

Chapter 1: The Physics of Waves

Yijie Chen

Contents

1	One-dim wave equation	2
1.1	Example	2
2	General solutions to the one-dim wave equation	3
2.1	Example	4
3	Harmonic traveling waves	4
3.1	Example	6
4	The three-dimensional wave equation	7
4.1	Three-Dimensional Plane Waves	7
4.2	Spherical Waves	8
A	Appendix: Complex Numbers and the Complex Representation	10

1 One-dim wave equation

All classical mechanical waves can be described by the same equation, the *differential wave equation*:

$$\frac{\partial^2 \Psi(x, t)}{\partial x^2} = \frac{1}{v^2} \frac{\partial^2 \Psi(x, t)}{\partial t^2} \quad (1.1)$$

1.1 Example

Show that

$$\Psi(x, t) = (x + vt)^2$$

is a solution to the differential wave equation. Assume that v is a constant.

Solution

We must show that $\Psi(x, t)$ solves Equation (1.1):

$$\frac{\partial^2 \Psi(x, t)}{\partial x^2} = \frac{1}{v^2} \frac{\partial^2 \Psi(x, t)}{\partial t^2}$$

To do this, we must take the indicated derivatives. Begin with the partial derivative with respect to x :

$$\frac{\partial \Psi(x, t)}{\partial x} = \frac{\partial}{\partial x} (x + vt)^2 = 2(x + vt) \frac{\partial}{\partial x} (x + vt) = 2(x + vt)(1) = 2(x + vt)$$

Now take the second partial derivative with respect to x :

$$\frac{\partial^2 \Psi}{\partial x^2} = \frac{\partial}{\partial x} \left[\frac{\partial \Psi}{\partial x} \right] = \frac{\partial}{\partial x} [2(x + vt)] = 2$$

Next, take the partial derivatives of $\Psi(x, t)$ with respect to t :

$$\frac{\partial \Psi}{\partial t} = 2(x + vt) \frac{\partial}{\partial t} (x + vt) = 2(x + vt)(v) = 2v(x + vt)$$

$$\frac{\partial^2 \Psi}{\partial t^2} = \frac{\partial}{\partial t} [2v(x + vt)] = \frac{\partial}{\partial t} [2vx + 2v^2t] = 2v^2$$

Substitute these results into the differential wave equation.

$$\frac{\partial^2 \Psi(x, t)}{\partial x^2} = 2$$

$$\frac{1}{v^2} \frac{\partial^2 \Psi(x, t)}{\partial t^2} = \frac{1}{v^2} (2v^2) = 2$$

Thus

$$\frac{\partial^2 \Psi(x, t)}{\partial x^2} = \frac{1}{v^2} \frac{\partial^2 \Psi(x, t)}{\partial t^2}$$

$\Psi(x, t)$ solves the differential wave equation, and thus we have shown that it is a traveling wave.

2 General solutions to the one-dim wave equation

We will show below that the most general solutions to Equation(1.1) may be expressed as follows:

$$\Psi(x, t) = f(x - vt) \quad (2.1)$$

$$\Psi(x, t) = g(x + vt) \quad (2.2)$$

where the function f and g represent any function that has finite second derivatives, and where the parameters x, v and t all occur explicitly within the function as $x - vt$ or $x + vt$.

As an example, consider the function

$$\Psi(x, t) = \frac{A}{1 + (x - vt)^2} \quad (2.3)$$

where A is a constant. Equation(2.3) represents a peaked function whose maximum is located at points given by $x = vt$. A plot of this wave function at two different times is shown in Figure 1.

