STRAPDOWN SYSTEM ALGORITHMS

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By

Paul G. Savage
President
Strapdown Associates, Inc.
Woodbridge Plaza, Suite 150
10201 Wayzata Blvd.
Minnetonka, Minnesota 55343

SUMMARY

This papar addresses the attitude datermination, acceleration transformation, and attitude/hadding output computational operations parformed in modern-day strapdown inertial navigation systams. Contemporary algorithms are described for implementing these operations in raal-time computers. The attitude determination and acceleration transformation algorithm discussions are based on the two-speed approach in which high frequency coming and sculling effects are calculated with simplified high speed algorithms, with resulte fed into lower speed higher order algorithms. This is the approach that is typically used in most modern-day strapdown systems. Design equations are included for evaluating the performance of the strapdown computer algorithms as a function of computer execution speed and sensor assembly vibration amplitude/frequency/phase environment.

Both direction cosine and quaternion based attitude algorithms are described and compared in light of modern-day algorithm accuracy capabilities. Orthogonality and normalization operations are addressed for potential attitude elgorithm accuracy enhancement. The section on attitude data output algorithms includes a discussion on roll/yaw Euler angle singularities near high/low pitch angle conditions.

1. INTRODUCTION

The concept of strapdown inertial navigation was originated more than thirty years ago, largely from an analytical stendpoint. The theoretical enalytical expressions for processing etrapdown inertial sensor data to develop attitude, velocity, and position information were reesonably well understood in the form of continuous matrix operations and differential equations. The implementation of these equations in a digitial computer, however, was invariably keyed to severe throughput limitations of original airborne digitial computer technology. As a result, many of the strapdown computational algorithms originated during these early periods were inherently limited in accuracy, particularly under high frequency dynamic motion. A classical test for elgorithm accuracy during this early period was how well the algorithm computed ettitude under cyclic coning motion as the coning frequency approached the computer updata cycle frequency.

In the late 1960's and early 1970's, several analytical efforts eddrassed the problem of splitting the strapdown computation process into low end high speed sections (7, 8, 10). The low speed section contained the bulk of the computational equations, end wes designed to accurately account for low fraquency lerge emplitude dynamic motion effects (a.g., vehicle maneuvering). The high speed computation section was designed with a small set of simple algorithms that would accurately account for high frequency small amplitude dynamic motion (e.g., vehicle vibrations). Splitting the computational process in this manner allowed the bulk of the strapdown algorithms to be iterated at reasonable speeds competible with computer throughput limitations. The high speed elgorithms were simple anough that they could be machenized individually with special purpose alactronics, or as a minor high speed loop in the main processor.

Over the past ten years, the structura of most etrapdown algorithms has evolved into this two speed structura. The techniques have been refined today so that fairly straight-forward enelytical design methods can be used to define elgorithm enalytical forms and computational rates to echieve required levels of parformance in specified dynamic environments.

This paper dascribes the algorithms used today in most modern-day strapdown inertial navigation systams to calculate ettitude end transform acceleration vector measuraments from sensor to navigation axes. The algorithms for integrating the transformed eccelerations into velocity and position data are not addressed because it is believed that these operations are generic to inertial navigation in general, not only strapdown inertial navigation.

For the algorithms discussed, the analytical basis is presented together with a discussion on general design methodology used to develop the algorithms for compatibility with particular user accuracy and anvironmental requirements.

STRAPDOWN COMPUTATION OPERATIONS

Figure 1 depicts tha computational elamants implemented by software algorithms in typical atrapdown inertial navigation systems. Input data to the algorithms is provided from a triad of strapdown gyroa and acceleromatera. The gyros provide pracision measurements of strapdown ensor coordinate frame ("body axea") angular rotation rate relative to nonrotating inertial space. The accelarometers provide precision measurements of 3-axis orthogonal apecific force acceleration along body axes.

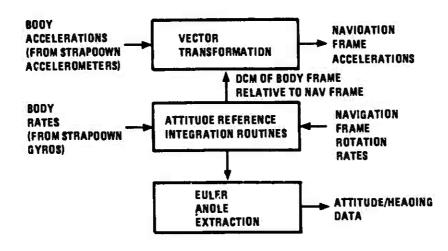


FIGURE 1 - STRAPDOWN ATTITUDE REFERENCE OPERATIONS

The strapdown gyro data is processed on an iterative basis by suitable integration algorithms to calculate the attitude of the body frame ralative to navigation coordinates. The rotation rate of the navigation frame is an input to tha calculation from the navigation section of the overall computation software. Typical navigation coordinate frames are oriented with the z-axis vartical and the x, y, axes horizontal.

The attitude information calculated from the gyro and navigation frama rate data is used to transform the accelerometer specific force vector measurements in body axas to their equivalent form in navigation coordinatas. The navigation frame spacific force accalerations are then integrated in the navigation software saction to calculate velocity and position. The velocity/position computational algorithms are not unique to the strapdown mechanization concept, hence, are not treated in this paper. Sevaral texts treat the velocity/position integration algorithms in datail (1, 2, 3, 4, 12).

Figure 1 also shows an Eular Angla Extraction function as part of the strapdown attitude reference operations. This algorithm is used to convert the calculated attitude data into an output format that is more compatible with typical user requiramants (e.g., roll, pitch, heading Euler angles).

STRAPDOWN ATTITUDE INTEGRATION ALGORITHMS

The attituda information in strapdown inertial navigation systems is typically calculated in the form of a direction cosine matrix or αs an attitude quaternion. The direction cosine matrix is a threa-by-three matrix whose rows represent unit vectors in navigation axes projected along body axas. As such, the alamant in the ith row and jth column represents the cosine of the angle between the navigation frame i-axis and body frame j-axis. The quaternion is a four-vector whose elements are defined analytically (5, 9) as follows:

 $a = (\alpha_{\chi}/\alpha) \sin (\alpha/2)$

b = $(\alpha_y/\alpha) \sin (\alpha/2)$ c = $(\alpha_z/\alpha) \sin (\alpha/2)$

cos (a/2)

where

$$\alpha_{X'}\alpha_{Y'}\alpha_{Z} = Componante of an engla vector $\underline{\alpha}$.

Magnitude of $\underline{\alpha}$.$$

Tha $\underline{\alpha}$ vector is defined to have direction and magnitude such that if the nevigetion frame was rotated about $\underline{\alpha}$ through an angle α , it would be rotated into elignment with the body frame. The $\underline{\alpha}$ rotation angla vector and its quaternion equivalent (a, b, c, d, from equations (1)), or the direction cosine matrix, uniquely define the attitude of the body axes relative to navigation exes.

3.1 Direction Cosina Updeting Algorithms

3.1.1 Direction Comina Updating Algorithm For Body Rotetions

The direction cosine matrix can be updeted for body frame gyro sensed motion in the strepdown computer by executing the following classical direction cosine matrix chain rule elgorithm on a repetative basis:

$$C(m+1) = C(m) \lambda(m)$$
 (2)

where

- C(m) = Direction coming matrix relating body to navigetion axaa at the mth computer cycle time
- A(m) = Direction cosine martix that transforms vectors from body coordinets et the (m+1)th computer cycle to body coordinetes et the mth computer cycle.

It is well known (9) thet:

$$A(m) = I + f_1(\phi x) + f_2(\phi x)^2$$
 (3)

where

$$f_1 = \frac{\sin \phi}{\phi} = 1 - \phi^2/31 + \phi^4/41 - \cdots$$

$$f_2 = \frac{1 - \cos \phi}{\phi^2} = 1/21 - \phi^2/41 + \phi^4/61 - \cdots$$

$$\phi^2 = \phi_x^2 + \phi_y^2 + \phi_z^2 \tag{4}$$

$$(\underline{\phi}\mathbf{x}) \stackrel{\Delta}{=} \begin{bmatrix} 0 & -\phi_{\mathbf{Z}} & \phi_{\mathbf{Y}} \\ \phi_{\mathbf{Z}} & 0 & -\phi_{\mathbf{X}} \\ -\phi_{\mathbf{Y}} & \phi_{\mathbf{X}} & 0 \end{bmatrix}$$

I = 3 x 3 unity matrix

 $\phi_{\mathbf{X}}, \phi_{\mathbf{Y}}, \phi_{\mathbf{Z}} = \text{Components of } \underline{\phi}.$

Angle vector with direction end magnitude such that e rotation of the body frame about \$\phi\$ through an angle equal to the magnitude of \$\phi\$ will rotate the body frame from its orientation at computer cycle m to its orientation at computer cycle m+1. The \$\phi\$ vector is computed for each computer cycle m by processing the dete from the strepdown gyros. The elgorithm for computing \$\phi\$ will be described subsequently.

The "order" of the algorithm defined by equations (2) through (4) is determined by the number of terms carried in ths f_1 , f_2 expansions. A fifth order algorithm, for example, retains sufficient terms in f_1 and f_2 such that A(m) contains all ϕ term products out to fifth order. Hence, f_1 would be truncated after the ϕ^4 term and f_2 would be truncated after the ϕ^2 term to retain fifth order accuracy in A(m). The order of accuracy required is determined by system accuracy requirements under maximum rate input conditions when ϕ is a maximum. The computation iteration rate is typically selected to assure that ϕ remains small at maximum rate (e.g., 0.1 radians). This assures that the number of terms required for accuracy in the f_1 , f_2 expansions will be reasonable.

3.1.2 Direction Coeine Updating Algorithm For Navigation Frame Rotations

Equation (2) is used to update the direction cosine matrix for gyro sensed body frame motion. In order to update the direction cosines for rotation of the navigation coordinate frame, the following classical direction cosine matrix chain rule algorithm is used:

$$C(n+1) = B(n) C(n)$$
(5)

where

B(n) = Direction cosine matrix that transforms vectors from navigation axes at computer cycle n to navigation axes at computer cycls (n+1).

The equation for B(n) parallels equation (3):

$$B(n) = I - (\theta x) + 0.5(\theta x)^2$$
 (6)

with

$$(\underline{\theta}\mathbf{x}) \stackrel{\Delta}{=} \begin{bmatrix} 0 & -\theta_{\mathbf{z}} & \theta_{\mathbf{y}} \\ \theta_{\mathbf{z}} & 0 & -\theta_{\mathbf{x}} \\ -\theta_{\mathbf{y}} & \theta_{\mathbf{x}} & 0 \end{bmatrix}$$
 (7)

where

 $\theta_{x}, \theta_{y}, \theta_{z} = Components of \theta$.

Angle vector with direction and magnitude euch that a rotation of the navigation frame about $\underline{\theta}$ through an angle equal to the magnitude of $\underline{\theta}$ will rotate the navigation frame from its orientation at computer cycle n to its orientation at computer cycle n+1. The $\underline{\theta}$ vector is computed for each computer cycle n by processing the navigation frame rotation rate data from the navigation software section (12).

It is important to note that the n cycle (for navigation frame rotation) and m cycle (for body frame rotation) are generally different, n typically being executed et a lower iteration rate than m. This is permissable because the navigation frame rotation rates are considerably smaller than the body rates, hence, high execution rates are not needed to maintain $\underline{\theta}$ small to reduce the order of the iteration algorithm. The algorithm represented by equations (5) and (6) is second order in $\underline{\theta}$. Generally, first order is of sufficient accuracy, and the $(\underline{\theta}\mathbf{x})^2$ term need not be carried in the actual software implementation.

3.2 Quaternion Updating Algorithms

3.2.1 Quaternion Transformation Properties

The updating algorithms for the attitude quaternion can be developed through an investigation of its vector transformation properties (5, 9). We first introduce nomenclature that is useful for describing quaternion algebraic operations. Referring to equation (1), the quaternion with components a, b, c, d, can be described as:

Whara

a,b,c = Components of the "vector" part of the quatarnion.

i,j,k = Qusternion vector operators analagous to unit vactors along orthogonal coordinata axes.

d = "Scalar" part of the quaternion.

