二维可压缩磁流体模拟(用 MATLAB 语言复原工作)

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摘要

本工作是围绕谢华生博士的著作《计算等离子体物理导论》^[1]第 5 章 5.42 节撕裂模及磁重联的内容展开讨论的。此书是等离子体物理数值计算与模拟的入门教程,通过具体的算例来帮助初学者理解相关的物理概念和物理图像,既有新意又有实用性,对于初学入门者来说,是一本不可多得的书。书中的算例均提供了相关的代码,基本上是以 MATLAB 语言为主,有些算例并没有完全用 MATLAB 语言编写,例如撕裂模及磁重联的那一节^[2]的算例提供的是 Fortran 代码,对于多种代码语言不能一一精通的初学者来说这是一个挑战。此外,作者提供的参考著作(傅竹风,1995《空间等离子体数值模拟》)^[3]网上缺货,电子版在全国磁约束核聚变专业群也未能求得。为此,我按照书籍提供的线索及其提供的 Fortran 代码注释说明,用MATLAB 语言对撕裂模及磁重联这一节的内容进行复原工作。

1 二维可压缩磁流体方程

平板模型,所有的变量都在(x,z)平面,各变量 y 方向是均匀的 $\partial/\partial y=0$ 。

1.1 原始方程

原始方程没有霍尔项, η, γ, ν 为常数。

$$\frac{\partial \rho}{\partial t} = -\mathbf{u} \cdot \nabla \rho - \rho \nabla \cdot \mathbf{u},$$

$$\frac{\partial p}{\partial t} = -\mathbf{u} \cdot \nabla p - \gamma p \nabla \cdot \mathbf{u},$$

$$\rho \frac{\partial \mathbf{u}}{\partial t} = -\rho \mathbf{u} \cdot \nabla \mathbf{u} - \nabla p - \frac{1}{\mu_0} \nabla \frac{B^2}{2} + \frac{1}{\mu_0} \mathbf{B} \nabla \cdot \mathbf{B} + \nu \nabla^2 \mathbf{u},$$

$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\mathbf{u} \times \mathbf{B}) - \nabla \times \left[\eta \left(\nabla \times \mathbf{B} \right) \right].$$
(1)

其中,使用条件

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t},$$

$$\mathbf{J} = \nabla \times \mathbf{B},$$

$$\mathbf{E} = -\mathbf{u} \times \mathbf{B} + \eta \mathbf{J}.$$
(2)

考虑到 $\nabla \cdot \mathbf{B} = 0$,我们使用 $A_{\mathbf{x}}$ 来表示 \mathbf{B} ,例如

$$\mathbf{B} = \nabla \times \mathbf{A} = \nabla \times (0, A_{v}, 0) = (-\partial A_{v} / \partial z, 0, \partial A_{v} / \partial x)$$
(3)

把上述方程写成显示的二维方程组(已经归一化)

$$\frac{\partial \rho}{\partial t} = -u_x \frac{\partial \rho}{\partial x} - u_z \frac{\partial \rho}{\partial z} - \rho \left(\frac{\partial u_x}{\partial x} + \frac{\partial u_z}{\partial z} \right),$$

$$\frac{\partial p}{\partial t} = -u_x \frac{\partial p}{\partial x} - u_z \frac{\partial p}{\partial z} - \gamma p \left(\frac{\partial u_x}{\partial x} + \frac{\partial u_z}{\partial z} \right),$$

$$\frac{\partial u_x}{\partial t} = -u_x \frac{\partial u_x}{\partial x} - u_z \frac{\partial u_x}{\partial z} - \frac{1}{\rho} \frac{\partial}{\partial x} \left(p + \frac{B^2}{2} \right) + \frac{1}{\rho} \left(B_x \frac{\partial B_x}{\partial x} + B_z \frac{\partial B_x}{\partial z} \right) + \frac{1}{\rho} v_m \left(\frac{\partial^2 u_x}{\partial x^2} + \frac{\partial^2 u_x}{\partial z^2} \right),$$

$$\frac{\partial u_z}{\partial t} = -u_x \frac{\partial u_z}{\partial x} - u_z \frac{\partial u_z}{\partial z} - \frac{1}{\rho} \frac{\partial}{\partial z} \left(p + \frac{B^2}{2} \right) + \frac{1}{\rho} \left(B_x \frac{\partial B_z}{\partial x} + B_z \frac{\partial B_z}{\partial z} \right) + \frac{1}{\rho} v_m \left(\frac{\partial^2 u_z}{\partial x^2} + \frac{\partial^2 u_z}{\partial z^2} \right),$$

$$\frac{\partial A_y}{\partial t} = -u_x \frac{\partial A_y}{\partial x} - u_z \frac{\partial A_y}{\partial z} + \eta_m \left(\frac{\partial^2 A_y}{\partial x^2} + \frac{\partial^2 A_y}{\partial z^2} \right).$$
(4)