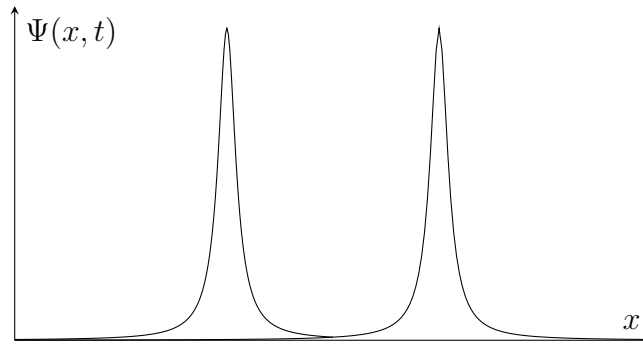


Figure 1: The traveling pulse of Equation (2.3), shown at two different times $t_1 < t_2$.

Velocity v is determined, x depends on time t . While $vt_1 = x$ or $vt_2 = x$, $\Psi(x, t) = A$, x is a coordinates of wave occurs pulse during traveling.

While v has a plus sign, such as $+v$, we have $x - (+v)t = x - vt$ in Equation(2.1), wave *traveling forward* and v has a minus sign, such as $-v$, $x - (-v)t = x + vt$ in Equation(2.2), wave *traveling backward*.

2.1 Example

Show that

$$\Psi(x, t) = Ae^{-(a^2x^2 + b^2t^2 + 2abxt)}$$

is a traveling wave, and find the wave speed and direction of propagation. Assume that A, a and b are all constants, and that a and b have units that make the quantity in the exponential function unitless.

Solution

$$\begin{aligned}\Psi(x, t) &= Ae^{-(a^2x^2 + b^2t^2 + 2abxt)} \\ &= Ae^{-(ax + bt)^2}\end{aligned}$$

While $ax + bt = 0$, thus $x = -\frac{bt}{a}$ and $v = -\frac{b}{a}$, this wave traveling backward.

3 Harmonic traveling waves

According to the results of the previous section, any function described by Equation(2.1) or (2.2) represents a traveling wave. In particular, *harmonic function* (i.e., sines and cosines) with the appropriate arguments solve the differential wave equation. Thus, the following function represents a traveling wave:

$$\Psi(x, t) = A \sin \frac{2\pi}{\lambda} (x \mp vt) \quad (3.1)$$

The $\Psi(x, t)$ given in Equation(3.1) is periodic in both space and time coordinates. The term λ represents the *spatial period*:

$$\Psi(x + \lambda, t) = A \sin \frac{2\pi}{\lambda} (x + \lambda \mp vt) = A \sin \left[\frac{2\pi}{\lambda} (x \mp vt) + 2\pi \right] = \Psi(x, t)$$

Let T represent the *temporal period*; the time required for one cycle. Since Equation (3.1) is periodic in T , we have

$$\Psi(x, t + T) = A \sin \frac{2\pi}{\lambda} [x \mp v(t + T)] = A \sin \left[\frac{2\pi}{\lambda} (x \mp vt) \mp \frac{2\pi}{\lambda} vT \right]$$

T represents the temporal period provided that

$$\frac{vT}{\lambda} = 1 \quad (3.2)$$

or

$$v = \frac{\lambda}{T} \quad (3.3)$$

Thus, a periodic classical wave travels one wavelength λ in one temporal period T . It is customary to define the *wave frequency* as

$$f = \frac{1}{T} \quad (3.4)$$

In terms of frequency, (3.3) becomes

$$f\lambda = v \quad (3.5)$$

A plot of $\Psi(x, t)$ vs. x is shown in Figure 2a. Figure 2b shows a plot of $\Psi(x, t)$ vs. t . It is customary to define the *propagation constant* as follows:

$$k = \frac{2\pi}{\lambda} \quad (3.6)$$

This quantity is also sometimes referred to as the *wave number*. Since k converts meters to radians, the units are rad/m . We may rewrite Equation (3.1) as

