

Architectures and Approaches in Scaling the Trapped-Ion Quantum Computer

Chen Huang
Department of Physics, Imperial College London

Contents

1 Overview

- Why compute with qubits?
- How to build a Quantum Computer?

2 Quantum Computing with Trapped Ions

- Ion Traps
- The Basic Hamiltonian
- Quantum Gates

3 Approaches to Scaling Trapped Ions

- Linear Arrays
- Two Dimensional Arrays
- Remote Entanglement

Why compute with qubits?

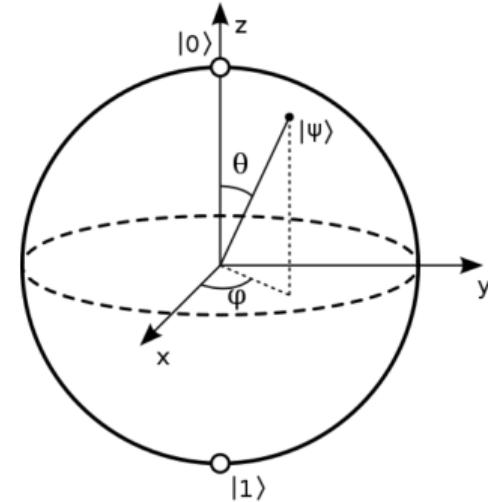
Quantum Rules

- ① Superposition $|\psi\rangle = \cos \frac{\theta}{2} |0\rangle + e^{i\varphi} \sin \frac{\theta}{2} |1\rangle$
- ② Coherence
- ③ Entanglement $|\psi\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$

Quantum parallel processing on 2^N inputs

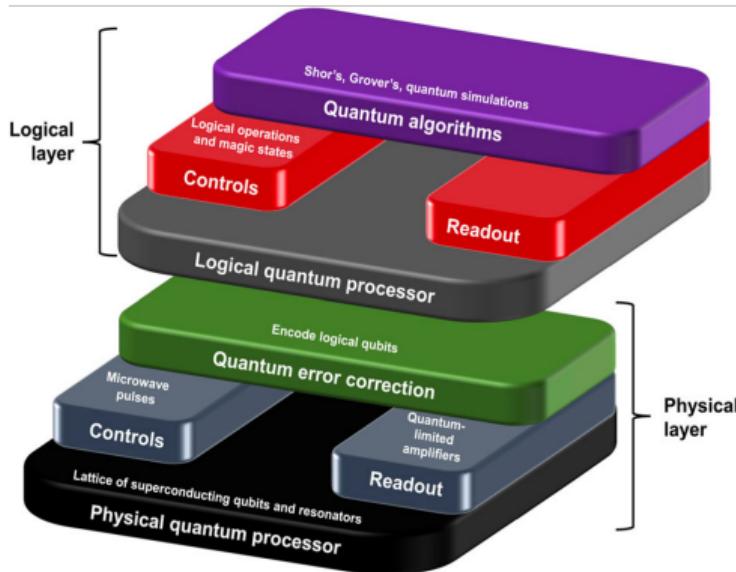
Example: $N = 3$

$$|\psi\rangle = a_0 |000\rangle + a_1 |001\rangle + a_2 |010\rangle + a_3 |011\rangle \\ + a_4 |100\rangle + a_5 |101\rangle + a_6 |110\rangle + a_7 |111\rangle$$



$N = 300$ qubits: more information than particles in the universe!

