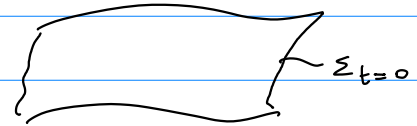


⑥

INITIAL DATA PROBLEM

$$\begin{cases} C_0 := R + K^2 - K_{ij}K^{ij} - 16\pi E = 0 \\ C_i := D_j K_i^j - D_i K - 8\pi P_i = 0 \end{cases}$$

Wart: γ_{ij}, K_{ij} on Σ_0 

(i) Constraints are satisfied

(ii) Physically meaningful

→ BH spacetime, N^m BH spacetime, NS spacetime ...

Here (today): Matter fields are zero or given.

4 equations for 12 unknowns!

→ Need to prescribe 8 quantities, and solve for the other 4.

- CONSTRAINED DATA (4)
- FREE DATA (8)

Q: How to?

A: "Principles":

- 1) Physical (astro) expectation for Σ_0
- 2) Euistic/intuition on the fields we prescribe/solve
- 3) Mathematical necessity
 - $C_a = 0$: solvable!
 - Linear equations, decouple
 - Well-posed BVP

Conformal decomposition allows us to implement 1-3)

Two main formalisms:

- Conformal Transverse Traceless formalism (CTT, York 1973)
- Conformal Thin Sandwich formalism (CTS, York 1993)

LICHNEROWICZ EQ (HAM. CONSTRAINT)

$$(L) \quad \boxed{C_0 := \underbrace{\tilde{D}_i \tilde{D}^i \psi}_{\sim \Delta \psi} - \frac{1}{8} \tilde{A}_{ij} \tilde{A}^{ij} \psi^{-7} + \left(-\frac{K^2}{12} + 2\pi E \right) \psi^5 = 0} \quad (p = -10)$$

$\sim \Delta \psi \rightsquigarrow$ Nonlinear eq for ψ !?

- An eq for the conf. factor ψ (elliptic type?)
- $K^2 \equiv 0$ the equation simplifies...

\hookrightarrow the BVP can be studied: several known results about well-posedness of (L) under the hypothesis $\boxed{K \equiv \text{const}}$

Constant Mean Curvature (CMC) spacetime

Example:

$\left. \begin{array}{l} \text{A.F. spacetimes} \\ \text{CMC} \quad -''- \quad w/ K \equiv 0 \\ E = 0 \end{array} \right\} \Rightarrow \text{BVP w/ (L) is solvable for a "large" class of } \tilde{\gamma}_{ij}$

Prototype equation for (L) w/ $K \equiv 0$:

$$\Delta u = +f^2 u^p \quad \text{in flat spacetime}$$

Theorems: (i) $p=1$ (linear): BVP w/ data on $\partial\Omega$ boundary of domain and $u|_{\partial\Omega} \equiv 0 \Rightarrow u \equiv 0$ unique

(ii) (nonlinear) unique solution iff $p > 0$

(same sign in front of the positive coefficient of u^p)

How about (L)? Linearize (L) and use (i) [set $K \equiv 0$]:

$$(L): \tilde{\Delta} \psi + H(\psi) = 0$$

$$\psi = \psi_0 + \epsilon \quad \psi_0: \text{solution}$$

$$H(\psi) = H(\psi_0) + \frac{\partial H}{\partial E} \Big|_0 \epsilon + \mathcal{O}(\epsilon^2)$$

$$\hookrightarrow \tilde{\Delta} \epsilon = f \epsilon \quad \text{with:}$$

$$f = \frac{1}{8} \tilde{R} + \frac{7}{8} \hat{A}_{ij} \hat{A}^{ij} \psi_0^{-8} - \underbrace{10\pi E \psi_0^4}$$

$$\text{Note: } K=0 \Rightarrow \tilde{R} > 0$$

f is not positive because of the "matter term"

\hookrightarrow not solvable!

We can still solve (L) w/ $K=0$ if we rescale the matter term:

$$\tilde{E} := \psi^s E$$

$$\hookrightarrow -5 \cdot 2\pi E \psi_0^4 \rightarrow \underbrace{+(s-5) 2\pi \tilde{E} \psi_0^{4-s}}_{>0 \quad \forall s>5}$$

This is not a trick:

(!) \Rightarrow impracticable nonlinear elliptic eqs are solved by iteration (linearization)

\hookrightarrow specifying the \tilde{E} r.h.s. is key to find a numerical solution

Q: What to use for s ?

A: $S=8$: this way the dominant energy condition can be expressed in conf. variables:

$$\tilde{E} := \psi^8 E \quad \tilde{P}^i := \psi^{10} P^i$$

$$\tilde{E}^2 \geq \tilde{P}^2 \Rightarrow E^2 \geq P^2$$

$$[\psi^{2s} E^2 \geq \psi^{-4} \psi^{10} \psi^{10} P^2]$$

MAXIMAL SLICING

Q: What is the meaning of $K \equiv 0$?