Ws also define rules for quaternion vector operator products as:

$$ii = -1$$
 $ij = k$ $ji = -k$
 $jj = -1$ $jk = i$ $kj = -i$
 $kk = -1$ $ki = j$ $ik = -j$

With the above dsfinitions, tha product w of two quatarnions (u and v) becomes:

or in "Four-vector" matrix form:

Ws also define tha "complex conjugats" of the general quaternion u in equation (8) as:

$$u^* = -ai - bi - ck + d$$

We now define a quatarnion operator h(m) for tha body angle change & over computar cycle m as:

$$h(m) = \begin{cases} (\phi_{X}/\phi) & \sin (\phi/2) \\ (\phi_{Y}/\phi) & \sin (\phi/2) \\ (\phi_{Z}/\psi) & \sin (\phi/2) \\ \cos (\phi/2) \end{cases}$$
(9)

where the alements in the above column matrix refer to the i, j, k, and scalar components of h. We also define a general vector \underline{v} with commonents v_x , v_y , v_z , and a corresponding quaternion v having the same vector components with a zero scalar component:

$$\mathbf{v} = \begin{bmatrix} \mathbf{v}_{\mathbf{X}} \\ \mathbf{v}_{\mathbf{Y}} \\ \mathbf{v}_{\mathbf{Z}} \\ \mathbf{0}^{\mathbf{Z}} \end{bmatrix}$$

Using the above definitions and the ganaral rules for quatarnion algabra, it is readily damonstrated by substitution and trigonometric manipulation that:

$$v' = h(m) \cdot v \cdot h(m) * = A'(m) \cdot v \tag{10}$$

whers

$$A'(m) \stackrel{\Delta}{=} \begin{bmatrix} A(m) & 0 \\ 0 & 0 \end{bmatrix}$$

$$\mathbf{v}' \qquad \stackrel{\Delta}{=} \qquad \begin{bmatrix} \mathbf{v}_{\mathbf{x}'} \\ \mathbf{v}_{\mathbf{y}'} \\ \mathbf{v}_{\mathbf{g}'} \end{bmatrix}$$

A(m) = As defined in (3).

Equation (10), therefore, is the quaternion form of the vector transformation equation that transforms a vector from body coordinates at computer cycle (m+1) to body coordinates at computer cycle m:

$$\mathbf{v}' = \mathbf{A}(\mathbf{m}) \ \underline{\mathbf{v}} \tag{11}$$

whera

 $\underline{v}',\underline{v}$ = "Thras-vector" form of v' and v (i.a., with componants v_x' , v_y' , v_z' and v_x , v_y , v_z).

v = The general vector v in body coordinates at computer cycle (m+1).

v' = The general vector v in body coordinates at computer cycle m.

3.2.2 Quaternion Updating Algorithm For Body Movion

Equation (10) with ita equation (11) dual can be used to dafine analagous vactor transformation operations between body coordinates and navigation coordinates at computer cycla m as:

$$v'' = q(m) v' q(m)*$$
 $v'' = C(m) v'$
(12)

whera

q(m) = Quatarnion relating body axss to navigation axas at computer cycla m.

v' = Tha vector v in navigation coordinates.

v" = The vector v in body coordinates at computer cycle m.

 $v',v'' = Quatarnion ("Four vsctor") form of <math>\underline{v}', \underline{v}''$.

The q qustarnion has four alements (i.e., a, b, c, d) that are updated for body motion at each computer cycle m. The updating equation is easily derived by substituting equation (10) into (12):

$$v'' = q(m) h(m) v h(m) * q(m) *$$

Using the dafinition for the quatarnion complax conjugata, it is raadily demonstrated that:

$$h(m)* q(m)* = (q(m) h(m))*$$

Thus,

$$v^{\alpha} = q(m) h(m) v (h(m) q(m))*$$

But we can also write the direct expression:

$$v'' = q(m+1) v q(m+1)*$$

Tharafore, by direct comparison of tha lattar two equations:

$$q(m+1) = q(m) h(m)$$
 (13)

Equation (13) is the quaternion equivalent to direction cosine updating equation (2). For computational purposes, h(m) as defined in equations (9) is equivalently:

$$h(m) = \begin{cases} f_3 & \phi x \\ f_3 & \phi y \\ f_3 & \phi z \\ f_4 & \end{cases}$$

$$f_3 = \frac{\sin (\phi/2)}{\phi} = 0.5(1 - (0.5\phi)^2/31 + (0.5\phi)^4/51 - \cdots)$$

(14)

$$f_4 = cos (\phi/2) = 1 - (0.5\phi)^2/2i + (0.5\phi)^4/4i - \cdots$$

$$(0.5)^2 = 0.25 (\phi_x^2 + \phi_y^2 + \phi_z^2)$$

The "order" of the equation (13) and (14) updating algorithm depends on the order of ϕ terms carried in h which depends on the truncation point used in f3 and f4. The rationala for selecting the algorithm order and associated algorithm iteration rate is directly analagous to selection of the direction cosins updating algorithm order (discussed previously).

3.2.3 Quaternion Updating Algorithm For Navigation Frams Rotation

Equation (13) with (14) is used to updata the quaternion for body frame motion sensed by gyros. In order to update the quatarnion for rotation of the navigation coordinata frame, an algorithm analogous to aquation (5) (for the direction cosine matrix) is used with a navigation frame rotation quaternion r:

$$q(n+1) = r(n) q(n)$$

$$\mathbf{r}(\mathbf{n}) = \begin{bmatrix} -0.5 & \theta_{\mathbf{x}} \\ -0.5 & \theta_{\mathbf{y}} \\ -0.5 & \theta_{\mathbf{z}} \\ 1-0.5(\theta/2)^2 \end{bmatrix}$$
(15)

$$(\theta/2)^2 = 0.25 (\theta_x^2 + \theta_y^2 + \theta_z^2)$$

where

$$\theta_x, \theta_y, \theta_z = \text{Components of } \frac{\theta}{2} \text{ as defined previously for equations}$$
 (6) and (7).

The development of equation (15) parsllels tha development of (13). The equation for r(n) is a truncated form of the thaoretical exact analytical exprassion (analagous to the sacond order truncated form of equation (14)). The θ^2 tarm in equation (15) ganarally is not required for accuracy (due to the smallness of $\underline{\theta}$ in typical spplications).

As for the direction cosine updating algorithm for navigation frame motion, the equivalent quaternion updating algorithm (equation (15)) updating cycle n naed not be procassed as fast as the body rate cycle m to maintain equivalent accuracy. This is due to the considerably smaller navigation frame rotation rates compared to body rotation rates.

3.2.4 Equivalancies Betwean Direction Cosins And Quaternion Elamants

The analytical equivalency between the elements of the diraction cosine matrix and the attitude quaternion can be derived by diract expansion of equations (12). If we define the elements of q as:

equation (12) becomes after expansion, factorization of v', and neglecting the scalar part of the v'' and v' quaternion vectors (i.e., carrying only the vector components v'' and v'):

$$\underline{v}^{"} = \begin{bmatrix} (d^{2} + a^{2} - b^{2} - c^{2}) & 2(ab - cd) & 2(ac + bd) \\ 2(ab + cd) & (d^{2} + b^{2} - c^{2} - a^{2}) & 2(bc - ad) \\ 2(ac - bd) & 2(bc + ad) & (d^{2} + c^{2} - a^{2} - b^{2}) \end{bmatrix} \underline{v}'$$
 (16)

Defining C in equation (12) as:

$$C = \begin{bmatrix} C_{11} & C_{12} & C_{13} \\ C_{21} & C_{22} & C_{23} \\ C_{31} & C_{32} & C_{33} \end{bmatrix}$$

equation (16) when compared with (12) ahowa that:

$$C_{11} = d^{2} + a^{2} - b^{2} - c^{2}$$

$$C_{12} = 2(ab - cd)$$

$$C_{13} = 2(ac + bd)$$

$$C_{21} = 2(ab + cd)$$

$$C_{22} = d^{2} + b^{2} - c^{2} - a^{2}$$

$$C_{23} = 2(bc - sd)$$

$$C_{31} = 2(ac - bd)$$

$$C_{32} = 2(bc + ed)$$

$$C_{33} = d^{2} + c^{2} - a^{2} - b^{2}$$

$$(17)$$

The converse of equation (17) is a omewhat more complicated. Using the property (from equation (1)) that:

$$a^2 + b^2 + c^2 + d^2 = 1$$

ths converse of equation (17) can be shown (11) to be computable from the following sequence of operations:

operations:

$$T_{r} = C_{11} + C_{22} + C_{33}$$

$$P_{1} = 1 + 2C_{11} - T^{2}$$

$$P_{2} = 1 + 2C_{23} - T^{2}$$

$$P_{3} = 1 + 2C_{23} - T^{2}$$

$$P_{3} = 1 + 2C_{33} - T^{2}$$

If $P_{1} = \max \{P_{1}, P_{2}, P_{3}, P_{0}\}$, then:
$$P_{2} = 1 + T^{2}$$

$$P_{3} = 1 + T^{2}$$

If $P_{2} = \max \{P_{1}, P_{2}, P_{3}, P_{0}\}$, then:
$$P_{3} = (C_{13} + C_{13})/4a$$

$$P_{4} = (C_{32} - C_{23})/4a$$

If $P_{2} = \max \{P_{1}, P_{2}, P_{3}, P_{0}\}$, then:
$$P_{3} = (C_{32} + C_{23})/4b$$

$$P_{4} = (C_{13} + C_{13})/4b$$

If $P_{3} = \max \{P_{1}, P_{2}, P_{3}, P_{0}\}$, then:
$$P_{3} = \max \{P_{1}, P_{2}, P_{3}, P_{0}\}$$
, then:
$$P_{4} = (C_{21} + C_{12})/4b$$

If $P_{3} = \max \{P_{1}, P_{2}, P_{3}, P_{0}\}$, then:
$$P_{4} = (C_{32} + C_{23})/4c$$

$$P_{5} = (C_{32} + C_{23})/4c$$

If $P_{6} = \max \{P_{1}, P_{2}, P_{3}, P_{0}\}$, then:
$$P_{6} = (C_{32} + C_{23})/4c$$

If $P_{6} = \max \{P_{1}, P_{2}, P_{3}, P_{0}\}$, then:
$$P_{6} = (C_{32} + C_{23})/4c$$

If $P_{6} = \max \{P_{1}, P_{2}, P_{3}, P_{0}\}$, then:
$$P_{6} = (C_{32} + C_{23})/4c$$

$$P_{6} = (C_{32} + C_{23})/4c$$

If $P_{6} = (C_{32} + C_{23})/4c$

$$P_{6} = (C_{32} + C_{23})/4c$$

$$P_{6} = (C_{32} + C_{23})/4c$$

$$P_{7} = (C_{13} + C_{23})/4c$$

$$P_{7} = (C_{13} + C_{13})/4c$$

3.3 The Computation Of •

3.3.1 Continous Form

The ϕ "body attitude change" vector is calculated by processing data from the strapdown giros. Under situations where the angular rotation rate vector (sensed by the gyros) lies along a fixed direction (i.e., is nonrotating in inertial space), the ϕ vector is equal to the simple integral of the angular rats vector over the time interval from computer cycle m to computer cycle (m+1):

$$\underline{\bullet} = \int_{t_m}^{t_{m+1}} \underline{\omega} dt$$
 for casse when $\underline{\omega}$ is nonrotating. (19)

where

 $\underline{\omega}$ = Angular rate vector eeneed by the strapdown gyros.