1.2 归一化

对于方程组(4), 归一化的密度, 压强, 长度, 速度和时间的归一化分别是 ho_0, p_0, L_0, B_0 ,

$$u_A = B_0 / \sqrt{\mu_0 \rho_0} , \quad \tau_A = L_0 / u_A \circ$$

参数估计

$$V_m \to \frac{V}{u_A L_0 \rho_0}, \eta_m \to \frac{\eta}{u_A L_0}, \tag{5}$$

磁 Lundquist 数 $S=1/\eta_m$ 。

1.3 离散化

对于方程组(4)右手边的空间微分,采用中心差分格式如下

$$\frac{\partial f}{\partial x} \to \frac{f_{i+1} - f_{i-1}}{2\Delta x},$$

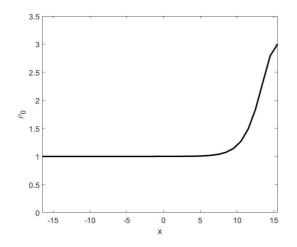
$$\frac{\partial f}{\partial z} \to \frac{f_{j+1} - f_{j-1}}{2\Delta z},$$

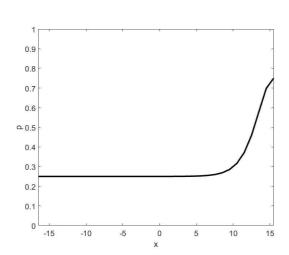
$$\frac{\partial^2 f}{\partial x^2} \to \frac{f_{i+1} + f_{i-1} - 2f_i}{\Delta x^2},$$

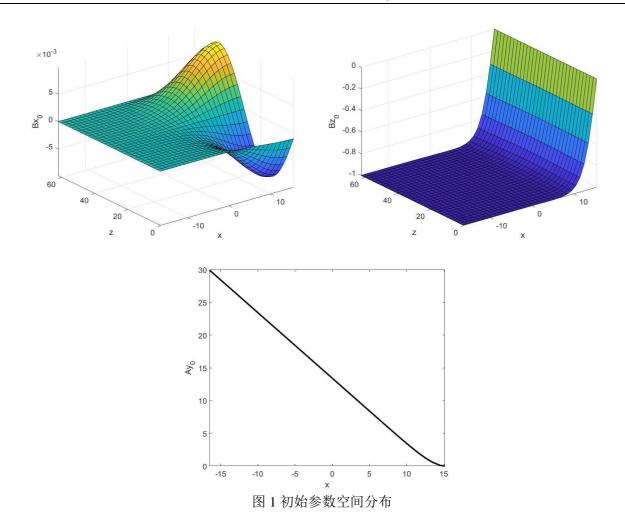
$$\frac{\partial^2 f}{\partial z^2} \to \frac{f_{j+1} + f_{j-1} - 2f_j}{\Delta z^2}.$$
(6)

对于方程组的左边的时间偏微分,采用的是四阶龙格库塔法格式求解。

1.4 初始条件







1.5 边界条件

在z方向上采用的是周期性边界条件,在x方向上用的是混合边界条件。在x轴的下边界采用条件为

$$u_x(x_d,:) = 0,$$

 $A_y(x_d,:) = 0,$
(7)

而 ρ , p, u_z 在 x 轴的下边界采用的是自由边界条件。在 x 轴的上边界采用条件为

$$u_{x}(x_{u},:) = 0,$$

$$\frac{\partial u_{z}}{\partial x} = 0,$$

$$\frac{\partial A_{y}}{\partial x} = 0.$$
(8)

2 MATLAB 求解代码结构

变量 U(ni,nj,5)存放了 5 个二维变量,分别代表 ρ , p, u_x , u_z , A_y 。本工作是基于 MATLAB 语言重新对著 作^[2, 4]提供的 Fortran 代码进行改写,结合 MATLAB 的矩阵运算优势,把原来单点循环赋值运算,全部改写 成矩阵赋值运算,使代码得到了很大程度上的简化。为擅于 MATLAB 语言而不会用 Fortran 的初学者提供 参考。另外,在 Fortran 代码的子程序 "Subroutine right(xo,xi)"中,求解 $\frac{\partial u_x}{\partial t}$ 的右边 pt 项本来是对 x 进行 求微分,但实际上输入的字母有误,具体是子程序中 "-rdz*(pt(i+1,j)-pt(i-1,j)))"正确应该是乘以 rdx。由于rdx 和 rdz 是个比较相接近的常数,没有对整体的趋势产生显著影响。

