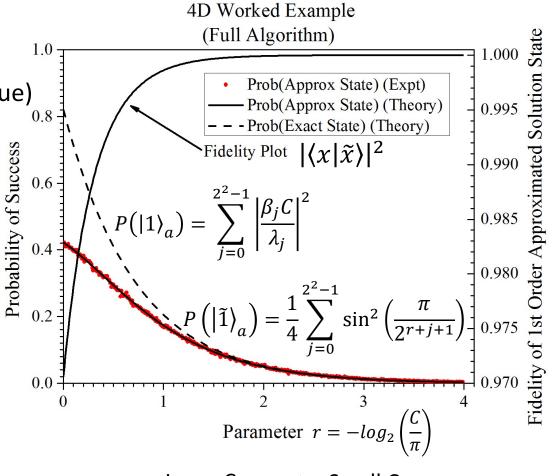
4D Case – Worked Example Results

IBM Qiskit QASM Simulator

Qubits: 9 (10000 runs/ r value)

Quantum Operations: 812

- Average Runtime : $\sim 10^0 s$
- Classical Runtime (Numpy Solver) : $\sim 10^{-5}s$
- Experimental value error :
 ± 9% (w.r.t theoretical value)



Large C → Small C

Analysis

Advantages:

- No need to know of Eigenvalues and eigenstates
- It can be used to solve for any 2^nD complex M and $|b\rangle$.

Disadvantages:

- Implementing Eigenvalue Inversion (EI) subroutine is complicated and very difficult.
- Higher n^{th} order AR requires $O(t^n)$ gates.
- Predicting probability of success needs eigenvalues.

Caveat:

• Full QLS Algorithm can be implemented, if controlled- $exp(2\pi i 2^{-t}M)$ quantum gates are given.

Crucial Gate for QPE!!!

Future Research

Full QLS Applications:

- Quantum Data Fitting.
- Quantum State Tomography.

QLS Extensions:

- Replace any subroutines in the QLS with another one as long the input-output are the same.
 - E.g. QPE, AR & EI.
- Add other quantum subroutines to increase the probability of success or improve computation time.
 - E.g. Grover search algorithm

Thank You Questions & Answers

More FYP Contents:

- Background on Quantum Computation
- Literature Review on Quantum Linear Solvers (QLS)
- Detailed descriptions and explanations on Simplified and Full QLS
 - More Worked Examples: 2D, 4D and 2^n D
- In-depth Discussions and Analysis.

Below are my experiment codes (python jupyter notebook):

Simplified QLS Algorithm



https://gist.github.com/cheechonghian/ea7ed914a264444b540aa57c428df497

Full QLS Algorithm



https://gist.github.com/cheechonghian/fc1c1a321f1e36b216eb7615d7edba65

Anaconda Environment



https://anaconda.org/chee0122/myQuantumCircuitSimulator2/files

Eigenvalue Marking Modification

Idea by Stefanie Barz et al:

- Modified AR: Apply $C[R_y(\theta_j)]$ on $|0\rangle_a$, conditioned on a partial sequence of the binary qubits of $|\lambda_j\rangle_e$.
- Modified QPE: Truncate QPE (Eigenvalue Marking, EM), to generate partial $|\lambda_j\rangle_{\rho}$ (Eigenvalue marking qubits).

Eigenvalue Assumption:

• Modification is possible if the all eigenvalues λ_j have the same binary sequence X up to the last n bits, where the last n bits are index j in binary.

t: No. of bits needed to specify the largest eigenvalue.

Assume in binary
$$\lambda_j \xrightarrow{t-n} 0.Xj$$
 bits

Also, assume knowing the eigenvalues.

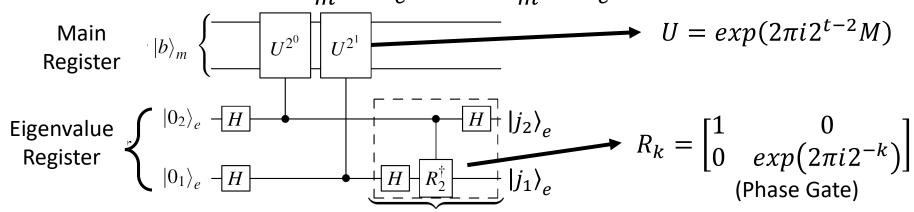
4D Case – Truncated QPE (EM)

Four eigenvalues are as follows,

$$\lambda_0 = 0.X00$$
, $\lambda_1 = 0.X01$ for some binary $\lambda_2 = 0.X10$, $\lambda_3 = 0.X11$ sequence X .