Under general motion conditions (when $\underline{\omega}$ may be rotating), equation (19) has the more complex form (as derived in (10) or alternatively, in Appendix A):

$$\underline{\alpha}(t) = \int_{t_{m}}^{t} (\underline{\omega} + 1/2 \underline{\alpha} \times \underline{\omega} + \frac{1}{\alpha^{2}} (1 - \frac{\alpha \sin \alpha}{(1 - \sin \alpha)}) \underline{\alpha} \times (\underline{\alpha} \times \underline{\omega})) dt$$
(20)

 $\underline{\phi} = \underline{\alpha}(t=t_{m+1})$

It can varified by power series expansion that to first order,

$$(1/\alpha^2)$$
 $(1-\frac{\alpha \text{ ein }\alpha}{(1-\cos\alpha)})$ = $\frac{1}{12}$

Hence, $\alpha(t)$ in equation (20), to third order accuracy in α can be approximated by:

$$\underline{\alpha}(t) \sim \int_{t_{m}}^{t} (\underline{\omega} + 1/2 \underline{\alpha} \times \underline{\omega} + \underline{1}_{2} \underline{\alpha} \times (\underline{\alpha} \times \underline{\omega})) dt$$
 (21)

A second order expression for $\alpha(t)$ can be obtained from (21) by dropping the 1/12 term. An even simpler expression for $\alpha(t)$ is obtained by dropping the 1/12 term, and approximating the α term in the integral by the direct integral of $\underline{\omega}$:

$$\underline{\beta}(t) = \int_{t_m}^{t} \underline{\omega} dt$$

$$\delta \underline{\rho}(t) = 1/2 \int_{t_{m}}^{t} \underline{\rho} \times \underline{\omega} dt$$
 (22)

$$\phi = \beta(t=t_{m+1}) + \delta\beta(t=t_{m+1})$$

An interesting characteristic about equation (22) is that its accuracy is in fact comparable to that of third order equation (21). In other words, the simplifying assumption of replacing α with β in the 1/2 α x ω term is in fact equivalent to introducing an error in equation (21) that to third order, equals the 1/12 α x $(\alpha$ x ω) term. This property can be verified by simulation as well as analytical expansion under hypothesized angular motion conditions.

Equation (22) is the equation that is mechanized in software in most modern-day strapdown inertial navigation systems to calcuste ϕ . It can be demonstrated analytically and by simulation that for representative vehicle angular motion and vibration, equation (22) faithfully calculates ϕ to accuracy levels that are compatible with high performance strapdown inertial navigation system requirements.

For situations where $\underline{\omega}$ is nonrotating, the $\underline{\delta \beta}$ term in (22) is zero and $\underline{\phi}$ equals the simple time integral or $\underline{\omega}$ over the computer interval m (i.e., the equation (19) approximation). For situations where $\underline{\omega}$ is rotating (s situation defined analytically as

"coning"), the $\underline{\delta \beta}$ tarm is nonzero and must be calculated and used as a correction to the $\underline{\omega}$ integral to properly calculate $\underline{\phi}$.

It is important to note that the accuracy by which equation (22) approximates (20) is dependant on ϕ being small (e.g., less than 0.1 radian). In order to protect the accuracy of this approximation, the computer iteration rats must be high anough that ϕ remains small under maximum vehicle rotation rate conditions.

3.3.2 Recursiva Algorithm Form

The implementation of equation (22) in a digital computar implies that a higher spaed integration summing operation be performed during each body motion attitude updata cycle. A computational algorithm for the integration function can be derived by first rewriting equation (22) in the equivalent incremental updating form:

$$\underline{\beta}(t) = \underline{\beta}(1) + \int_{t_{\underline{1}}}^{t} \underline{\omega} dt$$

$$\frac{\delta\beta(\bar{x}+1)}{\delta\beta(\bar{x})} = \frac{\delta\beta(\bar{x})}{t} + 1/2 \int_{t_2}^{t_{\bar{x}}+1} \underline{\beta}(t) \times \underline{\omega} dt \qquad (23)$$

$$\beta(1+1) = \beta(t=t_{\ell+1})$$

with initial conditions:

$$\underline{\beta}(t=t_m) = 0 \tag{24}$$

$$\underline{\delta\beta}(t=t_m) = 0$$

where

High spead computar cycla within the m body rate updats cycla.

The integrals in (23) can be replaced by analytical forms that are compatible with gyro input data processing if $\underline{\omega}$ is raplaced by a generalized time series expansion. For aquations (23), it is sufficient to approximate $\underline{\omega}$ over the $\frac{1}{2}$ to $\frac{1}{2}+1$ time interval as a constant plus a linear ramp:

$$\underline{\omega} \sim \underline{A} + \underline{B} (t - t_{1}) \tag{25}$$

where

 \underline{A} , \underline{B} = Constant vactors.

Substituting (25) in (23), and racognizing with tha equation (25) approximation that:

$$\underline{A}(t_{\underline{t+1}} - t_{\underline{t}}) = 1/2 \left(\underline{A\theta}(\underline{t}) + \underline{A\theta}(\underline{t-1})\right)$$

$$1/2 B(t_{t+1} - t_t)^2 = 1/2 (\Delta \theta(t) - \Delta \theta(t-1))$$

where by dafinition:

$$\Delta \theta(z) \triangleq \int_{t_z}^{t_{z+1}} \underline{\omega} dt$$

yialds the dasirad final form for tha & updating algorithm:

$$\frac{\delta \beta(\lambda+1)}{\delta \beta(\lambda)} = \frac{\delta \beta(\lambda)}{\delta \beta(\lambda)} + \frac{1}{2} \left(\frac{\beta(\lambda)}{\delta \beta(\lambda)} + \frac{1}{6} \frac{\delta \theta(\lambda-1)}{\delta \beta(\lambda)} \right) \times \frac{\delta \theta(\lambda)}{\delta \beta(\lambda)}$$

$$\underline{\Delta\theta}(1) \qquad \approx \int_{t_{\perp}}^{t_{\perp}+1} \underline{\omega} \, dt = \int_{t_{\perp}}^{t_{\perp}+1} \underline{d\theta}$$
 (26)

$$\underline{B}(\pm\pm1) = \underline{B}(\pm) + \underline{A}\underline{\theta}(\pm)$$

$$\underline{\bullet} = \underline{\beta}(t=t_{m+1}) + \underline{\delta\beta}(t=t_{m+1})$$

with initial conditions:

$$\underline{\beta}(t=t_m) \stackrel{\Delta}{=} \underline{\beta}(\ell=0) = 0$$

$$\underline{\delta\beta}(t=t_m) \stackrel{\Delta}{=} \underline{\delta\beta}(\ell=0) = 0$$

where

- de = Gyro output pulse vector. Each component (x,y,z) represents the occurance of a rotation through a specified fixed angle increment about the gyro input axis.
- Δθ = Gyro output pulse vactor count from 1 to 1+1.

The computational algorithm described by equation (26) is used on a recursive basic to calculata $\frac{1}{2}$ once each m cycle. After $\frac{1}{2}$ is calculated, the $\frac{1}{2}$ and $\frac{1}{2}$ functions are reset for the next m cycle $\frac{1}{2}$ calculation. The iteration rate for $\frac{1}{2}$ within m is maintained at a high enough rate to properly account for anticipated dynamic $\frac{1}{2}$ motion effects. Section 6. describes analytical techniques that can be used to access the adequacy of the $\frac{1}{2}$ iteration rate under dynamic angular rate conditions.

3.4 The Computation Of θ

The θ vactor in equations (6) and (15) is computed as a simple integral of navigation frame angular rate over the n cycle iteration period:

$$\underline{\theta} = \int_{t_n}^{t_{n+1}} \underline{\Omega} dt$$
 (27)

where

Navigation frame rotation rate as calculated in the navigation software section (12).

Standard recursive integration algorithms can be used to calculate θ in equation (27) (e.g., trapezoidal) over the time intervel from n to n+1. The update rate for the integration algorithm is celected to be compatible with software accuracy requirements in the anticipated dynamic maneuver environment for the user vehicle.

3.5 Orthogonality And Normalization Algorithms

Most strapdown attitude computation techniquas periodically employ self-consistancy correction algorithms as an outer-loop function for accuracy enhancement. If the basic attituda data is computed in the form of a direction cosine matrix, the self-consistancy check is that the rows should be orthogonal to each other and equal to unity in magnitude. This condition is besed on tha fact that the rows of the direction cosine matrix represent unit vectors along orthogonal navigation coordinate frame axea as projected in body axes. For the quaternion, the self-consistancy check is that the sum of the squares of the quaternion elemants be unity (this can be verified by operation on equation (1)).

3.5.1 Direction Cosine Orthogonalization And Normalization

The test for orthogonality between two direction cosine rows is that the dot product be zero. The error condition, than is:

$$E_{ij} = C_i C_j^T \tag{28}$$

where

Ci = ith row of C

C_i = jth row of C

T = Transpose

A calculated orhogonality error E_{ij} can be corrected by rotating C_i and C_j relative to each other about an axie perpendicular to both by the error angle E_{ij} . Since it is not known whether C_i or C_j is in error, it is assumed that each are equally likely to be generating the error, and each is rotated by half of E_{ij} to correct the error. Hence, the orthogonality correction algorithm is:

$$C_{i}(n+1) = C_{i}(n) - 1/2 E_{ij} C_{j}(n)$$

$$C_{j}(n+1) = C_{j}(n) - 1/2 E_{ij} C_{i}(n)$$
(29)

It is easily verified using (29) that an orthogonality error E_{ij} originally present in $C_i(n)$ and $C_j(n)$ is no longer present in $C_i(n+1)$ and $C_j(n+1)$ after application of equation (29).

The unity condition on C_i (i.e., normality) can be tested by comparing the magnitude squared of C_i with unity:

$$\mathbf{E_{ii}} = 1 - \mathbf{C_i} \ \mathbf{C_i}^{\mathrm{T}} \tag{30}$$

A measured normality error Eii can be corrected with:

$$C_{i}(n+1) = C_{i}(n) - 1/2 E_{i} C_{i}(n)$$
 (31)

Equations (28) through (31) can be used to measure and correct orthogonality and normalization errors in the direction cosine matrix. In combined matrix form, the overall measurement/correction operation is sometimes written as:

$$c_{n+1} = c_{n+1/2} (\bar{x} - c_n c_n^T) c_n$$
 (32)

3.5.1.1 Rows or Columns - The previous discussion addressed the problem of orthogonalizing and nomalizing the rows of a direction cosine matrix C. In combined form, equation (32) shows that the correction is:

$$\delta C = 1/2 (I - CC^{T}) C$$
 (33)

Equation (33) can be operated upon by premultiplication with C postmultiplication by \mathbf{C}^{T} , and combining terms. The result is:

$$\delta C = 1/2 C (I - C^{T}C)$$
 (34)

The (I - C^TC) term in (34) is the error matrix based on testing orthogonality and normality of the columns of C. Thus, if the rows of C are orthonormalized (i.e., δC is nulled), the columns of C will also be implicitly orthonormalized. The inverse applies if the columns are directly orthonormalized with (34). The question that remains is, which is preferred? The answer is related to the real time computing problem associated with the calculation and correction of orthogonalization and normalization errors.

Ideally, the orthogonalization and normalization operations are performed as an outer loop function in a strapdown navigation computer so as not to impact computer throughput requirements. A computational organization that facilities such an approach divides the orthonormalization operations into submodules that are executed on successive passes in the outer-loop software path. A logical division of the orthonormalization operations into submodules is as defined by equations (28), (29), (30), and (31).

This implies that measurement and correction of orthogonalization and normslization effects are performed at different times in the computing cycle. Such an approach is only valid if the orthogonality and normalizations errors (i.e., E_{ij} and E_{ii}) remain reasonably stable as a function of time.