$$\Psi(x, t) = A \sin k(x \mp vt)$$

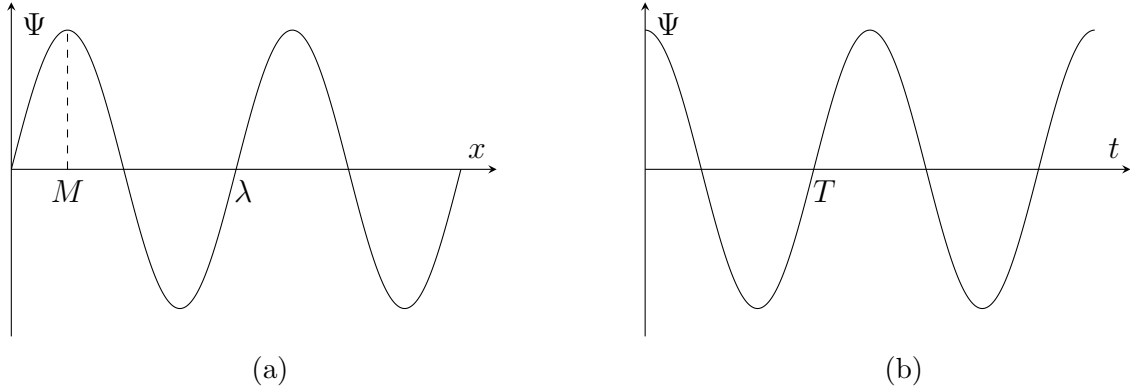


Figure 2: Plots of harmonic wavefunction. (a) A plot of $\Psi(x, t)$ vs. position x . (b) A plot of $\Psi(x, t)$ vs. time using data recorded by a single measuring device located at M in Figure (a)

Similarly, we can define the *angular frequency*

$$\omega = kv = \frac{2\pi}{\lambda}v = 2\pi f \quad (3.7)$$

Since ω converts time to radians, the units are *rad/s*. Note that according to the last result.

$$\frac{\omega}{k} = v \quad (3.8)$$

In terms of k and ω , Equation (3.1) becomes

$$\Psi(x, t) = A \sin(kx \mp \omega t)$$

Collectively, the terms $kx \mp \omega t$ are called the *phase* of the harmonic traveling wave. We may also include an explicit value of the phase when x and t are zero by specifying the *initial phase* ϕ :

$$\Psi(x, t) = A \sin(kx \mp \omega t + \phi) \quad (3.9)$$

3.1 Example

A harmonic traveling wave is given by $\Psi(z, t) = A \sin(50z + 3000t)$. Find the wave speed, frequency, angular frequency, and direction of propagation.

Solution

$$\Psi(z, t) = A \sin(50z + 3000t) = A \sin(kz + \omega t)$$

$$k = \frac{2\pi}{\lambda} = 50$$

$$\omega = kv = 2\pi f = 3000$$

wave speed:

$$v = \frac{3000}{k} = 60$$

frequency:

$$f = \frac{3000}{2\pi} = \frac{1500}{\pi}$$

angular frequency:

$$\omega = 3000$$

direction of propagation:backward

4 The three-dimensional wave equation

In Cartesian coordinates,the extension of Equation (1.1) to include three dimensions is made in the obvious way:

$$\frac{\partial^2 \Psi}{\partial x^2} + \frac{\partial^2 \Psi}{\partial y^2} + \frac{\partial^2 \Psi}{\partial z^2} = \frac{1}{v^2} \frac{\partial^2 \Psi}{\partial t^2} \quad (4.1)$$

where $\Psi = \Psi(x, y, z, t)$ is now understood to be a function of time and all three spatial coordinates.

Equation(4.1) can be written in a more general form using the Laplacian operator:

$$\nabla^2 \Psi = \frac{1}{v^2} \frac{\partial^2 \Psi}{\partial t^2} \quad (4.2)$$

where in Cartesian coordinates,the *Laplacian* is given by

$$\nabla^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \quad (4.3)$$

Equation (4.2) is the *coordinates independent* representation of the three-dimensional wave equation.