3 复原模拟结果

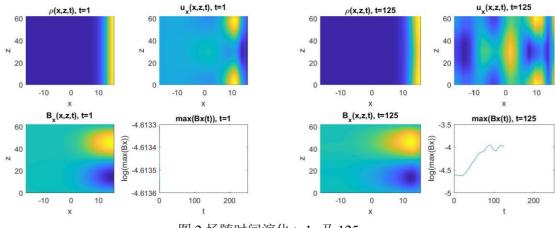


图 2 场随时间演化 t=1s 及 125s。

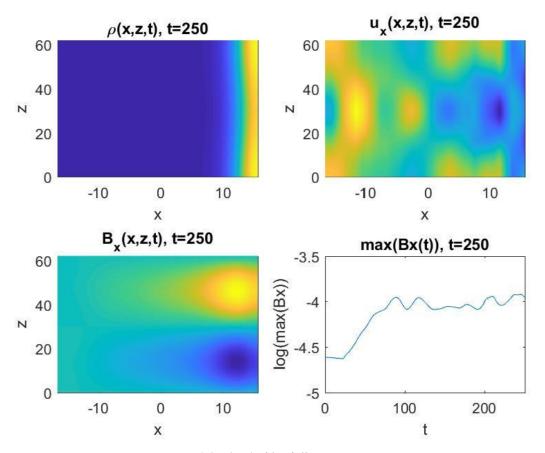


图 3 场随时间演化 t=250s.

图 2-3 为一组典型的非线性模拟结果,其中我们可以看到初始扰动调整后,开始进入一段指数增长的线性阶段,然后达到非线性饱和,饱和后基本稳态演化。

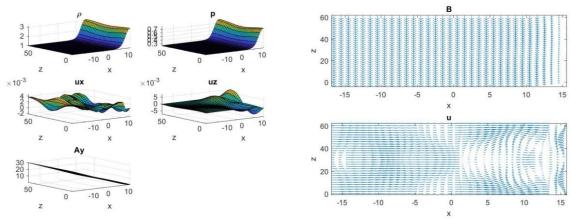


图 4 场随时间演化 t=250s.

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附录: MATLAB 程序

MHD2D main

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```
close; clear; clc;
global dx dz gamma nu eta dt Bx Bz
nx=33;
nz=32:
nt=5000;
dt=0.05; % 要满足 dt < min(dx, dz)/[sqrt(1.0+0.5*gamma*beta)*va]
dx=1.0; dz=2.0;
x1=dx*nx/2; z1=dz*(nz-1);
xx=-x1:dx:x1-dx;
zz=0:dz:z1:
[x,z]=meshgrid(xx,zz);
X=X'; Z=Z';
gamma=1.66667;
eta=0.01:
nu=0.05;
beta=0.5:
                                             % beta in x=Lx
rho0=1.0;
                       % rho0 -- mass density rho0 in x=Lx
B0=1.0;
                                                 % b0 -- B0
b1=3.0;
                            % bl -- width of current sheet
va=sqrt(B0*B0/rho0);% Alfven velocity
p0=0.5*beta*B0*B0; % p0 -- pressure p0 in x=Lx
t0=0.5*beta*va*va; % t0 -- temperature T0 in x=Lx
cfl=1.5; % Courant - Friedrichs - Lewy condition parameter
U=zeros(nx, nz, 5);
  %%初始场分布
s=((1:nx)-nx)*dx/b1;
b=B0*tanh(s);
p=p0+0.5*(B0.^2-b.^2);
rho=p/t0;
Ay=B0.*b1.*log(cosh(s));
U(:,:,1) = repmat(rho',1,nz);
U(:,:,2) = repmat(p',1,nz);
U(:,:,5) = repmat(Ay', 1, nz);
  %%初始场扰动
nw=1;% nw -- number of initial perturbation
abd=10; % abd -- perturbation width in x
am=[0.01, zeros(1,9)];% am -- prtb amplitude
s=((2:nx)-nx)*dx/abd;
for m=1:nw
    sii = \exp(-s.*s).*am(m)*B0;
    sij=sin((2*m*((1:nz)-0.5*nz)/(nz-1)+0.5)*pi)*dz*(nz-1)/(2.0*m*pi);
end
U(:,:,5)=U(:,:,5)+[zeros(1,nz);sii'*sij];
rho0=U(:,:,1); p0=U(:,:,2); ux0=U(:,:,3); uz0=U(:,:,4); Ay0=U(:,:,5); [Bx0,Bz0]=calcBxz(t,U);
for it=1:nt
    t=t+dt;
     [Bx, Bz] = calcBxz(t, U);
    Bxm(it) = max(max(Bx));
    U=RK4(t, U);
     if \mod(it, 20) == 0
         figure (2)
         subplot(3, 2, 1); surf(x, z, U(:,:,1)); xlabel('x'); ylabel('z'); title('hro'); axis tight
         subplot(3, 2, 2); surf(x, z, U(:, :, 2)); xlabel('x'); ylabel('z'); title('p'); axis tight
```