To assess the time stability of the orthogonality/normalization error is to investigate

the rate of changa of the breckatad terms in equations (33) and (34). For convenienca, these will be defined es:

$$E_{R} = (I - CC^{T})$$

$$E_{C} = (I - C^{T}C)$$
(35)

The time derivative of (35) is:

$$\dot{\mathbf{E}}_{R} = -\dot{\mathbf{C}}\mathbf{C}^{T} - \mathbf{C}\dot{\mathbf{C}}^{T}$$

$$\dot{\mathbf{E}}_{R} = -\dot{\mathbf{C}}^{T}\mathbf{C} - \mathbf{C}^{T}\dot{\mathbf{C}}$$
(36)

Expressions for \hat{C} end \hat{C}^T cen be davaloped by returning to equations (2), (3), (5), and (6). These equations can be rearranged to show that over a given time interval, the change in C is given by:

$$\Delta C = C(A - I) + (B - I)C$$

which with (3) end (4) becomes to first order:

$$\Delta C = C(\underline{\phi}x) - (\underline{\theta}x)C \tag{37}$$

Dividing by tha time interval for the change in C, recognizing that ϕ end θ ere approximately integrals of ω end Ω over the time interval, and letting the time interval go to zero in the limit, yields the classical equation for the rate of change of C:

$$\dot{C} = C(\underline{\omega}x) - (\underline{\Omega}x)C \tag{38}$$

where

 (ωx) , (Ωx) = Skew symmetric matrix form of vectors ω , Ω .

The trenspose of (38) is :

$$\dot{c}T = -(\underline{\omega}x) C^{T} + C^{T}(Qx)$$
 (39)

Wa now substituta (38) end (39) into (36). After combining terms end epplying equations (35), the finel result is:

$$\dot{\mathbf{E}}_{\mathbf{R}} = \mathbf{E}_{\mathbf{R}} \left(\underline{\mathbf{Q}} \mathbf{x} \right) - \left(\underline{\mathbf{Q}} \mathbf{x} \right) \mathbf{E}_{\mathbf{R}}$$

$$\dot{\mathbf{E}}_{\mathbf{C}} = \mathbf{E}_{\mathbf{C}} \left(\underline{\mathbf{w}} \mathbf{x} \right) - \left(\underline{\mathbf{w}} \mathbf{x} \right) \mathbf{E}_{\mathbf{C}}$$

$$(40)$$

Equations (40) show that the reta of changs of E_R is proportional to E_R and tha navigation frame rotation rate \underline{o} , whereas the rate of change of E_C is proportional to E_C and the body rotation rate \underline{o} . Since \underline{o} is generally much larger than \underline{o} , E_C is generally larger than \underline{E}_R . It can be concluded that E_R is more stable over time, hence, orthonormalizing the direction cosine matrix rows (based on the E_R measurement) is the preferred computational approach if the real time computing problem is taken into account.

3.5.2 Queternion Normalization

The quetarnion is normalized by measuring its magnitude squared compered to unity, end edjusting each element proportionally to correct the normalization error. The normalization error is given by:

$$\mathbf{E}_{\mathbf{q}} = \mathbf{q} \cdot \mathbf{q}^* - \mathbf{1} \tag{41}$$

It is aesily verified using the rulas for quatarnion algabric that $E_{\bf q}$ equals the sum of the squares of the elements of q minus 1. The correction algorithm is given by:

$$q_{(n+1)} = q_{(n)} - 1/2 E_q q_{(n)}$$
 (42)

3.6 Direction Cosine Versus The Quaternion For Body Attitude Referencing

The tradeoff between direction cosine versus quaternion parameters as the primary attitude reference data in strapdown inertial systems has been a popular area of debate between strapdown analysts over the past three decades. In its original form, the tradeoff centered on the relative accuracy between the two methods in accounting for body angular motion. These tradeoffs invariably svolved from the differential equation form of the direction cosine and quaternion updating equations and investigated the securacy of equivalent algorithms for integrating these equations in a digital computer under hypothscized body angular motion. Invariably, the body motion investigated was coning motion at various frequencies relative to the computer update frequency. For these early studies, the tradeoffs generally demonstrated that for comparable integration algorithms, the quaternion approach generated solutions that more accurately replicated the true coning motion for situations where the coning frequency was within a decade of the computer update frequency.

As presented in this paper, both the quaternion and direction cosine updating algorithms have been based on processing of a body engle motion vector ϕ which accounts for sll dynamic motion effects including coning. These updating algorithms (equation (2) and (3) for direction cosines and (13) and (14) for the quaternion) represent exact solutions for the attitude updating process for s given input angle vector ϕ . Consequently, the question of accuracy for different body motion can no longer be considered a viable tradeoff area. The principle tradeoffs that remain between the two approaches ere the computer memory and throughput requirements associated with each in a strapdown navigation system.

In order to assess the relative computer memory and throughput requirements for quaternion parameters versus direction cosines, the composite of all computer requirements for each must be assessed. In general, these can be grouped into three major computional areas:

- Basic updating algorithm
- Normalization and orthogonalization slgorithms
- Algorithms for conversion to the direction cosine matrix form needed for acceleration transformation and Eulsr angle extraction

Basic Updating Algorithms - The basic updating elgorithm for the quaternion parameters is somewhat simpler than for direction cosines as expansion of equations (2) and (3) compared with (13) and (14) would reveal. This results in both a throughput and memory advantage for the quaternion approach. Part of this advantage arises because only four quaternion elements have to be updated compared to nine for direction cosines. The advantage is somewhat diminished if it is recognized that only two rows of direction cosines (i.e., 6 elements) need actually be updated since the third row can then be easily derived from the other two by a cross-product operation (i.e., the third row represents a unit vector along the z-axis of the navigation frame as projected in body axes. The first two rows represent unit vectors along x and y navigation frame axes. The cross-product of unit vectors along x end y navigation axes equels the unit vector along ths z-navigation sxis).

Normalization And Orthogonalization Algorithms - The normalization and orthogonalization operations associated with direction cosines are given by equation (28) through (31). The quaternion normalization equation is given by equations (41) and (42).

The normalization equation for the quaternion is generally simpler to implement than the orthogonalization and normalization equations for the direction cosines. If only two rows of the direction cosine matrix are updated (as described in the previous paragraph) the direction cosine orthogonalization and normalization operations required are half that dictated by (2B) through (31), but are still more than required by (41) and (42) for the qusternion. Since the orthonormalization operations would in general be iterated at low rete, no throughput edvantage results for the quaternion. Some memory savings may be realized, however.

A key factor that must be addressed relative to orthonormalization tradeoffs is whether or not orthonormalization is actually needed at all. Clearly, if the direction cosine or quaternion updating algorithms were implemented perfectly, orthonormalization would not be required. It is the author's contention that, in fact, the accuracy requirements for strapdown systems dictate that strapdown attitude updating software cennot tolerate any errors whatsoever (compared to sensor error effects). Therefore, if the attitude updating software is designed for negligible drift and scale factor error (compared to sensor errors) it will also implicitly exhibit negligible orthogonalization and/or normalization errors.

The above argument is valid if the effect of orthonormalization errors in strapdown attitude data is no more detrimental to system performance than other software attitude error effects. This is in fact the case, as detailed error analyses would reveal. Since modern-day general purpose computers used in today's strapdown inertial navigation systems have the capability to implement attitude updating algorithms essentially perfectly within a reasonable throughput and memory requirement, it is the author's opinion that orthonormalization error correction should not be needed, hence, is not a viable tradeoff area relative to the use of quaternion parameters versus direction cosines.

attituds data is direction cosinea directly, no conversion processe is required. For cases where only two rows of direction cosines are updated, the third row must be generated by the cross-product between the two rows calculated. For example:

$$\begin{array}{rcl} C_{31} & = & C_{12} C_{23} - C_{13} C_{22} \\ C_{32} & = & C_{13} C_{21} - C_{11} C_{23} \\ C_{33} & = & C_{11} C_{22} - C_{12} C_{21} \end{array} \tag{43}$$

For quaternion parameters, equation (17) must be implemented to develop the direction cosine matrix, a significantly more complex operation compared with (43) for the two row direction cosins approach. Since direction cosins elementa are generally required at high rats (for acceleration transformation and Euler angls output extraction) both a throughput and memory penalty is accrued for the quaternion approach. The penalty is compounded if the calculated direction cosine outputs are required to greeter than single precision accuracy (including computational round-off error). For noise-free acceleration transformation operations (such as may be needed to effect an accurate system calibration) double-precision accuracy is needed. The result is that equation (17) for the quaternion versus (43) for direction cosines would have to be implemented in double-precision imposing a significant penalty for the more complex quaternion conversion process.

Tradeoff Conclusions - From the above qualitative discussion, it is difficult to draw hard conclusions regarding a preference for direction cosines vereus quaternion parametere for attitude referencing in strapdown inertial systems. Pros and cons sxist for each in the different tradeoff areas. Quantitative comparisons based on actual software sizing and computer loading studies have led to eimilar inconclusive results. Fortunately, today's computer technology is such that the slight advantage one attitude parameter approach may have over the other in any particular application is insignificant compared with composite total atrapdown insertial system throughput and memory software requirements. Hence, ultimate selection of the attitude approach can be safely made based on "analyst's choics".

4. STRAPDOWN ACCELERATION TRANSFORMATION ALGORITHMS

The acceleration vector measurement from the accelerometers in a etrapdown insrtial system is transformed from body to navigation axee through a mechanization of the classical vector transformation equation:

$$\underline{\mathbf{a}}^{\mathbf{N}} = \mathbf{C}\,\underline{\mathbf{a}} \tag{44}$$

whsrs

- Specific force acceleration measured in body axes by the strapdown accelerometers
- e^{N} = Specific forcs acceleration with componente svaluated along navigation axes.

The implementation of equation (44) is accomplished on a repetative basis as a recursive algorithm in a digital computer such that its integral properties are preserved at the computer cycle times. In this manner, the velocity which is formed from the integral of (44) will be accurate under dynamic conditions in which a may have stratic high frequency components. The recursive algorithm for (44) must account for the effects of body rotation (and secondarily, rotation of the navigation coordinate frame) as well as variations in a over the computer iteration period.

4.1 Acceleration Transformation Algorithm That Accounts For Body Rotation Effects

To develop an algorithm for squation (44) that preserves ite integral properties, we begin with its integral over a computer cycle:

$$\underline{\mathbf{u}}^{\mathbf{N}} = \int_{\mathbf{t}_{\mathbf{m}}}^{\mathbf{t}_{\mathbf{m}}+1} \mathbf{C} \underline{\mathbf{a}} d\mathbf{t} \qquad (45)$$

where

 \underline{u}^{N} = Change in the integral of equation (44) (or specific force valocity change) over a computer cycle m

The velocity vector in the newigation computer is generated by summing the \underline{u}^{N} 's corrected for Coriolia and gravity effects.

The C matrix in (45) is a continuous function of tims in the interval from t_m to t_{m+1} . An equivelent form for C in terms of its value at the computer update time. (m) is:

$$C = C(m) A(t)$$

whare

C(m) = Value of C at tm

A(t) = Direction cosine matrix that transform vectors from body axes at time t to the body attitude at the start time for the computation interval t_m .

Equation (46) with the dafinition for A(t) above accounte for the affect of gyro seneed body motion over the computer interval. The next section will discuss the correction used to account for the small rotation of the navigation frame over the computer interval.