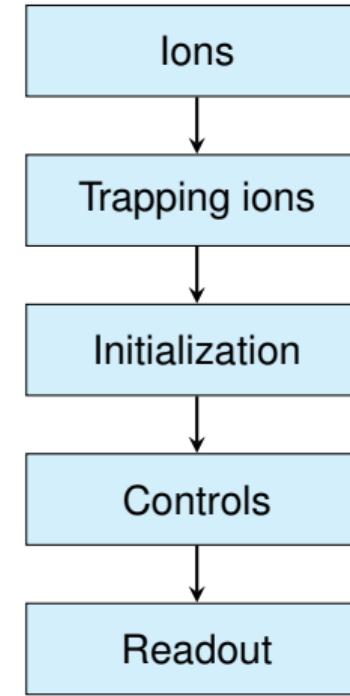
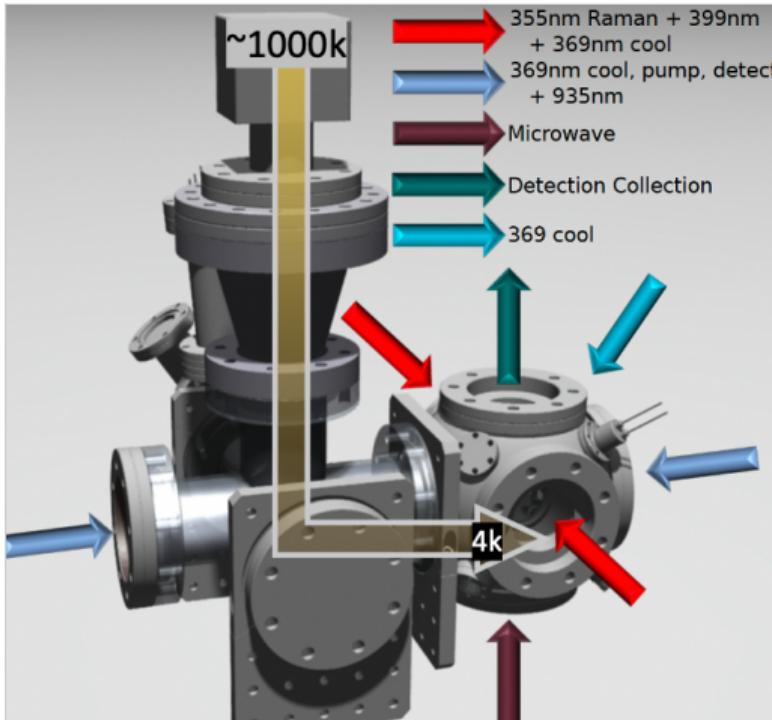
How to build a Quantum Computer?



DiVincenzo Criteria

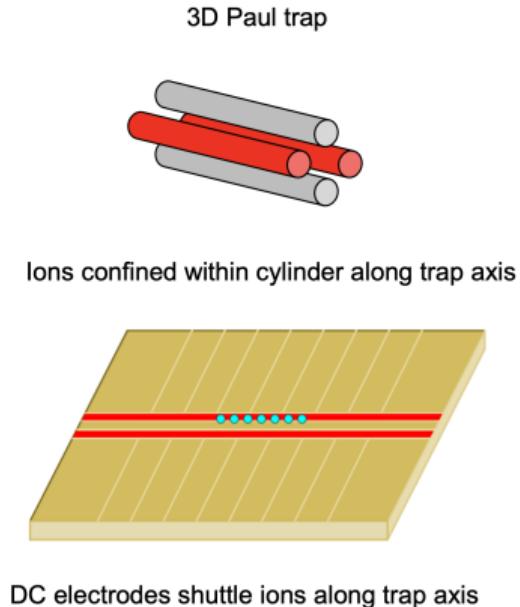
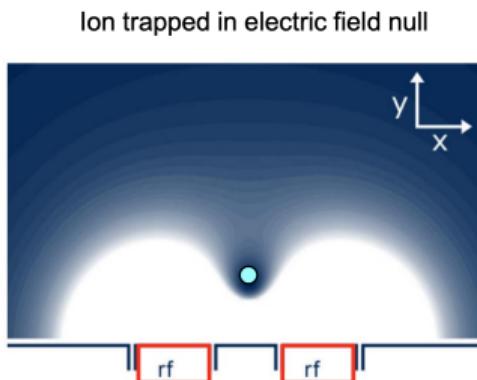
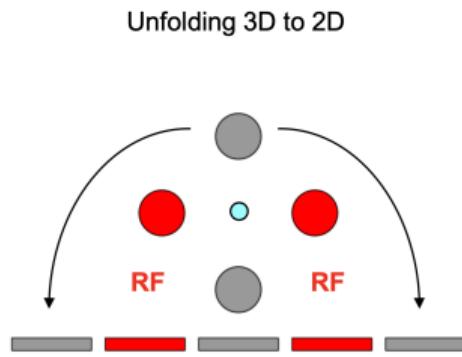
- ① A **scalable** physical system with well characterized qubits;
- ② The ability to initialize the qubits;
- ③ Long relevant coherence times;
- ④ A ‘universal’ set of quantum gates;
- ⑤ A qubit-specific measurement capability.