```
subplot(3, 2, 3); surf(x, z, U(:, :, 3)); xlabel('x'); ylabel('z'); title('ux'); axis tight
         subplot(3, 2, 4); surf(x, z, U(:,:,4)); xlabel('x'); ylabel('z'); title('uz'); axis tight
         subplot(3, 2, 5); surf(x, z, U(:,:,5)); xlabel('x'); ylabel('z'); title('Ay'); axis tight
         figure (3)
         subplot(2, 1, 1); quiver(x, z, Bx, Bz, 2); xlabel('x'); ylabel('z'); title('B'); axis tight
         subplot(2, 1, 2); quiver(x, z, U(:,:,3), U(:,:,4),2); xlabel('x'); ylabel('z'); title('u'); axis tight
         subplot(2, 2, 1); pcolor(x, z, U(:, :, 1)); shading('interp');
         xlabel('x'); ylabel('z'); title(['\rho(x, z, t), t=', num2str(t)])
         subplot(2, 2, 2); pcolor(x, z, U(:, :, 3)); shading('interp');
         xlabel('x');ylabel('z');title(['u_x(x,z,t), t=',num2str(t)])
         subplot(2, 2, 3);pcolor(x, z, Bx); shading('interp');
         xlabel('x');ylabel('z');title(['B_x(x,z,t), t=',num2str(t)])
         subplot(2, 2, 4);
         plot([1:it]*dt, log(Bxm)); xlabel('t'); ylabel('max(Bx(t))'); title(['max(Bx(t)), t=', num2str(t)])
         xlabel('t');ylabel('log(max(Bx))');title(['max(Bx(t)), t=', num2str(t)])
         xlim([0, nt*dt])
         drawnow
     end
end
function Uo=RK4(t, Ui)
 % U(:,:,i), 其中 i= 1 2 3 4 5 --> rho p ux uz ay
global dt
k1=rightEq(t,Ui);
U1=Ui+0.5*dt*k1;
k2 = rightEq(t, U1);
U2=Ui+0.5*dt*k2;
k3=rightEq(t,U2);
U3=Ui+dt*k3;
k4=rightEq(t,U3);
Uo=Ui+dt*(k1+2*k2+2*k3+k4)/6;
function k=rightEq(t,U)
 % U(:,:,i), 其中 i= 1 2 3 4 5 --> rho p ux uz ay
global gamma nu eta dx dz Bx Bz
rho=U(:,:,1);
p=U(:,:,2);
ux=U(:,:,3);
uz=U(:,:,4);
Ay=U(:,:,5);
pt=p+0.5*(Bx.^2+Bz.^2);% pt=p+b^2/2
[nx, nz, \sim] = size(U);
ic=1:nx:
ic=1:nz:
ip=ic+1; ip(nx)=1;
im=ic-1; im(1)=nx;
jp=jc+1; jp(nz)=1;
j_{m}=j_{c}-1; j_{m}(1)=nz;
k(:,:,1) = -ux.*(rho(ip,:)-rho(im,:))/(2*dx) \dots
     -uz.*(rho(:, jp)-rho(:, jm))/(2*dz) ...
     -rho.*((ux(ip, :)-ux(im, :))/(2*dx) ...
     +(uz(:, jp)-uz(:, jm))/(2*dz));
k(:,:,2) = -ux.*(p(ip,:)-p(im,:))/(2*dx) \dots
     -uz.*(p(:, jp)-p(:, jm))/(2*dz) ...
     -gamma.*p.*((ux(ip,:)-ux(im,:))/(2*dx) ...
     +(uz(:, jp)-uz(:, jm))/(2*dz));
k(:,:,3) = -ux.*(ux(ip,:)-ux(im,:))/(2*dx) \dots
     -uz.*(ux(:, jp)-ux(:, jm))/(2*dz) ...
     -1./rho.*(pt(ip,:)-pt(im,:))/(2*dz) ... % 原程序有错误(pt(ipx,:)-pt(imx,:))/(2*dz)中应该是除以 dx 的
     +1./\text{rho.}*(Bx.*(Bx(ip,:)-Bx(im,:))/(2*dx)+Bz.*(Bx(:,jp)-Bx(:,jm))/(2*dz)) ...
     +nu. /rho. *((ux(ip,:)+ux(im,:)-2*ux(ic,:))/(dx*dx)+(ux(:,jp)+ux(:,jm)-2*ux(:,jc))/(dz*dz));
k(:,:,4) = -ux.*(uz(ip,:)-uz(im,:))/(2*dx) \dots
```