Subatituting (46) in (45) and expanding:

$$\underline{u}^{N} = C(m) \int_{t_{m}}^{t_{m}+1} A(t) \underline{a} dt$$

We now use a first order approximation for A(t) as given by equation (3), with ϕ translated as a function of time in the interval as defined to first order in equation (22):

$$\underline{\bullet}(t) \sim \underline{\beta}(t) = \int_{t_m}^t \underline{\omega} dt$$

Thus,

$$h(t) = I + (\underline{\theta}(t)x) \tag{47}$$

and

$$\underline{u}^{N} \sim C(m) \int_{t_{m}}^{t_{m+1}} (I + (\underline{\varrho}(t)x)) \underline{a} dt$$

$$= C(m) \left(\int_{t_{m}}^{t_{m+1}} \underline{a} dt + \int_{t_{m}}^{t_{m+1}} (\underline{\varrho}(t) \times \underline{a}) dt \right)$$

We now dafine

$$\underline{\mathbf{u}} = \int_{\mathbf{t_m}}^{\mathbf{t_{m+1}}} \underline{\mathbf{a}} \, d\mathbf{t}$$

Hence,

$$\underline{\underline{u}}^{N} = C(m) \left(\underline{\underline{u}} + \int_{t_m}^{t_{m+1}} (\underline{\underline{\beta}}(t) \times \underline{\underline{a}}) dt\right)$$
 (48)

with

$$\underline{\beta(t)} = \int_{t_m}^{t} \underline{\omega} dt$$

$$\underline{u} = \int_{t_m}^{t_{m+1}} \underline{a} dt$$

An alternative form of (48) can also be derived through direct application of the integration by parts rule to the integral term in the equation (48) \underline{u}^N expression.:

$$\underline{\underline{u}}^{N} = C(\underline{m}) \left(\underline{\underline{u}} + 1/2 \underline{\underline{\beta}} \times \underline{\underline{u}} + 1/2\right)^{t}_{t_{\underline{m}}} \left(\underline{\underline{\beta}}(t) \times \underline{\underline{a}} + \underline{\underline{u}}(t) \times \underline{\underline{\omega}}\right) dt$$
(49)

with

$$\underline{\beta}(t) = \int_{t_m}^{t} \underline{\omega} dt$$

$$\underline{u}(t) = \int_{t_m}^{t} \underline{a} dt$$

$$\underline{\beta} = \underline{\beta}(t=t_{m+1})$$

$$\underline{u} = \underline{u}(t=t_{m+1})$$

Equations (48) and (49) are algorithmic forms of equation (44) that can be used to calculate \underline{u}^N in the strapdown computer exactly (within the approximation of equation (47)). These equations show that the specific force valocity change in navigation coordinates is approximately equal to the integrated output from the strapdown acceleromater (\underline{u}) over the computer cycle, timas the direction cosine matrix which was "lid at the pravious computar update time. Correction tarms are applied to account for body rotation. In general, the corraction tarm involves an integral of the interractive effects of angular ω and linear a motion over the update cycle. The integral terms have been coined "eculling" affects.

The equation (49) form of the \underline{u}^N equation includes a 1/2 \underline{g} \underline{x} \underline{u} term which can be evaluated at t_{m+1} as the aimpla cross-product of intagratad gyro and accalarometar measurements (i.a., without a dynamic intagral operation). Furthermora, it is easily damonetrated that for approximately constant angular rates and accalerations over the computer cycle, the integral term in (49) is identically zero. This forms the basis for an approximate form of (49) which is valid under benign flight conditions (i.a., using equation (49) without including the integral term). The 1/2 g x u term in (49) is sometimes denoted as "rotation compensation".

4.1.1 Incremental Form of Transformation Operationa and Sculling Tarms

In a severe dynamic anvironment, equations (48) or (49) would be implamented axplicitly with the integral terms mechanized as a high spead digital algorithmic operation within the t_m to t_{m+1} update cycle. The integral terms we are dealing with are from (48) and (49):

$$\underline{S}_{1} \stackrel{\triangle}{=} \int_{t_{m}}^{t_{m+1}} (\underline{s}(t) \times \underline{a}) dt$$

$$\underline{S}_{2} \stackrel{\triangle}{=} 1/2 \int_{t_{m}}^{t_{m+1}} (\underline{s}(t) \times \underline{a} + \underline{u}(t) \times \underline{\omega}) dt$$
(50)

With the aquation (50) definitions, (48) and (49) become:

$$\underline{u}^{N} = C(m) (\underline{u} + \underline{s}_{1})$$
 (51)

or $\underline{\mathbf{u}}^{\mathbf{N}} = \mathbf{C}(\mathbf{m}) \left(\underline{\mathbf{u}} + 1/2 \mathbf{g} \times \underline{\mathbf{u}} + \underline{\mathbf{S}}_{2}\right)$ (52)

Racureiva algorithms for \underline{s}_1 or \underline{s}_2 can be darivad by first rawriting (50) in that equivalent form:

$$\underline{\beta}(t) = \underline{\beta}(t) + \int_{t_{\underline{\beta}}}^{t} \underline{\omega} dt$$

$$\underline{u}(t) - \underline{u}(t) + \int_{t_{\underline{\beta}}}^{t} \underline{a} dt$$

$$\underline{x}_{1}(t+1) = \underline{x}_{1}(t) + \int_{t_{\underline{\beta}}}^{t} (\underline{\beta}(t) \times \underline{a}) dt$$

$$\underline{x}_{2}(t+1) = \underline{x}_{2}(t) + \frac{1}{2} \int_{t_{\underline{\beta}}}^{t} (\underline{\beta}(t) \times \underline{a} + \underline{u}(t) \times \underline{\omega}) dt$$

$$\underline{\beta}(t+1) = \underline{\beta}(t=t_{\underline{\beta}+1})$$

$$\underline{u}(t+1) = \underline{u}(t=t_{\underline{\beta}+1})$$

$$\underline{u}(t+1) = \underline{u}(t=t_{\underline{\beta}+1})$$
(53)

$$\underline{\mathbf{s}}_1 = \underline{\mathbf{y}}_1 (\mathbf{t} = \mathbf{t}_{m+1})$$

$$\underline{s}_2 = \underline{\gamma}_2 (t=t_{m+1})$$

with initial conditiona

$$\underline{\beta}(t=t_m) = 0$$

$$\underline{u}(t=t_m) = 0$$

$$\underline{\chi}_1(t=t_m) = 0$$

$$\underline{\chi}_2(t=t_m) = 0$$

(54)

where

= High spaed computer cycle within m lowar speed computation cycle.

The intagrals in (53) can be replaced by analytical forms that are compatible with gyro and accelerometer input data processing if $\underline{\omega}$ and \underline{a} are replaced by a generalized tima saries expansion. For equations (53), it is sufficient to approximate $\underline{\omega}$ and \underline{a} over the 1 to 11 time interval as constants. Using this approximation in (53) yialds the final algorithm forms. For \underline{S}_1 , the companion to equation (51), the algorithm is:

$$\underline{\gamma}_{1}(k+1) = \underline{\gamma}_{1}(k) + (\beta(k) + 1/2 \underline{\Delta\theta}(k)) \times \underline{\Delta\nu}(k)$$

$$\underline{\beta}(k+1) = \underline{\beta}(k) + \underline{\Delta\theta}(k)$$

where

$$\underline{\Delta\theta}(k) = \int_{t_k}^{t_{k+1}} \underline{\omega} \, dt = \sum_{t_k}^{t_{k+1}} \underline{d\theta}$$

$$\underline{\Delta v}(k) = \int_{t_k}^{t_{k+1}} \underline{a} \, dt = \sum_{t_k}^{t_{k+1}} \underline{dv}$$

and

$$\underline{\mathbf{S}}_{1} = \underline{\mathbf{Y}}_{1}(\mathsf{t} = \mathsf{t}_{m+1}) \tag{55}$$

For equation (51):

$$\underline{u}(k+1) = \underline{u}(k) + \Delta v(k)$$

$$\underline{\underline{u}} \stackrel{\Delta}{=} \underline{\underline{u}}(t=t_{m+1})$$

with initial conditions:

$$\underline{\beta}(t=t_m) \stackrel{\Delta}{=} \underline{\beta}(\ell=0) = 0$$

$$\underline{\gamma}_1(t=t_m) \stackrel{\Delta}{=} \underline{\gamma}_1(t=0) = 0$$

where

de, dv, = Gyro and accelerometer output pulse vectors. Each component (x, y, z) reprasents the occurance of a rotation through a specified angle about the gyro input axis (for de components) or an acceleration through a specific force velocity change along the acceleromater input axis (for dv components).

 $\Delta\theta$, Δv , = Gyro and accelaromater pulse vector counts from ℓ to $\ell+1$.

For the alternative S_2 form, the companion to equation (52), the algorithm is:

$$\underline{\gamma_2}(1+1) = \underline{\gamma_2}(1) + 1/2 \left(\underline{\beta}(1) \times \underline{\Delta \nu}(1) + \underline{u}(1) \times \underline{\Delta \theta}(1)\right) \\
\underline{\beta}(1+1) = \underline{\beta}(1) + \underline{\Delta \theta}(1) \\
\underline{u}(1+1) = \underline{u}(1) + \underline{\Delta \nu}(1)$$

where

$$\underline{\Delta\theta}(\lambda) = \int_{t_{1}}^{t_{1}+1} \underline{\omega} \ dt = \int_{t_{1}}^{t_{1}+1} \underline{d\theta}$$

$$\underline{\Delta u}(\lambda) = \int_{t_{\lambda}}^{t_{\lambda+1}} \underline{a} dt = \sum_{t_{\lambda}}^{t_{\lambda+1}} \underline{dv}$$

and

(56)

$$\underline{s}_2 = \underline{\gamma}_2(t=t_{m+1})$$

For equations (52):

$$\underline{\beta} = \underline{\beta}(t=t_{m+1})$$

$$\underline{\mathbf{u}} = \underline{\mathbf{u}}(\mathsf{t=t}_{m+1})$$

with initial conditions:

$$\underline{\beta}(t=t_m) \stackrel{\Delta}{=} \underline{\beta}(1=0) = 0$$

$$\underline{u}(t=t_m) \stackrel{\Delta}{=} \underline{u}(1=0) = 0$$

$$\underline{\Upsilon}_2(t=t_m) \stackrel{\Delta}{=} \underline{\Upsilon}_2(1=0) = 0$$

Equations (51) with (55), or (52) with (56) are computational algorithms that can be used to calculate the navigation frame specific force velocity changes. Two iteration rates are implied: a basic m cycle rate, and a higher speed I cycle rate within each m cycle.

The m cycle rata is selected to be high enough to protect the approximation of neglecting the $(\underline{\beta}(t)x)^2$ term in A(t) (contrast equation (47) with the equation (3) exact form for A). This design condition is typically avaluated under maximum expected linear acceleration/angular rats envalope conditions for the particular application. Typically, the m cycle rete required for accuracy in the attitude updating algorithms is also sufficient for accuracy requirements in the m cycle of the acceleration transformation algorithms.

That cycle rata within m is set high enough to proparly account for anticipated composite dynamic $\underline{\omega}$, \underline{a} effects. Section 6. dascribas analytical tachniques that can be used to assess that edequacy of the \underline{S} iteration rata for the sculling computation under dynamic input conditions.

4.1.3 Acceleration Transformation Algorithms Basad on Quaternion Attitude Dete

Equations (51) or (52) were based on the use of direction cosine data (C) in the strapdown computer. If the basic attitude data is calculated in the form of e quaternion, the equivalent C matrix for transformation can be calculated using equations (17). Alternatively, the quaternion data can be epplied directly in the implementation of the transformation operation through application of equations (12) to equations (51) and (52):

$$u^{N} = q(m) (u + S_{1}) q(m)^{*}$$
 (57)

or

$$u^{N} = q(m) (u + s_{2}^{'}) q(m)^{*}$$
 (58)

$$\underline{\mathbf{s}}_{2}^{\star} \stackrel{\Delta}{=} 1/2 \, \underline{\beta} \times \underline{\mathbf{u}} + \underline{\mathbf{s}}_{2}$$

where u and the terms in the middle breckets are the queternion form of the vector of the same nonmencleture defined as heving the first three terms (i.e., vector components) equal to the vector elements, and the fourth sceler term equal to zero. The $\underline{S_1}$ and $\underline{S_2}$ terms ere celculeted es defined by equations (55) end (56).