Trapped-Ion Quantum Computer Architecture



Ion Traps¹

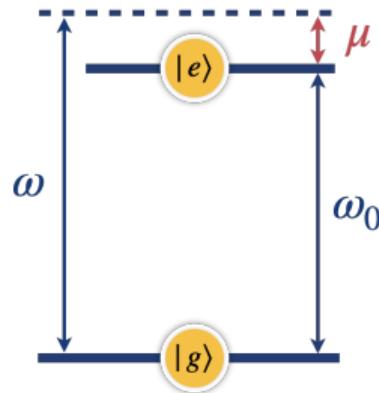
- **Penning trap:** static electric and magnetic field
- **Paul trap:** static and oscillating electric fields
 - Linear trap
 - Surface-electrode trap



¹Image source: MIT Lincoln Laboratory

The Basic Hamiltonian

The interaction between a two-level ion and a single-mode light.



$$\hat{H}_{int} = \frac{\hbar}{2}\Omega (|e\rangle\langle g| + |g\rangle\langle e|) [e^{i(kx - \omega t + \phi)} + e^{-i(kx - \omega t + \phi)}]$$

- ① Trans to Interaction picture
- ② $kx = kx_0(\hat{a}^\dagger + \hat{a}) = \eta(\hat{a}^\dagger + \hat{a})$
- ③ Rotational wave approximation (RWA)
- ④ Lamb-Dicke limit $\eta = k\sqrt{\hbar/(2m\omega_m)} \ll 1$

$$\hat{H} = \frac{\hbar}{2}\Omega\hat{\sigma}_+ [1 + i\eta (\hat{a}^\dagger e^{i\omega_m t} + \hat{a} e^{-i\omega_m t})] e^{i(\phi - \mu t)} + h.c.$$

Transitions

$$\begin{aligned}\hat{H} = & \frac{\hbar}{2}\Omega (\hat{\sigma}_+ e^{i\phi} + \hat{\sigma}_- e^{-i\phi}) & \Rightarrow \hat{H}_{car} \\ & + i\eta \frac{\hbar}{2} \Omega \left[\hat{\sigma}_+ \hat{a}^\dagger e^{i\phi} e^{-i(\mu - \omega_m)t} + \hat{\sigma}_- \hat{a} e^{-i\phi} e^{i(\mu - \omega_m)t} \right] & \Rightarrow \hat{H}_{bsb} \\ & + i\eta \frac{\hbar}{2} \Omega \left[\hat{\sigma}_+ \hat{a} e^{i\phi} e^{-i(\mu + \omega_m)t} - \hat{\sigma}_- \hat{a}^\dagger e^{-i\phi} e^{i(\mu + \omega_m)t} \right] & \Rightarrow \hat{H}_{rsb}\end{aligned}$$

- **Carrier transition** ($\mu = 0$)

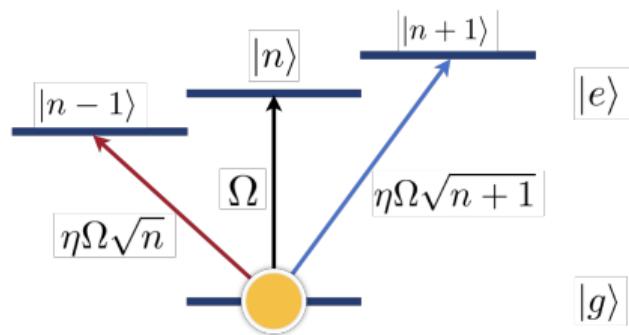
$$|g\rangle |n\rangle \leftrightarrow |e\rangle |n\rangle$$

- **Blue sideband transition** ($\mu = \omega_m + \delta$)

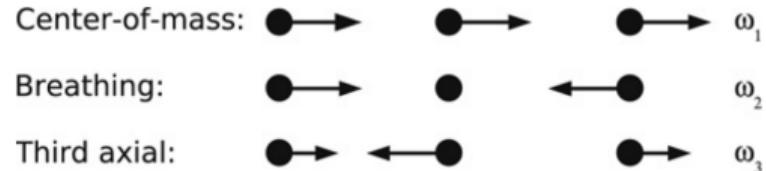
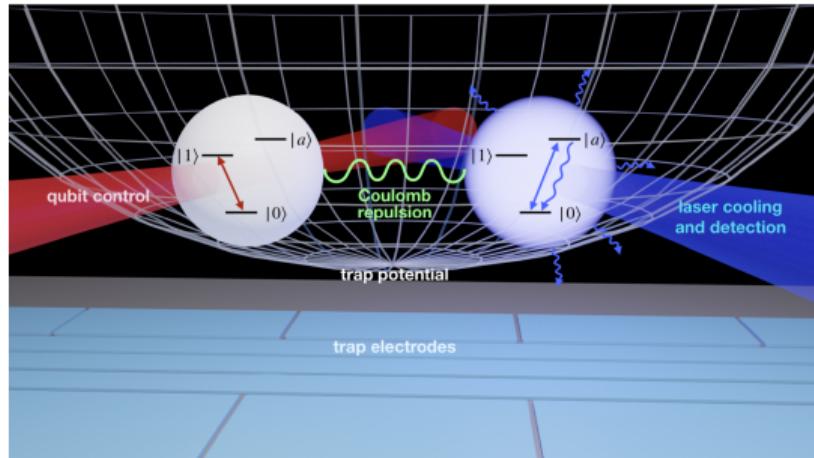
$$|g\rangle |n\rangle \leftrightarrow |e\rangle |n+1\rangle$$