4.1 Three-Dimensional Plane Waves

Consider an acoustic wave traveling through air.In a three-dimensional *plane wave*,the properties of the medium are constant over any *plane* oriented normal to the direction of propagation. To describe such a wave,we define the propagation vector \vec{k} with

magnitude $\frac{2\pi}{\lambda}$ and direction given by the wave propagation. The corresponding plane wave is given by

$$\Psi(x, y, z, t) = A \sin(\vec{k} \cdot \vec{r} - \omega t + \varphi) \quad (4.4)$$

A plane has Cartesian symmetry, and in Cartesian coordinates, \vec{k} is given by

$$\vec{k} = k_x \hat{i} + k_y \hat{j} + k_z \hat{k} \quad (4.5)$$

The position vector in Cartesian coordinates is given by

$$\vec{r} = x \hat{i} + y \hat{j} + z \hat{k} \quad (4.6)$$

For example, in a plane harmonic sound wave traveling in the positive x -direction, the value of the air pressure is constant over any plane that is parallel to the y - z plane:

$$\Psi(x, y, z, t) = A \sin(kx - \omega t + \varphi) \quad (4.7)$$

In Figure it is seen that \vec{k} is normal to planes defined by $\vec{k} \cdot \vec{r} = \text{const}$. It is convenient to define *wavefronts* located at points where the phase of Equation (4.4) is equal to integer multiples of 2π . Thus, a three-dimensional plane wave may be visualized as a train of wavefronts separated by one wavelength λ and moving with the wave speed v .

In the complex representation, harmonic plane waves are given by

$$\Psi(x, y, z, t) = A e^{i(\vec{k} \cdot \vec{r} - \omega t + \varphi)} \quad (4.8)$$

4.2 Spherical Waves

Spherical waves are waves with wavefronts that are spherical in shape. Examples include waves that emanate from an isotropic point source, or are converging to a point. An isotropic source emits waves symmetrically in all directions.

In spherical coordinates, spherical waves have no dependence on the angular coordinates θ and ϕ , so the derivatives with respect to these coordinates are zero, this means that the

$$\Psi(\vec{r}, t) = \Psi(r, t)$$

In other words, $\Psi(r, t)$ depends only on the spherical coordinates r and time t . Derivatives of such a function with respect to θ and ϕ given zero, leading to a simplified version of the Laplacian for spherical coordinates:

$$\nabla^2 = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) \quad (4.9)$$

Using this Laplacian, we obtain the *spherical symmetric differential wave equation*:

$$\frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial \Psi(r, t)}{\partial r} \right) = \frac{1}{v^2} \frac{\partial^2 \Psi(r, t)}{\partial t^2} \quad (4.10)$$

An important solution to this equation is the *harmonic spherical wave*

$$\Psi(r, t) = \frac{A}{r} \sin(kr \mp \omega t + \varphi) \quad (4.11)$$

If the minus sign is chosen, the wave emanates from an isotropic source located at $r = 0$. A plus sign represents a wave that is converging to the point $r = 0$. The amplitude $\frac{A}{r}$ of the wave is not constant, since it depends on the radial coordinate r . For a wave emanating from a point, the amplitude decreases as the wave travels. Waves converging to a point have amplitudes that increase as time increases.

In the complex representation, Equation (4.11) becomes

$$\Psi(r, t) = \frac{A}{r} e^{i(kr \mp \omega t + \varphi)} \quad (4.12)$$

A Appendix:Complex Numbers and the Complex Representation

Complex numbers can provide algebraic shortcuts that will prove very convenient as we continue our discussion of classical optics. In this section,we provide a quick overview of the properties of complex numbers,and a few of their algebraic features that we will find most useful.

Complex numbers include the concept of the *imaginary number* i :

$$i = \sqrt{-1} \quad (\text{A.1})$$

Clearly,the square root of a negative number has no counterpart within the set of all *real* numbers.A complex number z has both a real part and an imaginary part:

$$z = x + iy \quad (\text{A.2})$$

where x is the *real* part of z and y is the *imaginary* part of z .