Apply truncated QPE (EM),

$$\begin{aligned} |m_{0}\rangle_{m}|00\rangle_{e} &\rightarrow |m_{0}\rangle_{m}|00\rangle_{e} \\ |m_{1}\rangle_{m}|00\rangle_{e} &\rightarrow |m_{1}\rangle_{m}|01\rangle_{e} \\ |m_{2}\rangle_{m}|00\rangle_{e} &\rightarrow |m_{2}\rangle_{m}|10\rangle_{e} \\ |m_{3}\rangle_{m}|00\rangle_{e} &\rightarrow |m_{3}\rangle_{m}|11\rangle_{e} \end{aligned}$$



 QFT^{\dagger} : Inverse Quantum Fourier Transform

4D Case – AR

Apply AR:

Eigenvalue Register
$$\begin{cases} |k_2\rangle_e \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_2\rangle_e \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_2\rangle_e \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_2\rangle_e \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_2\rangle_e \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_2\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_2\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_2\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_2\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_2\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_2\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_2\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_2\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_2\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_2\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_2\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_2\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_2\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_2\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_2\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_2\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_2\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_2\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_2\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_2\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_2\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_2\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_1\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_1\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_1\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_1\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_1\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_1\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_1\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_1\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_1\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_1\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_1\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_1\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_1\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_1\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_1\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_1\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_1\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_1\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_1\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_1\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_1\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_1\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_1\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_1\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_1\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_1\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_1\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_1\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_1\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} = - \begin{cases} |k_1\rangle_e - AR \\ |k_1\rangle_e - AR \end{cases} =$$

Conditioned Qubits

$$|00\rangle_{e} |01\rangle_{e} |10\rangle_{e} |11\rangle_{e}$$

$$-R_{y}(\theta_{0}) - R_{y}(\theta_{1}) - R_{y}(\theta_{2}) - R_{y}(\theta_{3}) - R_{y}(\theta_{3})$$

Calculate angles
$$\theta_j = 2 \sin^{-1} \left(\frac{c}{\lambda_j} \right)$$

- After that, apply inverse EM to un-compute eigenvalue marking qubits.
- Finally, measure and select $|1\rangle_a$ quantum state to obtain

$$|x\rangle_m = \frac{1}{d} \sum_{j=0}^{2^2 - 1} \frac{\beta_j}{\lambda_j} |m_j\rangle$$
 with probability $\sum_{j=0}^{2^n - 1} \left| \frac{\beta_j C}{\lambda_j} \right|^2$.

4D Case – Worked Example

$$M = \frac{1}{16} \begin{pmatrix} 12 & 0 & -2 & 0 \\ 0 & 10 & 10 & -2 \\ -2 & 0 & 12 & 0 \\ 0 & -2 & 0 & 10 \end{pmatrix}$$

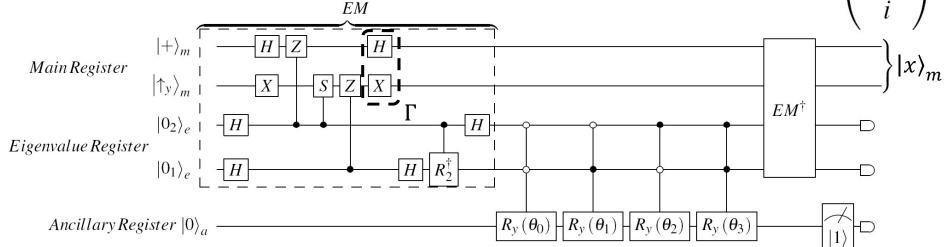
Let use a simplified case,

•
$$M = \Gamma \Lambda \Gamma^{\dagger}$$
, where $\Gamma = (H \otimes X)$.

$$\theta_j = 2\sin^{-1}\left(\frac{1}{2^r \cdot 0.1j}\right)$$

- Eigenvalues $\lambda_j=0.1j$, where j=00,01,10,11,
- Input $|b\rangle = |+\rangle \otimes |\uparrow_y\rangle$.
- $C = \frac{1}{2^r}$ (Re-parameterized)

and
$$X = 1$$
.
$$|x\rangle_m = \sqrt{\frac{25}{82}} \begin{pmatrix} \frac{8}{10} \\ i \\ \frac{8}{10} \\ i \end{pmatrix}$$



4D Case – Worked Example Results

IBM Qiskit QASM Simulator (10000 runs/ r value) 4D Worked Example (Simplified Algorithm)

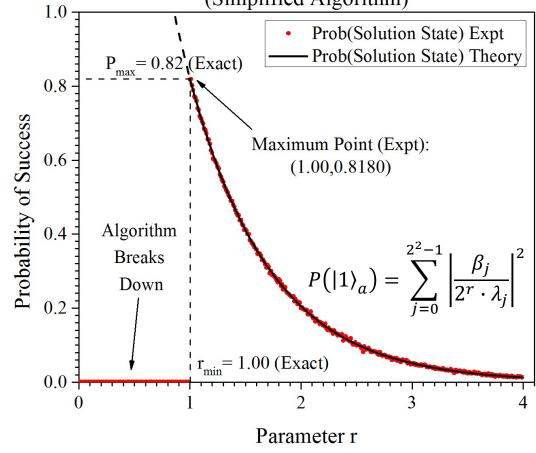
Qubits: 6

Quantum Operations: 51

- Average Runtime : $\sim 10^{-3} s$
- Classical Runtime (Numpy Solver) : $\sim 10^{-5}s$
- Experimental value error :
 ± 6% (w.r.t theoretical value)

Theoretical min parameter r:

$$r_{min} = -\log_2(C_{max})$$
$$= -\log_2(min(\lambda_j))$$



Simplified HHL QLS Algorithm Analysis

Advantages:

- Possible to maximize the probability by finding the minimum parameter r.
- Any parameter r below the minimum is unattainable as the corresponding quantum gate does not exist.

Disadvantages:

- Needs to know all eigenvalues so that angles can be calculated for AR.
- Physical examples of eigenvalues 0.Xj are unknown, applications are limited.