

Gathering more grid points near the boundary via grid transformation

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Abstract

To gather more grid points near the boundary of the 2D computational domain, grid transformation is applied. Here we only talk about stretching transformations with univariate functions, i.e. $y = y(\eta)$ and $z = z(\zeta)$. The final formula in this document is therefore not valid for general grid transformations $y = y(\eta, \zeta)$ and $z = z(\eta, \zeta)$.

To gather more grid points near the boundaries of a 2D rectangular physical domain, the following stretching transformation which maps the uniform computational grid system (η, ζ) to the non-uniform grid (y, z) is applied

$$y(\eta) = \frac{\tanh(a\eta)}{\tanh(a)}, \quad z(\zeta) = \frac{\tanh(a\zeta)}{\tanh(a)}, \quad (1)$$

where $-1 \leq \eta \leq 1$, $-1 \leq \zeta \leq 1$, and a is a parameter controlling the distribution of the grids.

When the univariate functions Eq. (1) are used, the derivatives with respect to (y, z) are given by [1, chapter 4]

$$\begin{bmatrix} \frac{\partial}{\partial y} \\ \frac{\partial}{\partial z} \end{bmatrix} = \frac{1}{|\mathbb{J}|} \begin{bmatrix} \frac{dz}{d\zeta} & 0 \\ 0 & \frac{dy}{d\eta} \end{bmatrix} \begin{bmatrix} \frac{\partial}{\partial \eta} \\ \frac{\partial}{\partial \zeta} \end{bmatrix} = \begin{bmatrix} \left(\frac{dy}{d\eta}\right)^{-1} \frac{\partial}{\partial \eta} \\ \left(\frac{dz}{d\zeta}\right)^{-1} \frac{\partial}{\partial \zeta} \end{bmatrix}, \quad (2a)$$

$$\begin{aligned} \frac{\partial^2}{\partial y^2} &= \frac{1}{|\mathbb{J}|^2} \left(\frac{dz}{d\zeta}\right)^2 \frac{\partial^2}{\partial \eta^2} - \frac{1}{|\mathbb{J}|^3} \left(\frac{dz}{d\zeta}\right)^3 \frac{\partial^2 y}{\partial \eta^2} \frac{\partial}{\partial \eta} \\ &= \left(\frac{dy}{d\eta}\right)^{-2} \frac{\partial^2}{\partial \eta^2} - \left(\frac{dy}{d\eta}\right)^{-3} \frac{\partial^2 y}{\partial \eta^2} \frac{\partial}{\partial \eta}, \end{aligned} \quad (2b)$$

$$\begin{aligned} \frac{\partial^2}{\partial z^2} &= \frac{1}{|\mathbb{J}|^2} \left(\frac{dy}{d\eta}\right)^2 \frac{\partial^2}{\partial \zeta^2} - \frac{1}{|\mathbb{J}|^3} \left(\frac{dy}{d\eta}\right)^3 \frac{\partial^2 z}{\partial \zeta^2} \frac{\partial}{\partial \zeta} \\ &= \left(\frac{dz}{d\zeta}\right)^{-2} \frac{\partial^2}{\partial \zeta^2} - \left(\frac{dz}{d\zeta}\right)^{-3} \frac{\partial^2 z}{\partial \zeta^2} \frac{\partial}{\partial \zeta}, \end{aligned} \quad (2c)$$

where

$$\mathbb{J} = \begin{bmatrix} \frac{dy}{d\eta} & 0 \\ 0 & \frac{dz}{d\zeta} \end{bmatrix} \quad (3)$$

is the Jacobian matrix, the determinant of which is

$$|\mathbb{J}| = \frac{dy}{d\eta} \frac{dz}{d\zeta}. \quad (4)$$

One should avoid $dy/d\eta = 0$ or $dz/d\zeta = 0$ (or $dy/d\eta = \infty$ or $dz/d\zeta = \infty$) otherwise the transformation given by Eq. (2) is singular.

The derivatives $dy/d\eta$, $d^2y/d\eta^2$, $dz/d\zeta$ and $d^2z/d\zeta^2$ are calculated by Eq. (1) as

$$\frac{dy}{d\eta} = \frac{a \operatorname{sech}^2(a\eta)}{\tanh(a)}, \quad \frac{d^2y}{d\eta^2} = -\frac{2a^2 \operatorname{sech}^2(a\eta) \tanh(a\eta)}{\tanh(a)}, \quad (5)$$

and

$$\frac{dz}{d\zeta} = \frac{a \operatorname{sech}^2(a\zeta)}{\tanh(a)}, \quad \frac{d^2z}{d\zeta^2} = -\frac{2a^2 \operatorname{sech}^2(a\zeta) \tanh(a\zeta)}{\tanh(a)}. \quad (6)$$

In Fig. 1, an example of the grid transformation via Eq. (2) is shown.

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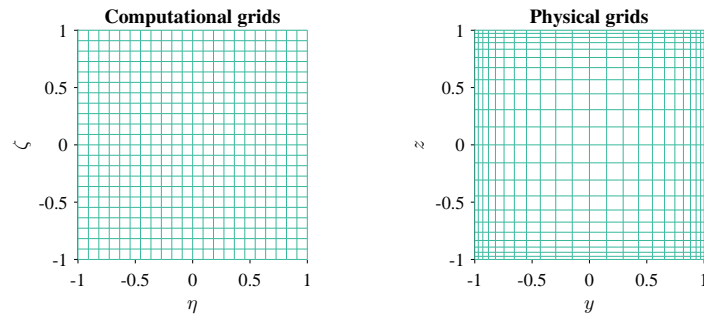


Figure 1. Grid transformation given by Eq. (2), with $a = 1.5$ and 1.6 for y and z respectively. See `testGridTransform.m` for the generation of the figure.

References

- [1] T.J. Chung. *Computational Fluid Dynamics*. Cambridge University Press, New York, USA, 2nd edition, 2010. ISBN 9780521769693.