4.2 Accelaration Transformation Algorithm Correction For Navigation Frame Rotations

The acceleration transformation elgorithms represented by equation (51), (52) or (57), (58) with (55), (56) neglects the effect of navigation frame rotation. In general, this is a minor correction term that can be easily accounted for at the nocycle update rate (i.e., the computer cycle rate used to update the attitude date for the effect of navigation frame rotations). It can be shown through a development similar to that leading to equation (52), that the correction algorithm for local navigation frame motion is given to first order by:

$$\underline{\Delta u}^{N}(n) = -1/2 \underline{\theta} \times \underline{v}(n) \tag{59}$$

where

 $\underline{\Delta u}^{N}(n)$ = Correction to the velue of \underline{u}^{N} computed in the m cycle that occurs at the current n cycle time. (Note: the m cycle is within the lower speed n cycle time freme).

y(n) = Summation of y(m) over the n cycle updete period.

 $\frac{\theta}{\theta}$ = Integral of the nevigation frame angular rotation reta over the n cycle psriod (as described in Sections 3.1.2 end 3.4)

5. EULER ANGLE EXTRACTION ALGORITHMS

If the body ettitude relative to navigation axas is defined in terms of three successive Euler angla rotetions ϕ , θ , ϕ about exes z, y, x reepectively (from ravigation to body exes), it can be readily demonstrated (9) that the relationship between the direction cosine elements end Euler engles is givan by:

$$C_{12} = -\cos\phi \sin\phi + \sin\phi \sin\theta \cos\phi$$

$$C_{22} = \cos\phi \cos\phi + \sin\phi \sin\phi \sin\phi$$
 (60)

$$C_{31} = - \sin\theta$$

$$C_{32} = sin\phi cos\theta$$

For conditions where $\theta \neq \pi/2$ the inversa of equations (60) can be used to evaluate the Eular engles from the direction cosinas:

$$\phi = \tan^{-1} \frac{C_{32}}{C_{33}}$$

$$\theta = - \tan^{-1} \frac{c_{31}}{\sqrt{(1-c_{31}^2)}}$$
 (61)

$$\phi = \tan^{-1} \frac{c_{21}}{c_{11}}$$

For situations where $/\theta$ / approaches $\pi/2$, the ϕ and ϕ equations in (61) become indeterminata because the numerator and denominator approach zero simultaneously (see

(63)

equations (60)). Under these conditions, an elternative equation for ϕ , ψ can be dsveloped by first applying trigonometric algebra to equations (61) to obtsin:

$$C_{23} + C_{12} = (\sin\theta - 1) \sin(\phi + \phi)$$

$$C_{13} - C_{22} = (\sin\theta - 1) \cos(\phi + \phi)$$

$$C_{23} - C_{12} = (\sin\theta + 1) \sin(\phi - \phi)$$

$$C_{13} + C_{22} = (\sin\theta + 1) \cos(\phi - \phi)$$
(62)

Taking appropriate reciprocals of sine, cosine terms in (62) and applying the inverse tangent function:

For
$$\theta$$
 near + $\pi/2$

$$\psi - \phi = \tan^{-1} \frac{c_{23} - c_{12}}{c_{13} + c_{22}}$$

For θ near - $\pi/2$

$$\psi + \phi = \tan^{-1} \frac{C_{23} + C_{12}}{C_{13} - C_{22}}$$

Equations (63) can be used to obtain expressions for the sum or difference of ψ and ϕ under conditions where $/\theta/$ is near $\pi/2$. Explicit separate solutions for ψ and ϕ cannot be found under the $/\theta/=\pi/2$ condition because ϕ and ϕ both become sngle measures about parsllel axes (about vertical), hence, measure the same angle (i.e., a degree of rotational freedom is lost, and only two Euler angles, $\theta=\pm\pi/2$ and ϕ or ϕ define the body to navigation frame attitude). Under $/\theta/$ near $\pi/2$ conditions, θ or ψ can be srbitrarily selected to satisfy another condition, with the unspecified variable calculated from (63). As an example, ϕ might be set to a constant at the value it had from equations (61) when the $/\theta/$ near $\pi/2$ region was entered. This selection avoids jumps in ψ as the solution equation is transitioned from the (61) to the (63) form.

6. ALGORITHM PERFORMANCE ASSESSMENT

The division of the ettitude updating end acceleration transformation algorithms into high and low speed loops for body motion effects (1 and m ratee) provides for flexibility in selection of the iteration rates to maintain overall algorithm accuracy at system specified performance levels. The 1 and m rate algorithms have been designed such that the high rate 1 loop consists of simple computations that can be iterated at the high rate needed to properly account for high frequency vibration effects. The m rate loop algorithms, on the other are more complicated, based on computationally exact solutions.

Iteration rates for the m loop are selected to maintain accuracy under maximum maneuver induced motion conditions. The m loop iteration rate to maintain accuracy under maximum maneuver conditions can be easily evaluated enelytically, or by simulation, through comparision of the actual algorithm solution with the Taylor series truncated forms selected for system mechanization. Iteration rates for the 1 loop are selected to maintain accuracy under anticipated vibratory environmental conditions.

6.1 Vibration Environment Assessment

A fundemental calculation that should be performed prior to the analysis of 1 loop algorithm iteration rate requirements is an assessment of the dynamic inputs that must be measured by the algorithms. In essence, this consists of an evaluation of the continuous (i.e., infinitely fast iteration rate) form of the elgorithms in question under dynamic input conditions. The specific continuous form equations of interest are equations (22) for $\underline{\delta}\beta$ and (50) for \underline{S}_1 or \underline{S}_2 .

6.1.1 δβ Dynamic Environment Assessment (Coning)

We repeat equations (22) for $\delta \beta$ evaluated at t = t_{m+1}:

$$\underline{\beta}(t) = \int_{t_m}^t \underline{\omega} \, dt$$

 $\underline{\delta\beta}(t=t_{m+1}) = 1/2 \int_{t_m}^{t_{m+1}} \underline{\beta}(t) \times \underline{\omega} dt$

(64)

end enalysa the solution for $\frac{\delta\beta}{\theta}(t=t_{m+1})$ under gamerel cyclic motion et frequancy f in axes x end y with enguler amplitudes $\frac{\delta\beta}{\theta}$, $\frac{\delta\gamma}{\theta}$ end relative phase angle ϕ such thet:

$$\int_{0}^{t} \underline{\omega} dt = (\theta_{x} \sin(2\pi f t), \theta_{y} \sin(2\pi f t + \phi), 0)^{T}$$

$$\underline{\omega} = 2\pi f (\theta_{x} \cos(2\pi f t), \theta_{y} \cos(2\pi f t + \phi), 0)^{T}$$
(65)

Substituting (65) in (64), expanding through epplication of epproprieta trigonometric identities, and carrying out the indicated integrals enelytically between the assigned limits, yields zero for the x, y components end the following for the z component of $\delta \theta$ (t=t_{m+1}):

$$\delta \beta_{z}(t=t_{m+1}) = \pi \theta_{x} \theta_{y} (\sin \phi) f \left((t_{m+1} - t_{m}) - \frac{\sin 2\pi f(t_{m+1} - t_{m})}{2\pi f} \right)$$

Defining the m cycle time interval es $T_{\mathbf{m}}$, the letter exression is equivalently:

$$\delta \beta_z = \pi \theta_x \theta_y \left(\sin \phi \right) f \gamma_m \left(1 - \frac{\sin 2\pi f T_m}{2\pi f T_m} \right)$$
 (66)

Hence, even though the ω rate is cyclic in two axas as defined by equation (65) in x end y, the value for $\delta\beta_z$ is a constant proportional to the sine of the phase engle between the x, y angular vibrations. Under conditions where $\phi=0$ (defined es "rocking" motion), $\delta\beta_z$ is zero. Under conditions where $\phi=\pi/2$, $\delta\beta_z$ is maximum. The equation (65) rata when $\phi=\pi/2$ has been tarmed "coning motion" due to the characteristic responsa of the z exis under this motion which describes e cone in inertial space.

Equation (66) cen be put into a "drift rete" form by dividing the $^{\delta\beta}_{Z}$ angla by the tima interval T_m ovar which it was evaluated:

$$\delta \dot{\beta}_{z} = \pi \theta_{x} \theta_{y} (\sin \phi) f \left(1 - \frac{\sin 2\pi f T_{m}}{2\pi f T_{m}}\right)$$
 (67)

Equation (67) is e fundamental equation that can be used to essess the magnitude of $\delta\beta_Z$ that must be accounted for by the $\delta\beta$ computer algorithm under discrete fraquency input conditions. If $\delta\beta_Z$ is small relative to system performance requirements, it can be neglected, and the 1 loop elgorithm for $\delta\beta$ need not be implemented.

Equation (67) dascribas how $\delta \hat{\beta}_Z$ cen be celculated for a discrete input vibration frequency f. In a more general case, the input rate is composed of e mixture of frequencies in x end y at different phese angles ϕ for each. If the source of the generalized angular vibration is random input noise to the strapdown system, the x, y motion is colored by the transmission cherecteristics of the noise input to the x, y angular rasponse. A more general development of equation (67) that accounts for the latter effects shows that the comparable equation for $\delta \hat{\beta}_Z$ is given by:

$$\delta \dot{\beta}_{\mathbf{z}} = \int_{0}^{\infty} \omega \, A_{\mathbf{x}}(\omega) \, A_{\mathbf{y}}(\omega) \, \sin \left(\phi_{\mathbf{A}\mathbf{y}}(\omega) - \phi_{\mathbf{A}\mathbf{x}}(\omega) \right) \, \left(1 - \frac{\sin \omega T_{\mathbf{m}}}{\omega T_{\mathbf{m}}} \right) \, P_{\mathbf{n}\mathbf{n}}(j\omega) \, d\omega \tag{68}$$

where

 $A_{x}(\omega)$, $A_{y}(\omega)$ = Amplitude of trensfar function relating system input vibration noisa to angular attituda rasponsa of sansor assambly about x, y axas.

 $\phi_{Ax}(\omega), \phi_{Ay}(\omega)$ = Phasa of transfer function relating system input vibration noisa to angular attitude responsa of sansor assembly about x, y axas.

 $P_{nn}(j_w) = Power epectral density of input vibration noise.$

ω = Fourier frequency (rad/eec)

Note: Mean squared vibration energy = $\int_0^{\infty} P_{nn}(j_{\omega}) d_{\omega}$

Equation (68) can be used to assess the extent of random spectrum dynamic angular environment to be measured by the $\delta\beta$ computational algorithm. The $\delta\beta_Z$ value calculated by (68) measuree the composite correlated coning drift in the eensor assembly that must be calculated to accurately account for the actual motion present. If the $\delta\beta_Z$ magnitude calculated from (68) is small compared to other eystems error budget effects, the mechanization of an algorithm to calculate $\delta\beta$ is not needed (i.e., can be approximated by zero).

The extension of equations (67) and (68) to y, z or z, x axis angular vibration motion should be obvious.

6.1.2 \underline{S}_1 , \underline{S}_2 Dynamic Environment Assessment (Sculling)

We repeat equations (50) with \underline{u} and \underline{s} from (48) and (49):

$$\underline{\beta}(t) = \int_{t_{\overline{M}}}^{t} dt$$

$$\underline{u}(t) = \int_{t_m}^{t} \underline{a} dt$$

 $\underline{s}_{1} = \int_{t_{m}}^{t_{m+1}} (\underline{s}(t) \times \underline{a}) dt$ (69)

$$\underline{s}_2 = 1/2 \int_{t_m}^{t_{m+1}} (\underline{s}(t) \times \underline{a} + \underline{u}(t) \times \underline{\omega}) dt$$

and analyse the $\underline{S_1}$, $\underline{S_2}$ solutions under general cycle motion at frequency f in axes x, y with angular amplitude θ_X about axis x and acceleration amplitude D_y along axis y at relative phase ϕ euch that:

$$\int_{0}^{t} \underline{\omega} dt = (\theta_{X} \sin(2\pi f t), 0, 0)^{T}$$

$$\underline{\omega} = (2\pi f \theta_{X} \cos(2\pi f t), 0, 0)^{T}$$

$$\underline{a} = (0, D_{Y} \sin(2\pi f t + \phi), 0)^{T}$$
(70)

Substituting (70) in (69), expanding through application of appropriate trigonometric identities, and cerrying out the indicated integrals analytically between the assigned limits, yields zero for the x, y components and the following for the z component of \underline{S}_1 and \underline{S}_2 :

$$S_{2z} = 1/2 T_m \theta_x D_y (coe_{\phi}) \left(1 - \frac{sin_{\pi}fT_m}{2\pi fT_m}\right)$$
 (71)

$$S_{1Z} = 1/2 (g \times u)_{z} + S_{2Z}$$
 (72)

where

 $(\underline{\beta} \times \underline{u})_z = z - \text{component of } \underline{\beta} \times \underline{u} \text{ evaluated at } t = t_{m+1}$.