- **Red sideband transition** ($\mu = -\omega_m - \delta$)

$$|g\rangle |n\rangle \leftrightarrow |e\rangle |n-1\rangle$$



More Ions



When N oscillators are coupled, N different canonical modes are obtained, which correspond to **phonons** of N **different frequencies**.

$$\hat{H} = \left[\frac{\hbar}{2} \sum_{j=1}^N \Omega_j \hat{\sigma}_+^j e^{i(\phi_j - \mu_j t)} \right] \left[\sum_{k=1}^N \left[1 + i\eta_{j,k} (\hat{a}_k^\dagger e^{i\omega_k t} + \hat{a}_k e^{-i\omega_k t}) \right] \right] + h.c.$$

internal state phonon state (ion vibration)

The Cirac-Zoller Gate²



- Based on **red** sideband transition $|g\rangle|1\rangle \leftrightarrow |e\rangle|0\rangle$
- Laser frequency
 $\omega = \omega_0 - \omega_m$

- Apply a π -pulse on ion A
- Apply a 2π -pulse on ion B
- Apply a π -pulse on ion A again

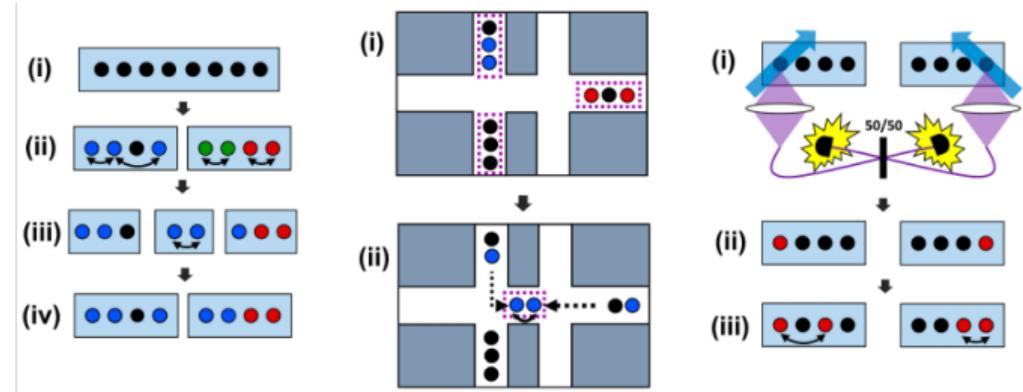
	(1)	(2)	(3)
$ gg0\rangle$	\rightarrow	$ gg0\rangle$	$ gg0\rangle$
$ ge0\rangle$	\rightarrow	$ ge0\rangle$	$ ge0\rangle$
$ eg0\rangle$	\rightarrow	$-i gg1\rangle$	$i gg1\rangle$
$ ee0\rangle$	\rightarrow	$-i ge1\rangle$	$-i ge1\rangle$
			$\rightarrow - ee0\rangle$

$$|\pm\rangle = (|g\rangle \pm |e\rangle)/\sqrt{2} \quad \Rightarrow \quad \begin{cases} |g\rangle |\pm\rangle \rightarrow |g\rangle |\pm\rangle \\ |e\rangle |\pm\rangle \rightarrow |e\rangle |\mp\rangle \end{cases} \quad \Rightarrow \quad \textbf{Controlled-Z Gate!}$$

²Cirac and Zoller 1995

Approaches to Scaling Trapped Ions³

- ① Linear arrays
- ② Two dimensional arrays
- ③ Remote entanglement



³Image source: Bruzewicz et al. 2019

Linear Arrays⁴

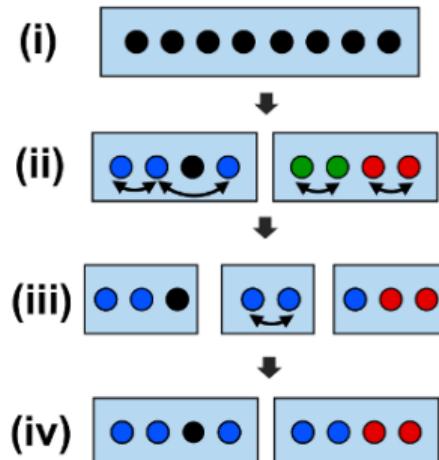
Problem: may prove more difficult as the chain size grows.