In the *Cartesian representation*, a complex number is plotted with coordinates (x,y) ; thus the horizontal axis is called the *real axis*, and the vertical axis is called *imaginary axis*.

Many of the features of complex numbers that we will find most useful result from the *Euler relation*:

$$e^{i\theta} = \cos\theta + i\sin\theta \quad (\text{A.3})$$

We may use the Euler relation to express any complex number z in polar form.Let

$$x = r\cos\theta \quad (\text{A.4})$$

$$y = r\sin\theta \quad (\text{A.5})$$

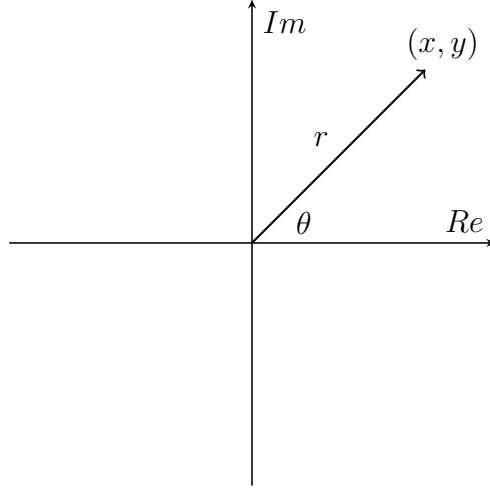


Figure 3: The complex plane

where

$$r = \sqrt{x^2 + y^2} \quad (\text{A.6})$$

$$\theta = \tan^{-1}\left(\frac{y}{x}\right) \quad (\text{A.7})$$

Using these relation, we may represent z as

$$z = x + iy = r\cos\theta + i(r\sin\theta) \quad (\text{A.8})$$

Thus, according to Equation (A.3),

$$z = re^{i\theta} \quad (\text{A.9})$$

In the polar representation, r is called the *magnitude* of z , and θ is called the *phase* of z .

The following examples are often useful:

$$e^{\pm i(2\pi)} = 1; e^{\pm i(\pi)} = -1; e^{i\frac{\pi}{2}} = i; e^{i\frac{3\pi}{2}} = -i \quad (\text{A.10})$$

The *complex conjugate* z^* of a complex number z is obtained by inverting the sign on each occurrence of i . For example, in the Cartesian representation where $z = x + iy$, z^* is given by

$$z^* = x - iy \quad (\text{A.11})$$

In the polar representation where $z = re^{i\theta}$, z^* is given by

$$z^* = re^{-i\theta} \quad (\text{A.12})$$

Let $Re[z]$ and $Im[z]$ denote the real and imaginary parts of z . In the Cartesian form, $Re[z]=x$ and $Im[z]=y$. In the polar form, $Re[z] = r\cos\theta$ and $Im[z] = r\sin\theta$. In either case,

$$Re[z] = \frac{z + z^*}{2} \quad (\text{A.13})$$

$$Im[z] = \frac{z - z^*}{2i} \quad (\text{A.14})$$

According to the Euler relation,

$$\cos\theta = \frac{e^{i\theta} + e^{-i\theta}}{2} \quad (\text{A.15})$$

$$\sin\theta = \frac{e^{i\theta} - e^{-i\theta}}{2i} \quad (\text{A.16})$$

Multiplication of a complex number by its complex conjugate gives the square of the magnitude. In the Cartesian form,

$$zz^* = (x + iy)(x - iy) = x^2 + y^2 \quad (\text{A.17})$$

and in the polar form,

$$zz^* = (re^{i\theta})(re^{-i\theta}) = r^2 \quad (\text{A.18})$$

which, according to Equation (A.6), agrees with the previous result. We will also refer to the magnitude of z as $|z|$:

$$zz^* = |z|^2 \quad (\text{A.19})$$