Hence, even though the $\underline{\omega}$ and \underline{a} inpute are cyclic in two sxes as defined in equations (70), the value for S_{2z} is a constant proportional to the cosine of the phase sngle between

the x angular vibretion and y lineer ecceleration vibration. Under conditions where $\phi=\pi/2$, S_{2z} is xaro. Under conditions where $\phi=0$, S_{2z} is a maximum. Equation (70) motion when $\phi=0$ has been termed "sculling motion" due to the analogy with the characteristic engular movement end acceleration forces imparted to an oar used to propel a boat from the stern. Note also that S_{1x} is equal to S_{2x} plus the correction term (rotation compensation) measured as the cross-product of the simple engular rate and linear acceleration integrals taken over the m computation cycle. (See equations (48) and (49) for definitions).

Equation (71) for $S_{2\chi}$ can be put into en "acceleration bies" form by dividing the velocity change correction $S_{2\chi}$ by the time intervel T_m over which it was evaluated:

$$\dot{s}_{2\pi} = 1/2 \,\theta_{\pi} \,D_{y} \,(\cos\phi) \,\left(1 - \frac{\sin \,2\pi f T_{m}}{2\pi f T_{m}}\right)$$
 (73)

Equation (73) (with (72) for S_{1Z}) is a fundemental equation that can be used to essess the magnitude of S_{2Z} that must be accounted for by the S_1 or S_2 computer algorithm under discrete frequency input conditions. If S_{2X} is small relative to system performance requirements, it can be neglected, and the \$\mathcal{L}\$ loop algorithm for calculating S_1 or S_2 need not be implemented. Under the latter conditions, S_1 would be set equal to the cross-product term in (72) which makes the basic equation (51) and (52) transformation algorithms identical.

Equation (73) describes how \hat{S}_{2x} can be calculated with a discreta input vibration frequency f for engular motion about x and linear action along y. In a more general case, the input rates and eccelerations are composed of mixtures of engular and linear motion about x and y at different frequencies and relative phase engles. If the source of the generalized vibration motion is random input noise to the strapdown system, the x, y angular and linear motion is colored by the transmission characteristics of the noise input to the x, y angular and linear response. A more general development of equation (73) that accounts for the latter effects show that the comparable equation for \hat{S}_{2x} is given by:

$$\dot{S}_{2z} = \int_{0}^{z} \left(A_{y}(\omega) B_{x}(\omega) \cos(\phi_{Ay}(\omega) - \phi_{Bx}(\omega)) - A_{x}(\omega) B_{y}(\omega) \cos(\phi_{Ax}(\omega) - \phi_{By}(\omega)) \right) \left(1 - \frac{\sin\omega T_{m}}{\omega T_{m}} \right) P_{nn}(j\omega) d\omega$$
(74)

whsrs

 $A_{\mathbf{X}}(\omega)$, $A_{\mathbf{Y}}(\omega)$, $\phi_{\mathbf{A}\mathbf{X}}(\omega)$, $\phi_{\mathbf{A}\mathbf{Y}}(\omega)$

 $B_{\mathbf{X}}(\omega)$, $B_{\mathbf{Y}}(\omega)$, = x, y, amplitude/phese linear acceleration reaponse of the sensor $\phi_{\mathbf{B}\mathbf{X}}(\omega)$, $\phi_{\mathbf{B}\mathbf{Y}}(\omega)$ assembly to the input vibration.

Equation (74) can be used to essess the extent of random spectrum dynamic motion environment to be measured by the S_1 or S_2 computational algorithms. The \dot{S}_{2z} value calculated by (74) measures the composite correlated sculling acceleration bias in the sensor essembly that must be calculated to accurately account for the actual motion present. If the \dot{S}_{2z} magnitude calculated from (74) is small compared to other eystem error budget effects, the mechanization of an algorithm to calculate S_1 or S_2 in the high rate 1 loop is not needed (i.e., S_2 can be approximated by zero in (52) or S_1 can be set equal to the cross-product term in (52)).

The extension of equetions (73) and (74) for y, z or z, x exis vibration motion should be obvious.

6.2 Algorithm Accuracy Assessment

The eccurecy of the computation elgorithm for $\underline{\delta\beta}$ or $\underline{S_1}$, $\underline{S_2}$ cen be assessed by compering their solutions to the comparable continuous form solutions developed in Section 6.1 under identical input conditions.

6.2.1 δβ Coning Algorithm Error Assessment

The computational algorithm for calculating $\delta\beta$ in a strapdown system is given by equation (26). A measure of the accuracy of the equation (26) algorithm can be obtained by enalytically calculating the solution generated from (26) under assumed cyclic motion and

comparing this result to the equivalent solution obtained from the idealizs continuous algorithm described in Section 6.1. For a discrete frequency vibration input, the equation (65) motion can be used analytically in equation (26) to calculate the algorithm solution for $\delta\beta$ at t = t_{m+1} (i.e., analogous to the equation (67) solution for the continuous (infinitely faet) algorithm. After much algebraic manipulation, it can be demonstrated that the algorithm solution for $\delta\beta$ as calculated from equation (26) under equation (65) input motion, has zero x, y components, with a z component rate given by:

$$\delta \hat{\beta}_{ZALG} = \pi \theta_{\chi} \theta_{\chi} \left(\sin \phi \right) \left((1 + 1/3 \left(1 - \cos 2\pi f T_{\chi} \right) \right) \frac{\sin 2\pi f T_{\chi}}{2\pi f T_{\chi}}$$

$$- \frac{\sin 2\pi f T_{m}}{2\pi f T_{m}} \right)$$
(75)

whers

 $\delta\beta_{ZALG}$ = Recursive algorithm solution for $\delta\beta_{Z}$ rate

Time interval for high speed & computer iteration cycle

Equation (75) for the $\underline{\delta \beta}$ discrete recursive algorithm solution of equation (26) is directly analogous to the equation (67) solution of the equation (22) continuous $\underline{\delta \beta}$ algorithm. It is easily verified that (75) reduces to (67) as $T_{\underline{\lambda}}$ approaches zero.

The error in the $\delta\beta$ algorithm is measured by the difference between (67) and (75); i.e.:

$$e(\delta \hat{\theta}_{Z}) = \pi f \theta_{X} \theta_{Y} (\sin \phi) ((1 + 1/3 (1 - \cos 2\pi f T_{\chi})) \frac{ein 2\pi f T_{\chi}}{2\pi f T_{\psi}} - 1)$$
 (76)

where

 $e(\delta\beta_Z)$ = Error rate in the equation (26) algorithm.

Equation (76) can be used to assess the error in the equation (26) $\underline{\delta g}$ algorithm caused by finite iteration rate (i.e., the effect of $T_{\underline{I}}$) under discrete frequency input conditions.

Under random vibration input conditions, the equation (26) algorithm can be analysed to obtain the more general solution for the $\delta \beta_{zALG}$ rate:

$$\delta\beta_{zALG} = \int_{O}^{\infty} \omega A_{x}(\omega) A_{y}(\omega) ein(\phi_{Ay}(\omega) - \phi_{Ax}(\omega)) ((1 + 1/3 (1 - \cos\omega T_{z}) \frac{\sin \omega T_{z}}{\omega T_{z}} - \frac{\sin \omega T_{m}}{\omega T_{m}}) P_{nn}(j\omega) d\omega$$
(77)

The $\delta\beta$ algorithm error under random inputs is the difference betwean the equetion (77) discrete solution and the equivalent continuous equation (68) solution form. The result is:

$$e(\delta \hat{\beta}_{\mathbf{Z}}) = \int_{0}^{\infty} \omega A_{\mathbf{X}}(\omega) A_{\mathbf{Y}}(\omega) \sin(\phi_{\mathbf{A}\mathbf{Y}}(\omega) - \phi_{\mathbf{A}\mathbf{X}}(\omega)) ((1 + 1/3 (1 - \cos\omega T_{\mathbf{Z}}) \frac{\sin \omega T_{\mathbf{Z}}}{\omega T_{\mathbf{Z}}} - 1) P_{\mathbf{n}\mathbf{n}}(j\omega) d\omega$$
(78)

Equations (76) and (78) can be used to assess the error in the equation (26) $\frac{\delta\beta}{2}$ algorithm caused by finite iteration rate under discrete or random vibration input conditions. The extension of equations (76) and (78) to y, z or z, x axis effects should be obvious.

6.2.2 Sculling Algorithm Error Assessment

The computational algorithm for calculating \underline{s}_1 or \underline{s}_2 is given by equations (55) and (56). A measure of the accuracy of these algorithms can be obtained by analytically

celculeting the solution generated from (55) or (56) under essumed cyclic motion end comparing the rasult to the equivelent solution obtained from the continuous algorithm es described in Saction 6.1.2. For a discrete frequency vibration input, the equation (70) motion can be used analyticelly in equation (55) and (56) to celculete the algorithm solution for S_1 , S_2 (i.s., enelogous to the equation (72) end (73) solution for the continuous (infinitaly fest) elgorithm). After much algebraic manipulation, it can be demonstrated that the algorithm solution for S_1 end S_2 as calculated from equations (55) and (56) under equation (70) input motion, hes zero z, y components, with a z component rete given by:

$$\dot{S}_{2zALG} = 1/2 \theta_z D_y \left(\cos\phi\right) \left(\frac{\sin 2\pi f T_{f}}{2\pi f T_{f}} - \frac{\sin 2\pi f T_{m}}{2\pi f T_{m}}\right)$$
(79)

$$S_{1zALG} = 1/2 \left(\underline{\beta} \times \underline{u} \right)_z + S_{2zALG}$$
 (80)

whera

 S_{1zALG} , S_{2zALG} = Racursive elgorithm solutions for S_{1z} , S_{2z} .

Equations (79) end (80) for the S_1 , S_2 discreta racursiva elgorithm addution is directly enalogous to the equetions (73) end (72) addution to the continuous S_1 , S_2 algorithm. It is easily varified that (79) and (80) reduce to (73) end (72) as T_1 approaches zero.

The error in the S_1 , S_2 elgorithm is measured by the difference between (79), (80) and (73), (72); i.e.,

$$a(\hat{s}_{1z}) = e(\hat{s}_{2z}) = 1/2 e_z D_y (\cos_{\theta}) (1 - \frac{\sin 2\pi f T_{\hat{z}}}{2\pi f T_{\hat{z}}})$$
 (81)

whers

 $e(\hat{s}_{1z})$, $e(\hat{s}_{2z})$ = Error reta in the equetion (55) and (56) algorithm solutions.

Equation (81) cen be used to assess the error in the equetion (55) end (56) algorithms caused by finite iteration rate (i.a., the effect of $T_{\underline{1}}$) under discrets frequency input conditions.