```
-uz.*(uz(:, jp)-uz(:, jm))/(2*dz) ...
           -1./\text{rho.}*(\text{pt}(:, jp)-\text{pt}(:, jm))/(2*dz) \dots
           +1./\text{rho.}*(Bx.*(Bz(ip,:)-Bz(im,:))/(2*dx)+Bz.*(Bz(:,jp)-Bz(:,jm))/(2*dz)) ...
           +nu. /rho. *((uz(ip,:)+uz(im,:)-2*uz(ic,:))/(dx*dx)+(uz(:,jp)+uz(:,jm)-2*uz(:,jc))/(dz*dz));
      k(:,:,5) = -ux.*(Ay(ip,:)-Ay(im,:))/(2*dx) \dots
           -uz.*(Ay(:, jp)-Ay(:, jm))/(2*dz) \dots
           + \text{eta.} * ((Ay(ip, :) + Ay(im, :) - 2*Ay(ic, :)) / (dx*dx) + (Ay(:, jp) + Ay(:, jm) - 2*Ay(:, jc)) / (dz*dz));
       % x 的下边界, i=1
      k(1, :, 1) = k(2, :, 1);
      k(1, :, 2) = k(2, :, 2);
      k(1, :, 3)=0;
      k(1, :, 4) = k(2, :, 4);
      k(1, :, 5)=0;
        % x 的上边界, i=end
      k (end, :, 1) = 0 \dots
           -uz (end, :).*(rho(end, jp)-rho(end, jm))/(2*dz) ...
           -\text{rho}(\text{end}, :).*((\text{ux}(\text{end}, :)-\text{ux}(\text{end}-1, :))/(1*\text{dx}) \dots
           +(uz (end, jp) - uz (end, jm)) / (2*dz));
      k \text{ (end, :, 2)} = 0 \dots
           -uz (end, :).*(p(end, jp)-p(end, jm))/(1*dz) ...
           -gamma.*p(end,:).*((ux(end,:)-ux(end-1,:))/(1*dx) ...
           +(uz (end, jp)-uz (end, jm))/(2*dz));
      k (end, :, 3) = 0;
      k \text{ (end, :, 4)} = 0 \dots
           -uz (end, :).*(uz (end, jp)-uz (end, jm))/(2*dz) ...
           -1./\text{rho}(\text{end}, :).*(\text{pt}(\text{end}, \text{jp})-\text{pt}(\text{end}, \text{jm}))/(2*\text{dz}) \dots
           +1./\text{rho}(\text{end}, :).*(Bx(\text{end}, :).*(Bz(\text{end}, :)-Bz(\text{end}-1, :))/(1*dx)+0) ...
           +nu. /rho (end, :). *((2*uz (end-1, :)-2*uz (end, :)))/(dx*dx) + (uz (end, jp) + uz (end, jm) - (end, index))
2*uz(end, jc))/(dz*dz));
      k \text{ (end, :, 5)} = 0 \dots
           -uz (end, :).*(Ay (end, jp) -Ay (end, jm))/(2*dz) ...
           +eta.*((2*Ay(end-1,:)-2*Ay(end,:))/(dx*dx)+(Ay(end,jp)+Ay(end,jm)-2*Ay(end,jc))/(dz*dz));
     function [Bx, Bz]=calcBxz(t, U)
     global dx dz
      [nx, nz, ^{\sim}] = size(U);
      ic=1:nx;
      jc=1:nz;
      ip=ic+1; ip(nx)=1;
      im=ic-1; im(1)=nx;
      jp=jc+1; jp(nz)=1;
      jm=jc-1; jm(1)=nz;
      Bx=-(U(:, jp, 5)-U(:, jm, 5))/(2*dz);
      Bz = (U(ip, :, 5) - U(im, :, 5)) / (2*dx);
      Bx (end, :) = -(U(end, jp, 5) - U(end, jm, 5)) / (2*dz);
      Bz (end, :) = 0;
     end
```

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