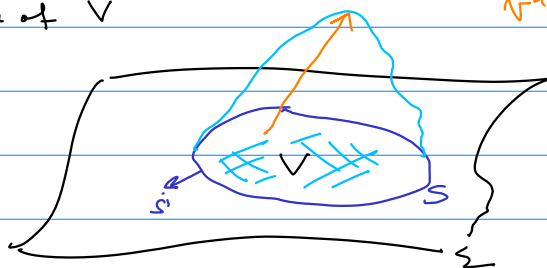
A: Gauge condition that extremize the volume of Σ .

$$K = \gamma^{ab} \underline{k_{ab}} = -\frac{1}{2\alpha} \gamma^{ij} \underline{\mathcal{L}_m \gamma_{ij}} = -\frac{1}{2\alpha} \mathcal{L}_m \ln \gamma = -\frac{1}{2\alpha} \mathcal{L}_m \ln \sqrt{\gamma}$$

↑ identity (see last lecture)
↑ determinant
↑ volume element! on Σ

Volume: $V = \int \sqrt{\gamma} d^3x$

Perform a variation of V



$$n^a = \text{St}(\alpha n^a + \beta^a) :$$

$$n^a|_S = 0$$

$$\frac{\delta V}{\delta t} = \int_V \underline{\partial_t \sqrt{\gamma} d^3x} = \int_V \underline{(-\alpha K + D_i \beta^i) \sqrt{\gamma} d^3x} = -\int_V \alpha K \sqrt{\gamma} d^3x + \underbrace{\oint_S \beta^i S_i}_{=0} = -\int_V \alpha K \sqrt{\gamma} d^3x$$

$K=0 \Rightarrow \delta V=0$

Example: Some diff.-geom. problem of film of soap on a ring S
 \hookrightarrow the soap film minimize the V (Euclidean geom).

But we are in a Lorentzian geometry \Rightarrow MAXIMUM.

CTT ($p = -10$ scaling of K_{ij})

- use (L) for $C_0 = 0$
- Derive an eq from $C_i = 0$ for a quantity X^i (vector) obtained by a further decomposition of

$$\hat{A}^{ij} := \hat{A}_L^{ij} + \hat{A}_\pi^{ij} \quad (1)$$

where:

$$\pi\pi : \tilde{\gamma}_{ij} \hat{A}^{ij} = 0 \quad \text{and} \quad \tilde{B}_j \hat{A}^{ij}_\pi = 0$$

$$L : (\tilde{L}X)^{ij} = \tilde{D}^i X^j + \tilde{D}^j X^i - \frac{2}{3} \tilde{D}_k X^k \tilde{\gamma}^{ij}$$

Conformal Killing operator on X^i

determined by taking the divergence of (1):

$$\tilde{D}_j \hat{A}^{ij} = \tilde{D}_j (\tilde{L}X)^{ij} = \tilde{B}_j \tilde{D}^j X^i + \frac{1}{3} \tilde{B}^i \tilde{D}_j X^j + \tilde{R}^i_j X^j =: \underline{\tilde{\Delta}_L X^i}$$

Conformal vector Laplacian operator

$\exists!$ L+TT decomposition (1) iff $\exists! X^i$ of the conf. Laplacian eq:

$$\boxed{\tilde{\Delta}_L X^i = \tilde{B}_j \hat{A}^{ij}}$$

Th (Contar 179):

$$\left. \begin{array}{l} \Sigma \text{ is K.F.} \\ \partial_k \partial_l \tilde{\gamma}_{ij} = \mathcal{O}(r^{-3}) \end{array} \right\} \Rightarrow \exists! \text{ is guaranteed.}$$

CTT

$$\left\{ \begin{array}{ll} (L) & (C_0 = 0) \\ \tilde{\Delta}_L X^i - \frac{2}{3} \tilde{B}^i_k \psi^k - 8\pi \tilde{P}^i = 0 & (C_i = 0) \end{array} \right.$$

Constrained data: ψ, X^i

Free data: $\tilde{\gamma}_{ij}, \hat{A}_{\pi}^{ij}, K$

Note: (i) under max. slicing $K \equiv 0$
the 2 equations decouple!

(ii) for CMC slices they partially decouple
(can solve one after the other)

(iii) $\tilde{\gamma}_{ij}, \hat{A}_{\pi}^{ij} \leadsto$ GW content of Σ

CTT: GNF. #LAT, ASYMPTOTICALLY #LAT, MAX. SLICING DATA, VACUUM

2) 1) 3) 4)

Simplest choice of free data:

$$\begin{aligned}\tilde{\gamma}_{ij} &= f_{ij} & 2) \\ K &\equiv 0 & 3) \\ E = 0 = P^i & & 4) \\ \hat{A}_{\pi}^{ij} &\equiv 0 & 5)\end{aligned}$$

$$\begin{aligned}L \rightarrow \tilde{D}_i &= D_i & \tilde{D}_i \tilde{D}^i &= D_i D^i = \Delta \\ \tilde{R} &= 0 & \tilde{L} &= L\end{aligned}$$

$$\begin{aligned}L \rightarrow \text{CTT: } & \begin{cases} \Delta \psi + \frac{1}{3} (LX)_{;i} (LX)^{;i} \psi^{-7} = 0 & (\text{CTT1}) \\ \Delta_L X^i = \Delta X^i + \frac{1}{3} D_j D^j X^i = 0 & (\text{CTT2}) \end{cases}\end{aligned}$$

These are "easy": flat (Euclidean) operators & decoupled.

BVP:

OUTER BCs: A.F. $\psi = 1, X^i = 0$ at i_0 ($r \rightarrow +\infty$)
INNER BCs (stray fields): fix the topology

(*)

CASE 1: $\textcircled{*}$ $\Sigma = \mathbb{R}^3$ (No inner BC)

$$(CTT_2) \Rightarrow \dot{x}^i \equiv 0$$

$$(CTR_1) \Rightarrow \begin{cases} \Delta \psi = 0 \\ \psi = 1 \end{cases} \quad r \rightarrow +\infty \Rightarrow \psi \equiv 1$$

\hookrightarrow flat spacetime.

If we want a nontrivial solution we need an inner BC.

... Next time ...