Under rendom vibration input conditions, the equation (55) end (56) algorithms can be enelyeed to obtain the more general solution for S_{1z} , S_{2z} :

$$\dot{s}_{2z} = \int_{0}^{\infty} \left(A_{y}(\omega) \, S_{x}(\omega) \, \cos \left(\phi_{Ay}(\omega) - \phi_{Sx}(\omega) \right) \right)$$

$$- A_{x}(\omega) \, B_{y}(\omega) \, \cos \left(\phi_{Ax}(\omega) - \phi_{Sy}(\omega) \right) \left(\frac{\sin \omega T_{t}}{\omega T_{t}} \right)$$

$$- \frac{\sin \omega T_{m}}{\omega T_{m}} \, P_{nn}(j\omega) \, d\omega$$
(82)

 $S_{1z} = 1/2 (\beta z u)_z + S_{2z}$

The S_{1z} , S_{2z} algorithm error under vibration is the difference between the equation (82) discrete solutions end the equivalent continuous equation (74) with (72) forms:

$$s(\hat{S}_{1z}) = a(\hat{S}_{2z}) = \int_{0}^{\infty} (A_{y}(\omega) B_{x}(\omega) \cos(\phi_{Ay}(\omega) - \phi_{Bx}(\omega))$$

$$- A_{x}(\omega) B_{y}(\omega) \cos(\phi_{Ax}(\omega) - \phi_{By}(\omega))) (1$$

$$- \frac{\sin \omega T_{x}}{\omega T_{x}}) P_{nn}(j\omega) d\omega$$
(83)

Equation (82) end (83) can be used to essass the arror in tha equetion (55) end (56) algorithms caused by finita iteration rate under discrete or random vibration input conditions. The extension of equetion (83) to y, z or z, x exis affacts should be obvious.

CONCLUDING REMARKS

The strapdown computational algorithms and associated design considerations presented in this paper are representative of the algorithms being used in most modern-day strapdown inartial navigation systems. The unique characteristic of the attituda and transformation algorithms presented is the separation of each into a complex low speed and simple high speed computation section. Due to the simplicity of the high speed calculations they can be executed at the high rates necessary to properly account for high frequery but generally low amplituda vibratory effects without posing an insurmountable throughput burdan on the computar. The lower speed calculations which contain the bulk of the computational equations can than be executed at a fairly modest update rate selected to properly account for lower frequency but larger magnitude maneuver induced motion effects. Perhaps the principal advantage of the algorithm forms presented, is their ability to be analyzed for accuracy using straight-forward analytical techniques. This allows the algorithms to be easily tailored and evaluated for given applications as a function of anticipated dynamic environments and user accuracy requirements.

REFERENCES

- 1. Pitman, George R. Jr., ed., Inertial Guidance, John Wiley and Sons, New York, 1962.
- Leondes, Cornelius T., ed., <u>Guidance and Control of Aeroapace Vehicles</u>, <u>McGraw-Hill</u>, 1963.
- Macomber, George R. and Fernandes, Manuel, <u>Inertial Guidance Engineering</u>, Prentice-Hall Englewood Cliffs, New Jersey, 1962.
- Britting, Kenneth R., <u>Inertial Navigation System Analysis</u>, John Wiley and Sons, New York, 1971.
- Morse, Philip M. and Feshback, Herman, <u>Methods of Theoretical Physics</u>, McGraw-Hill, 1953.
- A Study of Critical Computational Problems Associated with Strapdown Inertial Navigation Systems, NASA Report CR-968, April 1968.
- Savage, P.G., "A New Second-Order Solution for Strapped-Down Attitude Computation", AIAA/JACC Guidance & Control Conference, Seattle, Washington, August 15 - 17, 1966.
- Jordan, J.W., "An Accurate Strapdown Direction Cosine Algorithm", NASA TN D-5384, September 1969.
- Mckern, Richard A., A Study of Transformation Algorithms For Use In A Digital Computer, Massachusetts Institute of Technology Master's Thesis, Department of Aeronautics and Astronatics, January 1968.
- Bortz J.E., "A New Mathematical Formulation for Strapdown Inertial Navigation, "IEEE Transactions on Aerospace and Electronic Systems, Volume AES-7, No. 1, January 1971.
- Shepperd, Stanley W., "Quaternion from Rotation Matrix". AIAA Journal of Guidance and Control, Vol. 1, No. 3, May-June 1978.
- 12. Savage, P.G., Introduction To Strapdown Inertial Navigation Systems, June 1, 1983 (Third Printing); Third Strapdown Associates Open Seminar On Strapdown Inertial Navigation Systems, Marquette Inn, Minneapolis, Minnesota, November 14 18, 1983.

APPENDIX A

DERIVATION OF # EQUATION

A differential equation for the rate of change of the ϕ vector can be derived from the equivalent quaternion rate equation. The quaternion h in equations (13) and (14) is the quaternion equivalent to the ϕ rotation angle vector. A differential equation for h can be derived from the incremental equivalent to (13) that describes how h changes over a short time period Δ t (from t½ to t½+1) within the larger time interval from tm to tm+1:

$$p = \begin{pmatrix} g_3 & \alpha \\ g_3 & \alpha \\ g_3 & \alpha \\ g_4 \end{pmatrix}$$

(A2)

$$g_3 = \frac{\sin (\alpha/2)}{\alpha/2}$$

$$g_4 = \cos (\alpha/2)$$

Rotation angle vector associated with the small rotation of the body over the short computer time interval from \$\mathbf{t}\$ to \$\mathbf{t}+\mathbf{l}\$ within the larger interval from \$m\$ to \$m+1\$.

 $\alpha_{\mathbf{x}}, \alpha_{\mathbf{y}}, \alpha_{\mathbf{z}}, \alpha =$ Components and magnitude of $\underline{\alpha}$.

Equation (Al) is equivalently:

$$\frac{h(l+1) - h(l)}{\Delta t} = h(l) \frac{p(l)-1}{\Delta t}$$
(A3)

$$\Delta t = t_{l+1} - t_{l}$$

The basic definition of angular rate states that for small Δt ,

$$\frac{\alpha}{\alpha} \sim \frac{\omega}{\omega} \Delta t$$
 (A4)

Hence, for small Δt , α is small, and therefore, from (A2),

$$g_3 = 1/2$$
 (A5)

Using mixed vector/scalar notation, substitution of (A4) and (A5) in (A2) yields:

$$p = g_3 α + g_4$$
 $= 1/2 ω Δt + 1 - \frac{ω^2 Δt^2}{2}$

Substituting in (A3) obtains:

$$\frac{h(\ell+1) - h(\ell)}{\Delta t} \sim h(\ell) \left(1/2 \underline{\omega} + 1/2 \omega^2 \Lambda t\right)$$

In the limit as $\Delta t \neq 0$, the latter reduce to the derivative form:

$$\dot{h} = 1/2 h \omega \tag{A6}$$

We now return to (14) and express h as a function of ϕ in mixed vector/scaler notation:

$$h = f_3 \phi + f_4$$

$$f_3 = \frac{\sin (\phi/2)}{\phi} \tag{A7}$$

$$f_4 = \cos (\phi/2)$$

Substituting in (A6),

$$\dot{h} = 1/2 \, f_3 \, \phi \, \omega + 1/2 \, f_4 \, \omega$$
 (A8)

It is readily demonstrated by algebraic expansion and using the rules of quaternion algebra that ϕ ω in (A8) is equivalently:

$$\phi \omega = \phi \times \omega - \phi \cdot \omega$$

Differentiation of (A7) shows that:

$$\dot{f}_{3} = \dot{f}_{3} + \dot{f}_{3} + \dot{f}_{4}$$

$$\dot{f}_{3} = 1/2 \quad \frac{\cos \phi/2}{\phi} + \frac{\sin \phi/2}{\phi^{2}} + \frac{\sin \phi/2}{\phi^{2}} + \frac{\dot{\phi}}{\phi}$$

$$= \frac{\dot{\phi}}{\phi} \quad (1/2 \, f_{4} - f_{3})$$

$$\dot{f}_{4} = -1/2 \, (\sin \phi/2) + \frac{\dot{\phi}}{\phi} = -1/2 + \frac{\dot{\phi}}{\phi} + \frac{\dot{\phi}}{\phi}$$

Hence, with (A8),

$$\dot{h} = f_3 \dot{\phi} + \frac{\dot{\phi}}{\phi} (1/2 f_4 - f_3) \dot{\phi} - 1/2 \phi \dot{\phi} f_3$$
$$= 1/2 f_3 (\dot{\phi} \times \underline{\omega}) - 1/2 f_3 \dot{\phi} \cdot \underline{\omega} + 1/2 f_4 \omega$$

Dividing by f3 and solving for :

$$\frac{\dot{\phi}}{\dot{\phi}} = 1/2 \frac{f_4}{f_3} \underline{\omega} + 1/2 \left(\underline{\phi} \times \underline{\omega} \right)$$

$$- \frac{\dot{\phi}}{\dot{\phi}} \left(1/2 \frac{f_4}{f_3} - 1 \right) \underline{\phi} + 1/2 \phi \dot{\phi} - 1/2 \underline{\phi} \cdot \underline{\omega}$$
(A9)

Equation (A9) is now separated into its vector and scalar components:

$$\frac{\dot{\phi}}{\dot{\phi}} = 1/2 \frac{f_4}{f_3} \underline{\omega} + 1/2 \left(\frac{\dot{\phi}}{\dot{\phi}} \times \underline{\omega} \right) - \frac{\dot{\phi}}{\dot{\phi}} \left(1/2 \frac{f_4}{f_3} - 1 \right) \underline{\phi}$$

$$1/2 \phi \dot{\phi} = 1/2 \phi \cdot \omega$$
(A10)

The scalar equation is equivalently:

$$\frac{\phi}{\phi} = \frac{1}{\phi^2} \phi \cdot \omega$$

Substituting in the vector part of (A10) yields:

$$\frac{\dot{\Phi}}{\dot{\Phi}} = 1/2 \frac{f_4}{f_3} \omega + 1/2 \left(\frac{\dot{\Phi}}{\dot{\Phi}} \times \underline{\omega} \right) - \frac{1}{\dot{\Phi}^2} \left(1/2 \frac{f_4}{f_3} - 1 \right) \left(\underline{\Phi} \cdot \underline{\omega} \right) \underline{\Phi}$$

Using the vector triple product rule, it is easily demonstrated that:

$$(\phi \cdot \underline{\omega}) \phi = \phi \times (\phi \times \underline{\omega}) + \phi^2 \underline{\omega}$$

Substituting,

$$\frac{\dot{\bullet}}{\dot{\bullet}} = 1/2 \frac{f_4}{f_3} \underline{\omega} + 1/2 \underline{\phi} \times \underline{\omega} - (1/2 \frac{f_4}{f_3} - 1) \underline{\omega} + \frac{1}{\phi^2} (1 - \frac{f_4}{2f_3}) \underline{\phi} \times (\underline{\phi} \times \underline{\omega})$$

Combining terms:

$$\dot{\phi} = \underline{\omega} + 1/2 \phi \times \underline{\omega} + \frac{1}{\phi^2} (1 - \frac{f_4}{2f_3}) \phi \times (\phi \times \underline{\omega})$$

Using the definition for f_4 and f_3 from (A7), it can be shown by trigonometric manipulation that the bracketed coefficient in the latter expression is equivalently:

$$1 - \frac{f_4}{2f_3} = \frac{1}{\phi^2} \left(1 - \frac{\phi \sin \phi}{2(1 - \cos \phi)}\right)$$

Substitution yields the final expression for $\dot{\underline{\phi}}$:

$$\underline{\dot{\phi}} = \underline{\omega} + 1/2 \underline{\phi} \times \underline{\omega} + \frac{1}{\phi^2} \left(1 - \frac{\phi \sin \phi}{2(1 - \cos \phi)} \right) \underline{\phi} \times (\underline{\phi} \times \underline{\omega})$$
 (A11)

Equation (20) in the main text is the integral from of (All) over a computer cycle (from t_m to t_{m+1}).