- ➊ weak coupling to the ion motion
- ➋ an increased susceptibility

Modularity!

Strategy:

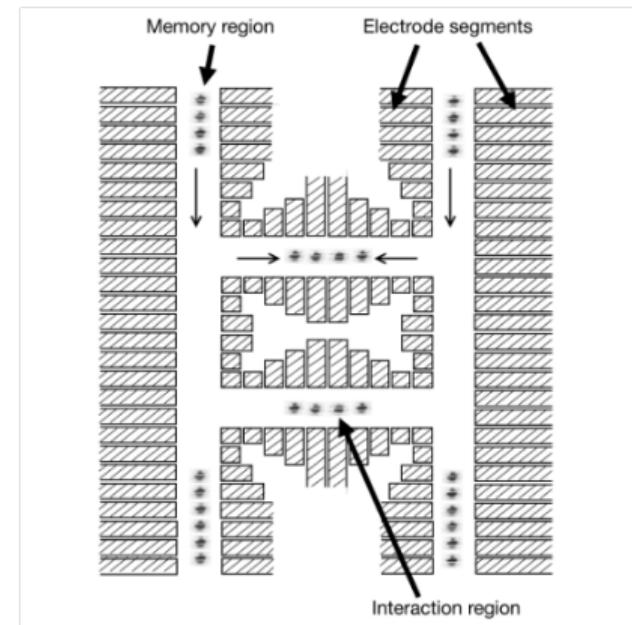
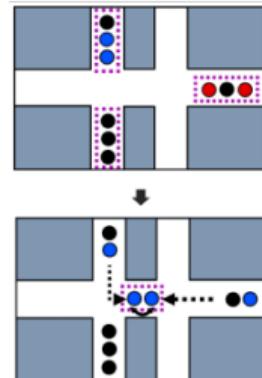
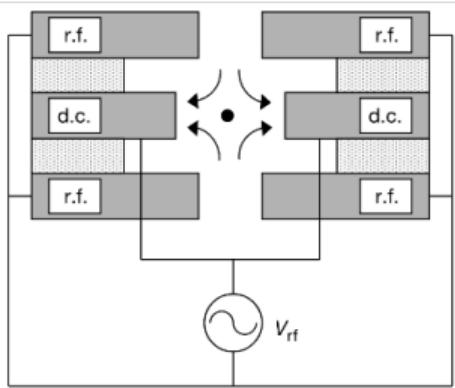
- ➊ Splitting the chain into smaller modules;
- ➋ Reconfiguring the modules allows for high-fidelity entangling gates between ions in different modules;
- ➌ Upon completion, the modules can revert to their original state.



⁴Bruzewicz et al. 2019

Two Dimensional Arrays - QCCD (Quantum Charge-Coupled Device)⁵

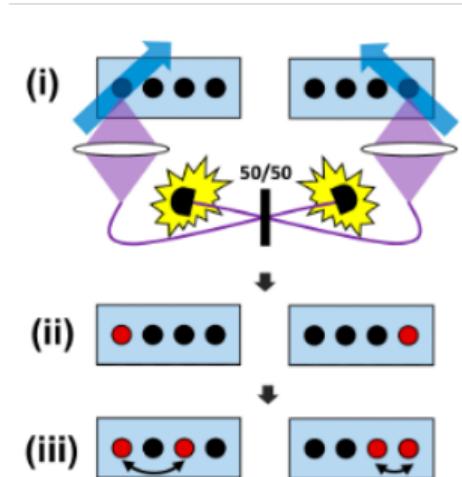
Ion transport is typically much faster than a two-ion gate or a reordering operation.



⁵Kielpinski, Monroe and Wineland 2002

Remote Entanglement - Photonic Interconnect

- ① Ions within each module initially exist in an **unentangled** state.
- ② A dedicated “communication ion” in each module is excited by a laser, resulting in the emission of photons **entangled** with the internal ion state.
- ③ There is a probability that these photons reach the 50/50 beam splitter and **interfere**.
- ④ Simultaneous detection of photons by single-photon detectors heralds **entanglement** between communication ions.



Summary

1 Overview

- Why compute with qubits?
- How to build a Quantum Computer?

2 Quantum Computing with Trapped Ions

- Ion Traps
- The Basic Hamiltonian
- Quantum Gates

3 Approaches to Scaling Trapped Ions

- Linear Arrays
- Two Dimensional Arrays
- Remote Entanglement

Thanks for Listening!